

Statistical Inference for Nanopore Sequencing

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ABSTRACT DNA translocation through a nanopore is modeled as a one-dimensional biased random walk. A Hidden Markov Model is used to infer the input DNA sequence from a set of output read sequences. Bounds on inference accuracy are set based on model parameters.

Received for publication 1 April 2013 and in final form 1 April 2013.

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Current generation sequencing platforms have led to an explosion in available sequence data, rapidly advancing our understanding of the genomic underpinnings of life. Making sense of this data has relied on the development of new algorithms in bioinformatics to process the raw short-read data into usable sequence assemblies. However, the shortness of the raw data makes some information, particularly phasing and haplotype structure, unavailable even at the highest levels of coverage. The development of longer read sequence technologies is critical for continued progress in genomics (cite?).

Nanopore sequencing has emerged as a potential candidate to supercede current generation sequencing and allow for theoretically unlimited read lengths [1]. A number of strategies for sequencing DNA with nanopores have been proposed, with the common basis of detecting individual nucleotides as they pass through a nanometer-scale aperture in a thin membrane separating two electrolytes. To date, a primary obstacle of this approach has been overcoming the fast stochastic motion of the individual molecules as they are driven through the pore [2, 3]. While recent methods have been proposed in order to controllably 'ratchet' DNA molecules through a nanopore one base at a time, experimental results show that motion of the molecule can still occur in both forward and backward directions within a single read [4–6]. Therefore, unidirectional stepwise motion remains difficult to reliably achieve, leading to a novel source of error in the sequence read due to this diffusive motion.

In this letter, we demonstrate a statistical method of combining multiple reads from a given input DNA sequence passing through a nanopore that is able to yield highly accurate sequencing, even in the presence of highly diffusive

molecular motion. We consider a simple physical model of DNA translocation as a one-dimensional biased random walk, and demonstrate the use of a Hidden Markov Model (HMM) to obtain a statistical estimate of the most likely input DNA sequence generating the observed series of reads. We then use this model to set bounds on the achievable inference accuracy as a function of experimental parameters. We conclude that multiple parallel reads are sufficient to compensate for both highly diffusive motion and high base-call error rates.

Translocation through the nanopore is modeled as a one-dimensional diffusion with driving force F and unit base step size a , which we treat as a biased random walk with fixed forward bias p [7], where p is given by

$$p = \frac{1}{2} + \frac{Fa}{4k_B T}. \quad (1)$$

Given an input DNA sequence of length L , we generate a single output read of the sequence by stepping through the sequence with a bias p , at each step making a base call with error probability e . Note that this error rate is strictly due to the base call and is independent of errors introduced due to backward motion. We treat the sequencing device itself as a black box, and assume that an appropriate method of making a base call from the raw signal has been implemented. From this simple physical model, we generate a set of N observed reads, each with length $T^n \geq L$ and denoted $X_{1:T^n}^n$. A representation of the experimental configuration is shown in Figure 1. In Figure 1B the distribution of obtained read lengths as a fraction of the input sequence length is shown for several different forward biases, and in Figure 1C an example of the stochastic nature of the transit is shown.

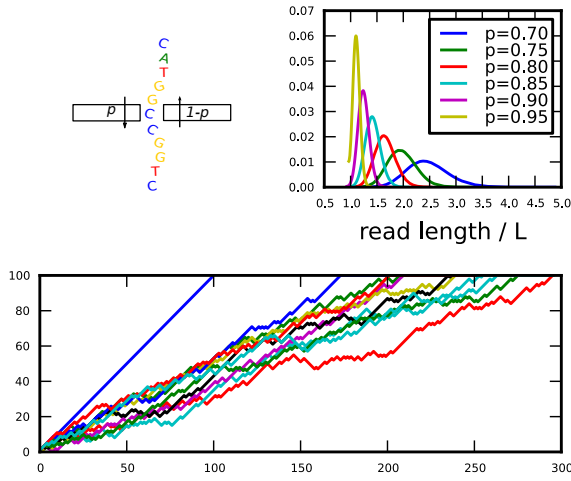


FIGURE 1 (A) Idealized example of a nanopore sequencing device. (B) Transit time distribution for different forward biases. (C)

Given this set of N read sequences, the statistical task is to infer the true sequence most likely to have generated the observed data. We model the data as a Hidden Markov Model (HMM) [8], where each output read is modeled as a discrete set of observed states, $\mathbf{x} = \{x_1 \dots x_T\}$, $x_i \in (A, G, C, T)$, a vector of observed bases, and a discrete set of hidden states, $\mathbf{z} = \{z_1 \dots z_T\}$, $z_i \in (1 \dots L)$, the unknown position along the sequence. For an HMM, we require three model parameters, the initial state distribution $\pi = p(\mathbf{z}_1)$, the state transition matrix $A = p(\mathbf{z}_t | \mathbf{z}_{t-1})$, and an emission distribution $S = p(\mathbf{x}_n | \mathbf{z}_n)$. π and A are taken as fixed given the experimental conditions, and our attention focuses on optimizing $p(\mathbf{x}_n | \mathbf{z}_n)$, which directly represents our true sequence inference. We maximize the complete data log likelihood, $p(X, Z | \theta)$, with respect to the model parameters using an implementation of expectation-maximization (EM). Because the likelihood factorizes over the independent output reads, $p(X, Z | \theta) = \prod_n p(X^n, Z^n | \theta)$, we can perform expectation updates on each read individually before averaging results in the maximization step. This key point provides us with a method of jointly using the all the information provided in the reads while still maintaining an efficient, parallel calculation. After a convergence criteria on the complete data log likelihood is satisfied, we recover an estimated emission distribution S , which can be converted to an estimate for the DNA sequence by taking $\max_d S_{dl}$. The algorithm has complexity $O(L^3)$. An example of the output of this algorithm showing the relationship between the true sequence and the inferred sequence distribution is shown in Figure X. Additionally, the column-wise entropy of the sequence inference is shown in Figure X(c). (to make)

