## Project 4 FYS4150

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#### Abstract

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#### 1 Introduction

## 2 Theory

#### 2.1 Introduction to statistics

$$P(a \le X \le b) = \int_{b}^{a} p(x) dx \tag{1}$$

$$\langle h \rangle_X = \int h(x) p(x) dx$$
 (2)

$$\langle x^n \rangle = \int x^n \, p(x) \, dx \tag{3}$$

$$\langle x \rangle = \int x \, p(x) \, dx$$
 (4)

#### 2.2 Magnet - properties

Energy, magnetic moment, susceptibility, heat capacity.

$$T_C(L) - T_C(L = \infty) = aL^{-1/\nu}$$
 (5)

Boltzman Markhow chain - convergance

L=2 case:

Table 2.1: text

No spin up	$\operatorname{Deg}$	Energy	Magnetization
0	1	-8J	-4
1	4	0	-2
2	4	0	0
2	2	8J	0
3	4	0	2
4	1	-8J	4

Error Random number

The partition function:

$$Z = \sum_{i}^{M} e^{-\beta E_i} = e^{-\beta 8J} + e^{-\beta 8J} + e^{\beta 8J} e^{\beta 8J} + 12$$

$$= 2e^{-\beta 8J} + 2e^{\beta 8J} + 12 = 4\left(\frac{e^{-\beta 8J} + e^{-\beta 8J}}{2}\right) + 12$$

$$= 4\cosh(\beta 8J) + 12$$
In

The energy:

$$\langle E \rangle = k_B T^2 \left( \frac{\partial Z}{\partial T} \right)_{V,N}$$

$$= k_B T^2 \frac{\partial}{\partial T} \left[ \ln \left( 4 \cosh \left( \frac{8J}{k_B T} \right) + 12 \right) \right]$$

$$\partial \ln Z = \partial Z \partial \beta = \partial \ln Z \left( -1 \right)$$

$$\frac{\partial \ln Z}{\partial T} = \frac{\partial Z}{\partial \beta} \frac{\partial \beta}{\partial T} = \frac{\partial \ln Z}{\partial \beta} \left( \frac{-1}{k_B T^2} \right)$$

$$\langle E \rangle = -\left(\frac{\partial Z}{\partial \beta}\right)_{V,N} = -\frac{\partial}{\partial \beta} \ln\left[4\cosh\left(8J\beta\right) + 12\right]$$

$$= \frac{-1}{4\cosh(8J\beta) + 12} 4\sinh(8J\beta)8J\beta$$

$$= \frac{-8J\sinh(8J\beta)}{3\cosh((J\beta) + 4}$$

Following the same method, we found that:

$$\langle |M| \rangle = \frac{1}{Z} \sum_{i=1}^{M} M_i e^{\beta E_i} = \frac{(8J)^2 \cosh(8J\beta)}{\cosh(8J\beta) + 3}$$

$$\langle M \rangle = 0$$

$$\langle E^2 \rangle = \frac{8 \left( e^{8J\beta} + 1 \right)}{\cosh(8J\beta) + 3}$$

$$\langle M^2 \rangle = \frac{1}{Z} \left( \sum_{i=1}^{M} M_i^2 e^{\beta E_i} \right) = \frac{2 \left( e^{8J\beta} + 2 \right)}{\cosh(8J\beta) + 3}$$

We can use these to calculate the rest:

$$C_V = k\beta^2 \left( \langle E^2 \rangle - \langle E \rangle^2 \right)$$

$$\chi = \beta \left( \left\langle M^2 \right\rangle - \left\langle M \right\rangle^2 \right)$$

## 3 Method

In this project we tried out many new concepts in our algorithm. We used Monte Carlo method.

Table 4.1: This table compares the analytical values for L=2 with the numerical ones after  $10^6$  Monte Carlo cycles. The values are in units per spin.

	Numerical:	Analytical:
$\langle E \rangle$	-1.9958	-1.9960
$\langle E^2 \rangle$	15.9664	15.9679
$\langle M \rangle$	0.0451	0
$\langle M^2 \rangle$	3.9930	3.9933
$\langle  M  \rangle$	0.9986	0.9987
$\chi$	3.9849	3.9933
$C_V$	0.0335	0.0321

#### 3.1 Monte Carlo cycles

In Monte Carlo methods, the goal is to

#### 3.2 Metropolis algorithm

#### 3.3 Random number generator

#### 3.4 Parallelizing

Metropolis (T,A,...) Stochastic matrix - convergences (forhold eigenvalue).

