

## Problem 5.17

Equations of motion

$$\begin{aligned} m_1 \ddot{x}_1 + (k_1 + k_2)x_1 - k_2 x_2 &= 0 \\ m_2 \ddot{x}_2 + k_2 x_2 - k_2 x_1 &= 0 \end{aligned} \quad (E_1)$$

With  $x_i(t) = X_i \cos(\omega t + \phi)$ ;  $i = 1, 2$ , Eqs. (E<sub>1</sub>) give the frequency equation

$$\begin{vmatrix} -\omega^2 m_1 + k_1 + k_2 & -k_2 \\ -k_2 & -\omega^2 m_2 + k_2 \end{vmatrix} = 0$$

or  $\omega^4 - \left( \frac{k_1 + k_2}{m_1} + \frac{k_2}{m_2} \right) \omega^2 + \frac{k_1 k_2}{m_1 m_2} = 0 \quad (E_2)$

Roots of Eq. (E<sub>2</sub>) are

$$\omega_1^2, \omega_2^2 = \frac{k_1 + k_2}{2m_1} + \frac{k_2}{2m_2} \mp \sqrt{\frac{1}{4} \left( \frac{k_1 + k_2}{m_1} + \frac{k_2}{m_2} \right)^2 - \frac{k_1 k_2}{m_1 m_2}} \quad (E_3)$$

If  $\vec{x}^{(1)} = \begin{Bmatrix} x_1^{(1)} \\ x_2^{(1)} = r_1 x_1^{(1)} \end{Bmatrix}$  and  $\vec{x}^{(2)} = \begin{Bmatrix} x_1^{(2)} \\ x_2^{(2)} = r_2 x_1^{(2)} \end{Bmatrix}$ ,

$$r_1 = \frac{x_2^{(1)}}{x_1^{(1)}} = \frac{-m_1 \omega_1^2 + k_1 + k_2}{k_2} = \frac{k_2}{-m_2 \omega_1^2 + k_2} \quad (E_4)$$

$$r_2 = \frac{x_2^{(2)}}{x_1^{(2)}} = \frac{-m_1 \omega_2^2 + k_1 + k_2}{k_2} = \frac{k_2}{-m_2 \omega_2^2 + k_2} \quad (E_5)$$

General solution of (E<sub>1</sub>) is

$$x_1(t) = X_1^{(1)} \cos(\omega_1 t + \phi_1) + X_1^{(2)} \cos(\omega_2 t + \phi_2) \quad (E_6)$$

$$x_2(t) = r_1 X_1^{(1)} \cos(\omega_1 t + \phi_1) + r_2 X_1^{(2)} \cos(\omega_2 t + \phi_2)$$

where  $X_1^{(1)}$ ,  $X_1^{(2)}$ ,  $\phi_1$  and  $\phi_2$  can be found using Eqs. (5.18).

For  $m_1 = m$ ,  $m_2 = 2m$ ,  $k_1 = k$  and  $k_2 = 2k$ , (E<sub>3</sub>) gives

$$\omega_1^2 = (2 - \sqrt{3}) \frac{k}{m}, \quad \omega_2^2 = (2 + \sqrt{3}) \frac{k}{m} \quad (E_7)$$

$$\omega_1^2, \omega_2^2 = \frac{8000}{2} + \frac{6000}{2} \mp \sqrt{\frac{1}{4} \left( \frac{8000}{1} + \frac{6000}{1} \right)^2 - \frac{12 \times 10^6}{1}} = 917.2, 13082.8$$

$$\omega_1 = 30.2853 \text{ rad/sec}, \quad \omega_2 = 114.3801 \text{ rad/s}$$

Eqs. (E<sub>4</sub>) and (E<sub>5</sub>) of solution of problem 5.5 give

$$r_1 = \frac{k_2}{-m_2 \omega_1^2 + k_2} = \frac{6000}{-917.2 + 6000} = 1.1805; \quad \vec{X}^{(1)} = \begin{Bmatrix} 1 \\ 1.1805 \end{Bmatrix} X_1^{(1)}$$

$$r_2 = \frac{k_2}{-m_2 \omega_2^2 + k_2} = \frac{6000}{-13082.8 + 6000} = -0.8471; \quad \vec{X}^{(2)} = \begin{Bmatrix} 1 \\ -0.8471 \end{Bmatrix} X_1^{(2)}$$

