

importance are collected in Table 26. When several states of the same kind result their number is given in parentheses after the term symbol. The number of  $\Sigma$  states that arise is always odd. There is always one  $\Sigma$  state (with  $M_{L_1} = 0$ ,  $M_{L_2} = 0$ ) in addition to pairs of  $\Sigma^+$  and  $\Sigma^-$  states similar to that discussed above for  $P + P$ . The symmetry of this single  $\Sigma$  state cannot be obtained by elementary means. According to Wigner and Witmer it is  $\Sigma^+$  when  $L_1 + L_2 + \Sigma l_{i1} + \Sigma l_{i2}$  is even and it is  $\Sigma^-$  when this sum is odd. It will be remembered that  $\Sigma l_i$  determines the *parity*, even or odd, of an atomic state. Thus the symmetry of the  $\Sigma$  state in question depends on the  $L$  values as well as the parities of the atomic states from which it results. It is for this reason that the parities are indicated by subscripts  $g$  and  $u$  in Table 26.

TABLE 26. MOLECULAR ELECTRONIC STATES RESULTING FROM GIVEN STATES OF THE SEPARATED (UNLIKE) ATOMS

[According to Wigner and Witmer (712); see also similar tables in Mulliken (514).]

States of the Separated Atoms	Molecular States
$S_g + S_g$ or $S_u + S_u$	$\Sigma^+$
$S_g + S_u$	$\Sigma^-$
$S_g + P_g$ or $S_u + P_u$	$\Sigma^-, \Pi$
$S_g + P_u$ or $S_u + P_g$	$\Sigma^+, \Pi$
$S_g + D_g$ or $S_u + D_u$	$\Sigma^+, \Pi, \Delta$
$S_g + D_u$ or $S_u + D_g$	$\Sigma^-, \Pi, \Delta$
$S_g + F_g$ or $S_u + F_u$	$\Sigma^-, \Pi, \Delta, \Phi$
$S_g + F_u$ or $S_u + F_g$	$\Sigma^+, \Pi, \Delta, \Phi$
$P_g + P_g$ or $P_u + P_u$	$\Sigma^+(2), \Sigma^-, \Pi(2), \Delta$
$P_g + P_u$	$\Sigma^+, \Sigma^-(2), \Pi(2), \Delta$
$P_g + D_g$ or $P_u + D_u$	$\Sigma^+, \Sigma^-(2), \Pi(3), \Delta(2), \Phi$
$P_g + D_u$ or $P_u + D_g$	$\Sigma^+(2), \Sigma^-, \Pi(3), \Delta(2), \Phi$
$P_g + F_g$ or $P_u + F_u$	$\Sigma^+(2), \Sigma^-, \Pi(3), \Delta(3), \Phi(2), \Gamma$
$P_g + F_u$ or $P_u + F_g$	$\Sigma^+, \Sigma^-(2), \Pi(3), \Delta(3), \Phi(2), \Gamma$
$D_g + D_g$ or $D_u + D_u$	$\Sigma^+(3), \Sigma^-(2), \Pi(4), \Delta(3), \Phi(2), \Gamma$
$D_g + D_u$	$\Sigma^+(2), \Sigma^-(3), \Pi(4), \Delta(3), \Phi(2), \Gamma$
$D_g + F_g$ or $D_u + F_u$	$\Sigma^+(2), \Sigma^-(3), \Pi(5), \Delta(4), \Phi(3), \Gamma(2), \text{H}$
$D_g + F_u$ or $D_u + F_g$	$\Sigma^+(3), \Sigma^-(2), \Pi(5), \Delta(4), \Phi(3), \Gamma(2), \text{H}$

We have yet to determine the *multiplicity* of the resulting molecular states. Let us assume that the coupling of the  $L_i$  to the field between the nuclei is strong compared to the coupling between  $L_i$  and  $S_i$ . Then, since the spin is not influenced by an electric field, the two spin vectors  $S_1$  and  $S_2$  of the separated atoms add together forming a resultant  $S$ , the resultant spin vector of the molecule. For the corresponding quantum number  $S$  we have (see p. 25)

$$S = (S_1 + S_2), (S_1 + S_2 - 1), (S_1 + S_2 - 2), \dots, |S_1 - S_2|. \quad (\text{VI}, 2)$$

For a given orientation of the  $L_i$ , each of the values of  $S$  in (VI, 2) is possible; that is, *each of the states given in Table 26 can occur with each of the multiplicities*