

Checking Imaginary Parts of the Compton Amplitudes by The Optical Theorem

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The trick is that imaginary parts in our formulation can only come from s-channel rescattering diagrams.

Onebody and Twobody are real below $W_{\text{rc}} \approx m_a$.

⇒ Use optical theorem to relate to $\sigma_{\text{tot}}^{\text{tot}}(\gamma^3\text{He} \rightarrow X)$ photodissociation.

Fortunately, I just can copy from [Höckel PhD (6.89)]

with minor modifications: There are 2 spin states in ${}^3\text{He}$ to be averaged over, not 3 as in the deuteron.

$$\Rightarrow \sigma_{\text{tot}}^{\text{tot}}[\gamma {}^3\text{He} \rightarrow X] = \frac{1}{4\pi} \frac{1}{4} \sum_{\substack{n_i = n_f = \pm \frac{1}{2} \\ \lambda_i = \lambda_f = \pm 1}} \gamma_m^{(1)} [M_{fi}^{(1)}(\delta=0)]$$

α_{EM} total interpreted [.]

$n_i = n_f = \pm \frac{1}{2}$
 $\lambda_i = \lambda_f = \pm 1$ with total = onebody + twobody

Johannes writes the optical theorem as

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$$\sigma_{\text{tot}}^{\text{tot}}[\gamma {}^3\text{He} \rightarrow X] = -\frac{2\pi \alpha_{\text{EM}}}{16\pi^2} \sum_{\substack{n_i = n_f = \pm \frac{1}{2} \\ \lambda_i = \lambda_f = \pm 1}} \gamma_m^{(2)} [T_{\lambda\lambda}^{fi}(\delta=0)]$$

$$\Rightarrow \boxed{\text{total}_g = -\frac{8\pi}{\sqrt{6!}} T_{\text{Johannes}}^{fi}}$$

as the relation between h_g & Johannes amplitudes.

Units: h_g amplitudes : $[\text{fm}]^4$

Johannes : ?