Hvilken random number engine

#### MPI:

- Develop codes locally, run with some few processes and test your codes. Do benchmarking, timing and so forth on local nodes, for example your laptop or PC. - When you are convinced that your codes run correctly, you can start your production runs on available supercomputers.

MPI functions:

#### 4 Result

#### 4.1 Matrix dimension L=2

OBS! Need an number of MC cycles necessary!

All calculations in this subsection are at T=1.0 K.

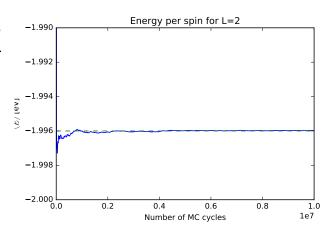


Figure 4.1: This is a plot of the expectation value of the energy per spin verus number of Monte Carlo cycles. The plot shows that at least  $9 \cdot 10^5$  MC cycles are necessary for a good argeement.

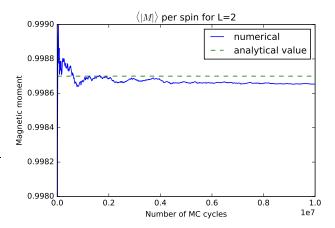


Figure 4.2: This is a plot of the expectation value of the mean absolute value of the magnetic moment per spin verus number of Monte Carlo cycles. The plot shows that at least  $8 \cdot 10^5$  MC cycles are necessary for a good argeement, but all the way to  $10^6$  the value is a bit low.

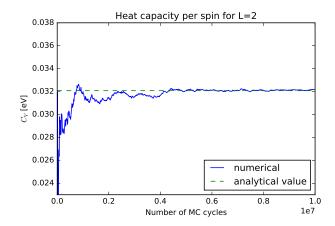


Figure 4.3: This is a plot of the heat capacity per spin verus number of Monte Carlo cycles. The plot shows that at least  $6 \cdot 10^5$  MC cycles are necessary for a good argeement.

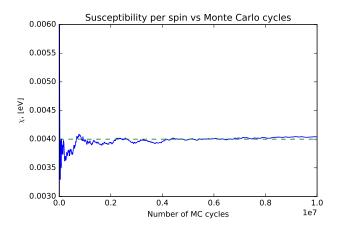


Figure 4.4: This is a plot of the susceptibility per spin verus number of Monte Carlo cycles. The plot shows that at least  $6 \cdot 10^5$  MC cycles are necessary for a good argeement.

#### 4.2 Matrix dimension L = 20

HMM: Should define an area that is enough for equilibrium!

OBS: Need the number of MC cycles to reach equilibrium!

OBS: Need equilibration time! (5 1e5?)

OBS: Comment accepted configs T dependency

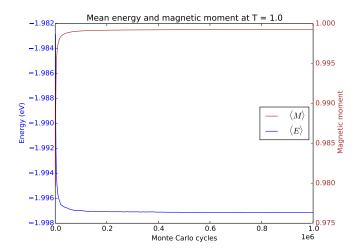


Figure 4.5: This is a plot of both the expectation value of the energy and absolute magnetic moment per spin verus number of Monte Carlo cycles at  $T=1.0~\mathrm{K}$ . The plot shows that at least ??C cycles are necessary to reach equilibrium.

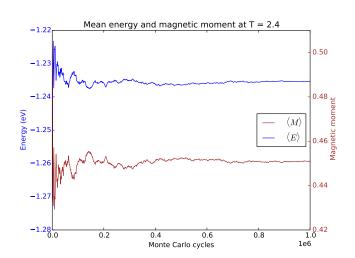


Figure 4.6: This is a plot of both the expectation value of the energy and absolute magnetic moment per spin verus number of Monte Carlo cycles at T = 2.4 K. The plot shows that at least ??C cycles are necessary to reach equilibrium.

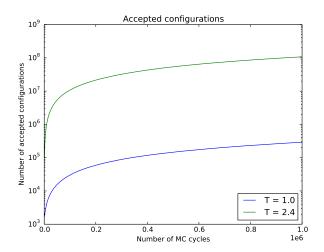


Figure 4.7: This is a plot of the total number of accepted configurations versus number of Monte Carlo cycles with random initial state.

