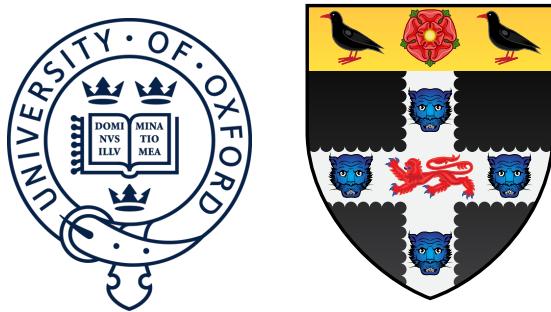


First Measurement of the Solar Neutrino Oscillation Parameters via Boron-8 Solar Neutrinos in SNO+



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A thesis submitted in fulfilment of the requirements for the degree of

Doctor of Philosophy

Trinity Term 2022

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Saturday 8th July, 2023 – 17:56

To my parents

and

To Oscar Jacobsson:

The best of us

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Acknowledgements

And I would like to acknowledge ...

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Abstract

Formal 1-page summary of the work completed in the thesis.

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Statement of Authorship

The nature of contemporary experimental particle physics is that the research work is highly collaborative. The SNO+ experiment, which this thesis is about, is an 80+ member Collaboration and as such is no exception. As a result, the work discussed in this thesis is in many places built upon the tireless work of many others within the Collaboration. This Statement, as well as throughout the thesis itself, endeavours to note where the author himself did the research, and where work was taken from others.

Chapters 1 and 2 cover the background knowledge of neutrino physics as well as the SNO+ detector needed to best understand the work performed in the rest of this thesis. This information has been compiled from a variety of books, journal articles, theses, and SNO+ internal technical reports. The author was trained as both a detector ‘operator’ and ‘expert’, taking dozens of shifts to monitor the detector whilst it was taking data, as well as being on-call to help if any issues arose with the detector electronics. The author also made numerous contributions to the Collaboration’s central software package, RAT.

Chapters 3–5 cover the work done by the author on the SMELLIE optical calibration system for SNO+. This calibration system historically been built and analysed by multiple people, including: Krishanu Majumdar, Esther Turner, Stephanie Langrock, and Jeff Lidgard. The author installed multiple pieces of new hardware for SMELLIE on-site, with help from Armin Reichold and Jeff Lidgard. With the same team, the author also helped to integrate the new hardware with the existing SMELLIE server

software so that the newly-installed hardware could actually be effectively used. The author has taken dozens of hours of calibration data using SMELLIE throughout the course of the DPhil, both by himself and alongside Armin Reichold, Jeff Lidgard, and Ana Sofia Inácio.

Chapter 4 considers improvements to the simulation of SMELLIE events; the work done by Esther Turner on this subject is considered as a starting point. The rest of the chapter covers a new simulation approach created and developed by the author. Similarly, the scattering analysis performed within Chapter 5 was inspired from the analyses done by Krishanu Majumdar, Esther Turner, and Stephanie Langrock. However, the work in this thesis uses a new method created by the author, and uses new scintillator phase data taken by the author. Also in this chapter is a separate analysis of this same data to measure the scintillator’s extinction length. This was a novel analysis for SMELLIE, designed and implemented entirely by the author.

Finally, Chapter 6 describes an analysis of scintillator phase data performed by the author to measure the solar neutrino oscillation parameters. An initial background-free sensitivity study of this topic was originally performed by Javier Caravaca. However, the work done by the author in this thesis analyses actual data, includes all relevant backgrounds appropriately, and considers systematics. Projections of expected sensitivity with greater livetime were also performed by the author. The analysis made by the author also uses a Bayesian MCMC approach, using the **OXO** signal extraction software framework initially built by Jack Dunger. A number of further people have since improved **OXO**, including the author: substantial improvements to the handling of systematics floated within the MCMC fit, as well as allowing for the floating of oscillation parameters within the fit.

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List of Acronyms

AmBe Americium-Beryllium radioactive source

AMELLIE Attenuation Module for ELLIE

AV Acrylic vessel

BD 1,2-Butanediol

BHT Butylated hydroxytoluene

BisMSB 1,4-Bis(2-methylstyryl)benzene

CI Credible Interval

CTC Crate Trigger Card

DAQ Data acquisition (system)

DB Daughter Board

DDA N,N-Dimethyldodecylamine

ECA Electronic Calibration

ELLIE Embedded LED/Laser Light Injection Entity

EXTA External Asynchronous (Trigger)

FEC Front-End Card**FV** Fiducial Volume**GTID** Global Trigger Identification number**GT** Global Trigger**HPD** (The point of) Highest Posterior Density**IW** In-window events**LAB** Linear alkylbenzene**MPU** Monitoring PMT Unit (for SMELLIE)**MTC/A+** Analogue Master Trigger Card**MTC/D** Digital Master Trigger Card**npe** Number of photoelectrons**OOW** Out-of-window events**OWLs** Outward-looking PMTs**PCA** PMT Calibration**PMTIC** PMT Interface Card**PMT** Photomultiplier Tube**PPO** 2,5-Diphenyloxazole**PQ** PicoQuant**PSUP** PMT support structure

QE Quantum efficiency (of a PMT)

QHL Charge with high gain over a ‘long’ integration time (390 ns)

QHS Charge with high gain over a ‘short’ integration time (60 ns)

QLX Charge with low gain over a ‘long’ integration time (390 ns)

SMELLIE Scattering Module for ELLIE

SNO Sudbury Neutrino Observatory

TAC Time-to-amplitude Converter

TeA Telluric acid, Te(OH)₆

TELLIE Timing subsystem for ELLIE

TeLS Tellurium-loaded liquid scintillator

TOF Time-of-flight (of a photon)

TTS Transit time spread (of a PMT)

TUBii Trigger Utility Board Mark ii

UPW Ultra-pure water

VFA Remotely-controllable Variable Fibre Attenuator

ZDAB Zebra Database (file format)

nhit Number of PMT hits in an event

RATDB RAT Database

RAT Reactor Analysis Tool

ADC Analogue-to-Digital Converter

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Introduction

Couple of pages outlining document's structure and contents (this is what each of the chapters is here for). Less formal than the abstract, also explaining what the expected audience of this thesis is: who will find this document useful!

1

2

3

4

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Chapter 1

The Theory of Neutrino Physics

Light

Light

The visible reminder of Invisible Light

The Rock

T. S. ELIOT

1.1 The Standard Model and Neutrinos

1.1.1 A Brief Introduction to the Standard Model

Covering how the SM works at the highest level, including:

- Quantum Field Theory and the Lagrangian dynamical framework
- The connection between symmetries of a QFT model and its gauge fields that describe the model's forces
- The SM's fundamental symmetries, and associated forces, but —

- ¹ • Not (exactly) what we see “normally”! The electromagnetic and weak forces
 - ² appear distinct, and the weak gauge bosons have mass. To explain this, we need
 - ³ a further component, the Brout-Englert-Higgs (BEH) Mechanism.
- ⁴ [2 pages total]

⁵ 1.1.2 Neutrinos within the Standard Model

- ⁶ • Basic description of where neutrinos fit into SM: 3 kinds of neutral fermion, the
- ⁷ counterparts to the charged fermions. Interacts with the weak force only.
- ⁸ • Summary of the experimental evidence for this picture: mainly, the discovery of
- ⁹ electron anti-neutrinos by Cowan and Reines, the muon neutrino by Lederman,
- ¹⁰ Schwartz, and Steinberger, and the tau neutrino by the DONUT Collaboration.
- ¹¹ Further critical experiments include the first measurement of a neutrino’s helicity
- ¹² by Goldhaber et al. as well as Danby et al.’s demonstration that ν_μ are distinct
- ¹³ from ν_e .
- ¹⁴ • More detailed description, via Feynman diagrams, of the two fundamental modes
- ¹⁵ of interaction by neutrinos with the weak force: charged- and neutral-current
- ¹⁶ interactions. A brief mention of the quantitative theory that underlies description:
- ¹⁷ Gashow, Salam, and Weinberg’s Electroweak Theory. This explains not only the
- ¹⁸ V–A structure of charged-current interactions, but also predicted accurately the
- ¹⁹ nature of neutral-current interactions. (Given space constraints, I see no reason
- ²⁰ to go into much of the details of the theory, or the many experimental tests of
- ²¹ its structure.)

$$\frac{d\sigma_{\nu_i}}{dT} = BLAH \quad (1.1)$$

²³ [4 pages]

1.2 Neutrino Oscillations and Neutrino Masses

1

1.2.1 The Evidence for Neutrino Oscillations

2

- Describe status quo ante of massless nature of neutrinos: BEH mechanism as exists cannot allow for neutrinos to have mass as only left-handed neutrinos have been observed.
3
4
5
- Furthermore, strong experimental limits on neutrino masses, from e.g. tritium-decay endpoint measurements by the KATRIN experiment and cosmological inferences from the CMB by the Planck satellite.
6
7
8
- But — then neutrino oscillations are observed over a variety of experiments and contexts. Summarise critical bits of evidence:
9
10
- Electron neutrino disappearance in solar neutrino experiments, including Ray Davis' Homestake experiment, the SAGE/GALLEX experiments, and SNO. For the latter, the comparison of charged-current and neutral-current modes of interaction was clear evidence of neutrino oscillations over other types of process (e.g. neutrino decay).
11
12
13
14
15
- Include in the above a brief description of Bahcall's Standard Solar Model.
16
- Muon neutrino disappearance in atmospheric and long-baseline accelerator neutrino experiments, such as Super-Kamiokande, T2K, and No ν a.
17
18
- A few further observations to note are: reactor electron anti-neutrino disappearance from both KamLAND and Daya Bay; tau neutrino appearance at the OPERA experiment; short-baseline neutrino anomaly within LSND and MiniBooNE (with recent contrary evidence from MicroBooNE).
19
20
21
22

1.2.2 The Phenomenology of Neutrino Oscillations

- Describe the current phenomenological model of 3-flavour neutrino oscillations that can explain all of this evidence: the PMNS mixing matrix.
- Describe also the MSW effect, which is critical for explaining solar neutrino oscillations.
- Show the formula for solar neutrino oscillations, given this MSW effect in both the Sun and Earth. Note the dependence of solar neutrino oscillations on only the “solar” oscillation parameters. This is all particularly useful for the solar analysis chapter.

$$P_{ee} \left(\tan 2\theta_{12}^M, \sin \theta_{13}^M, \Delta m_{21,M}^2 \right) = BLAH \quad (1.2)$$

[3 pages]

1.2.3 The Origins of Neutrino Mass

- Observed neutrino oscillations require at least two neutrino mass states to be non-zero. Given constraints of the current SM, two main ways of adding neutrino masses: a Dirac mass term (i.e. allowing for sterile neutrinos), and a Majorana mass term.
- For latter, briefly describe what a Majorana particle is, and how with the Seesaw Mechanism (just the simple Type 1 described in-text) one can not only get neutrino masses but also explain their lightness relative to the other massive SM particles. Note that there exist more elaborate versions of this theory.
- Furthermore, with reference to the Sakharov conditions, describe qualitatively how the Seesaw Mechanism also allows for possible leptogenesis/baryogenesis in the early Universe, and hence could explain its matter-antimatter asymmetry.

1.2 Neutrino Oscillations and Neutrino Masses

7

- Describe briefly the nuclear physics behind double-beta decay (i.e. why it can happen at all over just normal beta decay), and then how Majorana neutrinos allow for neutrinoless double beta decay, $0\nu\beta\beta$.
1
2
3
- Describe the experimental signature of $0\nu\beta\beta$: a spike of events of observed energy equal to the Q-value of the decay.
4
5
- Note Schecter-Valle Theorem ensures that any observation of $0\nu\beta\beta$ must be the result of neutrinos being Majorana. I.e. the Universe cannot conspire against us and have $0\nu\beta\beta$ without Majorana neutrinos.
6
7
8
- Very briefly note the current status of the search for $0\nu\beta\beta$, describing the main varieties of experimental setup seen, along with a nice canonical example of such an experiment and their best limit. In particular, the Germanium-crystal detectors such as GERDA, Xenon-TPC detectors like EXO-200, and large-scale liquid scintillators such as KamLAND-Zen.
9
10
11
12
13

[3 pages]

14

[CHAPTER TOTAL: 17 pages]

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Chapter 2

The SNO+ Detector

The light-soaked days are coming.

JOHN GREEN

2.1 Detector Geometry

The SNO+ detector is a large, multi-purpose neutrino detector built in the SNOLAB underground laboratory near Sudbury, Canada. Its main detector structure is taken from the Sudbury Neutrino Observatory (SNO) [1], which can be seen in Fig. 2.1. At the heart of the detector lies the main detector medium, which changes depending on the phase of the experiment — more on the specifics of this in Section 2.2. This medium is held within a 12 m diameter sphere known as the Acrylic Vessel (AV). The AV floats within a body of ultra-pure water (UPW), beyond which is a stainless steel support structure (PSUP) that holds 9362 inward-facing Photomultiplier Tubes (PMTs). It is these PMTs that detect the light generated from physics events that occur within the detector medium. The AV is kept in place relative to the PSUP through a series of ‘hold-up’ and ‘hold-down’ tensylon ropes. All of these components are suspended within a large cylindrical cavity also filled with UPW. 91 outward-looking PMTs (OWLs) are

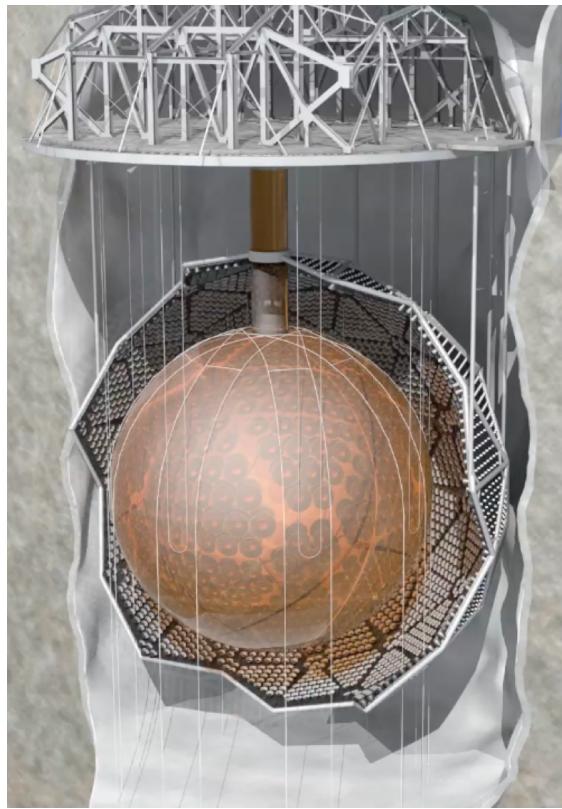


Fig. 2.1: 3D model of the SNO+ detector [2].

¹ also affixed to the outside of the PSUP, allowing for the effective vetoing of cosmic ray
² muons.

³ Directly above the detector is the Deck, within which all the detector electronics
⁴ are kept. Access within the AV for calibration tools and filling is possible only through
⁵ the acrylic ‘neck’ on top of the AV. Full details of the design of the current detector
⁶ can be found in [2].

⁷ 2.2 Experimental Phases

⁸ As mentioned earlier, SNO+ was designed to fulfil a number of physics goals over
⁹ multiple ‘phases’ of the detector’s lifetime. The phases are distinguished by the medium
¹⁰ that fills the AV. The first main phase (after a brief **Air Fill Phase** used only for

2.2 Experimental Phases**11**

detector commissioning) was that of the **Water Fill Phase**, with data taken between 1 May 2017 and July 2019. This was used to perform fundamental optical calibrations 2 of the detector [3], measurements of the solar neutrino flux [4], observation of neutrino 3 oscillations in reactor anti-neutrinos [5], and searches for nucleon decay [6, 7]. 4

After this, the detector was filled with 780 tonnes of liquid scintillator known as 5 linear alkylbenzene (LAB), mixed with the fluor 2,5-diphenyloxazole (PPO). More 6 information on the physics of scintillators can be found in Section 2.3.1. Filling of the 7 LABPPO cocktail had to be paused in March 2020 due to the COVID-19 pandemic, 8 leading to the detector having its bottom half still filled with UPW, and the top 9 half filled with LAB and PPO at 0.5 g/L. This impromptu phase became known as 10 the **Partial Fill**, and allowed for some creative analyses to be performed: an initial 11 neutrino oscillation analysis from reactor anti-neutrinos [8], as well as the first ever 12 observation of directionality in a high light yield scintillator [9, 10]. Eventually, filling 13 of the detector with liquid scintillator completed in May 2021. At that point, the 14 concentration of PPO in the detector was at 0.6 g/L, markedly below the target level 15 of 2.0 g/L. A further ‘PPO top-up’ campaign then proceeded, finishing in April 2022 16 with a final concentration of 2.2 g/L PPO. Thus began the **Scintillator Phase** of 17 the experiment, which continues on during the time of writing. The main goals for 18 this phase include a number of solar neutrino analyses (including the one described in 19 Chapter 6), a precision measurement of the neutrino oscillation parameter Δm_{21}^2 using 20 reactor anti-neutrinos [8], and further calibrations of the detector and its backgrounds. 21

Finally, in the near future the detector will be loaded with Tellurium for the 22 **Tellurium Phase**, allowing for the flagship analysis of the experiment to begin: 23 neutrinoless double beta decay. In order to load Te within the liquid scintillator 24 in a stable manner, a chemical loading process has been developed, as described 25 in [11]. The Te starts within Te(OH)_6 (telluric acid, otherwise known as TeA), which 26

¹ after purification will be reacted with 1,2-butanediol (BD) via heating and addition
² of N,N-Dimethyldodecylamine (DDA), which acts as a stabiliser. What results is
³ tellurium-loaded scintillator, TeLS.

⁴ Two further chemicals are planned to be added to the scintillator cocktail. The
⁵ antioxidant butylated hydroxytoluene (BHT) will be added to capture any free-radicals
⁶ within the liquid scintillator, hopefully preventing any oxidation reactions that could
⁷ lead to the ‘yellowing’ of the scintillator, a degradation of its optical properties. The
⁸ addition of BHT is not expected to impact the detector’s optics in any substantial
⁹ way. However, the other substance to also be added, 1,4-Bis(2-methylstyryl)benzene
¹⁰ (bisMSB), will impact the optics. BisMSB acts as a ‘wavelength-shifter’ which enables
¹¹ the scintillator cocktail to transmit light with a greater overall detection efficiency —
¹² more on the details of this in Section 2.3.1.

¹³ **2.3 Detecting and Recording an Event in SNO+:
14 A Journey**

¹⁵ To understand the SNO+ detector well, it is worth thinking about how the information
¹⁶ contained in a physics event, e.g. a solar neutrino interaction, gets observed. This
¹⁷ section follows the journey of such an event.

¹⁸ **2.3.1 Particle Interactions with Matter**

¹⁹ All observable physics events within the detector begin by the generation of some
²⁰ form of ionising radiation: α , β^\pm , γ , p or n . These can be created via numerous
²¹ processes, both exciting (e.g. $0\nu\beta\beta$ or interactions of neutrinos) and annoying (e.g.
²² decay of background radioisotopes): see Section 6.1.2 for some of them. Regardless of
²³ their origin, these particles begin propagating through the detector, and interacting

2.3 Detecting and Recording an Event in SNO+: A Journey

13

with the detector medium. A number of mechanisms then allow for the generation of optical-wavelength light as a result of these interactions.

Cherenkov Light Emission

Whenever a charged particle travels through a dielectric medium at speeds faster than the speed of light in that medium, light is generated from the ‘wake’ of induced dipoles. This is known as **Cherenkov light**, a process much akin to the ‘sonic boom’ that occurs when an object travels at supersonic speeds through a medium. This light emanates outwards in a cone along the direction of the charge’s travel; the angle of the cone θ_γ being purely a function of the speed of the charged particle relative to the speed of light in vacuum, β , and the refractive index of the medium $n(\omega)$ at a given frequency ω : $\cos \theta_\gamma(\omega) = \frac{1}{n(\omega)\beta}$. There is then a minimum speed necessary for Cherenkov light to be generated: $\beta_{\min}(\omega) = 1/n(\omega)$.

