

# Enlarged unit cell, decomposition of $\kappa$ , and Bloch Hamiltonian

## 1 Enlarged unit cell, decomposition of $\kappa$ , and Bloch Hamiltonian

### 1.1 Lattice and supercell

We work on the honeycomb in brick-wall coordinates with original integer coordinates  $(i, j) \in \mathbb{Z}^2$ . We enlarge to a  $2 \times 2$  supercell with

$$\mathbf{a} = 2\hat{i}, \quad \mathbf{b} = 2\hat{j},$$

and four internal offsets

$$r_0 = (0, 0), \quad r_1 = (1, 0), \quad r_2 = (0, 1), \quad r_3 = (1, 1).$$

A supercell position is  $\mathbf{R} = m\mathbf{a} + n\mathbf{b}$ . Sites are labeled  $(\mathbf{R}, r, s)$  with  $r \in \{0, 1, 2, 3\}$  and sublattice  $s \in \{A, B\}$ . We use the Bloch basis

$$\Psi_{\mathbf{K}} = (A_0, A_1, A_2, A_3; B_0, B_1, B_2, B_3)^T,$$

and define

$$\mathbf{K} \cdot \mathbf{a} = 2K_i, \quad \mathbf{K} \cdot \mathbf{b} = 2K_j, \quad \mathbf{K} \cdot (\mathbf{a} - \mathbf{b}) = 2(K_i - K_j).$$

### 1.2 Nearest-neighbor (Kitaev) map in real space

Using the convention

$$J_x : B_{i,j} \rightarrow A_{i,j}, \quad J_y : B_{i,j} \rightarrow A_{i+1,j}, \quad J_z : B_{i,j} \rightarrow A_{i,j+1},$$

the Majorana hopping Hamiltonian (itinerant  $c^z \equiv c$ ; flux-free gauge) is

$$H_J = i \sum_{\mathbf{R}} \left[ J_x \sum_{r=0}^3 c_{B,\mathbf{R},r} c_{A,\mathbf{R},r} + J_y (c_{B,\mathbf{R},r_0} c_{A,\mathbf{R},r_1} + c_{B,\mathbf{R},r_2} c_{A,\mathbf{R},r_3} + c_{B,\mathbf{R},r_1} c_{A,\mathbf{R}+\mathbf{a},r_0} + c_{B,\mathbf{R},r_3} c_{A,\mathbf{R}+\mathbf{a},r_2}) + J_z (c_{B,\mathbf{R},r_0} c_{A,\mathbf{R},r_2} + c_{B,\mathbf{R},r_1} c_{A,\mathbf{R},r_3} + c_{B,\mathbf{R},r_2} c_{A,\mathbf{R}+\mathbf{b},r_0} + c_{B,\mathbf{R},r_3} c_{A,\mathbf{R}+\mathbf{b},r_1}) \right].$$

### 1.3 Twelve vertex-disjoint $\kappa$ subsets in real space

We decompose the chiral term into 12 commuting blocks indexed by sublattice class  $S \in \{A, B\}$  (A–A or B–B), next-nearest-neighbor direction  $d \in \{+i, -j, j-i\}$ , and parity  $p \in \{\text{even, odd}\}$ . Each block is a *matching* (no shared sites), so  $\exp(-i\theta H_{\kappa}^{S,d,p})$  is exact per block. With a fixed orientation  $J \rightarrow K = J + d$ , the A–A blocks carry signs  $(+, -, +)$  for  $d = (+i, -j, j-i)$ , while B–B blocks carry the opposite  $(-, +, -)$ .

