

Department of Physics, NTNU

Homework 2 TFY4210/FY8302 Quantum theory of solids Spring 2021.

Problem 1

Consider a tightbinding Hamiltonian with nearest-neighbor hopping on a general two dimensional Bravais lattice with N lattice points, in contact with a particle reservoir. The chemical potential of the system is denoted μ . The nearest-neighbor hopping element is given by t . The Hamiltonian is given by

$$H = -t \sum_{\langle i,j \rangle, \sigma} c_{i,\sigma}^\dagger c_{j,\sigma} - \mu \sum_{i,\sigma} c_{i,\sigma}^\dagger c_{i,\sigma}$$

Here, $(c_{i,\sigma}^\dagger, c_{i,\sigma})$ create and destroy particles in spin-state σ on site i . Introduce Fourier-transformed operators

$$\begin{aligned} c_{\mathbf{k},\sigma}^\dagger &= \frac{1}{\sqrt{N}} \sum_i c_{i,\sigma}^\dagger e^{-i\mathbf{k}\cdot\mathbf{r}_i} \\ c_{\mathbf{k},\sigma} &= \frac{1}{\sqrt{N}} \sum_i c_{i,\sigma} e^{i\mathbf{k}\cdot\mathbf{r}_i} \end{aligned}$$

where \mathbf{r}_i is the position at lattice site i .

a) Show that Hamiltonian may be written on form

$$H = \sum_{\mathbf{k},\sigma} E_{\mathbf{k}} c_{\mathbf{k},\sigma}^\dagger c_{\mathbf{k},\sigma}$$

and give an expression for $E_{\mathbf{k}}$ for a general two-dimensional Bravais lattice.

c) Imagine that we now consider a model on a *honeycomb* lattice, see Figure 1. Explain what the *principal* difference between this lattice and the square lattice is.

d) Write down a tight-binding model for this system, considering only nearest neighbor interactions, and find $E_{\mathbf{k}}$. (Hint: You need to introduce two distinct types of fermions, one type for the red atoms and one type for the blue atoms).

Problem 2

Consider a tight-binding model in a uniform external magnetic field h directed along the z-direction. The Hamiltonian is given by

$$H = -t \sum_{\langle i,j \rangle, \sigma} c_{i,\sigma}^\dagger c_{j,\sigma} - \mu \sum_{i,\sigma} c_{i,\sigma}^\dagger c_{i,\sigma} - h \sum_{i,\sigma} \sigma c_{i,\sigma}^\dagger c_{i,\sigma}$$

in the same notation as in Problem 1. The last term is the second-quantized version of the Zeeman-term $-\mathbf{h} \cdot \mathbf{S} = -h \sum_i S_{iz} = -h \sum_i \left[c_{i,\uparrow}^\dagger c_{i,\uparrow} - c_{i,\downarrow}^\dagger c_{i,\downarrow} \right] = -h \sum_{i,\sigma} \sigma c_{i,\sigma}^\dagger c_{i,\sigma}$.

a) Introduce the same Fourier-transform as in Problem 1, and show that the Hamiltonian may be written on the form

$$H = \sum_{\mathbf{k}, \sigma} E_{\mathbf{k}, \sigma} c_{\mathbf{k}, \sigma}^\dagger c_{\mathbf{k}, \sigma}$$

and give an expression for $E_{\mathbf{k}, \sigma}$ for a general two-dimensional Bravais lattice.

b) Consider instead another variant of the tight-binding model, now in the absence of a magnetic field, but with a spin-dependent hopping matrix element. (In lectures Week 4, we discuss what the origin of such spin-dependent hopping could be). The Hamiltonian is given by, with a general spin-dependent hopping matrix element $t_{ij}^{\sigma\sigma'}$

$$H = - \sum_{\langle i,j \rangle, \sigma, \sigma'} t_{ij}^{\sigma\sigma'} c_{i,\sigma'}^\dagger c_{j,\sigma} - \mu \sum_{i,\sigma} c_{i,\sigma}^\dagger c_{i,\sigma}$$

Consider a special case $t_{ij}^{\sigma\sigma'} = t_{ij}^{\sigma\sigma} \delta_{\sigma,\sigma'}$, with $t_{ij}^{\uparrow\uparrow} \neq t_{ij}^{\downarrow\downarrow}$. Introduce the same Fourier-transform as in Problem 1, and show that the Hamiltonian may be written on the form

$$H = \sum_{\mathbf{k}, \sigma} E_{\mathbf{k}, \sigma} c_{\mathbf{k}, \sigma}^\dagger c_{\mathbf{k}, \sigma}$$

and give an expression for $E_{\mathbf{k}, \sigma}$ for a general two-dimensional Bravais lattice.

c) Compare with the results in Problem 2 a), and show that this sort of spin-dependent hopping may be interpreted as particles moving in a \mathbf{k} -dependent external magnetic field along the z -axis.