In addition to the characteristic cone shape of the light, the spectrum of the light generated is also distinctive. Igor Tamm and Ilya Frank determined the expected energy dE emitted per unit length travelled by the charged particle, dx , as [12]:

$$\frac{dE}{dx} = \frac{q^2}{c^2} \int_{\beta n(\omega) > 1} \omega \left(1 - \frac{1}{\beta^2 n^2(\omega)} \right) d\omega. \quad (2.1)$$

Here, q is the charge of the moving particle.

All SNO+ detection media allow Cherenkov light to be generated, as long as sufficiently high energy particles traverse it. In the water fill phase of the detector, Cherenkov light was the only means by which light could be generated. Light from Cherenkov emission can still be created in liquid scintillator, but it tends to be swamped by another form of light generation: scintillation.

1 Scintillation

2 For certain special classes of material, the excitation and ionisation of atomic electrons
 3 nearby a moving charged particle can lead to the generation of optical-wavelength light,
 4 in a process known as **scintillation** (often generally referred to as ‘luminescence’ or
 5 ‘fluorescence’). Although multiple varieties of scintillator exist, the one used in SNO+
 6 is that of an organic liquid scintillator. For such liquids, scintillation light is generated
 7 from the de-excitation of delocalised electrons within carbon–carbon ‘ π -bonds’ [13]. A
 8 major example of these π -bonds are found in benzene rings, which are present in LAB,
 9 PPO, and bisMSB.

10 Because of this delocalised structure, excited atomic π -electrons can stay in what
 11 is typically the first-excited state for somewhat longer than typical excited states:
 12 lifetimes of $\mathcal{O}(10^{-9}$ s) as opposed to $\mathcal{O}(10^{-12}$ s). This is what gives scintillation light
 13 its characteristic ‘slow’ response relative to the instantaneous light generated by the
 14 Cherenkov process. Moreover, decays from this state can emit light typically in the
 15 optical-wavelength range. In addition to excited electrons, ionised electrons can also
 16 recombine — that is, re-enter atomic orbitals — into various excited states, and then
 17 decay back to the ground state, also allowing for the possibility of scintillation light to
 18 be generated.

19 Because of atomic spin selection rules [13], scintillation light typically has, at the
 20 very least, a ‘fast’ and ‘slow’ time component. In SNO+, we currently model emission of
 21 scintillation light from LAB with 4 time components, following the timing distribution
 22 $f(t)$ given by:

$$23 \quad f(t) = \sum_i A_i \left(\frac{e^{-t/\tau_i} - e^{-t/\tau_{\text{rise}}}}{\tau_i - \tau_{\text{rise}}} \right), \quad t > 0. \quad (2.2)$$

24 Here, A_i and τ_i correspond to the fraction of light emitted and decay constant for each
 25 component respectively, and τ_{rise} is a common rise time. The current fitted values for

Component	A_i	τ_i [ns]
1	0.665	7.35
2	0.218	5.45
3	0.083	117.5
4	0.0346	425
Rise	–	0.8

Table 2.1: Current values used to model scintillator emission from electrons in 2.2 g/L LABPPO [14, 15].

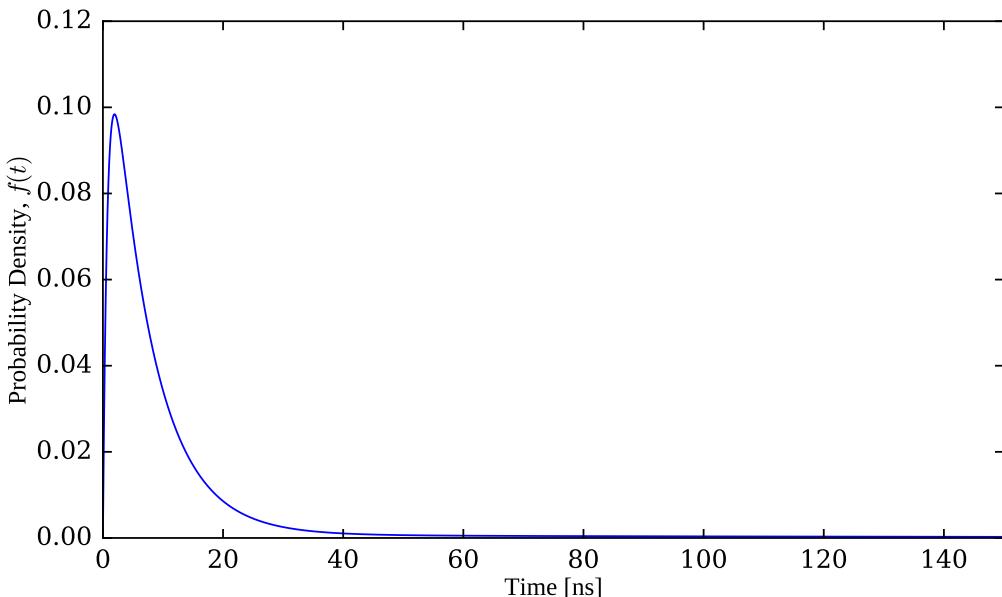


Fig. 2.2: The current modelled distribution for the emission timing of the SNO+ scintillator, for electron tracks [14, 15].

these parameters for the emission from electron tracks can be seen in Table 2.1. A plot of $f(t)$ with these values can be seen in Fig. 2.2.

When using just a single scintillating compound, the very same energy levels that can generate scintillation light are those that can absorb it. This can be a problem for large-scale detectors like SNO+, which depend on scintillation light being unobstructed in its path to the PMTs. Conveniently, this problem can be addressed with the addition of another scintillating component, known (somewhat confusingly) as the primary fluor. In SNO+, this is the PPO added to the LAB.

When an LAB molecule is excited, that energy can be transferred to a PPO molecule through what is known as a ‘non-radiative transfer’. In short, this transfer of energy occurs not through the emission and absorption of optical photons, but through the coupling of the molecules’ electric dipoles.¹ When the now-excited PPO molecule de-excites to emit scintillation light, the different molecular structure it has generates a different emission spectrum to that of LAB. These longer wavelengths of light are no longer able to be absorbed by the LAB, allowing for a scintillator with less optical absorption.

Adding in one additional component doesn’t have to be the end, either. In SNO+ we plan on adding in the compound BisMSB to the scintillator cocktail. This is a ‘wavelength-shifter’: scintillation light at short wavelengths is absorbed, and then re-emitted at longer wavelengths, where the detection efficiency of the PMTs is greater. More on the properties of the PMTs in SNO+ can be found in Section 2.3.3. The net effect of the three scintillating components within SNO+ can be seen in Fig. 2.3. Note how, as energy is transferred from one scintillation component to another, the wavelength of light emitted gets necessarily longer as energy is lost to heat.

The light yield of a scintillator, i.e. the amount of optical photons generated per unit of energy deposited into the scintillator, is a function not just of the scintillator but also the incident particle’s ionisation strength. In particular, α particles are far more effective at exciting and ionising nearby atoms, and so can deposit far more of its energy into the scintillator per unit volume. However, the strength of this ionisation for α s can actually become at detriment to the generation of scintillation light. Empirically, scintillators follow to first order Birks’ Law for their scintillation light yield [18]:

$$\frac{dL}{dx} = S \frac{\frac{dE}{dx}}{1 + k_{\text{Birks}} \frac{dE}{dx}}, \quad (2.3)$$

¹To be pedantic, photons are still transferred in this energy exchange, but they are virtual instead of real.

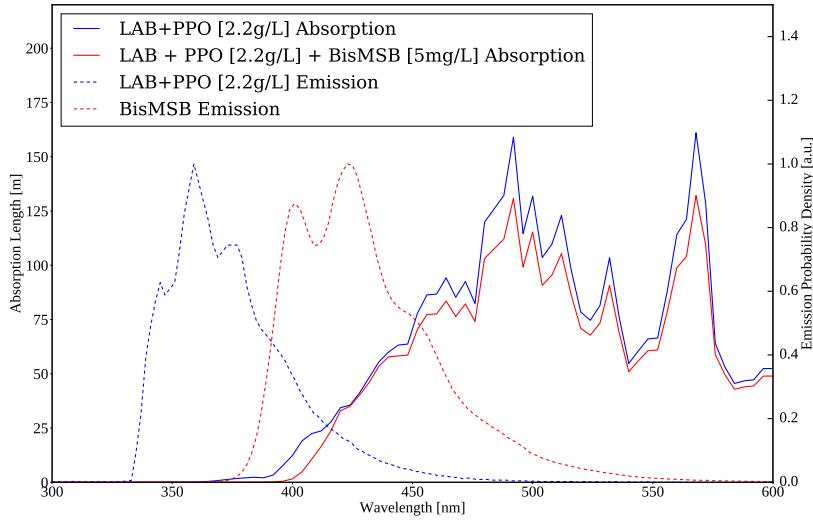


Fig. 2.3: Comparison of the SNO+ liquid scintillator cocktail’s absorption and emission properties, with/without BisMSB [16, 17].

where $\frac{dL}{dx}$ is the number of photons emitted per unit track length, $\frac{dE}{dx}$ is the energy loss of the incident particle per unit track length, S is the scintillator’s characteristic light yield constant, and k_{Birks} is the scintillator’s “Birks’ Constant”². For minimum-ionising particles such as a 6 MeV electron, the denominator of this equation is close to 1, and so the amount of scintillation light generated is just $\frac{dL}{dx} = S \cdot \frac{dE}{dx}$. However, for α -particles generated in radioactive decays, this denominator can become substantial, and in the limiting case we have merely $\frac{dL}{dx} = \frac{S}{k_{\text{Birks}}}$. In the current phase of SNO+, S and k_{Birks} are measured to be 14,000 γ/MeV and 0.077 mm MeV⁻¹, respectively [19].

2.3.2 Optical Processes

Once optical-wavelength photons have been created within the detector, various processes can then occur that can hinder its path towards a PMT, and therefore modify the observed signal. This subsection covers the main optical processes, with a fo-

²Birks’ Constant is often just written as k_B , but this is easily confused with the far-better known Boltzmann Constant, which is completely different!

¹ cus on Rayleigh scattering, as an understanding of this phenomenon is critical for
² Chapters 3–5.

³ Rayleigh Scattering

⁴ Optical scattering is the general process of how light is scattered by particles within a
⁵ medium. This is fundamentally an electrodynamical process: an electromagnetic wave
⁶ is incident on the set of particles within the medium, which induces these particles to
⁷ oscillate within the field, and therefore generating their own electromagnetic radiation
⁸ in response. Usually, this ‘scattered’ radiation has the same frequency as that of the
⁹ incident radiation, and therefore the scattering is said to be *elastic*. It is possible under
¹⁰ certain circumstances for this scattered radiation to be of a longer wavelength than
¹¹ the incident radiation: in which case, energy was absorbed by the particles and so the
¹² scattering was *inelastic*. However, this latter type of scattering, also known as Raman
¹³ scattering, is not relevant for SNO+.

¹⁴ The general solution to elastic optical scattering was first described by Gustav
¹⁵ Mie [20] and Ludvig Lorenz [21] in what is now known as *Mie Theory*. In this theory, it
¹⁶ is assumed that a plane wave of wavelength λ is incident on a dielectric sphere of radius
¹⁷ a . While the general solution to the problem of Mie scattering is somewhat complicated
¹⁸ (if tractable), in certain regimes one can make further simplifying assumptions that
¹⁹ substantially reduce the complexity of the result. In particular, if one assumes that
²⁰ the size of the particle is much smaller than the wavelength of light, and that any
²¹ induced dipole moment can actually be established in the time window allowed by
²² the oscillation period of the electromagnetic field [22], then one can obtain *Rayleigh*
²³ *scattering*. This simpler case is so-called because of its initial formulation by Lord
²⁴ Rayleigh [23].

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One can show that the differential cross-section associated with Rayleigh scattering of unpolarised light off of a single particle, $\frac{d^2\sigma_{\text{Ray}}}{d\theta d\phi}(\theta, \phi)$, is given by [24]:

$$\frac{d^2\sigma_{\text{Ray}}}{d\theta d\phi}(\theta, \phi) = \frac{8\pi a^6}{\lambda^4} \left(\frac{n_{\text{par}}^2 - 1}{n_{\text{par}}^2 + 2} \right)^2 (1 + \cos^2 \theta). \quad (2.4)$$

Here, θ and ϕ correspond respectively to the polar and azimuthal angles of the scattered waves, and n_{par} is the refractive index of the scattering particle. Most important to notice about this equation is that the cross-section follows a strong $1/\lambda^4$ dependence, meaning that short wavelengths of light will be scattered to far greater extents than that of longer wavelengths. Secondly, the light is not scattered isotropically, but according to a $1 + \cos^2 \theta$ dependence. This means that most light is either scattered directly forwards or backwards (known as *backscattering*), and little gets scattered orthogonally to the direction of the incident light. This is useful when it comes to trying to measure scattering in the SNO+ detector, as it provides a handle upon which to distinguish scattered light from isotropically-emitted scintillation light.

Of course, we care about the scattering that occurs within an entire bulk medium, not just the scattering off of a single molecule. From a macroscopic perspective, the key quantity of interest is a material's *Rayleigh scattering length*, l_{Ray} : the mean distance a photon is expected to travel before Rayleigh scattering. One can show that, assuming the above differential scattering cross-section, the Rayleigh scattering length is given by [25]:

$$l_{\text{Ray}} = \left[\frac{16\pi}{3} R \right]^{-1}. \quad (2.5)$$

R is the *Rayleigh ratio*, $R = \frac{1}{V} \frac{d^2\sigma_{\text{Ray}}(90^\circ)}{d\theta d\phi}$, equivalent to the power of the scattered light per unit volume of the scattering medium per unit incident intensity at $\theta = 90^\circ$.

This can lead to a few changes to Rayleigh scattering that are worth noting. Firstly, unlike for a single particle, the electric polarisability of a material can be

¹ *anisotropic.* Anisotropic materials have a modified angular dependence on their
² differential cross-section, governed by the *depolarisation ratio*, δ . In particular, the
³ $(1 + \cos^2 \theta)$ dependence becomes $(1 + \frac{1-\delta}{1+\delta} \cos^2 \theta)$. For isotropic materials, $\delta = 0$, and
⁴ so the angular dependence reduces to the original form.

⁵ Secondly, the above model has been shown to be insufficient to describe liquids
⁶ or solids [26], because of the non-negligible strength of their inter-molecular forces.
⁷ Fortunately, Einstein [27], Smoluchowski [28], and Cabannes [29] developed a theory
⁸ for describing how photons can scatter off of the local charge density fluctuations that
⁹ naturally are present in a medium because of the thermal motion of molecules. The
¹⁰ theory shows that the Rayleigh ratio of a medium is related to the medium's dielectric
¹¹ constant, ε , by:

$$R = \frac{\pi^2}{2\lambda^4} \left[\rho \left(\frac{\partial \varepsilon}{\partial \rho} \right)_T \right]^2 k_B T \kappa_T \frac{6 + 6\delta}{6 - 7\delta}, \quad (2.6)$$

¹² where ρ is the density of the medium, $\left(\frac{\partial \varepsilon}{\partial \rho} \right)_T$ is the partial derivative of the dielectric
¹³ constant with respect to a changing density assuming a constant temperature T , k_B
¹⁴ is the Boltzmann Constant, and κ_T is the medium's isothermal compressibility. This
¹⁵ latter quantity is given by the rate of change of volume given a changing pressure of
¹⁶ the medium, all at a constant temperature.
¹⁷

¹⁸ Furthermore, the Eykman Equation [25, 30] has been shown to be an effective
¹⁹ empirical formula relating how ε is impacted by density fluctuations to the medium's
²⁰ refractive index, n_{med} :

$$\rho \left(\frac{\partial \varepsilon}{\partial \rho} \right)_T = \frac{(n_{\text{med}}^2 - 1)(2n_{\text{med}}^2 + 0.8n_{\text{med}})}{n_{\text{med}}^2 + 0.8n_{\text{med}} + 1}. \quad (2.7)$$

²¹ This leads to a final formula for the Rayleigh scattering length:

$$l_{\text{Ray}} = \left[\frac{8\pi^3}{3\lambda^4} \left(\frac{(n_{\text{med}}^2 - 1)(2n_{\text{med}}^2 + 0.8n_{\text{med}})}{n_{\text{med}}^2 + 0.8n_{\text{med}} + 1} \right)^2 k_B T \kappa_T \frac{6 + 3\delta}{6 - 7\delta} \right]^{-1}. \quad (2.8)$$

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Discussions of the scattering lengths currently assumed within SNO+’s optical model for UPW and LABPPO can be found within the theses of Krishanu Majumdar [31] and Esther Turner [32]. In particular, whilst the scattering length of the UPW in the water phase was measured by Esther, major systematics in the measurement remained. Measurements of the scattering lengths in scintillator are a major focus of Chapters 3–5.

Absorption and Re-emission

In addition to scattering, an optical medium is also able to absorb light that propagates through it. For a given medium, the *absorption length* l_{abs} is analogous to l_{Ray} described above, and is typically strongly a function of wavelength. For most materials, absorbed light is forever lost, converted into heat. However, for the special case of scintillators, re-emission of absorbed light is possible: this is because of the physics described in Section 2.3.1.

Because both scattering and absorption impede a photon’s ability to propagate through a medium directly, it is often possible to measure their combined impact through what is known as the attenuation/extinction length, l_{ext} :

$$\frac{1}{l_{\text{ext}}} = \frac{1}{l_{\text{abs}}} + \frac{1}{l_{\text{Ray}}}. \quad (2.9)$$

In the water phase, the ‘Laserball’ calibration system was used to measure various optical properties of the detector, including the extinction lengths of the UPW and acrylic as a function of wavelength [3]. Using the water phase scattering measurements made by Esther, Eq. 2.9 allowed for the estimation of the absorption lengths of these two materials, shown in Figure 2.4. Measurements of the extinction length in the scintillator phase is discussed in detail in Chapter 5.

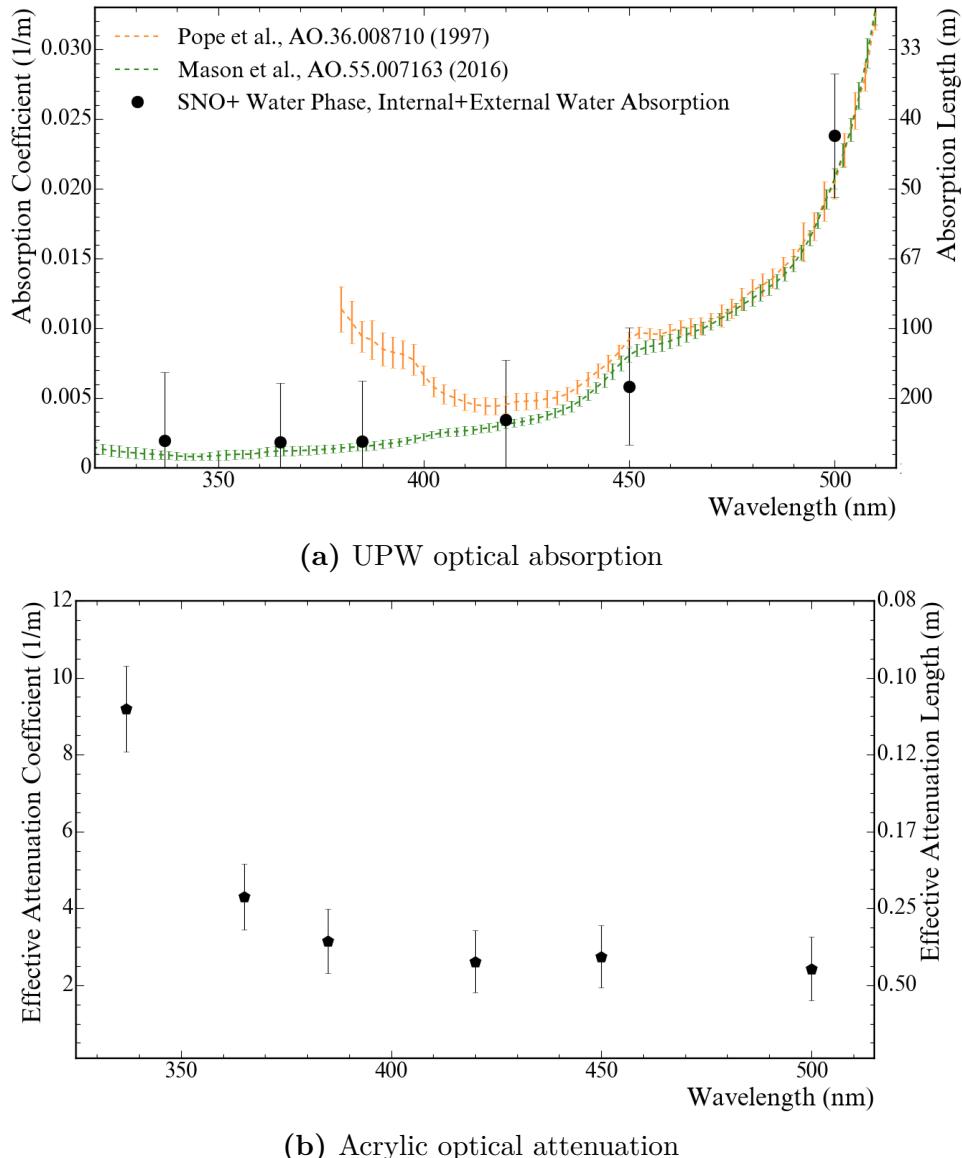


Fig. 2.4: Measured properties of the UPW and acrylic in the water phase, from [3].