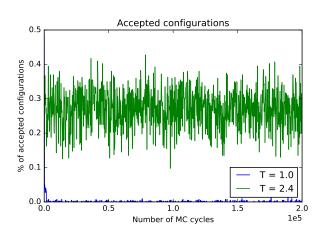


Figure 4.8: This is a plot of the percentage accepted of attempted configurations versus Monte Carlo cycles with random initial state.

### 4.3 Energy probability

OBS: Compare result with computed variance!

OBS: Discuss behavior (In Discussion - maybe just merge result and discussion?)

Computed variance (from same dataset?):

$$\sigma_E^2 = \left\langle E^2 \right\rangle - \left\langle E \right\rangle^2$$

T = 1.0 K:

$$\sigma_E^2 = 638181 - (-798.855)^2 = 11.69$$
$$\sigma = 3.42$$

T = 2.4 K:

$$\sigma_E^2 = 247886 - (-494.628)^2 = 3229.14$$

$$\sigma = 56.8$$

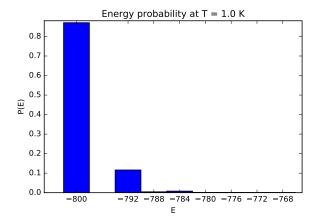


Figure 4.9: This is a plot of the energy probability when T = 1.0 K. The energy is the total energy of the 2D lattice with  $20 \times 20$  spins.

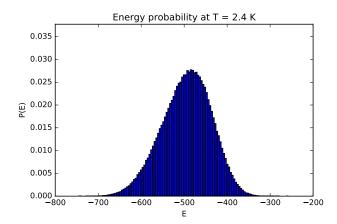


Figure 4.10: This is a plot of the energy probability when T = 2.4 K.The energy is the total energy of the 2D lattice with  $20 \times 20$  spins.

# 4.4 Increasing dimensionality/ Critical temperature

OBS: Plot of E, M, Cv, X as functions of T (put L as legend and plot together)

OBS: Indication of phase transition? (Peak - at least for Cv and X)

OBS: Use Equation 5 to extract  $T_C$ .

Getting these equations from 5 where  $\nu = 1$ :

$$T_C(40) - T_C(\infty) = a40^{-1}$$

$$T_C(60) - T_C(\infty) = a60^{-1}$$

$$T_C(80) - T_C(\infty) = a80^{-1}$$

$$T_C(100) - T_C(\infty) = a100^{-1}$$

## References

[1] Lars Onsager. Crystal statistics. i. a twodimensional model with an order-disorder transition. Phys. Rev., 65:117-149, Feb 1944.

## Appendix

	State	Spinn	Energi	Magnetization
$T_C(\infty) = -a40^{-1} + T_C(40)$	0	$\downarrow\downarrow\downarrow\downarrow$	-8J	-4
$T_C(\infty) = -a60^{-1} + T_C(60)$	1	$\downarrow\downarrow\downarrow\uparrow\uparrow$	0	-2
$T_C(\infty) = -a80^{-1} + T_C(80)$	2	$\downarrow\downarrow\uparrow\uparrow\downarrow$	0	-2
$T_C(\infty) = -a100^{-1} + T_C(100)$	3	$\downarrow\uparrow\downarrow\downarrow$	0	-2
	4	$\uparrow\downarrow\downarrow\downarrow$	0	-2
	5	$\downarrow\downarrow\uparrow\uparrow\uparrow$	0	0
	6	$\downarrow\uparrow\downarrow\uparrow$	0	0
$= kTC/J = 2/\ln(1+\sqrt{2}) \approx 2.269$ [1]	7	$\downarrow\uparrow\uparrow\downarrow$	8J	0
, , , , , , , , , , , , , , , , , , , ,	8	$\uparrow\downarrow\downarrow\uparrow\uparrow$	8J	0
	9	$\uparrow\downarrow\uparrow\downarrow$	0	0
scussion	10	$\uparrow\uparrow\downarrow\downarrow$	0	0
	11	$\downarrow\uparrow\uparrow\uparrow$	0	2

12

13

14

15

 $\uparrow\downarrow\uparrow\uparrow$ 

 $\uparrow\uparrow\downarrow\uparrow$ 

 $\uparrow\uparrow\uparrow\downarrow$ 

 $\uparrow\uparrow\uparrow\uparrow$ 

0

0

0

-8J

2

2 2

4

- $T_{i}$  $T_{i}$  $T_{i}$
- $T_{i}$

Exact  $T_C =$ 

- Dis 5
- Conclusion 6