

HW 3 solutions

Pb. 4.8

1 Problem 4.8

1.1 Decomposition of $y(t)$ into Fourier series

$y(t)$ is the following function of time:

$$\begin{aligned} y(t) &= \frac{Y}{\tau}t \\ y(t + \tau) &= y(t) \end{aligned} \quad (1)$$

Because y is periodic with a period τ ,

$$y(t) = \frac{a_0}{2} + \sum_{j=1}^{\infty} a_j \cos(j\omega_0 t) + \sum_{j=1}^{\infty} b_j \sin(j\omega_0 t) \quad (2)$$

where $\omega_0 = \frac{2\pi}{\tau}$. The Fourier series coefficient of $y(t)$, a_j ($j = 0, \dots, +\infty$) and b_j ($j = 1, \dots, +\infty$), are given by:

$$\begin{aligned} a_j &= \frac{2}{\tau} \int_0^{\tau} y(t) \cos(j\omega_0 t) dt \\ &= \frac{2}{\tau} \int_0^{\tau} \frac{Y}{\tau} t \cos(j\omega_0 t) dt \\ &= \frac{2Y}{\tau^2} \int_0^{\tau} t \cos(j\omega_0 t) dt \end{aligned} \quad (3)$$

For $j = 0$:

$$\begin{aligned} a_0 &= \frac{2Y}{\tau^2} \int_0^{\tau} t dt \\ &= \frac{2Y}{\tau^2} \frac{\tau^2}{2} \\ &= Y \end{aligned} \quad (4)$$

For $j \neq 0$, integration by part (let $u = t$, $v' = \cos(j\omega_0 t)$)

$$\begin{aligned} a_j &= \frac{2Y}{\tau^2} \left[\frac{t}{j\omega_0} \sin(j\omega_0 t) \Big|_0^{\tau} - \int_0^{\tau} \frac{1}{j\omega_0} \sin(j\omega_0 t) dt \right] \\ &= \frac{2Y}{\tau^2} \left[0 + \left(\frac{1}{j\omega_0} \right)^2 \cos(j\omega_0 t) \Big|_0^{\tau} \right] \\ &= \frac{2Y}{\tau^2} \left(\frac{\tau}{j2\pi} \right)^2 [\cos(j2\pi) - \cos 0] \\ &= 0 \end{aligned} \quad (5)$$

$$\begin{aligned} b_j &= \frac{2}{\tau} \int_0^{\tau} y(t) \sin(j\omega_0 t) dt \\ &= \frac{2}{\tau} \int_0^{\tau} \frac{Y}{\tau} t \sin(j\omega_0 t) dt \\ &= \frac{2Y}{\tau^2} \int_0^{\tau} t \sin(j\omega_0 t) dt \end{aligned} \quad (6)$$

Integration by part (let $u = t$, $v' = \sin(j\omega_0 t)$)

$$\begin{aligned}
 b_j &= \frac{2Y}{\tau^2} \left[-\frac{t}{j\omega_0} \cos(j\omega_0 t) \right]_0^\tau + \int_0^\tau \frac{1}{j\omega_0} \cos(j\omega_0 t) dt \\
 &= \frac{2Y}{\tau^2} \left[-\frac{\tau}{j\omega_0} \cos(2\pi j) + \left(\frac{1}{j\omega_0} \right)^2 \sin(j\omega_0 t) \right]_0^\tau \\
 &= \frac{2Y}{\tau^2} \left[-\frac{\tau^2}{j2\pi} + 0 \right] \\
 &= -\frac{Y}{j\pi}
 \end{aligned} \tag{7}$$

such that:

$$y(t) = \frac{Y}{2} - \sum_{j=1}^{\infty} \frac{Y}{j\pi} \sin(j\omega_0 t) \tag{8}$$

which can be written as a complex Fourier series:

$$y(t) = \frac{Y}{2} + \sum_{j=1}^{\infty} \frac{iY}{2j\pi} e^{ij\omega_0 t} + \sum_{j=-\infty}^{-1} \frac{-iY}{2j\pi} e^{-ij\omega_0 t} \tag{9}$$