Problem 3: 5.20

(5.20) (a)  $\omega_1^2 = 917.2, \quad \omega_2^2 = 13082.8, \quad r_1 = 1.1805, \quad r_2 = -0.8471$   
 $x_1(0) = 0.2, \quad x_2(0) = \dot{x}_1(0) = \dot{x}_2(0) = 0$

Eg. (5.18) gives

$$X_1^{(1)} = \frac{1}{(-0.8471 - 1.1805)} [(-0.8471)(0.2)] = 0.08356$$

$$X_1^{(2)} = \frac{1}{(-0.8471 - 1.1805)} [(-1.1805)(0.2)] = 0.11644$$

$$\phi_1 = \phi_2 = \tan^{-1}(0) = 0$$

$$x_1(t) = 0.08356 \cos 30.2853t + 0.11644 \cos 114.3801t$$

$$x_2(t) = (1.1805)(0.08356) \cos 30.2853t + (-0.8471)(0.11644) \cos 114.3801t \\ = 0.09864 \cos 30.2853t - 0.09864 \cos 114.3801t$$

Problem 4: 5.31

### 5.31 (a) Equations of motion:

Assume:  $\theta_1, \theta_2$  are small.

Moment equilibrium equations of the two masses about P and Q:

$$m l^2 \ddot{\theta}_1 + m g l \theta_1 + k d^2 (\theta_1 - \theta_2) = 0 \quad (1)$$

$$m l^2 \ddot{\theta}_2 + m g l \theta_2 - k d^2 (\theta_1 - \theta_2) = 0 \quad (2)$$

### (b) Natural frequencies and mode shapes:

Assume: Harmonic motion with

$$\theta_i(t) = \Theta_i \cos(\omega t - \phi); \quad i = 1, 2 \quad (3)$$

where  $\Theta_1$  and  $\Theta_2$  are amplitudes of  $\theta_1$  and  $\theta_2$ , respectively,  $\omega$  is the natural frequency, and  $\phi$  is the phase angle.

Using Eq. (3), Eqs. (1) and (2) can be expressed in matrix form as

$$-\omega^2 m l^2 \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{Bmatrix} \Theta_1 \\ \Theta_2 \end{Bmatrix} + \begin{bmatrix} m g l + k d^2 & -k d^2 \\ -k d^2 & m g l + k d^2 \end{bmatrix} \begin{Bmatrix} \Theta_1 \\ \Theta_2 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix} \quad (4)$$

Frequency equation:

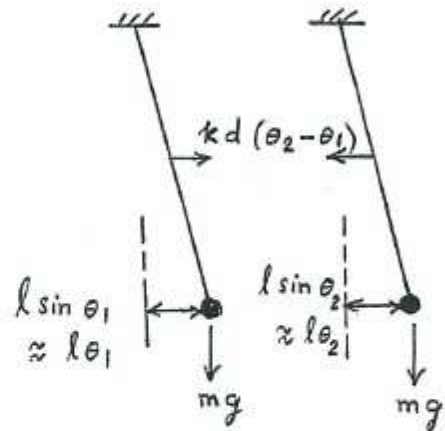
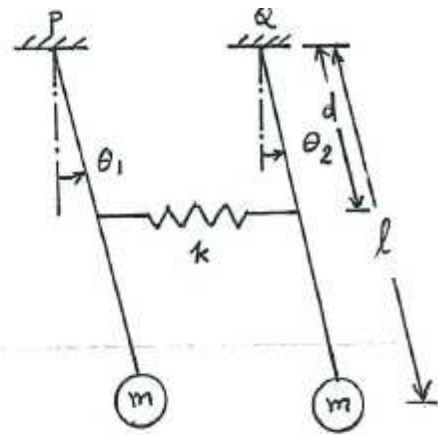
$$\begin{vmatrix} -\omega^2 m l^2 + m g l + k d^2 & -k d^2 \\ -k d^2 & -\omega^2 m l^2 + m g l + k d^2 \end{vmatrix} = 0$$

$$\text{or } \omega^4 - \omega^2 \left( \frac{2g}{l} + \frac{2kd^2}{ml^2} \right) + \left( \frac{g^2}{l^2} + \frac{2gkd^2}{ml^3} \right) = 0 \quad (5)$$