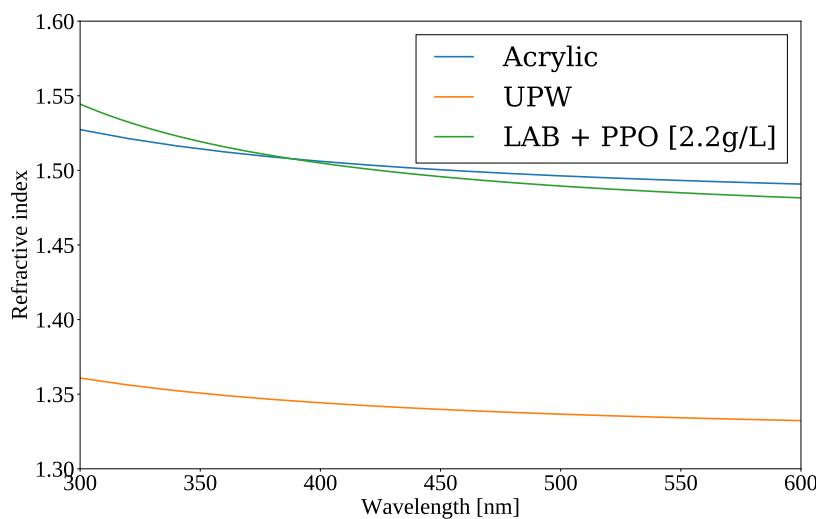


Fig. 2.5: Refractive indices of acrylic, UPW, and LABPPO as a function of wavelength [3, 33, 34].

Surface reflection and refraction

When light travels through the boundary of one medium to another, both reflection and refraction can be possible, depending on the relative refractive indices of the two media. The refractive indices of the UPW, acrylic, and LABPPO are shown as a function of wavelength in Figure 2.5. Note that, for most optical wavelengths, LABPPO has a very close refractive index to acrylic, whereas UPW is somewhat farther away. By consequence, negligible refraction is expected in most cases for light travelling between the liquid scintillator and the acrylic; however, substantial refraction is possible for light travelling between acrylic and UPW. Because of this, isotropically-emitting point-like physics events within the AV that are close enough to the acrylic will have some of their light undergo Total Internal Reflection (TIR) at the AV, reflecting back into the AV instead of continuing outward into the outer water.

Even when not undergoing TIR, some light at a boundary can still reflect. The fraction of light that reflects is known as the *reflectance* R , compared to that which is able to transmit through the boundary, the *transmittance* $T = 1 - R$. The *Fresnel*

¹ Equations determine the reflectance of an interface [35]:

$$\text{where } R_s = \left| \frac{n_1 \cos \theta_i - n_2 \cos \theta_t}{n_1 \cos \theta_i + n_2 \cos \theta_t} \right|^2, R_p = \left| \frac{n_1 \cos \theta_t - n_2 \cos \theta_i}{n_1 \cos \theta_t + n_2 \cos \theta_i} \right|^2, \quad (2.10)$$

³ where R_s and R_p are the reflectances of *s*- and *p*-polarised light, n_1 and n_2 are the
⁴ refractive indices of the first and second optical media, and θ_i and θ_t are the angles of
⁵ incidence and refraction, respectively. For SNO+, we are only interested in unpolarised
⁶ light, so the total reflectance $R = (R_s + R_p) / 2$.

⁷ 2.3.3 Detection by PMTs

⁸ The final step for photons in our journey is detection by a PMT. Almost all PMTs in
⁹ SNO+ are of the Hamamatsu R1408 design [1]. These PMTs within SNO+ are housed
¹⁰ within a reflecting cone known as a ‘concentrator’. The combined PMT–concentrator
¹¹ ‘bucket’, shown in Fig. 2.6a, is designed to maximise the collection efficiency of light
¹² emanating from within the AV, whilst minimising the collection efficiency of light
¹³ outside the AV [36]. The so-called ‘angular response’ of the PMT buckets has been
¹⁴ measured in both SNO and SNO+ using the Laserball, which describes the relative
¹⁵ collection efficiency as a function of the polar angle of the incident light ray relative
¹⁶ to the direction in which the PMT bucket points. The results of this can be seen in
¹⁷ Fig. 2.6b.

¹⁸ Once a photon is incident on the PMT’s photocathode, it is possible for that photon
¹⁹ to be absorbed and generate a photoelectron. The probability of this happening is
²⁰ governed by the photocathode’s Quantum Efficiency (QE) at the photon’s wavelength,
²¹ as well as the collection efficiency of a photoelectron onto the first dynode of the
²² PMT. The combined measured efficiency of PMTs tested *ex-situ* for SNO can be
²³ seen in Fig 2.7. Once this photoelectron has been created, the dynodes within the
²⁴ PMT generate a cascade of electrons that allow for an observable voltage signal to

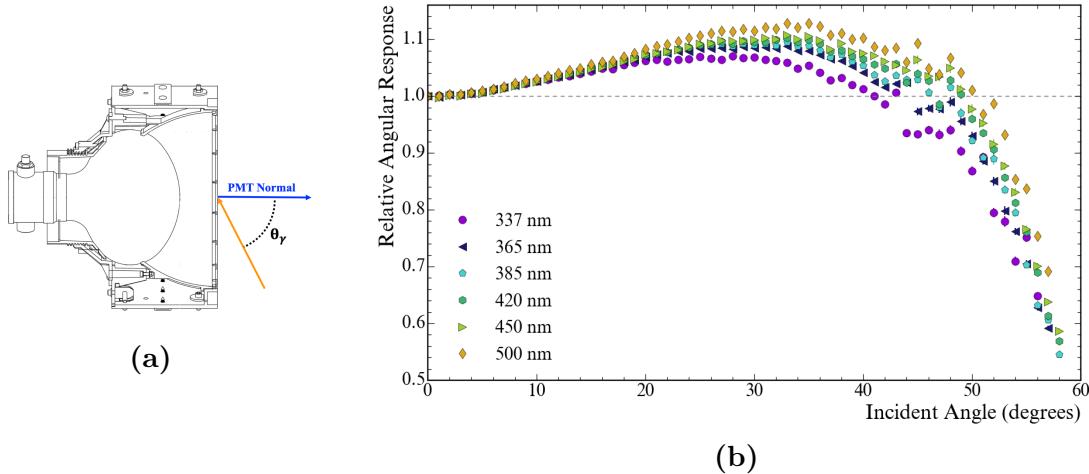


Fig. 2.6: (a): Diagram of the PMT and concentrator ‘bucket’ used within SNO+, showing also the definition of the incidence angle. (b): Plot of the measured relative angular response of the PMTs in SNO+, as a function of both incidence angle and wavelength. Both figures taken from [3].

be produced. The dynamics of this cascade are such that there is a natural spread of possible times between the creation of a photoelectron and the generation of the voltage signal pulse in the PMT’s wire. This is known as the ‘Transit Time Spread’ (TTS) of the PMTs: for SNO+, the RMS timing resolution of the TTS for the R1408-type PMTs is 1.7 ns [1].

Finally, if multiple photons generate photoelectrons on the same PMT close enough in time, the detector’s front-end electronics (described in Section 2.3.4) is unable to tell. However, the amount of charge generated increases in proportion to the number of photoelectrons (npe). Much like with the transit time, the strength of the signal observed by the PMT is governed by a distribution, a function of the npe generated. Examples of these distributions can be seen in Fig. 2.8. The relatively large widths of these charge distributions precludes the ability to straightforwardly determine the npe purely from charge when the npe is small. To work around this, various techniques can be employed to try and estimate the npe in a given PMT — an example of one such method can be seen in Section 4.2.

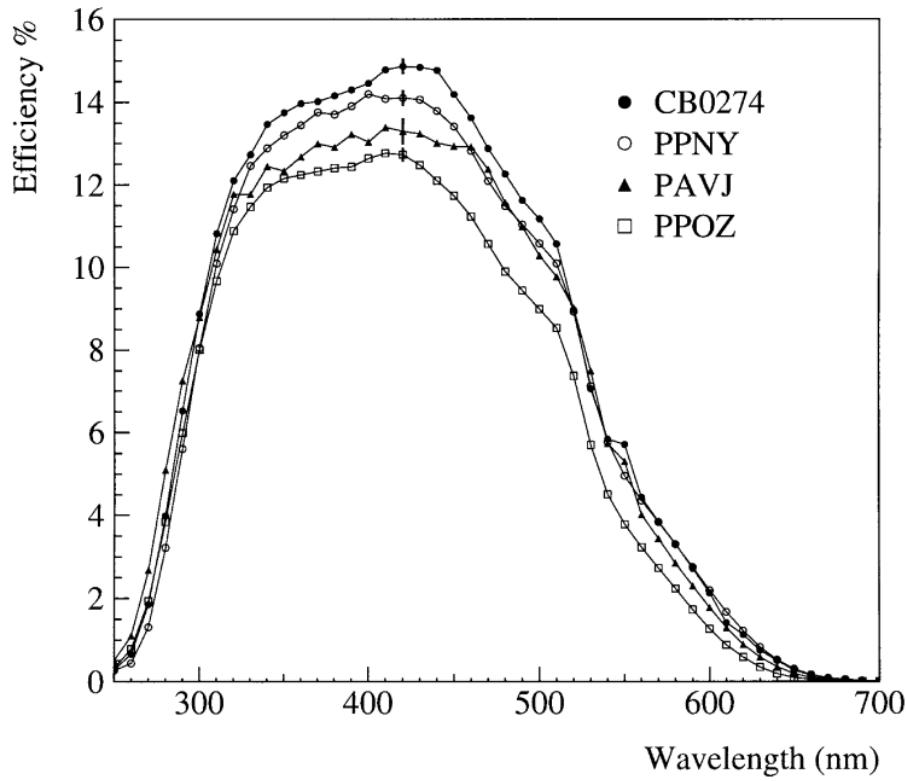


Fig. 2.7: Efficiencies of the R1408-type PMTs used as standard within SNO+ [37].

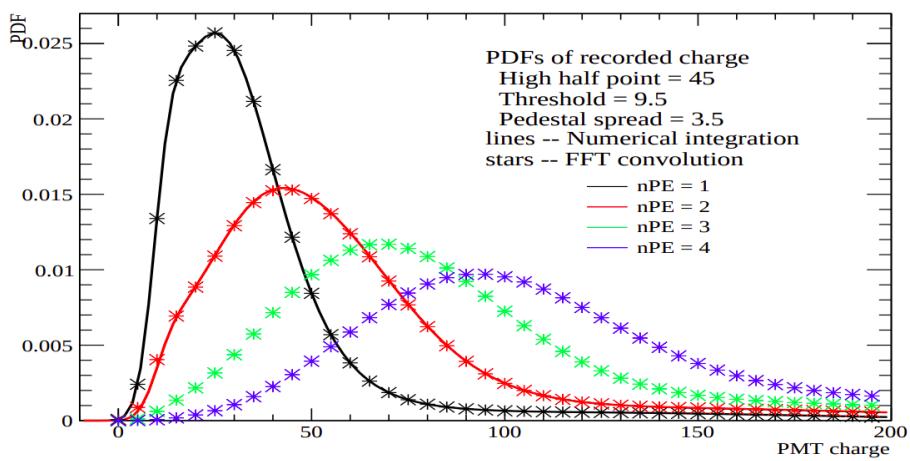


Fig. 2.8: Example charge spectra for a PMT as a function of the true npe generated [38].

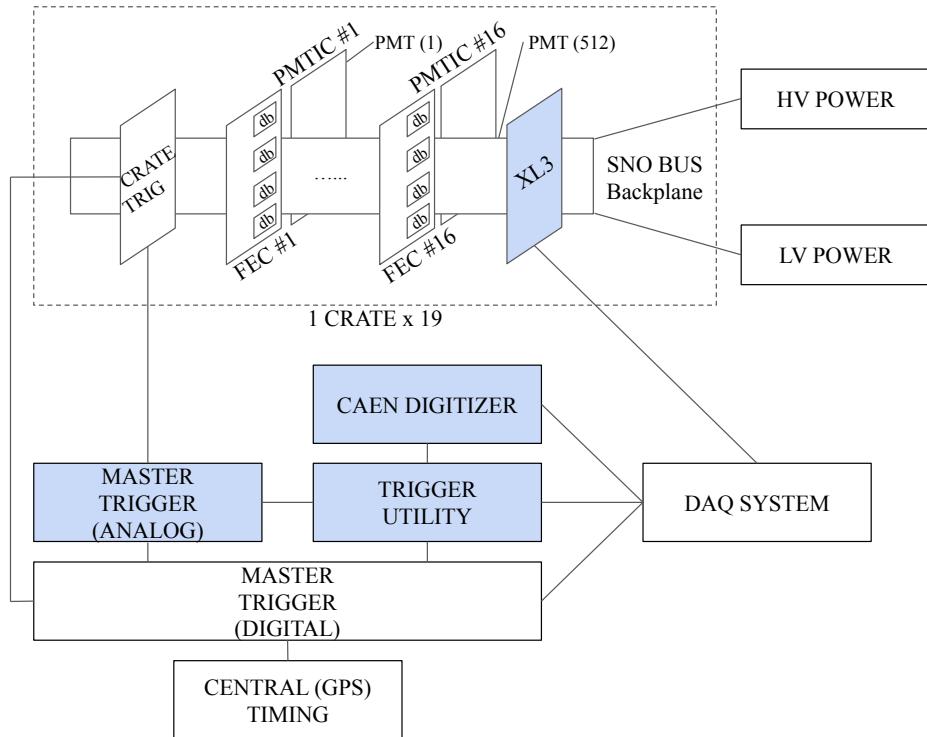


Fig. 2.9: Schematic of the front-end electronics used for data acquisition and triggering in SNO+, taken from and discussed in [2].

2.3.4 Data Acquisition and Triggering

Once a signal is generated in the cable attached to a PMT, it travels along up to the front-end electronics on the deck above the detector. The job of these electronics, known as the data acquisition (DAQ) and triggering system, is to convert raw electronic signals from the PMTs into recorded digital ‘events’ that can be used for analysis. A schematic showing the setup of the electronics is shown in Fig 2.9, with full details in [2].

A signal passes first through the PMT Interface Card (PMTIC), which then sends it through to one of the Daughter Boards (DBs) which are stored on Front-End Cards (FECs) within one of 19 electronic crates on deck. The DBs determine if the analogue signal along a given PMT channel has crossed a pre-defined charge threshold, at which

¹ point we say that we have detected a ‘hit’ on that PMT’s channel. When this occurs,
² the DB performs a set of important actions:

- ³ 1. Begins a timer for that channel, in the form of a Time-to-Amplitude Converter
⁴ (TAC). TAC also corresponds to the resulting quantity being measured.
- ⁵ 2. Begins integrating the total charge signal for that channel in three ways, known
⁶ as QHS, QHL, and QLX. These correspond to using different integration times
⁷ and gain settings.
- ⁸ 3. Generates trigger pulses for that channel for each available trigger type. Three
⁹ main trigger signals are the ‘N20’ (a square pulse for 20 ns), ‘N100’ (a square
¹⁰ pulse for 100 ns), and ‘ESUMHI’ (a pulse copying the shape of the voltage signal
¹¹ for that channel). The reason for these names will be explained shortly.

¹² Whilst the TAC, QHS, QHL, and QLX are being calculated, the trigger signals
¹³ from each channel are sent over to the Crate’s Trigger Card (CTC), where the signals
¹⁴ are then summed for each trigger type. These crate-level trigger signals are then sent
¹⁵ over to the 7 detector-level Analogue Master Trigger Cards (MTC/A+), which further
¹⁶ sum the signals by trigger type from all the crates in the detector. It is at this point of
¹⁷ the process where the names of the trigger types becomes clear: the combined N20
¹⁸ and N100 signals are proportional to the total number of hit PMTs within a 20 ns and
¹⁹ 100 ns time window, respectively, whilst the total ESUMHI signal corresponds to the
²⁰ total charge seen over all the PMTs.

²¹ If these trigger signals go above certain pre-defined thresholds, then a signal is
²² sent for that trigger type to the Digital Master Trigger Card (MTC/D). The MTC/D
²³ receives all trigger signals from the detector, and if a given trigger type has been
²⁴ ‘masked in’ the Card will generate a Global Trigger (GT) for the detector according to
²⁵ its 50 MHz clock. Under certain circumstances, such as calibrations, a trigger signal can

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be generated externally and asynchronous to the MTC/D: these are ‘EXTA’ triggers.
1
Any such EXTA trigger signal is first handled by an electronics box named ‘TUBii’
2
(Trigger Utility Board Mark ii), before then being passed onto the MTC/D. This
3
becomes relevant when discussing the calibration electronics described in Chapter 3.
4

Once a GT signal is generated, it is then sent back to all the CTCs, which then
5
orders the integration of time and charge to be stopped on all channels for that crate.
6
The time and charge information that has been temporarily stored on each channel’s
7
CMOS chip on the FEC is then sent to the crate’s ‘XL3’ card, which packages the
8
crate’s raw information via ethernet over to a set of computers. The trigger signals for
9
a triggered event are also digitised by a CAEN brand Analogue-to-Digital Converter
10
(ADC), and sent to the same DAQ computers. The total window of time in which
11
data is gathered from one GT signal is 400 ns, with data from up to 180 ns before and
12
220 ns after the GT has arrived. There is then a necessary ‘dead time’ of 420 ns after a
13
given GT has been given in which no further GTs can be made.
14

Finally, the raw data from the crates and trigger system arrives in a set of computers,
15
which organise all of this into an individually-packaged ‘event’, stored on disk in the
16
'Zebra Database' (ZDAB) format. A given built event contains the TAC and charge
17
information from each hit PMT, the CAEN digitised waveforms, a unique identifying
18
number for that triggered event (the GTID), as well as the times from both the
19
MTC/D’s 50 MHz clock and a GPS-calibrated 10 MHz clock. The former time is used
20
for measuring relative times between events whilst the latter is used for knowing the
21
time of day of an event: both are used in Chapter 6.
22

2.3.5 Operation of the Detector

Control of the detector’s DAQ system is handled through a custom-built GUI known
24
as ORCA [39]. This program allows operators of the detector to modify settings in the
25

¹ detector electronics at both a low- and high-level. It also allows operators to monitor
² the current status of the detector, such as voltage and current levels within each crate.

³ ORCA allows for the detector to have its data split into ‘runs’ of different types.

⁴ For the majority of the time, the detector is run in the ‘Physics’ mode, with individual
⁵ runs split into 1 hour periods. It is this data that is used for almost all high-level
⁶ physics analyses, such as the one described in Chapter 6. Other modes include ones
⁷ for detector maintenance, as well as for calibrations of various kinds. During certain
⁸ calibration runs, data can be further split into ‘subruns’ where necessary. Operation of
⁹ specific calibration sources, including the one described in Chapter 3, can be performed
¹⁰ through the ORCA GUI.

¹¹ **2.4 Calibrations and Detector Modelling**

¹² Once the raw data from triggered events has been stored in files, certain extra steps
¹³ must be taken before effective analysis of that data can be achieved. This section
¹⁴ covers those steps.