1.2 Equation of motion

Applying Newton's 2nd law:

$$m\ddot{x} + c\dot{x} + kx = c\dot{y} + ky \tag{10}$$

1.3 Solution

The steady-state solution is written as a complex Fourier series:

$$x(t) = \sum_{j=-\infty}^{\infty} X_j e^{ij\omega_0 t} \tag{11}$$

Value of X_0 :

$$kX_0 = k\frac{Y}{2} \tag{12}$$

such that $X_0 = \frac{Y}{2}$.

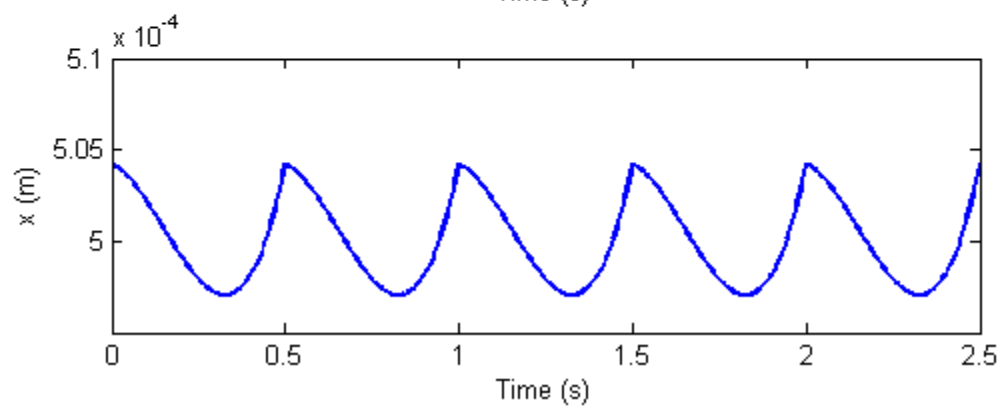
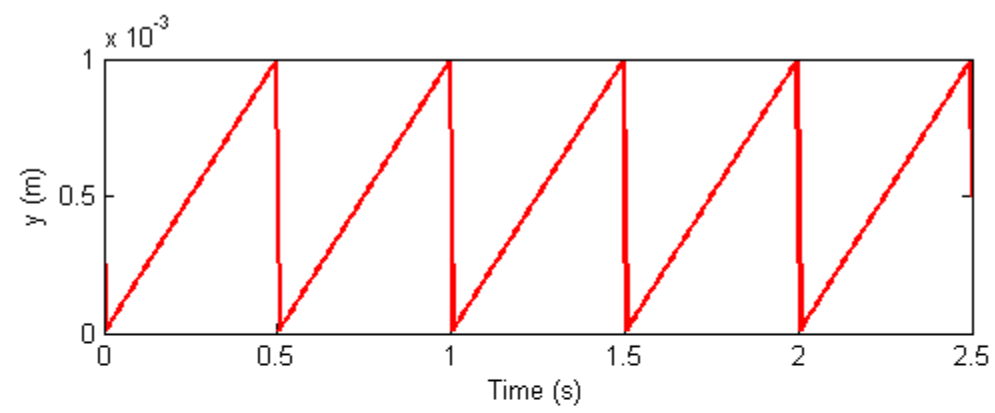
Value of X_j (for $j > 0$)

$$\left[-m(j\omega_0)^2 + Cij\omega_0 + k \right] X_j = \frac{-iY}{2j\pi} [cij\omega_0 + k] \tag{13}$$

$$X_j = \frac{-iY}{2j\pi} \frac{cij\omega_0 + k}{k - m(j\omega_0)^2 + Cij\omega_0} \tag{14}$$

$$X_j = \frac{-iY}{2j\pi} \frac{2i\xi \frac{j\omega_0}{\omega_n} + 1}{1 - \left(\frac{j\omega_0}{\omega_n} \right)^2 + 2i\xi \frac{j\omega_0}{\omega_n}} \tag{15}$$