Solution of Eq. (5) gives

$$\omega_1^2 = \frac{g}{l}, \quad \omega_2^2 = \frac{g}{l} + \frac{2kd^2}{ml^2} \quad (6)$$

END



Free body diagram

By substituting for  $\omega_1^2$  and  $\omega_2^2$  into Eq. (4), we obtain

$$\left(\frac{\Theta_2}{\Theta_1}\right)^{(1)} = 1 \quad \text{or} \quad \left\{\begin{matrix} \Theta_1 \\ \Theta_2 \end{matrix}\right\}^{(1)} = \begin{Bmatrix} 1 \\ 1 \end{Bmatrix} \Theta_1^{(1)}$$

and

$$\left(\frac{\Theta_2}{\Theta_1}\right)^{(2)} = -1 \quad \text{or} \quad \left\{\begin{matrix} \Theta_1 \\ \Theta_2 \end{matrix}\right\}^{(2)} = \begin{Bmatrix} 1 \\ -1 \end{Bmatrix} \Theta_1^{(2)}$$

Thus the motion of the masses in the two modes is given by

$$\vec{\theta}^{(1)}(t) = \begin{Bmatrix} \theta_1^{(1)}(t) \\ \theta_2^{(1)}(t) \end{Bmatrix} = \Theta_1^{(1)} \begin{Bmatrix} 1 \\ 1 \end{Bmatrix} \cos(\omega_1 t + \phi_1) \quad (7)$$

$$\vec{\theta}^{(2)}(t) = \begin{Bmatrix} \theta_1^{(2)}(t) \\ \theta_2^{(2)}(t) \end{Bmatrix} = \Theta_1^{(2)} \begin{Bmatrix} 1 \\ -1 \end{Bmatrix} \cos(\omega_2 t + \phi_2) \quad (8)$$

### (c) Free vibration response:

Using linear superposition of natural modes, the free vibration response of the system is given by

$$\vec{\theta}(t) = c_1 \vec{\theta}^{(1)}(t) + c_2 \vec{\theta}^{(2)}(t) \quad (9)$$

By choosing  $c_1 = c_2 = 1$ , with no loss of generality, Eqs.

(7) to (9) lead to

$$\theta_1(t) = \Theta_1^{(1)} \cos(\omega_1 t + \phi_1) + \Theta_1^{(2)} \cos(\omega_2 t + \phi_2) \quad (10)$$

$$\theta_2(t) = \Theta_1^{(1)} \cos(\omega_1 t + \phi_1) - \Theta_1^{(2)} \cos(\omega_2 t + \phi_2) \quad (11)$$

where  $\Theta_1^{(1)}$ ,  $\phi_1$ ,  $\Theta_1^{(2)}$  and  $\phi_2$  are constants to be determined from the initial conditions. When  $\theta_1(0) = w$ ,  $\theta_2(0) = 0$ ,

$\dot{\theta}_1(0) = 0$  and  $\dot{\theta}_2(0) = 0$ , Eqs. (10) and (11) yield

$$\left. \begin{aligned} w &= \Theta_1^{(1)} \cos \phi_1 + \Theta_1^{(2)} \cos \phi_2 \\ 0 &= \Theta_1^{(1)} \cos \phi_1 - \Theta_1^{(2)} \cos \phi_2 \\ 0 &= -\omega_1 \Theta_1^{(1)} \sin \phi_1 - \omega_2 \Theta_1^{(2)} \sin \phi_2 \\ 0 &= -\omega_1 \Theta_1^{(1)} \sin \phi_1 + \omega_2 \Theta_1^{(2)} \sin \phi_2 \end{aligned} \right\} \quad (12)$$