¹⁵ **2.4.1 Detector Monitoring**

¹⁶ No data taken from the detector can reasonably be used for analysis unless its quality
¹⁷ has been approved. This is done in a number of ways on SNO+. Firstly, a number of
¹⁸ automated systems monitor all aspects of the detector, including voltage levels in the
¹⁹ crates, trigger rates, as well as ‘slower’ quantities such as the tensions on the ropes
²⁰ holding the AV in place. Problems in any of these measured parameters trigger an
²¹ automatic alarm system, which notifies a human detector operator. A human detector
²² operator monitors the detector 24/7 whilst the detector is live.

In addition to systems that monitor whether anything has gone wrong, information about the state of the detector during each run is stored in a database known as RATDB. This information includes, amongst other things, a recording of which PMT channels have actually been raised to high voltage for that run, as well as any channels/cards/crates that have been flagged for having a known poor data quality (e.g. being overly noisy).

2.4.2 Electronic and PMT Calibrations

The lowest level of calibrations performed in SNO+ are the Electronic and PMT Calibrations: ECAs and PCAs, respectively. These calibrations convert the raw time and charge values recorded by the DAQ into quantities that can actually be used in analysis.

During an ECA, two main quantities are measured. Firstly, because of noise the integrated charge measured on each channel is offset by some amount. This offset, known as the ‘pedestal’, is recorded for each channel. The other quantity is the ‘time slope’ for each channel, which allows one to convert from the ADC TAC counts into an uncalibrated hit time of that channel’s PMT in ns. Both of these quantities are measured by manually sending external signals to channels in the crates, forcing them to start measuring TAC and charge even though no PMTs were actually hit. Running ECAs also enable us to spot any channels with unusual behaviour, so that they are not used during analysis.

Using ECAs alone is not enough to have fully-calibrated time and charge data. The lengths of cables between PMTs and PMTICs are all slightly different, leading to differences in the so-called ‘cable delay’ of each channel. This means that two PMTs that have a photoelectron generated at the same time can generate slightly different TAC values. Furthermore, because the start time of the TAC is determined by when

¹ the channel’s signal goes above a constant threshold, if a signal is very large (e.g. when
² numerous photoelectrons have been generated on one PMT) then the start time of the
³ TAC will be systematically earlier. This is known as the ‘time walk’. Both of these
⁴ quantities get measured during PCAs.

⁵ PCAs can be performed by either the Laserball or by the TELLIE calibration
⁶ system. The latter is a series of 92 optical fibres attached at various points of the
⁷ PSUP, through which optical-wavelength light can be fired from LEDs. TELLIE is the
⁸ Timing subsystem of the ELLIE calibration system: the Embedded LED/Laser Light
⁹ Injection Entity. The other two subsystems, AMELLIE and SMELLIE, are introduced
¹⁰ in Section 2.4.3. For both the Laserball and TELLIE calibration systems, the cable
¹¹ delay and time walk are measured by firing light from the source at a known time, and
¹² observing when the signal arrives in each PMT channel.

¹³ On top of calibrating the PMT hit times, PCAs also further calibrate the charge
¹⁴ information. In particular, the charge spectrum generated by a single photoelectron is
¹⁵ determined for each channel. This allows us to convert the pedestal-corrected charge
¹⁶ ADC counts into a number of photoelectrons.

¹⁷ Using the data gathered from both ECAs and PCAs, the raw data stored in
¹⁸ ZDABs is processed into a new file format known as RATDS files. These files contain all
¹⁹ the information of an event, but now the timing and charge information have been
²⁰ calibrated. It is this file type used in the optical calibration work of Chapters 3–5.

²¹ 2.4.3 Energy and Optical Calibrations

²² The next stage of calibrating the detector is modelling its optical properties. These
²³ properties include all the processes covered in Section 2.3.2, such as scintillator emission,
²⁴ optical absorption, re-emission, and Rayleigh scattering. This is crucial, as it allows us to

reconstruct information about events within the detector: more on event reconstruction
shortly.

In addition to deployments of the Laserball (discussed in Section 2.3.2), two further
calibration sources are used in SNO+ to measure properties of light propagation:
AMELLIE and SMELLIE. These are the ‘Attenuation Module’ and ‘Scattering Module’
for the ELLIE calibration system. Like TELLIE, AMELLIE and SMELLIE consist of
optical light sources that shine through optical fibres into the detector. The former
uses LEDs from TELLIE, whilst the latter uses optical wavelength lasers. Despite the
names both subsystems are similar enough that they are both capable of measuring
attenuation and scattering within the detector. More details about the SMELLIE
hardware can be read in Chapter 3.

Another critical component of the detector to calibrate well is the energy response:
given a specific amount of energy deposited in the water/scintillator, how many hits are
observed? For this, a number of radioactive sources are used at a variety of energies. In
the scintillator phase, there are three main sources. The first is an americium-beryllium
(AmBe) source inherited from SNO [40], which contains ^{241}Am that α -decays, which
can be captured by the ^9Be within the source. This capture leads to the emission of a
neutron as well as production of a ^{12}C nucleus, which 60% of the time is in an excited
state. When this excited state decays, a 4.4 MeV γ is emitted promptly. Eventually the
neutron is captured by hydrogen in the detector, leading to a characteristic 2.2 MeV
delayed γ being generated [2]. Both the prompt and delayed peak energies from the
AmBe source can be used for energy calibration. In addition to this, one can calibrate
the neutron detection efficiency with the AmBe source [40], which is important for the
analysis of antineutrino IBD events.

Another deployable radioactive source is the ^{16}N source, also originally used for
SNO [41]. The ^{16}N isotope β -decays to ^{16}O , with a distinctive 6.1 MeV γ also being

¹ generated 66% of the time. It is the γ that can make it out to the detector, whilst the
² β can be tagged by a block of scintillator and PMT held within the source container.

³ Neither of the above two calibration sources have been deployed internally within
⁴ the scintillator during the scintillator phase because of concerns over contaminating
⁵ the scintillator. However, there have been a number of deployments in the external
⁶ water. Alongside this, a different kind of radioactive source has been used during the
⁷ scintillator phase for energy calibration: the existing radioactive background spectra
⁸ within the detector. Backgrounds such as ^{14}C , ^{210}Po , $^{214}\text{BiPo}$, and ^{208}Tl all have
⁹ distinctive peaks in the energy spectrum of SNO+ during the scintillator phase, and
¹⁰ can be used to calibrate the scintillator’s energy response. Using $^{214}\text{BiPo}$ events in
¹¹ particular has been used for energy scale calibration with the solar oscillation analysis,
¹² as discussed in Section 6.2.3.

¹³ 2.4.4 Event Reconstruction

¹⁴ Once the detector has been calibrated, event ‘reconstruction’ becomes possible. This is
¹⁵ the process of deriving high-level physics quantities about a triggered event within the
¹⁶ detector, based upon the calibrated hit information. In SNO+, our base assumption in
¹⁷ most cases of event reconstruction is that a triggered event was due to a single point-like
¹⁸ electron track. Reconstructing an event involves running a number of algorithms, which
¹⁹ in the scintillator phase are together called the **ScintFitter**.

²⁰ The first critical pieces of information that gets determined by **ScintFitter** is the
²¹ event’s position and time. The reconstructed position corresponds to the point in the
²² detector where the triggered event most likely came from (assuming the event was
²³ point-like in extent), whilst the reconstructed time is the starting emission time of the
²⁴ event, relative to the event’s trigger time. The position of an event is critical to know,
²⁵ as far fewer background events occur near the centre of the detector compared to the

edges. It is also important to know the emission time of an event, as this allows us to build the so-called ‘time residual’ (t_{res}) distribution of an event. For a point-like physics event in the detector, t_{res} for a given PMT hit is defined as:

$$t_{res} = t_{\text{hit}} - t_{\text{TOF}} - t_{\text{emm}}, \quad (2.11)$$

where t_{hit} is the calibrated hit time of the PMT, t_{TOF} is the time one expects for light to travel directly from the reconstructed position to that PMT (the time-of-flight), and t_{emm} is the reconstructed emission time.

Whilst a number of algorithms have been developed for reconstructing position and time in SNO+, they all work on the same basic principle. Because of the spherical symmetry of the detector, if an event occurs at the centre of the detector one expects direct light to hit PMTs throughout the detector at the same time³. However, if an event happens some distance away from the detector’s centre then direct light will arrive at the PMTs it is closer to sooner. Therefore, by looking at the distribution of hit times for PMTs that were hit earliest as a function of the PMTs’ positions in the detector (ignoring PMT hits that arrived much later, presumably because the photon paths were not direct) one can try and estimate where the position of the event was. Reconstructed positions and reconstructed times are deeply linked by the time residual equation described above.

Currently on SNO+, a likelihood-based approach is used to reconstruct position and time. The algorithm endeavours to maximise the combined likelihood of the observed calibrated hit times of the hit PMTs, given proposed points in the four-dimensional (position, time) parameter space [42, 43, 44]. However, regardless of algorithm there are two factors that limit the position and timing reconstruction of an event. The

³It is possible for the centre of the AV and PSUP to not be completely aligned, i.e. there is an ‘AV offset’. Then, the spherical symmetry can be broken by refraction through the AV. Fortunately, we account for this when coordinating our position fitters.

¹ TTS of the PMTs used as well as the speed of the scintillator emission timing defines
² the fundamental timescale — and hence also length scale — by which events can
³ be reconstructed. Secondly, if more photons are able to generate prompt hits in the
⁴ detector from a given event, then more information can be used to determine the
⁵ position and time. Under current conditions, a 2.5 MeV event in the centre of the
⁶ detector will have a position resolution of 100 mm [45].

⁷ The other critical piece of reconstructed event information is the event’s energy.
⁸ More precisely, this is the kinetic energy of an electron that has been assumed to have
⁹ generated the event. By consequence, events due to α -decay will obtain reconstructed
¹⁰ energies well below the actual energy of the α particle, because of scintillator quenching.

¹¹ At its simplest, assuming that an event is from an electron of moderate energy,
¹² then we can expect the number of PMT hits observed to be directly proportional to
¹³ the energy of the event. Given that the number of hits observed in an event (called the
¹⁴ `nhit`) is governed by a Poisson distribution, then the uncertainty in energy will just be
¹⁵ proportional to the square root of the number of hits. As a result, the reconstructed
¹⁶ energy resolution in SNO+ is determined by the scintillator’s light yield and absorption
¹⁷ length, as well as the coverage and QE of the PMTs.

¹⁸ There are second-order corrections to the energy reconstruction that need to be
¹⁹ considered to minimise bias. At low energies, scintillator quenching becomes non-
²⁰ negligible, so an understanding of the scintillator’s Birks’ constant is needed. At
²¹ high energies, many PMTs will have had multiple photoelectrons generated, so merely
²² using the `nhit` will give an underestimate of the true energy. Finally, the detection
²³ efficiency of photons is non-uniform as a function of position in the detector. The
²⁴ current energy reconstruction algorithm used within SNO+ attempts to deal with all
²⁵ of these effects [46, 47].

After position, time, and energy reconstruction, **ScintFitter** calculates a number of additional quantities from what are known as classifiers. These describe a wide number of properties about an event, often using the derived t_{res} distribution of an event as the basis for classification. Some examples of classifiers used in analysis are discussed in Section 6.1.2.

All Physics data runs, as well as certain calibration data such as AmBe and ^{16}N , has the **ScintFitter** algorithm run over it after having been processed for time and charge calibration. This results in what is known as a fully-processed RATDS file, as well as a new file type known as an **ntuple**. This latter file type has much of the hit-level information removed, and contains only event-level information such as the reconstructed energy and position. Because these files are much smaller, they are the ones typically used in the high-level physics analyses on the experiment, such as the one described in Chapter 6.

2.4.5 Event Simulation

Simulations of events in SNO+ are performed using the software **RAT** [2]. Built on the **GEANT4** particle physics software framework, **RAT** is capable of simulating all aspects of the physics of an event within the detector via a Monte Carlo (MC) approach. This includes any particle physics that defines an event’s generation, propagation and interactions of those particles in the detector media, the generation of light by both scintillation and Cherenkov processes, the propagation of that light, as well as the detection of that light by PMTs and simulation of the expected DAQ response. **RAT** is then used to process both simulation and data in the same way. In addition to being highly customisable, **RAT** can use the **RATDB** tables generated from a given data run when simulating to try and match those particular run conditions as closely as possible. **RAT** also offers a suite of tools to assist with analysis of data. The software is

- ¹ constantly being updated with new features — the work done in this thesis uses RAT
- ² versions between 6.18.8 and 7.0.9, inclusive.

Chapter 3

The SMELLIE Calibration System

There's a certain Slant of light,

Winter Afternoons —

That oppresses, like the Heft

Of Cathedral Tunes —

EMILY DICKINSON

As mentioned in Section 2.4.3, one of the principal systems for calibrating the optics of the SNO+ detector is SMELLIE. This calibration device consists of 5 different optical wavelength lasers able to be fired through 15 optical fibres, whose endpoints are attached to the PSUP. A collimator is attached at the end of each fibre, ensuring the emitted light forms a narrow beam across the detector. A diagram of SMELLIE in the detector is shown in Fig. 3.1.

The primary goal of SMELLIE is the measurement and monitoring of optical scattering within the detector over the lifetime of the experiment. By firing light from SMELLIE into the detector, some fraction of the photons will be scattered by the detector medium, a fraction of those scattering at large angles relative to the

Fig. 3.1

1 direction of the SMELLIE beam. This strongly scattered light can be detected by
2 PMTs far from the ‘beamspot’, and will also arrive substantially later than light which
3 travelled directly from the fibre to those PMTs. By isolating this scattered light signal,
4 and comparing the quantity observed in data to equivalent simulations with varying
5 scattering lengths, in principle one can measure the detector medium’s scattering
6 length. If one takes SMELLIE data with various wavelengths of light at various points
7 in time, we can get a dynamical picture of the optical scattering in SNO+. An analysis
8 of optical scattering in the scintillator phase is made in Section 5.2.

9 Another substantial measurement that can be made with SMELLIE is the extinction
10 length of the detection media as a function of wavelength and time. This can be done by
11 comparing the fraction of light emitted by the fibre that gets observed in the beamspot
12 after passing through the detector. Section 5.1 covers this analysis in the scintillator
13 phase. Once measurements of both the scattering length and extinction length have
14 been made, it is then possible to derive the absorption length from Eq. 2.9.

15 Why is measuring the scattering and absorption lengths of the detector medium
16 important? Both optical characteristics of the detector impact the propagation of light
17 from physics events, and hence which PMTs get hit along with the times of those
18 hits. If these lengths are systematically off within simulation, this can lead to negative
19 consequences for reconstructing events. In particular, for the scintillator phase, if there
20 is more optical absorption occurring than expected, then because not all absorbed
21 light is re-emitted a larger fraction of photons are lost. Therefore, because energy
22 reconstruction is strongly dependent on the number of PMT hits observed in an event,
23 the energy of events will be systematically under-estimated. Alongside this, light that
24 is re-emitted will only do so after some time delay, and the direction of this re-emission
25 unlikely to be in the same direction as before. This leads to systematic changes in the

Fig. 3.2

observed time residual distributions, impacting position reconstruction, as well as any classifiers that use the time residual distribution.

If there is more optical scattering than expected within the scintillator, this also indirectly leads to a greater loss of light because of the increased path length that a photon will typically have to travel before being detected. By consequence, there will be a second-order impact on the energy reconstruction from systematics in the scattering length. Much like changes in the quantity of re-emitted light, increasing the amount of scattering will also systematically effect the position reconstruction and many classifiers.

This chapter covers the hardware used to take SMELLIE data, as well as the steps taken to ensure high quality data was taken throughout the scintillator phase of the experiment.

3.1 The SMELLIE Hardware

A full description of the initial hardware setup that was used during the air fill and early parts of the water fill phase can be read in [1]. Since then, a series of hardware upgrades have been made, with [2] covering the hardware status used in data taken throughout the water phase. Fig. 3.2 shows a timeline of the hardware upgrades as well as the data taking campaigns performed throughout the current lifetime of the experiment. The current layout of the SMELLIE hardware, showing the connections between each of the devices within the system, can be seen in Fig. 3.3. The rest of this section briefly summarises the current contents of the calibration system, along with descriptions of the major hardware changes made since the water phase.

Fig. 3.3

(a)
(b)

Fig. 3.4

3.1.1 Lasers

Fundamental to the SMELLIE calibration system are 5 optical-wavelength lasers. Four of these are fixed-wavelength pulsed-diode laser heads from the company PicoQuant. These ‘PQ’ laser heads each emit with a different narrow wavelength spectrum, peaking at 375 nm, 407 nm, 446 nm, and 495 nm. These are referred to as the PQ375, PQ407, PQ446, and PQ495 lasers, respectively. In addition to these lasers, a SuperK Compact laser made by NKT Photonics (hereafter referred to as the SuperK laser) is also used¹. Unlike the PQ laser heads, the SuperK is a super-continuum laser able to produce laser light over the whole optical wavelength spectrum. Because we are almost always interested in determining optical properties at specific wavelengths, a variable bandpass filter also built by NKT Photonics known as the SuperK Varia has been included. This allows the user to select any wavelength interval between 400–700 nm, with a minimum bandwidth of 10 nm. The wavelength and emission timing characteristics of all five lasers are shown in Fig. 3.4. The $\mathcal{O}(1\text{ ns})$ time widths for the lasers are essential for scattering analyses as they allow for much greater precision in knowing when photons should arrive at PMTs in the detector than with an LED source. This is why SMELLIE uses lasers for its light generation, unlike the LEDs used for AMELLIE and TELLIE.

¹ Apologies to those more familiar with ‘SuperK’ referring to the SuperKamiokande experiment based in Japan: this laser bears no relation.

3.1.2 Controlling Laser Intensities

It is important to be able to control the quantity of light that enters the detector from a given pulse of one of the lasers. Too much light, and because of the detector's DAQ analysis of data becomes difficult, and in the extreme case permanent damage could be done. Too little, and insufficient statistics are collected by the detector to actually do any analysis. Controlling the pulse intensity of light into the detector is done in two parts. For the SuperK laser, the raw power of the beam in a pulse can be set as a percentage of the maximum possible power for that wavelength. Once a light pulse has been generated, it is then attenuated by a neutral density filter contained within the SuperK laser hardware.

For the PQ lasers, a PicoQuant-brand SEPIA II laser driver is used to set the raw pulse intensity and firing rate. Unlike the SuperK, the intensity control is given by the laser driver's driving voltage, as a fraction of the maximum possible voltage. Problematically, the dependence on the raw output intensity of the PQ laser heads are highly nonlinear with respect to the amplitude of the driving voltage. To demonstrate, consider the observed npe per shot within the detector as a function of driving voltage for SMELLIE events using the PQ lasers: see Fig. 3.5. For low driving voltages, the resulting npe is very small, and rises slowly. However, near a certain ‘lasing threshold’ the npe observed increases super-exponentially, rapidly climbing many orders of magnitude in intensity. Above this threshold, increasing the driving voltage once again changes the observed intensity relatively slowly. A result of this is that if one requires an npe in the detector associated with a driving voltage that lies near the threshold, then any uncertainty in that driving voltage value can lead to dramatic changes in the npe observed. Furthermore, when driving a laser head near its lasing threshold the shot-to-shot variation in intensity can also become substantial: see Fig. 3.6 for an example of this occurring.

Fig. 3.5**Fig. 3.6**

1 During the water phase and for much of the scintillator phase, much of the data
2 taken, especially using the PQ407 and PQ446 lasers, suffered from large shot-to-shot
3 intensity variations. Throughout this period, after the light was generated by a PQ
4 laser head it would then be passed through an attenuator, fixed to some nominal
5 attenuation setting for each laser. In theory, one could solve the intensity variation
6 problem by deliberately setting the intensity well beyond the lasing threshold, and
7 then changing the attenuation of the attenuator to obtain the npe within the detector
8 one is interested in. However, under the original hardware this was untenable as this
9 would require someone to manually change the attenuations in-person every time a
10 different set of SMELLIE run conditions were proposed.

