

1 Hamiltonian Formulation

The Hamiltonian $H(\mathbf{R}_\theta)$ is a 4x4 arrowhead matrix defined as:

$$H(\mathbf{R}_\theta) = \begin{pmatrix} E_0(R_{\theta,i}) & c_{10} & c_{20} & c_{30} \\ c_{10} & E_1(R_{\theta,i}) & 0 & 0 \\ c_{20} & 0 & E_2(R_{\theta,i}) & 0 \\ c_{30} & 0 & 0 & E_3(R_{\theta,i}) \end{pmatrix} \quad (1)$$

where $E_0(R_{\theta,i}) = \hbar\omega + \sum_{i=0}^2 V_x^{(i)}(R_{\theta,i})$ and $E_{i+1}(R_{\theta,i}) = E_0(R_{\theta,i}) + V_a^{(i)}(R_{\theta,i}) - V_x^{(i)}(R_{\theta,i})$ for $i = 0, 1, 2$.

Here ω is a frequency parameter, and V_x and V_a are potential terms that depend on the parameter vector $R(\theta)$. These potential terms are defined as follows:

$$V_x^{(i)}(R_{\theta,i}) = a \cdot (R_{\theta,i})^2 + c \quad (2)$$

$$V_a^{(i)}(R_{\theta,i}) = a \cdot (R_{\theta,i} - x_{\text{shift}})^2 + c \quad (3)$$

The coupling constants c_{10} , c_{20} , and c_{30} are calculated using the transitional dipole moment between the ground state and excited states. The transitional dipole moment is computed as:

$$c_{i0} = c_{0i} = \langle \psi_i | \hat{\mathbf{r}} | \psi_0 \rangle \quad (4)$$

where ψ_0 is the ground state eigenvector, ψ_i is the i^{th} excited state eigenvector, and $\hat{\mathbf{r}}$ is the position operator. These couplings are then used in the Hamiltonian matrix:

$$H(\mathbf{R}_\theta) = \begin{pmatrix} E_0(R_{\theta,i}) & c_{10} & c_{20} & c_{30} \\ c_{10} & E_1(R_{\theta,i}) & 0 & 0 \\ c_{20} & 0 & E_2(R_{\theta,i}) & 0 \\ c_{30} & 0 & 0 & E_3(R_{\theta,i}) \end{pmatrix} \quad (5)$$

2 \mathbf{R}_θ generation

The \mathbf{R}_θ vector traces a perfect circle orthogonal to the $x = y = z$ line using the `create_perfect_orthogonal_vectors` function from the Arrowhead/generalized package.

3 Berry Connection

The Berry connection $A(\mathbf{R}_\theta)$ is calculated using:

$$A_n(\mathbf{R}_\theta) = \langle n(\mathbf{R}_\theta) | i \partial_{\mathbf{R}_\theta} | n(\mathbf{R}_\theta) \rangle \quad (6)$$

where $|n(\mathbf{R}_\theta)\rangle$ are the eigenstates of $H(\mathbf{R}_\theta)$.

4 Berry Phase

The Berry phase γ_n for state n is obtained by integrating the Berry connection:

$$\gamma_n = \int_0^{2\pi} A_n(\mathbf{R}_\theta) d\mathbf{R}_\theta \quad (7)$$

5 Verification

We verify the eigenvalue equation $H(\mathbf{R}_\theta)|n(\mathbf{R}_\theta)\rangle = E_n(\mathbf{R}_\theta)|n(\mathbf{R}_\theta)\rangle$ by comparing:

$$H(\mathbf{R}_\theta)|n(\mathbf{R}_\theta)\rangle \quad \text{vs} \quad E_n(\mathbf{R}_\theta)|n(\mathbf{R}_\theta)\rangle \quad (8)$$

6 Visualization

For each state n , we plot:

- Magnitude of $H(\mathbf{R}_\theta)|n(\mathbf{R}_\theta)\rangle$ and $E_n(\mathbf{R}_\theta)|n(\mathbf{R}_\theta)\rangle$
- Real and imaginary components separately
- All four vector components

7 References

- M. V. Berry, *Quantal phase factors accompanying adiabatic changes*, Proc. R. Soc. Lond. A 392, 45-57 (1984)
- D. J. Thouless, *Topological Quantum Numbers in Nonrelativistic Physics*, World Scientific (1998)
- B. Simon, *Holonomy, the Quantum Adiabatic Theorem, and Berry's Phase*, Phys. Rev. Lett. 51, 2167 (1983)
- J. E. Avron, R. Seiler, L. G. Yaffe, *Adiabatic Theorems and Applications to the Quantum Hall Effect*, Commun. Math. Phys. 110, 33-49 (1987)