And for $j < 0$, $X_j = X_{-j}^*$



- 4.36 (a) Unit impulse response function for undamped case:
Use $\gamma = 0$ and $\omega_d = \omega_n$ in Eq. (4.25):

$$x(t) = \frac{1}{m \omega_n} \sin \omega_n t \quad (E.1)$$

- (b) Unit impulse response function for underdamped case: Eq. (4.25):

$$x(t) = \frac{1}{m \omega_d} e^{-\gamma \omega_n t} \sin \omega_d t \quad (E.2)$$

- (c) Unit impulse response function for critically damped case:

Free vibration response of a critically damped system is given by Eq. (2.80):

$$x(t) = \{x_0 + (\dot{x}_0 + \omega_n x_0) t\} e^{-\omega_n t} \quad (E.3)$$

Using the initial conditions $x_0 = 0$ and $\dot{x}_0 = \frac{1}{m}$,

Eq. (E.3) gives

$$x(t) = \frac{t}{m} e^{-\omega_n t} \quad (E.4)$$

- (d) Unit impulse response function for an overdamped case:

Free vibration response of an overdamped system is given by Eq. (2.81):

$$x(t) = C_1 e^{(-\gamma + \sqrt{\gamma^2 - 1}) \omega_n t} + C_2 e^{(-\gamma - \sqrt{\gamma^2 - 1}) \omega_n t} \quad (E.5)$$

where

$$C_1 = \frac{x_0 \omega_n (\gamma + \sqrt{\gamma^2 - 1}) \dot{x}_0}{2 \omega_n \sqrt{\gamma^2 - 1}}; \quad C_2 = \frac{-x_0 \omega_n (\gamma - \sqrt{\gamma^2 - 1}) \dot{x}_0}{2 \omega_n \sqrt{\gamma^2 - 1}}$$

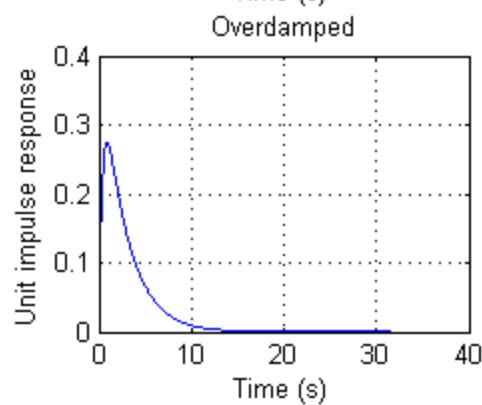
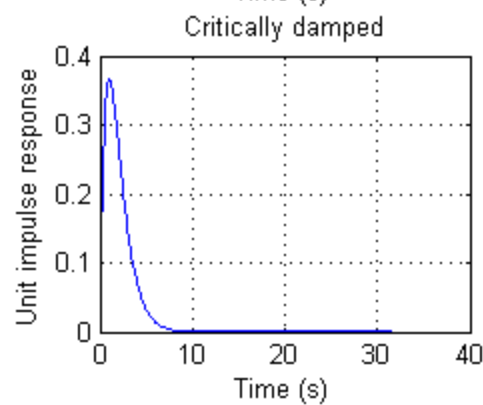
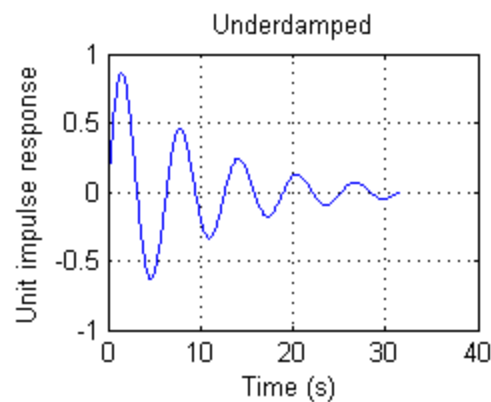
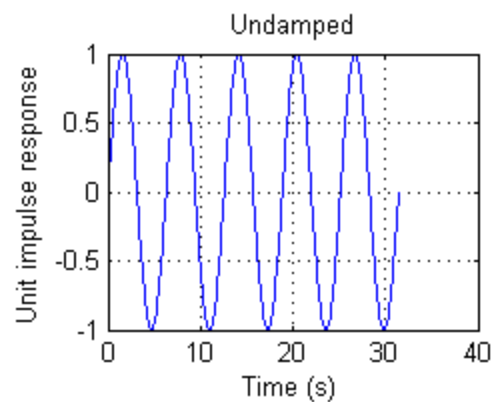
For the initial conditions $x_0 = 0$ and $\dot{x}_0 = \frac{1}{m}$,