11 Instead, Jeff Lidgard built a piece of hardware called the remotely-controllable
12 Variable Fibre Attenuator (VFA), shown in Fig. 3.7. Contained within a metal housing
13 were a ‘precision variable attenuator’ from DiCon Fibreoptics [] for each PQ laser,
14 along with an Arduino running firmware written by Jeff to enable communication with
15 each of the attenuators. Commands could be sent to a given attenuator asking for a
16 specific attenuation between 0–3000 dB. Following ex-situ testing by Jeff Lidgard and
17 Jasmine Simms, the VFA was installed underground by myself and Armin Reichold in
18 July 2022, with some assistance from Jeff in integration of the hardware and SMELLIE
19 server software.

20 During testing of the VFA in-situ, it was discovered that the inherent attenuation
21 of the variable attenuator at the minimum setting of 0 dB for the PQ375 laser was

Fig. 3.7

Fig. 3.8

so strong that negligible light was ever observed in the detector. Because of this, the PQ375 was not hooked up to the VFA, and kept its original attenuator setup. After fixing the driving voltage settings for PQ407, PQ446, and PQ495 to be XXX, XXX, and XXX respectively, the observed npe in the detector was once again compared to the input intensity setting. For PQ375 this input parameter remained the driving voltage; for the others, the attenuation setting was now used. The results can be seen in Fig 3.8. As hoped for, the three PQ lasers hooked up to the VFA can now have stable intensities of light observed in the detector over multiple orders of magnitude of observed intensity.

3.1.3 Propagation of Light into the Detector

Once a pulse of optical light has been generated by the lasers and attenuated to the desired intensity, the next step is to navigate that light into the detector. This is achieved through a network of Corning-brand “InfiniCor SXi” multimode optical fibres []. These fibres were chosen in part for their low intrinsic radioactivity [], as well as having a graded index as a function of radius. This latter property enables lower dispersion between different modes of the light, so that the initial sharpness of any given light pulse in time is maintained. However, because these fibres were mainly designed for telecommunication purposes, their nominal operating wavelengths are out in the near-infrared, 750–1450 nm. As SMELLIE only fires wavelengths in the range 375–700 nm, there is a small but non-negligible amount of light lost when propagating through the fibres.

After some light is split off by a beamsplitter to allow for ex-situ monitoring of the light pulse (see Section 3.1.4), it is sent to the Fibre Switch, two boxes manufactured

Fig. 3.9

1 by Laser Components UK [] that allows a user to remotely-control which of the fibres
2 to send the light down into the detector.

3 Finally, the light that passes through the fibre switch propagates along one of the 15
4 optical fibres that have been submerged in the SNO+ cavity, whose ends are mounted
5 to the PSUP. Specifically, sets of three fibres are mounted to a given node of the PSUP,
6 with associated node numberings: nominally 07, 25, 37, 55, and 21. These provide
7 for a variety of positions within the detector from which light can be emitted. Each
8 mounting which holds three of the optical fibres also contains collimators, designed to
9 reduce the possible range of angles with which the light can be emitted from. This is
10 particularly important for SMELLIE, because unlike the other ELLIE systems, a thin
11 ‘pencil’ beam of light across the detector is ideal for measuring scattering [31]. One last
12 thing the mounting achieves is to point each of the three fibres in different directions:
13 0°, 10°, and 20° from the direction radially towards the centre of the detector.

14 Each fibre is given a name that nominally refers to both its mounting position and
15 its pointing direction. For example, the label ‘FS107’ corresponds to the SMELLIE
16 fibre mounted at node 07 with a pointing direction of 10°. Unfortunately, during
17 installation some fibres were mislabelled, leading to the node mounting points and
18 pointing directions of some fibres being inconsistent with the labelling convention. The
19 actual pointing directions for each fibre can be seen in Table 3.1. The 3D positions
20 and pointing directions of all the fibres are shown in Fig. 3.9.

21 **3.1.4 The Monitoring PMT Unit**

22 As mentioned in the previous section, part of the light generated by the lasers gets split
23 off from the main fibre path down into the detector, and instead is used for monitoring

Fibre	Node	Pointing direction
FS007	07	0°
FS107	07	10°
FS207	07	20°
FS025	25	0°
FS125	25	10°
FS225	25	20°
FS037	37	10°
FS137	37	0°
FS237	37	20°
FS055	55	10°
FS155	55	20°
FS255	55	0°
FS093	21	0°
FS193	21	10°
FS293	21	20°

Table 3.1: SMELLIE fibre names, their associated mounting nodes on the PSUP, and their pointing direction. Taken from [32].

purposes. This is achieved with a box known as the Monitoring PMT Unit (MPU). As
1
the name suggests, the MPU contains a small PMT that generates an electronic signal
2
pulse from the laser light. This signal is then shaped by electronics, and passed to the
3
detector’s central CAEN digitiser to have that pulse digitised.
4

One problem with the existing MPU within SMELLIE was that the pulse it produced
5
was so broad that 300 ADC samples were needed to capture the full shape (the CAEN
6
samples at a rate of 1 every 4 ns). This led to a large fraction of data being generated
7
by SMELLIE events coming not from the PMT hit information, but simply the MPU’s
8
signal digitisation. A natural consequence of this was the rate at which the lasers could
9
be fired had to be limited to typically 1 kHz, otherwise the detector was not able to
10
handle the rate of data being generated.
11

Because of this, a new MPU was commissioned. Built by Adam Baird and Johan
12
Fopma from the Oxford Physics Central Electronics Group, this MPU had updated
13
electronics such that the pulse was shaped shorter. In addition, the rise time of the
14

Fig. 3.10

¹ pulse was made faster, in the hopes that the emission time of the light pulse for a given
² event could be captured more accurately. The new MPU was installed by myself and
³ Armin Reichold at the same time as the VFA, in Summer 2022.

⁴ Alongside the installation of the new hardware, the settings in ORCA for the
⁵ CAEN digitisation of the MPU signal were updated. In particular, the number of
⁶ samples made by the CAEN was shortened from 300 down to 124. The timing of the
⁷ CAEN sampling and delay on the TUBii trigger (more on the trigger shortly) was
⁸ also modified. As a result of these changes, the shortest possible trace was now being
⁹ generated, without missing any part of the MPU pulse or moving the observed TAC for
¹⁰ hits in the detector outside the trigger window. Fig. 3.10 shows a comparison between
¹¹ typical MPU pulses taken before and after the upgrades.

¹² 3.1.5 Event Triggering and Data Acquisition

¹³ As mentioned in Section 2.3.4, it is possible to trigger the SNO+ detector electronics via
¹⁴ an external asynchronous trigger, EXTA. Taking data with SMELLIE takes advantage
¹⁵ of this capability: instead of waiting for the normal ‘physics’ triggers such as N100 to
¹⁶ pick up the event, because we know when we are firing the laser we can send an EXTA
¹⁷ signal to trigger the detector precisely when a light pulse is within the detector.

- ¹⁸ • Describe the existing hardware, post-upgrade made in Summer 2022. For pre-
¹⁹ upgrade hardware, can simply cite previous SMELLIE theses. This includes the
²⁰ path of light into the detector, as well as the path of the trigger signal.
- ²¹ • Make sure to mention explicitly these upgrades: Tony Zummo’s fix to the TUBii
²² trigger logic, as well as the addition of the VFA, updated MPU, and modified
²³ trigger window. Make sure to motivate why these updates were made.

3.2 Software for SMELLIE Data-taking**49**

[7 pages]

1

3.2 Software for SMELLIE Data-taking

2

- Can be brief here! Little has changed since previous theses, so can mostly just summarise and cite.
- Server running on SNODROP machine, which converts high-level commands into low-level ones that the hardware can interpret.
- Run plan files written in JSON handed to ORCA which then sends relevant commands to SNODROP which fires as appropriate.
- Operator interacts with ORCA to perform SMELLIE calibration runs.
- After SMELLIE data taken, run description file created, containing metadata about the run conditions, used in analysis.

[2 pages]

12

3.3 Commissioning SMELLIE in the

13

Scintillator Phase

14

- Explain why commissioning of SMELLIE is needed: Need to confirm that SMELLIE is working as expected; determine intensity "set-points" for different use cases.
- Commissioning originally performed by Esther and JeffL back in the water phase; explain why this needed to be re-done for both the scintillator phase and after the hardware upgrades.

20

- ¹ • No need to describe the Tesseract in detail here - that can be in Jeff L's thesis. But,
² I do want to show the results of both commissioning campaigns in scintillator-fill,
³ one before the new hardware was added, and one after.

⁴ [5 pages] [15 PAGES TOTAL]

Chapter 4

1

Simulating SMELLIE Events

2

Max Power : *Kids. From now on there are three ways of doing things: the right way, the wrong way, and the Max Power way.*

Bart Simpson : *Isn't that just the wrong way?*

Max Power : *Yes, but faster!*

THE SIMPSONS

Critical to extraction of scattering information from SMELLIE data is an accurate Monte Carlo (MC) simulation of the SMELLIE system. By modelling the laser light emission into the detector correctly, we can simulate how SMELLIE light will be impacted by changing scattering lengths in the detector. Because of the complexity of the optics of the optical fibres used to direct the laser light into the detector, a given SMELLIE event is simulated as a partially-collimated “flash” of visible photons emanating from the emission point of the fibre into the detector. This flash then requires a number of parameters to be correctly described. In particular:

4

5

6

7

8

9

10

11

- **Fibre emission positions** were recorded during the installation of the fibres.

12

- **Wavelength and emission timing distributions** of light pulses were taken from measurements of the laser heads by their manufacturers [], or by colleague Jeff Lidgard in the case of the SuperK wavelength distribution [].
- **The “pulse magnitude”**, defined as the mean number of photons simulated per event, is determined on a subrun-by-subrun basis, and is assumed to fluctuate as a Poisson distribution.
- **The beam profiles**, which describe the angular emission distributions of each fibre, is the focus of this chapter. These are necessary because unlike scintillation light, the light emitted from SMELLIE fibres is not isotropic.
- **Nominal fibre emission directions** attempt to define the centre of the beam for a given fibre.

This chapter is split into three sections. Improvements to the existing simulation algorithm for the beam profiles are first made, and then the beam profiles themselves are updated. Finally, comparisons between data and simulation are made after the upgrades to investigate any remaining discrepancies.

4.1 Improving the SMELLIE Generator Algorithm

4.1.1 Previous Attempts at SMELLIE Event Simulation

Before we can determine the beam profiles, we must first decide how to specify them. Previous observations show that different fibres can have notably different beam profiles [31], so we let each fibre’s beam profiles be unique. We assume for now that a given fibre’s beam profile is stable over time, and independent of the wavelength of light fired. A straightforward, naïve approach to parameterising a beam profile would be as follows: specify some nominal fibre direction, corresponding to the direction light

4.1 Improving the SMELLIE Generator Algorithm

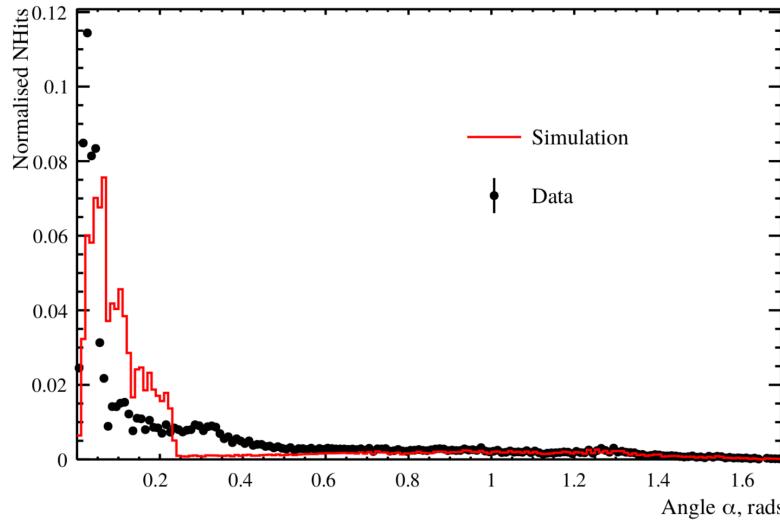


Fig. 4.1: Comparison between a simulation of one of the fibres, made from the 1D beam profile generator (red), with the associated data subrun that was used to create that beam profile (in black). For both MC and data, what is plotted is the PDF of observed PMT hits, as a function of the α angle. Poissonian errors have been added to the data points, but are too small to see. Clearly, this 1D generator does not replicate the observed beam profile correctly. Figure taken from [32].

takes travelling from the fibre to the centre of the “beamspot” observed on the other side of the detector. Then, specify a 1D beam profile, corresponding to the probability density of firing a photon at a given polar angle α relative to the nominal direction. One might even assume this distribution is Gaussian. The distribution in azimuthal direction, ϕ , is assumed to be uniform.

This 1D beam profile approach was used initially for SMELLIE, and remains in use for the other ELLIE sub-systems within SNO+. However, when SMELLIE data was taken in the water-phase of the experiment, simulations using these beam profiles failed to match them well at all - see figure 4.1 for an example. Not only was the distribution in α not Gaussian, a distinct speckle-pattern can be observed within the beamspot that is not uniform in ϕ . This fact led to colleague Esther Turner building a SMELLIE generator that could handle 2D beam profiles: dependent on both α and ϕ . The distribution was stored as a map from each inward-pointing PMT in the detector

¹ to a relative intensity value. This was chosen because the beam profile shapes were
² calibrated from existing SMELLIE data — more on this in section 4.2.

³ This original 2D generator then sampled the beam profile via a rejection sampling
⁴ approach, outlined as follows:

- ⁵ 1. Propose a test direction (α, ϕ) , by generating ϕ uniformly in the interval $[0, 2\pi]$,
⁶ and α according to some pre-determined Gaussian distribution, known as the
⁷ Gaussian envelope.
- ⁸ 2. Given this test direction, calculate where a line following this direction from the
⁹ fibre of interest will hit the PSUP on the other side of the detector. Find the 3
¹⁰ closest PMTs to that point.
- ¹¹ 3. From those PMTs, obtain their relative intensity values from the beam profile
¹² mapping, and perform an interpolation based on how close each PMT is to the
¹³ PSUP intersection point. This gives an interpolated relative intensity value for
¹⁴ this test direction.
- ¹⁵ 4. Because we are sampling using the angular coordinates (α, ϕ) , differential area
¹⁶ elements over this space of directions do not have the same size. We can correct
¹⁷ for this fact by multiplying our interpolated relative intensity by $\sin \alpha$, which
¹⁸ corresponds to the Jacobian of the direction-space.
- ¹⁹ 5. Calculate the value for the Gaussian envelope along this test direction.
- ²⁰ 6. Throw a random number uniformly between 0 and the Gaussian envelope value. If
²¹ the random number is less than the interpolated intensity, then this test direction
²² is accepted, and a photon is generated with that direction. Otherwise, we reject
²³ the direction and try the whole process again.

4.1 Improving the SMELLIE Generator Algorithm

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This generator certainly works, but has a key problem: efficiency. The 1D generator was able to generate a SMELLIE event (that is, to fully specify the starting parameters of all the photons emitted from a fibre) at a speed of ~ 1 ms. However, the 2D generator specified here could take upwards of ~ 50 s *per event* to generate. Because a typical SMELLIE analysis requires simulating many millions of events, the CPU time taken to perform this quickly became unfeasible. Fixing this generator speed problem was a high priority for the SMELLIE analysis.

4.1.2 The new generator

On careful inspection of the existing 2D generator, the main reason for the slowness of the algorithm is the use of a rejection approach. Even with use of the Gaussian envelope, which was included to help with speed, the vast majority of proposed directions are never selected. Figure 4.2 shows a histogram of number of attempts per event it took for a valid direction to be chosen for a representative SMELLIE simulation. Moreover, the calculations needing to be done for every proposed direction are relatively complex, notably trying to find the 3 nearest PMTs to some point on the PSUP.

A new 2D generator was built with these thoughts in mind. Firstly, the rejection method would no longer be used, given its inefficiency. We would also endeavour to try and “pre-calculate” as much as possible before run-time. Starting with the existing PMT relative intensity maps, we plot these in the 2D direction-space ($1 - \cos \alpha, \phi$): see Figure 4.3a. In a toy-MC simulation, 500,000 directions are then thrown uniformly in this 2D space per fibre. For each direction, the same method of obtaining an interpolated intensity value from the nearest PMTs to the corresponding point on the PSUP as from the original 2D generator was performed, the only difference being that these calculations were done well before any actual SMELLIE simulation. Figure 4.3b shows the interpolated intensities obtained for one fibre.

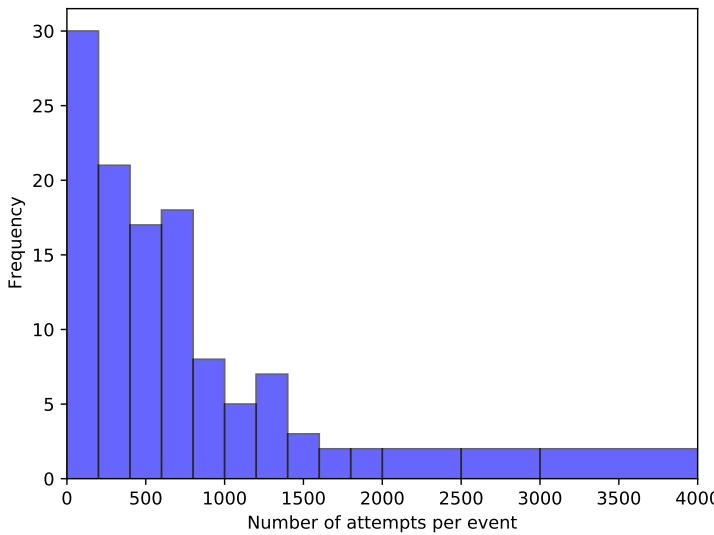


Fig. 4.2: Typical distribution of the number of attempts it takes for the existing 2D generator before the test direction gets accepted, per event.

Following this, the sampled intensities were then binned into a 2D histogram, where the bin value corresponds to the sum of all intensities for all directions found within this bin. Choosing a sensible binning procedure is important: too few bins, and necessary information about the shape of the beam is lost, whilst too many bins can oversample the data and capture statistical artefacts in the sampling process instead of just the beam profile. As a balance, 15 bins were chosen along the ϕ direction, and 60 in $r = 1 - \cos \alpha$. This was chosen to ensure that a reasonable number of PMTs were located within each bin, lessening the impact of any statistical fluctuations. Although the bins in ϕ were chosen to have uniform width, this was decided to be not the case for the other axis, as there is far more important information near $r = 0$ (the beamspot). Instead, the width of the bins in r were calculated so that roughly the same total probability was contained in each r -strip. By consequence, bins near the beamspot typically are of significantly smaller size than ones much further out. This allows us to both capture any rapid changes in intensity near the beamspot, where this

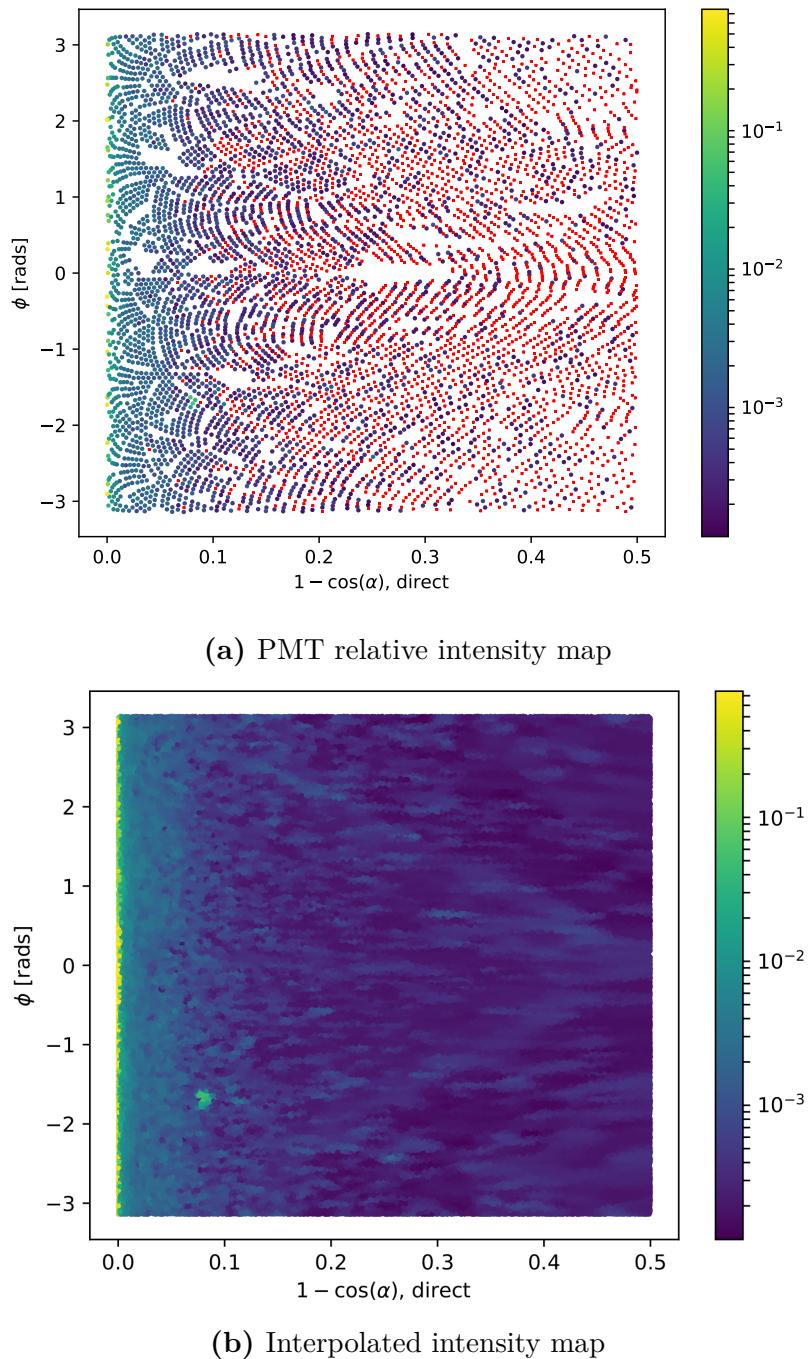


Fig. 4.3: The first step in the new method for preparing the new generator. In (a), the relative intensities used for the existing beam profile of fibre labelled FS055 are shown for each PMT, the position on the plot indicating the location of that PMT in the fibre coordinates. The colour indicates the relative intensity; PMTs marked red have an intensity of zero. Figure (b) shows the result of throwing 500,000 directions uniformly over this 2D space, the intensity of each point given by interpolating the intensities of nearby PMTs.