For the  $2 \times 2$  supercell:

$$\begin{aligned} H_{\kappa}^{A,+i,\text{even}} &= i \sum_{\mathbf{R}} \kappa_{A,+i,\text{e}} (c_{A,\mathbf{R},r_0} c_{A,\mathbf{R},r_1} + c_{A,\mathbf{R},r_2} c_{A,\mathbf{R},r_3}), \\ H_{\kappa}^{A,+i,\text{odd}} &= -i \sum_{\mathbf{R}} \kappa_{A,+i,\text{o}} (c_{A,\mathbf{R},r_0} c_{A,\mathbf{R}-\mathbf{a},r_1} + c_{A,\mathbf{R},r_2} c_{A,\mathbf{R}-\mathbf{a},r_3}), \\ H_{\kappa}^{A,-j,\text{even}} &= -i \sum_{\mathbf{R}} \kappa_{A,-j,\text{e}} (c_{A,\mathbf{R},r_0} c_{A,\mathbf{R},r_2} + c_{A,\mathbf{R},r_1} c_{A,\mathbf{R},r_3}), \\ H_{\kappa}^{A,-j,\text{odd}} &= +i \sum_{\mathbf{R}} \kappa_{A,-j,\text{o}} (c_{A,\mathbf{R},r_0} c_{A,\mathbf{R}-\mathbf{b},r_2} + c_{A,\mathbf{R},r_1} c_{A,\mathbf{R}-\mathbf{b},r_3}), \\ H_{\kappa}^{A,j-i,\text{even}} &= i \sum_{\mathbf{R}} \kappa_{A,j-i,\text{e}} (c_{A,\mathbf{R},r_1} c_{A,\mathbf{R},r_2} + c_{A,\mathbf{R},r_3} c_{A,\mathbf{R}+\mathbf{b},r_0}), \\ H_{\kappa}^{A,j-i,\text{odd}} &= -i \sum_{\mathbf{R}} \kappa_{A,j-i,\text{o}} (c_{A,\mathbf{R},r_1} c_{A,\mathbf{R}+\mathbf{a}-\mathbf{b},r_2} + c_{A,\mathbf{R},r_3} c_{A,\mathbf{R}+\mathbf{a},r_0}), \end{aligned}$$

and analogously for B–B with the opposite signs:

$$\begin{aligned} H_{\kappa}^{B,+i,\text{even}} &= -i \sum_{\mathbf{R}} \kappa_{B,+i,\text{e}} (c_{B,\mathbf{R},r_0} c_{B,\mathbf{R},r_1} + c_{B,\mathbf{R},r_2} c_{B,\mathbf{R},r_3}), \\ H_{\kappa}^{B,+i,\text{odd}} &= +i \sum_{\mathbf{R}} \kappa_{B,+i,\text{o}} (c_{B,\mathbf{R},r_0} c_{B,\mathbf{R}-\mathbf{a},r_1} + c_{B,\mathbf{R},r_2} c_{B,\mathbf{R}-\mathbf{a},r_3}), \\ H_{\kappa}^{B,-j,\text{even}} &= +i \sum_{\mathbf{R}} \kappa_{B,-j,\text{e}} (c_{B,\mathbf{R},r_0} c_{B,\mathbf{R},r_2} + c_{B,\mathbf{R},r_1} c_{B,\mathbf{R},r_3}), \\ H_{\kappa}^{B,-j,\text{odd}} &= -i \sum_{\mathbf{R}} \kappa_{B,-j,\text{o}} (c_{B,\mathbf{R},r_0} c_{B,\mathbf{R}-\mathbf{b},r_2} + c_{B,\mathbf{R},r_1} c_{B,\mathbf{R}-\mathbf{b},r_3}), \\ H_{\kappa}^{B,j-i,\text{even}} &= -i \sum_{\mathbf{R}} \kappa_{B,j-i,\text{e}} (c_{B,\mathbf{R},r_1} c_{B,\mathbf{R},r_2} + c_{B,\mathbf{R},r_3} c_{B,\mathbf{R}+\mathbf{b},r_0}), \\ H_{\kappa}^{B,j-i,\text{odd}} &= +i \sum_{\mathbf{R}} \kappa_{B,j-i,\text{o}} (c_{B,\mathbf{R},r_1} c_{B,\mathbf{R}+\mathbf{a}-\mathbf{b},r_2} + c_{B,\mathbf{R},r_3} c_{B,\mathbf{R}+\mathbf{a},r_0}). \end{aligned}$$

Summing all 12 blocks yields  $H_{\kappa} = \sum_{S,d,p} H_{\kappa}^{S,d,p}$ .