We examined the performance of the algorithm over a range of parameter values in order to identify a minimal experimental configuration capable of sequence inference at a

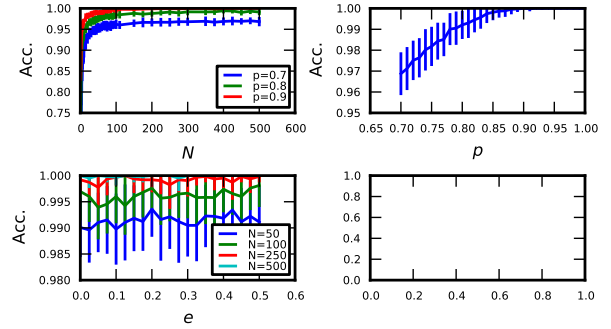


FIGURE 2 Parameter Sweeps. We plot inference accuracy (measured as 1 - edit distance/length) as a function of experimental parameters. (A) N vs P . (B) N vs P . (C) e vs N . (D) N vs e .

given accuracy. First, we consider how to select realistic values for our parameters. $p = 0.5$ corresponds to completely symmetric diffusion and is unlikely to reach the end of the sequence without exiting from the front (?) side (we discard these reads in the data generation). $p = 1.0$ corresponds to nondiffusive motion and is a trivial case in this model. Real experiments will operate somewhere in between. Lu et al. report an experimental bias term of $Fa/4k_B T = 0.2$, for which they claim accurate sequence recovery is impossible [3]. We take this as a starting point for investigation, exploring the range $p = 0.7$ to $p = 1.0$. An appropriate base-call error rate will be device specific and is presently unknown, so we examine the complete range, $e = 0.0$ to $e = 0.5$. The number of reads, N , is potentially unlimited. A sweep across these parameters is shown in Figure 2. We observe that, as expected, the strongest determinant of achievable sequence accuracy is the number of reads, N . This is expected because each read contributes an independent observation of the input DNA sequence. We are most interested in the relationship between forward bias and the number of reads. We observe an inference error rate that falls rapidly with increasing N , even at highly diffusive motion. Generally, the achievable accuracy converges at a fixed p even with the addition of more reads (Figure 2A). We compare this accuracy in Figure 2B and see that a threshold accuracy of 99% is achievable for p as low as 0.8. Finally, we observe that even high base-call error rates have minimal effect on inference accuracy, provided sufficient multiple reads are provided for averaging.

We next applied the method to fragments of λ -phage DNA. We observe that, contrary to random DNA, the addition of more reads was often unable to bring the accuracy to 100%. We determined that this is due to specific sequence patterns that are more difficult for the algorithm to correctly parse, suggesting that a more specific method of identifying backwards motion may be required.

Next, we observe that errors in the sequence inference due

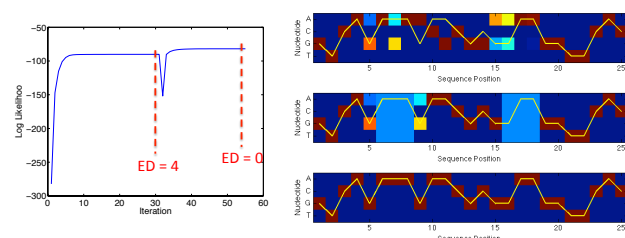


FIGURE 3 Restarts using entropy heuristic

to diffusive motion can take the form of extraneous insertions and deletion along the true sequence. These errors are not random, but have a predictable pattern that makes it possible to identify and correct them during the inference algorithm. This type of error occurs when there are multiple paths that may explain the observed data, and manifest themselves in the inference estimate as regions of higher entropy. We used this observation to develop a heuristic that, acting after an initial maximum likelihood convergence, locates regions of high entropy, randomly introduces shifts in the sequence estimate, and restarts EM inference. If the shifted inference yields a higher log likelihood, it is kept as explaining the data better than the previous estimate. The procedure is repeated until no high entropy regions remain. The procedure is detailed in Figure 3 on a segment of λ -phage DNA.

We have demonstrated a statistical method for inferring the DNA sequence passing through a nanopore sequencer under diffusive motion. The strength of the method lies in the way in which multiple reads of the input sequence are independently modeled and then iteratively combined to yield a joint estimate of the true sequence. Under this model, accurate sequence inference is achievable even within the experimental constraints of today's nanopore sequencers. ($p \approx 0.7$). Designers of nanopore sequencers may find that there is a tradeoff between sequencing speed and maintaining unidirectional motion of the molecules. In this scenario, techniques such as the random walk model may allow increased throughput without a loss of accuracy.

SUPPLEMENTARY MATERIAL

A full derivation of the model and additional supplementary material can be found by visiting BPJ Online at <http://www.biophysj.org>.

ACKNOWLEDGEMENTS

We acknowledge useful discussions with David Blei and David Pfau. Funding blurb.

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