C_1 and C_2 become

$$C_1 = \frac{1}{2m\omega_n\sqrt{\zeta^2-1}} \quad ; \quad C_2 = -\frac{1}{2m\omega_n\sqrt{\zeta^2-1}}$$

and hence Eq. (E.5) yields

$$x(t) = \frac{-1}{2m\omega_n\sqrt{\zeta^2-1}} \left\{ e^{-(\zeta+\sqrt{\zeta^2-1})\omega_n t} - e^{-(\zeta-\sqrt{\zeta^2-1})\omega_n t} \right\} \quad (E.6)$$



Pb. 4.17

4.17 From Eq. (4.31), $x(t) = \frac{1}{m\omega_d} \int_0^t F_0 e^{-\alpha\tau} e^{-\gamma\omega_n(t-\tau)} \sin \omega_d(t-\tau) d\tau$

$$= \frac{F_0}{m\omega_d} e^{-\gamma\omega_n t} \int_0^t e^{-(\alpha-\gamma\omega_n)\tau} \sin \omega_d(t-\tau) d\tau$$

$$x(t) = \frac{F_0 e^{-\gamma\omega_n t}}{m\omega_d} \int_0^t e^{-(\alpha-\gamma\omega_n)\tau} \{ \sin \omega_d t \cdot \cos \omega_d \tau - \cos \omega_d t \cdot \sin \omega_d \tau \} d\tau$$

$$= \frac{F_0 e^{-\gamma\omega_n t}}{m\omega_d} \sin \omega_d t \left[\frac{e^{-(\alpha-\gamma\omega_n)\tau}}{(\alpha-\gamma\omega_n)^2 + \omega_d^2} \{ -(\alpha-\gamma\omega_n) \cos \omega_d \tau + \omega_d \sin \omega_d \tau \} \right]_0^t$$

$$- \frac{F_0 e^{-\gamma\omega_n t}}{m\omega_d} \cos \omega_d t \left[\frac{e^{-(\alpha-\gamma\omega_n)\tau}}{(\alpha-\gamma\omega_n)^2 + \omega_d^2} \{ -(\alpha-\gamma\omega_n) \sin \omega_d \tau - \omega_d \cos \omega_d \tau \} \right]_0^t$$

$$= \frac{F_0 e^{-\gamma\omega_n t}}{m\omega_d [(\alpha-\gamma\omega_n)^2 + \omega_d^2]} \left\{ \omega_d e^{-(\alpha-\gamma\omega_n)t} + (\alpha-\gamma\omega_n) \sin \omega_d t - \omega_d \cos \omega_d t \right\}$$

$$= \frac{F_0 e^{-\alpha t}}{m [(\alpha-\gamma\omega_n)^2 + \omega_d^2]} + \frac{F_0 e^{-\gamma\omega_n t}}{m\omega_d [(\alpha-\gamma\omega_n)^2 + \omega_d^2]} \sin(\omega_d t - \phi)$$

where $\phi = \tan^{-1} \left(\frac{\omega_d}{\alpha - \gamma\omega_n} \right)$.