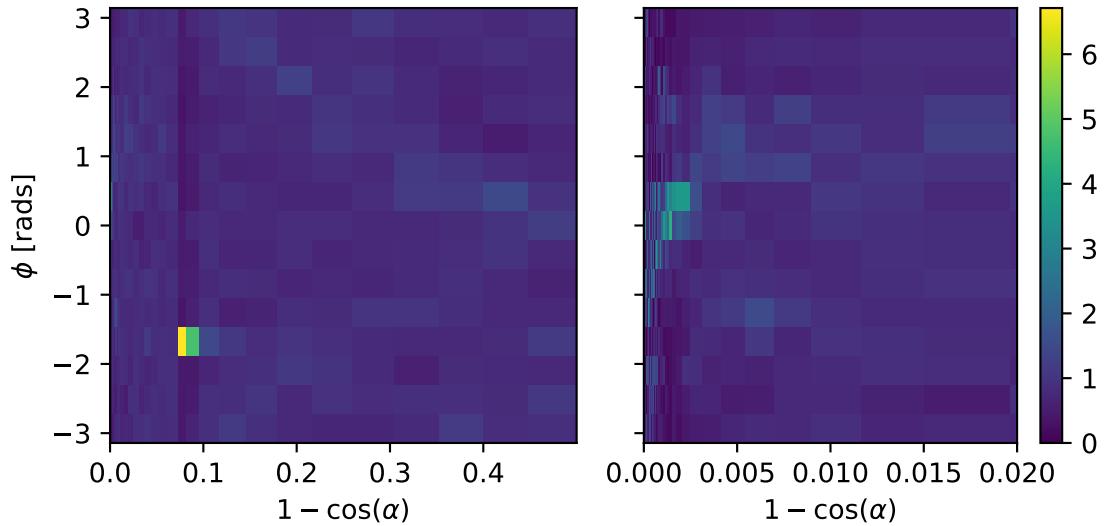


Fig. 4.4: Histogram of interpolated intensities within the 2D direction-space. The left view shows the full histogram; the right is a zoomed-in version near the beamspot. Unlike the binning in ϕ , the bin widths in r are not at all uniform. Instead, they have been determined such that the area summed over a given “strip” of bins of constant r will be the same.

¹ matters greatly, and smooth out the very-low intensities seen at larger polar angles.

² One of these histograms can be seen in Figure 4.4: the large change in bin widths as a
³ function of r is clear. One can also see that near the beamspot notable dependence on
⁴ the intensity as a function of ϕ . The mysterious “spot” at $r = 0.08$, well out of the
⁵ beamspot, is an indication that the underlying beam profile data being used requires
⁶ improvement: more on this in section 4.2.

⁷ The Cumulative Density Function (CDF) of this intensity histogram as a function
⁸ of bin was then produced, where the bins were ordered through a raster-scan: scanning
⁹ first over ϕ , and then r . The CDF was then normalised to 1 so that it was well-defined.

¹⁰ It is this CDF object that is then loaded in and sampled from during event generation.
¹¹ To do this, an “inverse-CDF” approach was used, which has the major benefit over
¹² rejection sampling of always producing a valid direction for every sample made. The
¹³ algorithm works as follows:

- ¹⁴ 1. Throw a random number uniformly in $[0, 1]$.

4.2 Improving the beam profiles

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2. Perform a binary search to find the bin that has the largest CDF value below
this random number. 1
3. Look at the bin edges in ϕ of this selected bin: use linear interpolation of the
random number to obtain a ϕ value located between these two ϕ -values. 2 3
4. Look at the selected bin's r -bin edges, and select a value of r by throwing a
second random number uniformly between the two edges. Convert this r into a
polar angle α . 4 5 6 7
5. The photon's direction is defined by the (α, ϕ) chosen by this process. 8

Because of the relative simplicity of this algorithm compared to the previous 2D generator, the speed improvement was very large: generation now took ~ 1 ms per SMELLIE event, a speed improvement of nearly 50,000. Event generation became as fast as it was when the 1D generator was being used. Furthermore, because of the approach taken, this major speed improvement comes at no sacrifice in accuracy. Figure 4.5 shows a comparison of the average number of photoelectrons (npe) per event per PMT between water-phase SMELLIE data and simulations with both the old and new 2D generator. One can see clearly that both generators are as accurate as one another. Note that this plot uses the updated beam profiles as explained in the next section. 9 10 11 12 13 14 15 16 17 18

4.2 Improving the beam profiles

Even with the new 2D profile generator, a problem remains: the simulation fails to reasonably recreate data, and much of this appears to be because of the poor beam profile data being used. The curious “spot” for one of the fibres was already noted in the previous section that doesn't seem to be physical, and more broadly at large angles 20 21 22 23

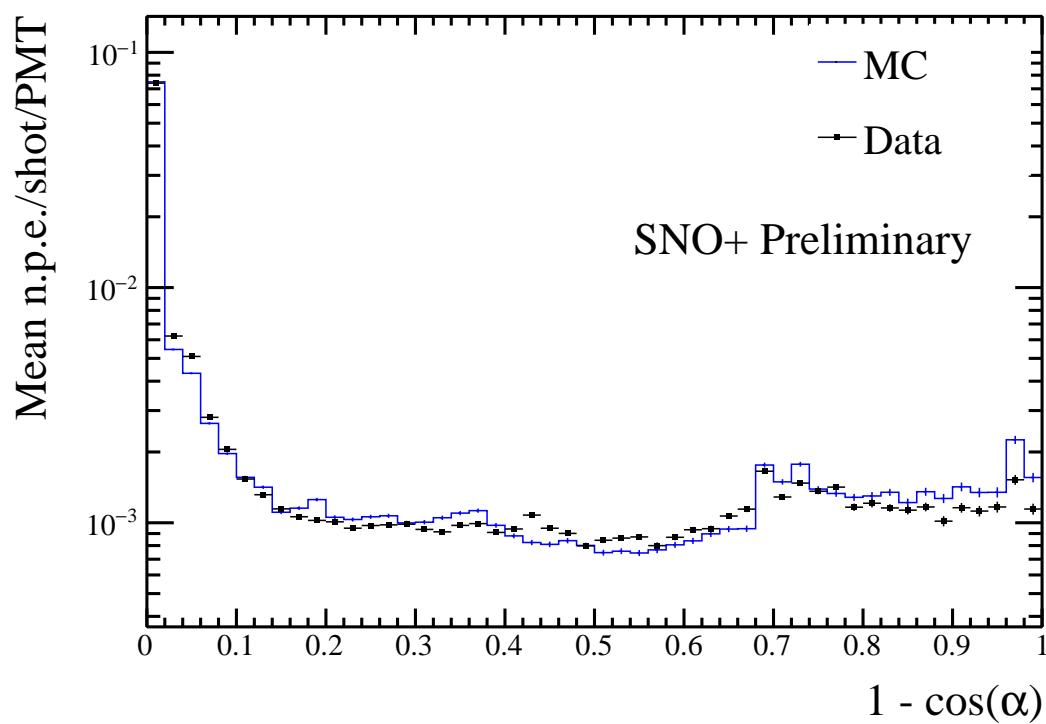


Fig. 4.5: Comparison of water-phase data to MC generated using both the old and new 2D beam profile generator approaches, with the updated beam profiles. Both versions of the generator are consistent with one another, but the new generator is many times faster.

4.2 Improving the beam profiles**61**

for all the fibres there are large swathes of PMTs with an intensity of zero, providing
 little useful information about the beam shape. It was shown in [32] that with the
 old 2D generator, the systematic uncertainty on the beam profiles was the dominant
 source of error in the main SMELLIE analysis. To help improve this situation, it was
 decided to update the existing beam profiles.

These old beam profiles were originally determined by looking at SMELLIE data
 taken during the water-phase. Specifically, a “medium”-intensity subrun with one
 of the lasers firing at a wavelength of 495 nm, was chosen for each fibre. “Medium”-
 intensity corresponds to firing the relevant laser at a set intensity determined during
 an earlier commissioning process, for which the maximum occupancy of PMT hits at
 that intensity, i.e. the proportion of hits per event, corresponded to roughly 80%. This
 value was chosen as it allowed for high statistics in a relatively short run-time, but not
 so intense that the occupancy of any given PMT in the beamspot was 100%. Because
 Rayleigh scattering is strongly-dependent on wavelength, the long wavelength of light
 was chosen so that impacts from this scattering were small in the data.

SNO+ PMTs are unable to distinguish the exact number of photoelectrons being
 generated. One is typically only able to know if a PMT has been triggered at all, by
 any number of photoelectrons. As a result, the occupancy of a PMT over a number
 of SMELLIE events, o , is a biased estimator of the mean number of photoelectrons
 generated, μ . Assuming the number of photoelectrons generated in a given event
 follows Poisson statistics, the probability of generating k photoelectrons is:

$$P(k|\mu) = \frac{\mu^k e^{-\mu}}{k!}. \quad (4.1)$$

The probability of observing a “hit” in a given PMT corresponds to generating at
 least one photoelectron:

$$P(\text{hit}|\mu) = P(k \geq 1|\mu) = 1 - P(k = 0|\mu) = 1 - e^{-\mu}, \quad (4.2)$$

which implies after rearrangement that one can determine the mean number of photoelectrons per event from the occupancy by:

$$\mu = \ln(1 - o). \quad (4.3)$$

This is the reason why we want to avoid PMTs with occupancies of 100%: they preclude one's ability to convert into a value for μ by looking at occupancy alone. We call this conversion from occupancy into npe the “multi-hit correction”. The impact of this correction is typically small for most PMTs, but can become very significant in a fibre's beamspot.

Once the npe mapping from data was obtained, a correction was then made for the detector's optics: even ignoring a fibre's beam profile, we still expect certain PMTs to be illuminated more than others because of e.g. reflections off the AV, or the solid angle subtended by the PMT bucket opening. For each fibre, a simulation was made where the beam profile was set as isotropic, and the corresponding npe mapping obtained: this map held information about the detector optics only. The beam profile mapping was then derived by simply dividing each fibre's npe mapping from data to its associated isotropic MC npe map. It is these maps that were first used in section 4.1.2.

4.2.1 Combining beam profile datasets

Fortunately, much more SMELLIE data was taken during the water-phase than was used for the original beam-profiling analysis. This additional data can be combined with that which was already used to far better constrain the beam profiles. In particular, given the existing assumption that scattering effects are minimal above wavelengths

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Run Number	Run Type	Comments
114,018	All PQ lasers; SuperK laser in 400–500 nm range	Only PQ495 laser and SuperK at 495 nm is used
114,023	SuperK laser in 500–600 nm range	Part 1 of this wavelength range; crash occurred on last subrun, so that subrun is ignored
114,034	SuperK laser in 500–600 nm range	Part 2 of this wavelength range

Table 4.1: Water-phase runs used for new beam profiling.

of ~ 490 nm, all data taken with wavelengths above this can also be used. The specific runs (and associated comments about their specifics) are described in Table 4.1. Because high-intensity runs require a different analysis approach (PMTs with high occupancies must use charge, not occupancy, to estimate npe), for this analysis we only considered subruns that used low or medium intensity set-points.

For each subrun j of data per fibre, we look only at PMT hits for each PMT i that has been identified as “good” for that subrun¹, $i \in G_j$. G_j here represents the set of good PMTs in subrun j . In particular, a “good” PMT must have valid electronic and timing calibrations, be at high voltage and masked into the detector’s trigger system for that subrun. In addition, an angular cut of $\alpha < 60^\circ$ was made to remove PMTs that are well outside any reasonable beam direction. To isolates the hits arriving directly from the fibre without reflecting, scattering, or being noise, a time cut was also made. Because what matters is the time relative to emission from the fibre, and the expected time-of-flight from fibre to different PMTs varies, a quantity known as the time residual was used. Starting with the calibrated hit time of a given PMT relative to the event’s trigger time, t_{hit} , the expected time-of-flight t_{TOF} from the fibre to the PMT was subtracted, estimated with the collaboration’s “Light Path Calculator”. Then, the emission time was also subtracted, t_{emm} , estimated by looking at the second-earliest value of $t_{hit} - t_{TOF}$ within the fibre’s central beamspot, defined as the PMTs for which

¹Strictly speaking, a PMT’s “goodness” is only determined on a run-by-run, not a subrun-by-subrun level, but this has no impact on the analysis.

¹ $\alpha < 3^\circ$. It was found that a “loose” time residual cut of $t_{res} \in [-10, +12]\text{ns}$ was
² sufficient to remove the vast majority of non-direct light with little signal sacrifice. In
³ the situation where a subrun with intensity was very small, it would not regularly have
⁴ at least two hits in the beamspot, and so the time residuals calculated would not be
⁵ valid for many events. To avoid this situation, a cut was made on any subruns with
⁶ mean intensities below 9 within their beamspot. This value was chosen as it would
⁷ mean a 2σ fluctuation downwards of $2 \cdot \sqrt{9} = 2 \cdot 3 = 6 \text{ npe}$ would still have more than
⁸ the 2 hits necessary for timing reconstruction. One fibre, FS207, has no data subruns
⁹ that satisfy this condition, and as such will have to be dealt with separately. For the
¹⁰ time being, this fibre was ignored.

¹¹ Extracting the underlying beam profiles from these data required some careful
¹² thought, especially because different subruns could have different intensities. Consider-
¹³ ing a PMT i in subrun j , the mean number of photoelectrons generated per event in
¹⁴ that PMT for that subrun, μ_{ij} can be decomposed as follows:

$$\mu_{ij} = I_j k_i = I_j b_i f_i. \quad (4.4)$$

¹⁵ I_j is the intensity of the subrun, i.e. the mean number of photons generated from the
¹⁶ fibre in that subrun per event. k_i is the probability that a given photon generated at
¹⁷ the fibre source ends up generating a photoelectron in PMT i . This itself can be further
¹⁸ split into two components: b_i , the probability that a given photon at the fibre source
¹⁹ points in the direction of PMT i ; and f_i , the probability that a given correctly-pointed
²⁰ photon actually makes it to the PMT and successfully generates a photoelectron. It is
²¹ b_i that is the actual beam profile we would like to measure.
²²

²³ Letting p_{ij} be the probability of observing a hit for a given event on a given
²⁴ PMT, the probability of observing m_{ij} hits out of N_j events in the subrun will be

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binomially-distributed:

$$P(m_{ij}|\mu_{ij}) = L(\mu_{ij}|m_{ij}) = \binom{N_j}{m_{ij}} p_{ij}^{m_{ij}} (1-p_{ij})^{N_j-m_{ij}} = \binom{N_j}{m_{ij}} (1-e^{-\mu_{ij}})^{m_{ij}} e^{-\mu_{ij}(N_j-m_{ij})}. \quad (4.5)$$

Here we have used equation 4.2, and noted that this probability distribution in m can be re-framed as a likelihood function for the parameter μ_{ij} . Considering only a single subrun of data, the maximum likelihood estimate of the parameter μ_{ij} can be shown to be:

$$\langle \mu_{ij} \rangle = -\ln \left(1 - \frac{m_{ij}}{N_j} \right) = \ln (1 - o_{ij}) \quad (m_{ij} \neq N_j), \quad (4.6)$$

where o_{ij} is just the occupancy of PMT i in subrun j . This is just the multi-hit correction formula seen in equation 4.3, which makes sense.

When looking at multiple subruns for the same fibre, the total likelihood function for a given PMT when considering all the data for a given fibre will be the product of the likelihoods from each dataset,

$$L(\{I_j\}, k_i | \{m_{ij}\}) = \prod_j L(I_j, k_i | m_{ij}) = \prod_j \binom{N_j}{m_{ij}} (1-e^{-I_j k_i})^{m_{ij}} e^{-I_j k_i (N_j-m_{ij})}. \quad (4.7)$$

This leads to a log-likelihood distribution of

$$\mathcal{L}(\{I_j\}, k_i | \{m_{ij}\}) = \sum_j \left[\ln \left(\binom{N_j}{m_{ij}} \right) + m_{ij} \ln \left(1 - e^{-I_j k_i} \right) - I_j k_i (N_j - m_{ij}) \right]. \quad (4.8)$$

Formally, one could combine the likelihoods of all the PMTs together, and by looking at the maximum likelihood estimates for each of the parameters measure the parameter values this way. However, the set of equations one obtains through this approach quickly become analytically intractable, because the PMTs are coupled by the intensity values I_j . Even a direct numerical approach would be liable to fail: for a given fibre

¹ there can be dozens of subruns, and many thousands of PMTs of relevance, so the
² dimensionality of the system of equations would be far too large.

³ Because of this, a different approach was taken. It is expected that in a subrun the
⁴ total npe, summed over all good PMTs, should be proportional to the intensity value
⁵ I_j . One must be careful about this construction — different subruns can have different
⁶ sets of good PMTs, so two subruns with identical I_j values could have a larger summed
⁷ npe merely because more PMTs were good in that subrun. To counter-act this effect,
⁸ only PMTs that were classified as good in *all* subruns being analysed for that fibre
⁹ would be used for the npe summation. In other words, we use data from PMT i for
¹⁰ summing only if:

$$\text{11} \quad i \in \mathcal{I} = \bigcap_j G_j. \quad (4.9)$$

¹² We can then define the summed npe for a given subrun as $S_j = \sum_{i \in \mathcal{I}} \text{npe}_{ij}$, and assert
¹³ that $I_j = cS_j$. By finding a value proportional to I_j , there is now enough information
¹⁴ to maximise the log-likelihood $\mathcal{L}(k_i | \{m_{ij}\}, \{I_j\})$ with respect to k_i for each PMT
¹⁵ independently, and hence obtain estimates for these k_i parameters.

¹⁶ Of course, what is actually wanted are the underlying b_i values, not k_i . This is
¹⁷ where isotropic simulations come in. For each run of data used, a matching isotropic
¹⁸ MC was produced. As an example, a simulation for run 114,023 contained 200,000
¹⁹ events for each fibre using an isotropic beam profile, over the full wavelength range
²⁰ considered in this run, 500–600 nm, using the same run conditions as in data (which
²¹ PMTs were at high voltage, etc.).

