

Strong magnetic fields and contact interactions in few-fermion systems

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Technical manual detailing the implementation of a variational solution of the non-relativistic few-body problem in an external, *i.e.*, static magnetic field.

a. The symmetric Gauge

$$\mathbf{A}_i = \frac{B_0}{2} (-y_i, x_i, 0) \quad (1)$$

b. The Hamiltonian

$$\hat{H} = -\frac{\hbar^2}{2m} \sum_{i=1}^N \left\{ \nabla_i^2 + i \left(\frac{\hbar^2}{2m} \right) \left(\frac{q_i B_0}{\hbar} \right) L_i^z + \left(\frac{\hbar^2}{2m} \right) \left(\frac{q_i B_0}{\hbar} \right)^2 \frac{1}{4} (x_i^2 + y_i^2) - g_i \left(\frac{\hbar^2}{2m} \right) \left(\frac{q_i B_0}{\hbar} \right) \sigma_{z_i} \right\} \quad (2)$$

$$+ \sum_{i < j}^N [C_a + C_b (\sigma_i^+ \sigma_j^- + \sigma_i^- \sigma_j^+ - \sigma_i^z \sigma_j^z)] e^{-\frac{\Lambda^2}{4} (\mathbf{r}_i - \mathbf{r}_j)^2} + \sum_{\text{cyc. } i < j < k} D \cdot e^{-\frac{\Lambda^2}{4} ((\mathbf{r}_i - \mathbf{r}_j)^2 + (\mathbf{r}_i - \mathbf{r}_k)^2)} \quad (3)$$

c. The variational basis

$$|A, \boldsymbol{\lambda}\rangle := e^{-\frac{1}{2} \mathbf{x}^T A_x \mathbf{x}} e^{-\frac{1}{2} \mathbf{y}^T A_y \mathbf{y}} e^{-\frac{1}{2} \mathbf{z}^T A_z \mathbf{z}} \cdot \sum_{\alpha} \lambda_{\alpha} \sum_{n=1}^{N_{\text{int}}} C_{\alpha}^n |s_1^n, \dots, s_N^n; t_1^n, \dots, t_N^n\rangle \quad (4)$$

input together with $|\lambda\rangle$

d. The generic matrix element

$$I_{\mathcal{O}}(A', \boldsymbol{\lambda}', A, \boldsymbol{\lambda}; P) := \langle A', \boldsymbol{\lambda}' | \hat{\mathcal{O}} \otimes \hat{\mathcal{G}} | \hat{P}(A), \hat{P}(\boldsymbol{\lambda}) \rangle = \langle A' | \hat{\mathcal{O}} | \hat{P}(A) \rangle \cdot \langle \boldsymbol{\lambda}' | \hat{\mathcal{G}} | \hat{P}(\boldsymbol{\lambda}) \rangle \quad (5)$$

with $\hat{P} \in \mathcal{A}$, hence,

$$\hat{P}(A) = T_P^{\dagger} A T_P := A^P \quad (6)$$

$$\hat{\mathcal{O}} \in \left\{ \mathbb{1} ; \mathbf{p}^T \mathbb{1}_{(3N \times 3N)} \mathbf{p} ; \sum_{i=1}^N q_i L_i^z ; \sum_{i=1}^N q_i (x_i^2 + y_i^2 + z_i^2) ; \sum_{i=1}^N q_i \sigma_i^z ; \sum_{i < j}^N e^{-\frac{\Lambda^2}{4} (\mathbf{r}_i - \mathbf{r}_j)^2} \right\} \quad (7)$$

e. The matrix elements

$\hat{\odot}$	$\langle A' \mid \hat{\odot} \mid \hat{P}(A) \rangle$	$\langle \lambda' \mid \hat{\odot} \mid \hat{P}(\lambda) \rangle$
$\mathbb{1} := \mathbb{1}_{\mathbf{r}}^P \cdot \mathbb{1}_{\mathbf{s}}^P$	$\left(\frac{(2\pi)^{3N}}{\det \mathbb{A}_x \det \mathbb{A}_y \det \mathbb{A}_z} \right)^{\frac{1}{2}}$	$\sum_{\alpha, n}^{ \lambda , N_{\text{int}}} \lambda_{\alpha} C_{\alpha}^n \langle \mathbf{s}^n; \mathbf{t}^n \mid \hat{P}(\mathbf{s}^n); \hat{P}(\mathbf{t}^n) \rangle$
$\frac{1}{2} \mathbf{p}^T \mathbb{1}_{3N} \mathbf{p} = -\frac{\hbar^2}{2} \nabla^T \mathbb{1}_{3N} \nabla$	$\frac{\hbar^2}{2} \mathbb{1}_{\mathbf{r}}^P \prod_{c=x,y,z} (A_c)_{im} (\mathbb{A}_c^{-1})_{mn} (A_c^P)_{ni}$	$\mathbb{1}_{\mathbf{s}}^P$
$\sum_{i=1}^N q_i L_i^z = q_i (x_i \partial_{y_i} - y_i \partial_{x_i})$	0	$\mathbb{Q}_{\mathbf{s}}^P := \sum_{i=1}^N \sum_{\alpha, n}^{ \lambda , N_{\text{int}}} \underbrace{(\hat{P}[\mathbf{t}^n])_i}_{=q_{P(i)}} \lambda_{\alpha} C_{\alpha}^n \langle \mathbf{s}^n; \mathbf{t}^n \mid \hat{P}(\mathbf{s}^n); \hat{P}(\mathbf{t}^n) \rangle$
$\sum_{i=1}^N q_i (\omega_x x_i^2 + \omega_y y_i^2 + \omega_z z_i^2)$	$\mathbb{1}_{\mathbf{r}}^P \prod_{c=x,y,z} \omega_c \sum_{i=1}^N (\mathbb{A}_c^{-1})_{ii}$	$\mathbb{Q}_{\mathbf{s}}^P$
$\sum_{i=1}^N q_i \sigma_i^z$	$\mathbb{1}_{\mathbf{r}}^P$	$\sum_{i=1}^N \sum_{\alpha, n}^{ \lambda , N_{\text{int}}} \underbrace{(\hat{P}[\mathbf{t}^n])_i}_{=q_{P(i)}} \underbrace{(\hat{P}[\mathbf{s}^n])_i}_{=\underbrace{s_{P(i)}}_z} \lambda_{\alpha} C_{\alpha}^n \langle \mathbf{s}^n; \mathbf{t}^n \mid \hat{P}(\mathbf{s}^n); \hat{P}(\mathbf{t}^n) \rangle$
$\sum_{i < j}^N e^{-\frac{\Lambda^2}{4} (\mathbf{r}_i - \mathbf{r}_j)^2}$		
$\sum_{\text{cyc.}} \sum_{i < j < k} e^{-\frac{\Lambda^2}{4} ((\mathbf{r}_i - \mathbf{r}_j)^2 + (\mathbf{r}_i - \mathbf{r}_k)^2)}$		

with

$$\mathbb{A}_x = A'_x + A_x^P \quad (9)$$