Eqs. (12) can be solved for  $\Theta_1^{(1)}$ ,  $\phi_1$ ,  $\Theta_1^{(2)}$  and  $\phi_2$  to obtain

$$\left. \begin{aligned} \theta_1(t) &= w \cos \frac{\omega_2 - \omega_1}{2} t \cdot \cos \frac{\omega_2 + \omega_1}{2} t \\ \theta_2(t) &= w \sin \frac{\omega_2 - \omega_1}{2} t \cdot \sin \frac{\omega_2 + \omega_1}{2} t \end{aligned} \right\} \quad (13)$$

(d) conditions for beating:

$$\text{When } \frac{2 \kappa d^2}{m l^2} \ll \frac{g}{l} \quad \text{or} \quad \kappa \ll \frac{m g l}{2 d^2}, \quad (14)$$

the two frequency components in Eqs. (13), namely,  $\frac{\omega_2 - \omega_1}{2}$  and  $\frac{\omega_2 + \omega_1}{2}$ , can be approximated as

$$\Omega_1 = \frac{\omega_2 - \omega_1}{2} \simeq \frac{\kappa}{2m} \frac{d^2}{\sqrt{g l^3}} \quad (15)$$

and

$$\Omega_2 = \frac{\omega_2 + \omega_1}{2} \simeq \sqrt{\frac{g}{l}} + \frac{\kappa}{2m} \frac{d^2}{\sqrt{g l^3}} \quad (16)$$

This implies that the motions of the pendulums are given by

$$\left. \begin{aligned} \theta_1(t) &\simeq a \cos \Omega_1 t \cdot \cos \Omega_2 t \\ \theta_2(t) &\simeq a \sin \Omega_1 t \cdot \sin \Omega_2 t \end{aligned} \right\} \quad (17)$$

This motion, Eqs. (17), denotes beating phenomenon.

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Problem 5.61

5.61

Equations of motion:  $m_1 \ddot{x}_1 + (k_1 + k_2)x_1 - k_2 x_2 = F_{10} \cos \omega t = \operatorname{Re}(F_{10} e^{i\omega t})$

$$m_2 \ddot{x}_2 + (k_2 + k_3)x_2 - k_2 x_1 = F_{20} \cos \omega t = \operatorname{Re}(F_{20} e^{i\omega t})$$

Assuming  $x_j(t) = X_j e^{i\omega t}$ ,  $j=1,2$  along with  $F_j(t) = F_{j0} e^{i\omega t}$ ,  $j=1,2$ , the equations of motion can be expressed as

$$(-\omega^2 m_1 + k_1 + k_2) X_1 - k_2 X_2 = F_{10}$$

$$-k_2 X_1 + (-\omega^2 m_2 + k_2 + k_3) X_2 = F_{20}$$

$$\text{i.e. } [Z(i\omega)] \vec{X} = \vec{F}_0 \quad \text{---- (E}_1\text{)}$$

$$\text{where } Z_{11}(i\omega) = -\omega^2 m_1 + k_1 + k_2, \quad Z_{12}(i\omega) = Z_{21}(i\omega) = -k_2,$$

$$Z_{22}(i\omega) = -\omega^2 m_2 + k_2 + k_3,$$

$$\vec{X} = \begin{Bmatrix} X_1 \\ X_2 \end{Bmatrix}, \quad \vec{F}_0 = \begin{Bmatrix} F_{10} \\ F_{20} \end{Bmatrix}$$

Solution of (E<sub>1</sub>) can be expressed, using Eqs. (5.35), as

$$X_1 = \frac{(-\omega^2 m_2 + k_2 + k_3) F_{10} + k_2 F_{20}}{(-\omega^2 m_1 + k_1 + k_2)(-\omega^2 m_2 + k_2 + k_3) - k_2^2} \quad \text{---- (E}_2\text{)}$$

$$X_2 = \frac{k_2 F_{10} + (-\omega^2 m_1 + k_1 + k_2) F_{20}}{(-\omega^2 m_1 + k_1 + k_2)(-\omega^2 m_2 + k_2 + k_3) - k_2^2} \quad \text{---- (E}_3\text{)}$$

Since  $X_1$  and  $X_2$  are real (since there is no damping), the final solution is given by

$$x_1(t) = X_1 \cos \omega t$$

$$x_2(t) = X_2 \cos \omega t$$

where  $X_1$  and  $X_2$  are given by (E<sub>2</sub>) and (E<sub>3</sub>).