²² For each isotropic MC run, both I_j^{MC} and k_i^{MC} were calculated via the method
²³ described above. Because the simulations were isotropic, the underlying value for b_i
²⁴ was constant across all the PMTs, and so $ak_i^{MC} = f_i$. By doing some rearranging of

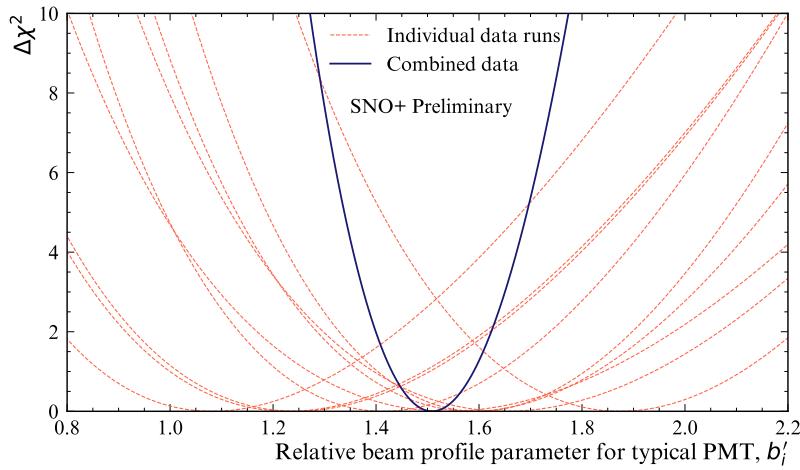


Fig. 4.6: Plot of $\Delta\chi^2 \simeq X_i$, twice the negative log-likelihood ratio, for both single subruns of a typical PMT, and when all relevant subruns are combined.

equation 4.4, we find that:

$$\mu_{ij} = I_j b_i f_i = c S_j b_i a k_i^{MC} = (acb_i) S_j k_i^{MC}. \quad (4.10)$$

As a result of this, given the set $\{S_j\}$ and k_i^{MC} , one can maximise the log-likelihood \mathcal{L} with respect to $b'_i = acb_i$ numerically, to obtain the maximum likelihood estimate of b'_i . Because a and c were global constants of proportionality, they would become irrelevant as soon as the beam profile was normalised in the CDF-creation process outlined in 4.1.2.

Figure 4.6 shows the shape of this log-likelihood distribution for a particular PMT when considering fibre FS007's beam profile. One can see how individual subruns provide much more information when combined than if one looked at a single subrun alone.

Another benefit of using this log-likelihood approach is that the resulting distribution's shape can be used for uncertainty estimation. In almost all cases, Wilks Theorem [48] allows us to produce 1σ confidence intervals about the maximum likeli-

hood estimate for b'_i , $\langle b'_i \rangle$, because

$$X(b'_i) = -2 [\mathcal{L}(b'_i) - \mathcal{L}(\langle b'_i \rangle)]$$

¹ approximates a χ^2 -distribution. As a result, the error bounds on our parameter estimate
² are given by when $X = 1$. The fact that the shape of X can be well-approximated by
³ a quadratic in the region near $X = 0$ indicates the validity of Wilks' Theorem being
⁴ used here.

⁵ Only a couple of exceptions to this approach of parameter estimation are possible.
⁶ In the case where $m_{ij} = N_j$, i.e. a PMT has 100% occupancy, no maximum likelihood
⁷ estimate exists: we need not worry about this, as subruns where this occurs have not
⁸ been used. On the other end, however, there are some PMTs for certain fibres where
⁹ after all subruns of data have been included, there remains no hits. In this scenario,
¹⁰ one can show that the log-likelihood becomes linear in the beam profile parameter:

$$\mathcal{L}(b'_i | \{m_{ij} = 0\}) = b'_i k_i^{MC} \cdot \sum_j [I_j N_j]. \quad (4.11)$$

¹² This scenario is very much reminiscent of rare-decay searches, and a similar approach
¹³ can be used. A 1σ upper limit on the possible value for b'_i can be analytically-calculated
¹⁴ to be:

$$b'_{i,ulim} = -\frac{k_i^{MC} \sum_j [I_j N_j]}{\ln [1 - \text{erf}(1/\sqrt{2})]}, \quad (4.12)$$

¹⁶ where $\text{erf}(x)$ is the error function. [18 pages for above two sections]

¹⁷ 4.2.2 Results & Discussion

¹⁸ WARNING: contents of this subsection will be gutted, focusing merely on impact of
¹⁹ combining data sets. Details about discrepancies will be covered in the next section.

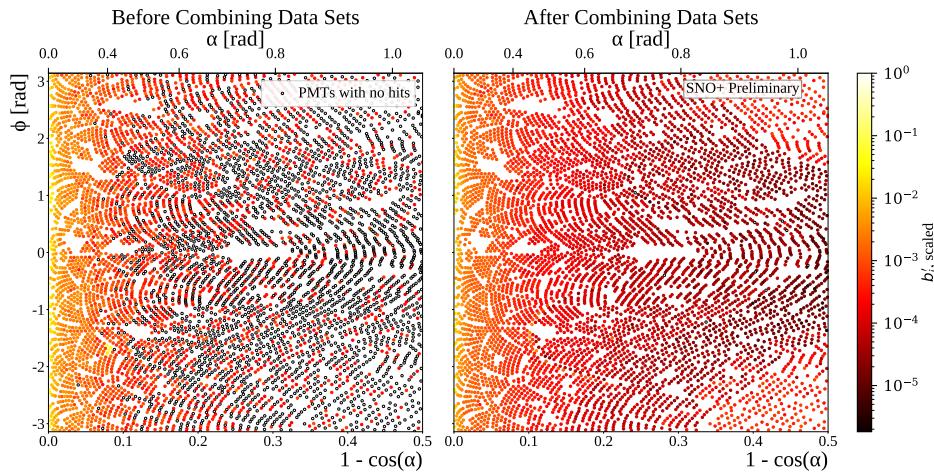


Fig. 4.7: Comparison between old and updated beam profiles for fibre FS055, after combining multiple data sets. Once again, the relative intensities (b'_i) for each PMT are given by the colour of each point, the position of each plotted in the 2D (r, ϕ) -space. The relative intensities have been both scaled here so that the largest value equals 1. Hollowed-out points are PMTs that, even after all relevant subruns have been combined, have no PMT hits.

[Will just be 1 page] Figure 4.7 shows the impact of using additional subruns of data on a typical beam profile. One can clearly see the great reduction in the number of PMTs with no hits in data. That many more data sets were included allowed for the major increase in dynamic range available for measuring these b'_i values. One can also note that by including additional data the curious spot that was seen in the old beam profile our at $r \approx 0.08$ has gone, further indicating that it was an artefact of that single data set.

Further details can be gathered from the interpolated intensity maps, one of which can be seen in figure 4.8. There are two curious stand-out features that can be seen here: firstly, there are multiple distinct parabolic arcs. These correspond to the shadows of the ropes that hold up/down the AV. More precisely, they are the mismodelling of those shadows — if the shadows were in the right place in the isotropic MC, then they would correctly cancel out any decreased intensity seen in the data of shadowed PMTs. These shadows could be mismodelled either because the positions of the ropes in the MC are in the wrong place, or the fibre's emission position is wrong. Note that any

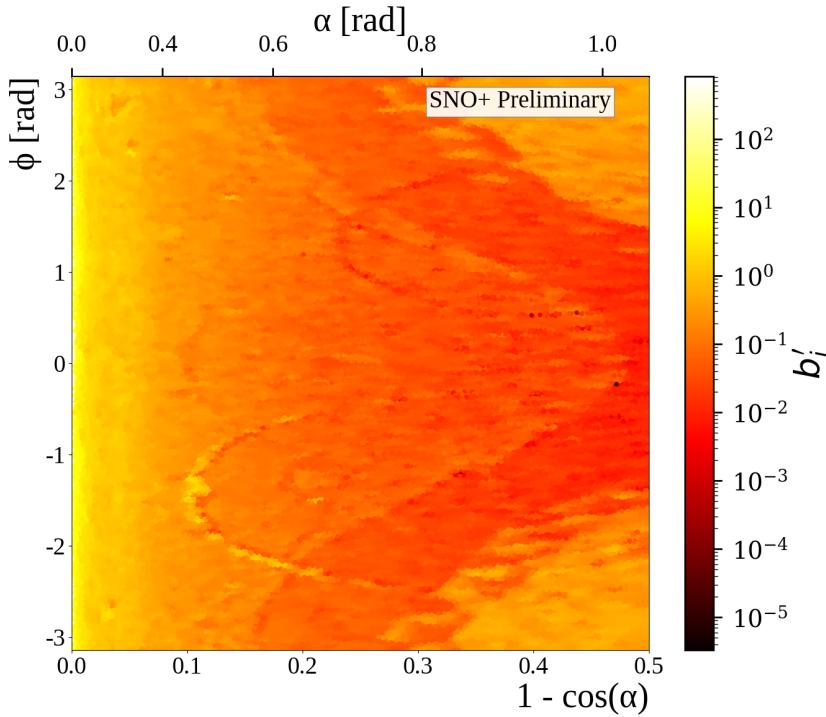


Fig. 4.8: Interpolated intensity map for the new updated beam profile of fibre FS055. The misalignment of rope shadows and AV effects, can both be seen.

1 mismodelling of the fibre's nominal emission direction has no impact on this shadowing
 2 problem, as changing that direction merely causes a change of basis in the (r, ϕ) -space.
 3 The latter possibility of incorrect fibre positions are more likely, and in fact these arcs
 4 in the beam profiles could be used as an effective way to correct for this problem.

5 The second distinctive feature of this intensity map is the large band of lower
 6 intensity varying between $r \approx 0.2 - 0.5$, followed by larger intensity out at large
 7 r values. This feature comes from light reflecting off the AV surface, or internally-
 8 reflecting. The reason for this band's functional dependence on ϕ is that this particular
 9 fibre, FS055, has a nominal fibre direction $\sim 10^\circ$ from pointing radially-towards the
 10 detector's centre. This feature appears in the updated beam profiles of all fibres, but
 11 its shape depends on the particular fibre's direction — for fibres pointing directly
 12 towards the detector's centre, there is little ϕ -dependence observed. Like the ropes,

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this feature must come from some form of mismodelling of the optics of the AV. A de-facto shadowing of PMTs in line with tangents from the AV surface which intersect the fibre position is to be expected. One also expects PMTs at polar angles larger than this to have their observed intensities boosted from reflected light off the AV. However, the discontinuities seen in the beam profiles indicate that for whatever reason this effect has been over-emphasised in the simulation.

There is a further phenomenon that can be seen, by comparing beam profile values obtained from a single subrun to the updated combined beam profile. This can be done by calculating the residuals corresponding to the single subrun, relative to the combined data set. The residual is negative if the combined data sets have a b'_i below the equivalent for a given single subrun; that is, the combined model underestimates this subrun for that PMT.

This information was plotted for two different subruns from the same fibre, seen in figure 4.9. One subrun was the same one used by Esther Turner for the original 2D beam profiling, with a wavelength of 495 nm; the latter was at the longer wavelength of 595 nm. For both subruns, most PMTs are seen to have intensities well-modelled by the combined model. However, there appears to be a significant amount of mismodelling within the beamspot. There also appears to be some systematic shift between data and model at somewhat larger polar angles. Moreover, this mismodelling seems not to be merely random, but a function of wavelength: at shorter wavelengths the beamspot tends towards being overestimated and then underestimated at larger values of α . At longer wavelengths, the beamspot becomes underestimated, with larger angles getting overestimated. This indicates that there appears to be a wavelength-dependence on the beam profiles, contradicting one of the main assumptions which we used to combine the water-phase data in the first place! All three of these features — rope shadows, AV reflections, and wavelength dependence — add systematic uncertainty to the beam

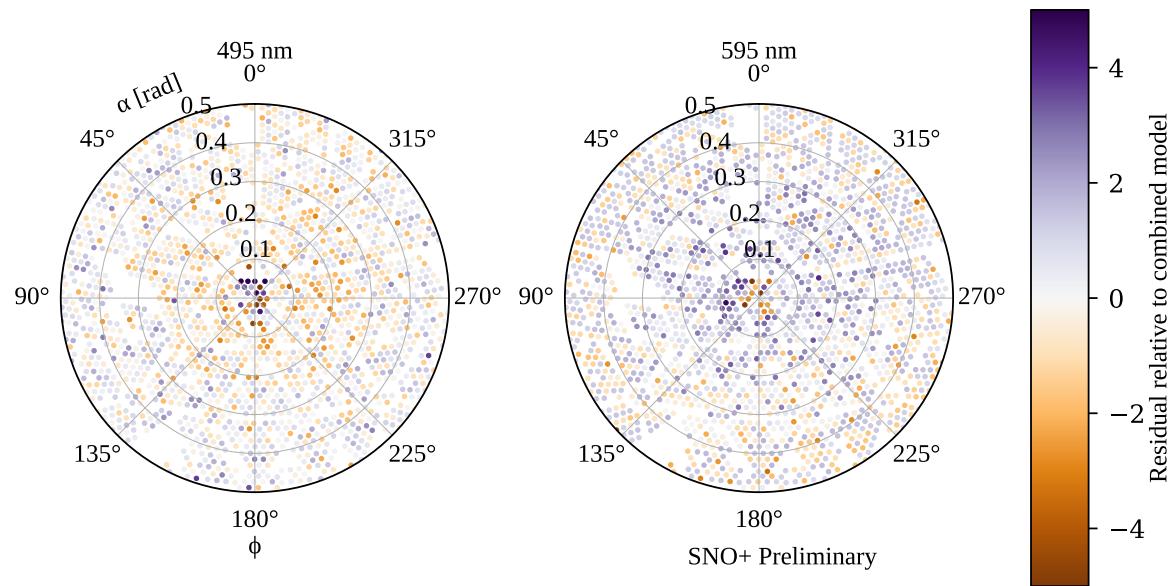


Fig. 4.9: Residuals from subruns at two different wavelengths, both compared to the combined beam profile model for fibre FS055. A negative sign, and hence bluer colours, indicate that the combined model underestimates the observed intensity for that particular subrun. Values with a magnitude beyond 5 are shown capped at this maximal value for the purposes of this plot. These PMTs are plotted in the polar fibre coordinates (α, ϕ) .

1 profiles, beyond the statistical uncertainty as measured by the width of the likelihood
 2 distribution. Certainly if one wanted to further improve the uncertainties in the beam
 3 profiles, tackling these challenges would be key.

4 4.3 Comparisons between Data and Simulation

- 5 • This focuses on disagreements noticed between data and MC even after the new
 6 generator and beam profiles have been used.
- 7 • Important to mention that none of these are necessarily game-ending, they are
 8 just systematics that may or may not be substantial in a given analysis with
 9 SMELLIE..

4.3.1 Forward Hemisphere Discrepancies

- The continued disagreement between data and MC when it comes to measured npe in various parts of the “forward” hemisphere. This includes:
 - The central beamspot,
 - The TIR region,
 - Rope shadows,
 - A noticeable wavelength-dependence.
- This is pretty much most of the contents of Section 4.2.2.

[4 pages]

4.3.2 Emission Time Discrepancies

- For certain lasers, a strong mismatch in the observed hit time residuals for prompt light.
- A mysterious +18 ns bump seen for the PQ495 laser.
- A trigger jitter in the SuperK laser.

[3 pages]

4.3.3 Backward Hemisphere Discrepancies

- The observed distribution of hits vs time and angle in MC does not match data in a number of ways for PMTs near the fibre emission point.
- Includes the outer-water scattering length, rope reflections, and investigations into whether certain modifications to the optics could plausibly fix things (so far, no).

¹ [6 pages]

² [32 PAGES TOTAL]

Chapter 5

Analysis of SMELLIE Data in the Scintillator Phase

This chapter contains two sets of analyses: measurements of the extinction lengths of the scintillator as a function of wavelength and time, as well as monitoring the Rayleigh scattering length over time.

5.1 Extinction Length Analysis

5.1.1 Motivation

- Explain motivating observations for this analysis: a substantial discrepancy between MC and data seen in the radial profiles of nhitsCleaned for ^{210}Po after the PPO top-up campaign. Performed by Serena.
- Hypothesis of a shortening of the absorption/scattering length proposed, further strengthened by Ben Tam’s ex-situ absorption measurements with scintillator taken from the detector, as well as knowledge about a likely “cooking” of PPO during the PPO-fill.

76

Analysis of SMELLIE Data in the Scintillator Phase

- Describe the provisional optics model decided on based on these measurements, which includes an additional non-re-emitting component of the scintillator.
- As a further cross-check, SMELLIE should be sensitive to changes in the overall extinction length of the scintillator, especially for short extinction lengths relative to the size of the detector.
- More straightforward in measuring extinction length compared to scattering length — no need to distinguish between scintillator re-emission and scattering.
- Further uses: can be used to monitor the extinction length over time!

[4 pages]

10 5.1.2 Analysis Approach**11 The 2-PMT Case**

- Outline theoretical approach for how one could measure the extinction length of scintillator through a comparison of SMELLIE data between the scintillator and water phases, in the simplified 1 dimensional case with only 2 PMTs.

[3 pages]

16 Combining Results Between PMTs

- Not doing analysis with just 2 PMTs, of course! Can combine results from multiple PMTs within a beamspot: I explain how here.

[2 pages]

Corrections Between the Water and Scintillator Phases

- Note the complications that we have to deal with. Namely, the differing refractive indices of the media bending the beamspot differently in the phases, as well as the method used to estimate t_{emm} .
- Explain how we deal with these, the former through MC simulation.

[2 pages]

5.1.3 Validation of the Analysis in Simulation

- Show results of this approach being used to measure the extinction length in simulation. How well does it do?

[3 pages]

5.1.4 Results in Data

- Describe the data used in this analysis, both water and scintillator, which can be shown in a table.
- Show examples of analysis of data in action for 375 nm data: typical t_{res} distributions of backscattered and beamspot PMTs; calculation of that particular extinction length measurement, followed by the graph for extinction length in 375 nm over all fibres and time periods.
- Discuss what results can be seen in this plot: consistency between fibres, the expected change as a function of PPO concentration, and stability of the extinction length during the main 2.2 g/L scintillator phase.

- Compare results to those made by Ben ex-situ: are they in agreement? If not, what possible systematics could there be? The main one for my analysis is likely to be uncertainties in the simulated beam profile that leak through into the refractive index correction of the beamspot. For the ex-situ analysis, the value of the extinction length obtained is achieved through background subtraction at some long wavelength, and the particular choice of this wavelength can lead to systematic changes in the obtained extinction length.
- Look at results at longer wavelengths: can anything reasonably be said at these longer wavelengths? Why/why not?
- Finally: describe any conclusions that can be reached, in particular whether we can affirm the optics model we use in RAT.

[8 pages]

5.2 Scattering Analysis

5.2.1 Historical Approaches and the Problem of Systematics

- Comparison to MC is necessary in scattering analysis, compared to merely being needed as a correction factor. This is because of the angular dependence of scattering. As a result, we can be far more susceptible to systematics from poor modelling!
- As a warning, show how Krish's/Esther's approach to the SMELLIE scattering analysis suffers majorly from these systematic effects. Requires describing their analysis approach briefly, and then explaining how the systematics described in Section 4.3 lead to major problems with this approach.

5.2 Scattering Analysis

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- Motivates the need for either reduced systematics, or an alternative analysis approach that is more robust to them! 1
2

[2 pages] 3

5.2.2 New Methodology 4

Signal Region Selection 5

- Propose the new analysis approach: looking at light in the “bad light-path” PMT region. Define what this region is. 6
7
- Give qualitative argument for why we expect this region to be robust to the beam profile systematics: dominated by the scattered signal as no direct light can make it here, and changes to beam profile should get “smeared out” after scattering. 8
9
10
11
- Show how simulations indicate this should be a region with a very high purity of scattered light, and (assuming all else being equal) robust to beam profile uncertainties. 12
13
14
- Confirm robustness of selected PMT region to uncertainties in AV offset and fibre position. 15
16

[5 pages] 17

Measuring the Emission Intensity 18

- Remaining systematics is now in the calculation of an average absolute emission intensity. 19
20

- 1 • Show how various methods don't work particularly well: whole detector npe,
- 2 beamspot npe, backscattered light npe, "bad light-path" PMTs but at later times.
- 3 Explain why it goes wrong for each method.
- 4 • Look at "beamspot but excepting the central bit": if that works well, then we
- 5 can continue!
- 6 • Otherwise, we'll have to live with measuring relative scattering lengths instead
- 7 of absolute amounts, using the outer water back-scattering as a measure of the
- 8 relative emission intensity.