## 1.4 Bloch Hamiltonian $H(\mathbf{K})$ as an explicit $8 \times 8$ matrix

In the Bloch basis, the total Hamiltonian is

$$H(\mathbf{K}) = \begin{pmatrix} H_{AA} & H_{AB} \\ H_{AB}^\dagger & H_{BB} \end{pmatrix}, \quad H_{ij}(\mathbf{K}) = H_{ji}(\mathbf{K})^*.$$

**A–A block (kappa, ABA).** Nonzero entries:

$$\begin{aligned} H_{AA}[0, 1] &= i(\kappa_{A,+i,e} - \kappa_{A,+i,o} e^{-2iK_i}), & H_{AA}[2, 3] &= i(\kappa_{A,+i,e} - \kappa_{A,+i,o} e^{-2iK_i}), \\ H_{AA}[0, 2] &= i(-\kappa_{A,-j,e} + \kappa_{A,-j,o} e^{-2iK_j}), & H_{AA}[1, 3] &= i(-\kappa_{A,-j,e} + \kappa_{A,-j,o} e^{-2iK_j}), \\ H_{AA}[1, 2] &= i(\kappa_{A,j-i,e} - \kappa_{A,j-i,o} e^{+2i(K_i-K_j)}), \\ H_{AA}[3, 0] &= i(\kappa_{A,j-i,e} e^{+2iK_j} - \kappa_{A,j-i,o} e^{+2iK_i}), \end{aligned}$$

with  $H_{AA}[j, i] = H_{AA}[i, j]^*$  and all diagonal elements zero.

**B–B block (kappa, BAB).** Nonzero entries:

$$\begin{aligned} H_{BB}[0, 1] &= i(-\kappa_{B,+i,e} + \kappa_{B,+i,o} e^{-2iK_i}), & H_{BB}[2, 3] &= i(-\kappa_{B,+i,e} + \kappa_{B,+i,o} e^{-2iK_i}), \\ H_{BB}[0, 2] &= i(\kappa_{B,-j,e} - \kappa_{B,-j,o} e^{-2iK_j}), & H_{BB}[1, 3] &= i(\kappa_{B,-j,e} - \kappa_{B,-j,o} e^{-2iK_j}), \\ H_{BB}[1, 2] &= i(-\kappa_{B,j-i,e} + \kappa_{B,j-i,o} e^{+2i(K_i-K_j)}), \\ H_{BB}[3, 0] &= i(-\kappa_{B,j-i,e} e^{+2iK_j} + \kappa_{B,j-i,o} e^{+2iK_i}), \end{aligned}$$

with  $H_{BB}[j, i] = H_{BB}[i, j]^*$  and zero diagonals.

**A–B block (nearest-neighbor  $J$ ).** Nonzero entries  $H_{AB}[\text{A}, \text{B}]$ :

$$\begin{aligned} J_x : \quad H_{AB}[0, 0] &= iJ_x, & H_{AB}[1, 1] &= iJ_x, & H_{AB}[2, 2] &= iJ_x, & H_{AB}[3, 3] &= iJ_x, \\ J_y : \quad H_{AB}[1, 0] &= iJ_y, & H_{AB}[0, 1] &= iJ_y e^{+2iK_i}, & H_{AB}[3, 2] &= iJ_y, & H_{AB}[2, 3] &= iJ_y e^{+2iK_i}, \\ J_z : \quad H_{AB}[2, 0] &= iJ_z, & H_{AB}[0, 2] &= iJ_z e^{+2iK_j}, & H_{AB}[3, 1] &= iJ_z, & H_{AB}[1, 3] &= iJ_z e^{+2iK_j}. \end{aligned}$$

The lower-left block is  $H_{BA} = H_{AB}^\dagger$ .

## 1.5 Remarks

- Each of the 12  $\kappa$ -blocks is a vertex-disjoint matching  $\Rightarrow$  exact exponentiation per subset (no intra-block Trotter error).

- Summing even + odd for fixed  $(S, d)$  produces the expected sine structure factor along  $d$ ; summing the three directions reproduces the usual real  $\Delta(\mathbf{k})$  sitting with opposite signs on the A/B diagonals.
- Hermiticity is enforced everywhere by  $H_{ij}(\mathbf{K}) = H_{ji}(\mathbf{K})^*$ .