9 [4 pages]

10 5.2.3 Results

- 11 • Actually do the proposed analysis on data, versus time and wavelength. Do the
- 12 results seem consistent between fibres? Are they sensible values?

13 [5 pages] [33 PAGES TOTAL]

Chapter 6

Solar Oscillation Analysis

Driving out into the Sun

Let the ultraviolet cover me up

Looking for a Creation Myth

Ended up with a pair of black lips

This is the End

PHOEBE BRIDGERS

Measuring the “solar” neutrino oscillation parameters Δm_{21}^2 and θ_{12} is one of the principal aims of the SNO+ detector during the scintillator-phase. There are, in fact, two complementary methods of measuring these parameters: the oscillations of anti-neutrinos from terrestrial nuclear reactors, and the oscillations of neutrinos from the Sun.

This chapter focuses on the latter approach, using ${}^8\text{B}$ neutrinos coming from the Sun to measure the solar oscillation parameters. An initial background-free study was performed by Javi Caravaca [49], which demonstrated that it was indeed possible to make such a measurement in the detector. The work in this chapter builds on substantially from that analysis. This chapter also draws on the associated reactor anti-

¹ neutrino analysis built by Iwan Morton-Blake [8], and more broadly from the general
² techniques used in the $0\nu\beta\beta$ analysis of Tereza Kroupova [47] and Jack Dunger [50].

³ This chapter begins by explaining how it is possible to measure the solar oscillation
⁴ parameters via ^8B events. Then, the framework used to perform the analysis is then
⁵ explained: that of a *Bayesian Analysis using Markov Chain Monte Carlo techniques*.
⁶ After the method has been described, the dataset upon which the analysis is performed
⁷ is then introduced. The results and associated validation are then given. Given these
⁸ results, a projection is then made for the expected sensitivity to θ_{12} as a function of
⁹ livetime.

¹⁰ 6.1 Analysis Methodology

¹¹ 6.1.1 Observational Principle

¹² How can we measure neutrino oscillation parameters via solar neutrinos in the SNO+
¹³ detector? As discussed in Chapter 1, it is possible to detect all flavours of neutrino
¹⁴ through elastic scattering with electrons in the detector. If this interaction was purely
¹⁵ neutral-current, then there would be no way of telling the flavour-state of an interacting
¹⁶ neutrino. However, electron neutrinos are able to interact through an additional
¹⁷ charged-current mode. This modifies the cross-section for electron neutrinos, and
¹⁸ means that as the survival probability for electron neutrinos generated from the Sun,
¹⁹ P_{ee} , is modified, the interaction probability of neutrinos with the detector will also.
²⁰ The resulting differential interaction rate R for solar neutrinos as a function of the
²¹ neutrino's energy, E_ν , is given by:

$$\frac{dR}{dE_\nu} = \Phi_{^8\text{B}} S_\nu(E_\nu) n_e \left[P_{ee}(E_\nu) \sigma_{\nu_e}(E_\nu) + (1 - P_{ee})(E_\nu) \sigma_{\nu_\mu, \tau}(E_\nu) \right], \quad (6.1)$$

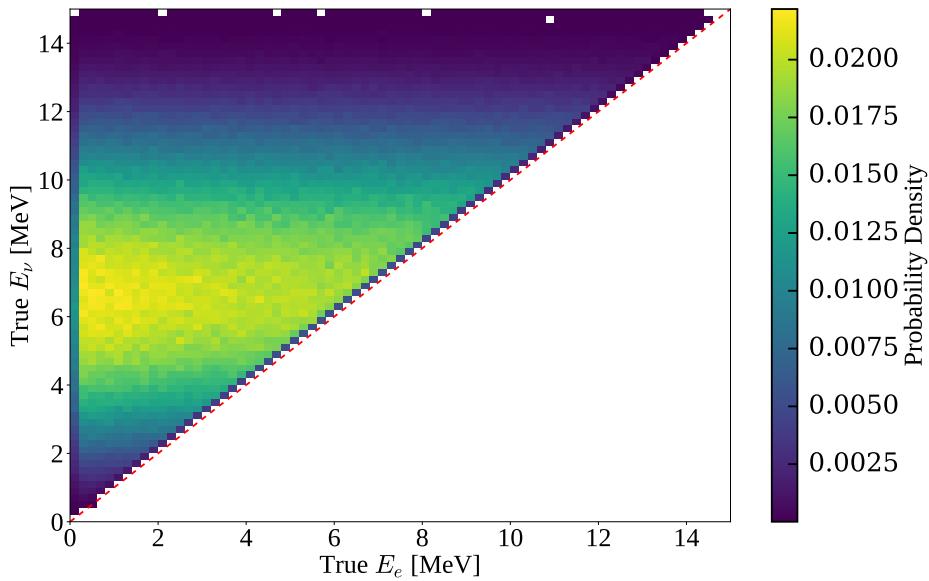


Fig. 6.1: 2D probability distribution comparing the true neutrino energy from a ${}^8\text{B}$ ν_e to the true energy of the scattered electron. Also shown in red is the line $E_\nu = E_e$.

where $\Phi_{{}^8\text{B}}$ is the total flux of ${}^8\text{B}$ solar neutrinos, S_ν is the normalised incident energy spectrum of the solar neutrinos, n_e is the number of electron targets in the detector detection medium, and σ_{ν_i} is the neutrino-electron elastic scattering cross-section for flavour i . 1
2
3
4

Of course, we do not directly measure neutrino energies in the detector — only the associated scattered electron. If there were no correlation between the observed electron energy and its associated neutrino, then the only effect of neutrino oscillations would be to change the overall observed rate of events due to this process. There would be no change in the shape of the event's energy spectrum, even though neutrino oscillations are a function of neutrino energy. Fortunately, there is some dependence of E_ν on that of the scattered electron, E_e . This dependence can be seen in Fig. 6.1 for ${}^8\text{B}$ electron neutrinos interacting in SNO+. As can be seen, the dependence is weak, and comes mostly from basic energy conservation: If one observes a 10 MeV electron event in the detector, it can't reasonably have come from a 5 MeV neutrino. 5
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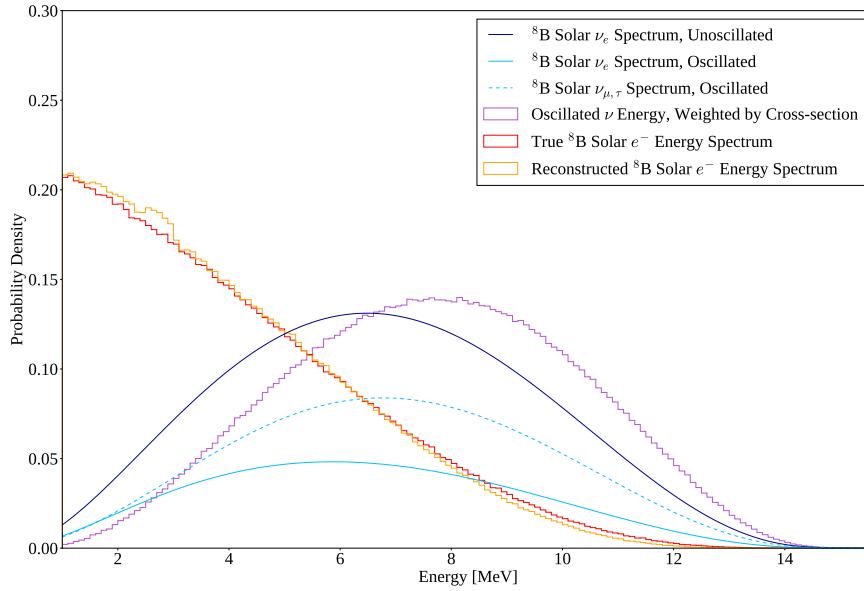


Fig. 6.2: The evolution of energy distributions related to ${}^8\text{B}$ solar neutrino detection. The unoscillated neutrino spectrum is taken from [51]; neutrino oscillations assume oscillation parameters from the current global fit results [52] and ν_e survival probabilities calculated via the method described in Section 6.1.8. The latter three distributions were obtained from MC production as described in Section 6.2.1, with the cross-section formula coming from [53].

In Fig. 6.2 we can see the impact each physical process has on the energy spectrum that we eventually observe. We start with a broad energy distribution of ${}^8\text{B}$ electron neutrinos generated in the Sun. These neutrinos then oscillate their flavour state as they propagate to the detector, in an energy-dependent manner. When a (tiny) fraction of these neutrinos interact with the electrons in our detector, there is both an energy- and flavour-dependence on the cross-section. The scattered electrons gain a kinetic energy with some mild dependence on the inciting neutrino's energy, which is then measured by the detector to within some energy resolution.

Let us now consider the dependence of P_{ee} on the individual neutrino oscillation parameters. Recall from Eq. 1.2 that, after considering matter-induced oscillations due to neutrinos passing through the Sun and possibly the Earth, $P_{ee} = P_{ee}(\tan 2\theta_{12}^M, \sin \theta_{13}^M, \Delta m_{21,M}^2) = P_{ee}(\theta_{12}, \theta_{13}, \Delta m_{12}^2, \Delta m_{13}^2)$. Fig. 6.3 shows the dependence of each of these four oscillation parameters on $P_{ee}(E)$. We can see that in reality

Fig. 6.3

only the two parameters Δm_{21}^2 and θ_{12} have a substantial impact on $P_{ee}(E)$ and hence the observed electron energy spectrum. Because of this, for this analysis we will only ever vary these two oscillation parameters, and keep θ_{13} and Δm_{13}^2 at their current global fit values¹ of $\sin^2 \theta_{13} = 0.0222$ and $\Delta m_{13}^2 = +2.515 \times 10^{-3} \text{ eV}^2$ [52].

6.1.2 Background Processes

Sadly, elastically-scattered electrons from ${}^8\text{B}$ neutrinos are not the only events we see in the SNO+ detector during the scintillator phase. There are a number of background processes that our signal must compete against. Below a reconstructed energy of $\sim 2.5 \text{ MeV}$, it is known that various backgrounds completely swamp any possible ${}^8\text{B}$ signal, and so for this analysis we only consider processes that can generate reconstructed energies of at least $E_{\min} = 2.5 \text{ MeV}$. The following subsections explain each of these backgrounds, as well as methods that have been used to mitigate them as much as possible.

Internal Uranium- and Thorium-Chain Backgrounds

Although every effort has been made to make the scintillator cocktail that fills SNO+ to be as radio-pure as possible, there inevitably remain trace amounts of the radioactive isotopes that derive from the decay chains of the ${}^{238}\text{U}$ and ${}^{232}\text{Th}$ isotopes. Fig. 6.4 shows these two decay chains. Fortunately, only a fraction of the radioactive isotopes in these chains actually are capable of generating events in the detector with energies above E_{\min} : these have been highlighted in Fig. 6.4 in gold.

¹We use the global fit results excluding Super-Kamiokande’s atmospheric data, and assuming normal ordering of the neutrino mass hierarchy. This choice has a tiny impact on the magnitudes of these two fixed parameters, the main impact being the sign of Δm_{13}^2 .

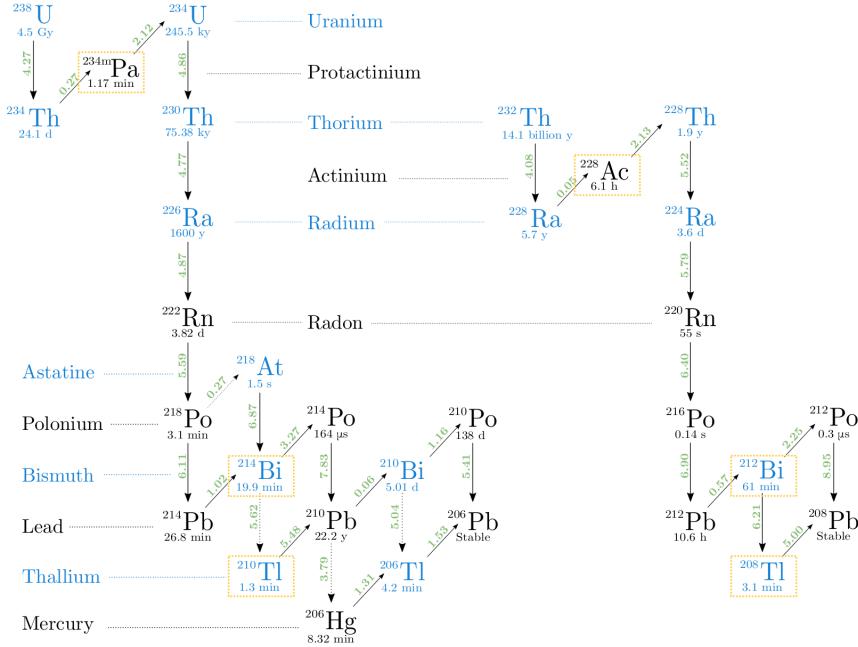


Fig. 6.4: The ^{238}U and ^{232}Th decay chains, taken from [47]. Isotopic half-lives are given below their symbol; the Q-values for each decay, in MeV, is given in green. Downward arrows indicate an α -decay; diagonal arrows indicate β -decay. Isotopes highlighted in gold are potential backgrounds for this solar analysis.

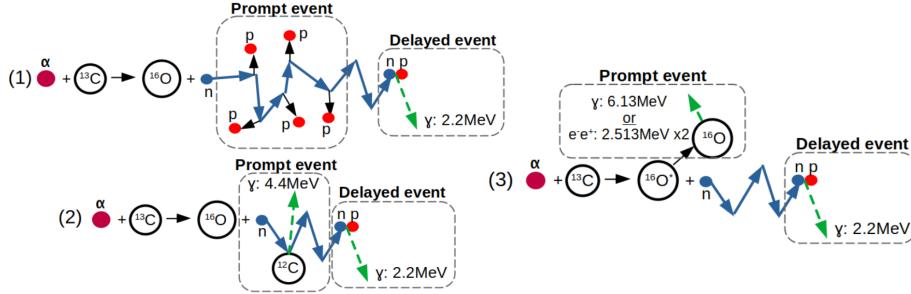
1 Of particular note are the decays of ^{212}Bi and ^{214}Bi . Both are capable of either α - or
 2 β -decays to Tl or Po isotopes, respectively. For the former, it is the subsequent β -decay
 3 of the Tl that can have a reconstructed energy above E_{\min} . For the latter, the Bi
 4 decay is the part of the pair of decays that can lie above E_{\min} . Although the α -decays
 5 here certainly have Q-values well above 2.5 MeV, the liquid scintillator quenches the
 6 observed energy well below E_{\min} . The so-called “Bi–Po” decays are particularly special
 7 because the lifetimes of ^{212}Po and ^{214}Po are 300 ns and 164 μs , respectively, which are
 8 short enough to allow for highly-effective coincidence tagging.

9 There are two classes of Bi–Po event in the detector: “out-of-window” (OOW)
 10 events for which the Bi and Po occur in separate event windows, and “in-window” (IW)
 11 events whereby the Bi and Po occur within the same event window. These lead to two
 12 distinct strategies for tagging these kinds of events. For out-of-window Bi–Pos, we
 13 look for a delayed coincidence of two events. Using the tagging algorithm suggested

in [] as a starting point, the chosen procedure was as follows. There must be two events that trigger the detector within $4\mu\text{s}$ of one another, and both have a valid `scintFit` position reconstruction within 2 m of one another. The delayed candidate event must also have at least 100 cleaned PMT hits. This very broad coincidence tagging procedure was designed to ensure that the cut was as *efficient* in tagging (and hence, rejecting) Bi–Po as possible, whilst negligibly impacting the solar signal. This is in contrast to the cuts chosen by Rafael Hunt-Stokes in [], which try and obtain a highly *pure* sample of Bi–Po tags.

Of course, the above delayed coincidence procedure cannot catch any of the in-window Bi–Po events. For these, we use a different approach. Because two decays happened in the same event window, we expect to see two distinct peaks in the event’s time residual spectrum. In order to look for this event topology, a likelihood-ratio classifier was run over events, first developed by Eric Marzec [] and re-coordinated for the 2.2 g l^{-1} LABPPO scintillator optics by Ziping Ye []. This classifier calculates the likelihood ratio between the null hypothesis of a $0\nu\beta\beta$ event (a proxy in this analysis for single-site events such as our ${}^8\text{B}$ signal) and the alternative hypothesis of an in-window Bi–Po event. The more negative the value of the result, `alphabeta212`, the greater the evidence there is for rejecting the null hypothesis of a single-site event. Events with $\text{alphabeta212} < 0$ were then rejected.

Combining both out-of-window and in-window Bi–Po tagging, the impact on ${}^{212}\text{Bi-Po}$, ${}^{214}\text{Bi-Po}$, and ${}^8\text{B } \nu_e$ events can be seen in Fig. 6.5. We consider here only events that pass all other cuts used in this analysis: see Section 6.2.2 for the specifics of the cuts used. Because of the different lifetimes of the decays, ${}^{214}\text{Bi-Po}$ decays predominantly fall out-of-window whilst ${}^{212}\text{Bi-Po}$ events are typically in-window. This explains why the out-of-window tagging is substantially better at cutting ${}^{214}\text{Bi-Po}$ decays, whereas the in-window tagging far better tags ${}^{212}\text{Bi-Po}$ decays. Overall, within

Fig. 6.5**Fig. 6.6:** Schematic of the three dominant modes of (α, n) interaction, taken from [8].

the analysis region of interest (ROI), the two combined cuts are able to tag TODO% of $^{214}\text{Bi}-\text{Po}$ triggered events, TODO% of $^{212}\text{Bi}-\text{Po}$ triggered events, whilst retaining TODO% of $^8\text{B} \nu_e$ signal events.

4 (α, n) Reactions

The impact of ^{238}U - and ^{232}Th -chain isotopes does not simply end at their direct decays. It is possible for the α s generated during these decays to undergo their own interactions with nuclei in the detector. Within the organic scintillator of SNO+, the dominant interaction of this type is when an α collides with a ^{13}C nucleus, emitting a neutron: $\alpha + ^{13}\text{C} \longrightarrow ^{16}\text{O} + n$. This is known as an (α, n) reaction.

The topology of this reaction in the detector is a delayed coincidence, as shown in Fig. 6.6. For the prompt signal, there is the light emitted from the α just before, and the n just after the (α, n) . The neutron then thermalises and gets captured by another nucleus — usually hydrogen in SNO+ — which creates an excited state that then eventually decays, creating a γ that creates the delayed signal in the detector [].

As can be seen in Fig. 6.7, (α, n) interactions can lead to events reconstructed at a wide variety of energies, which could be an issue for this analysis. However, because they are delayed coincidence events with a typical decay time of ~ 100 ns,

Fig. 6.7

the aforementioned out-of-window Bi–Po tagging algorithm also efficiently tags (α, n) events. Looking again at Fig. 6.7, simply by using the out-of-window Bi–Po tagger without any further modifications TODO% of events in the ROI are cut.

External Backgrounds

All materials within the SNO+ detector are radioactive, not just the liquid scintillator cocktail. This includes the acrylic, ropes, external water, and PMTs. These components have had their radiopurity “assayed” (that is, measured) throughout the detector’s lifetime, often back to the construction of the original SNO detector itself. The materials other than the liquid scintillator are known to have far higher background levels, especially in the important ^{238}U - and ^{232}Th -chain backgrounds [54]. To distinguish between the inherent backgrounds within the scintillator, and the backgrounds from materials at larger radii, we use the terminology “internal” and “external”, respectively.

Although there are numerous external backgrounds, with a suitably accurate and precise position reconstruction algorithm they can be suitably handled. The simplest approach is with a so-called “fiducial volume” (FV) cut: just throw out all events that reconstruct beyond some radius. The only external background events that will reach within the FV are those that have reconstructed very poorly, or have some long-distance radiation that manages to deposit radiation close to the centre of the AV. Because α and β radiation can only travel short distances through the detector, it is only γ radiation that can realistically travel far enough into the detector to be able to reconstruct anywhere near the centre. Moreover, the intensity of this γ radiation attenuates exponentially towards the centre of the detector, meaning only a tiny fraction of the total number of external events reconstruct within a 3.5 m FV, say. This strong radial-dependence can be seen in Fig. 6.8.

Fig. 6.8

What this figure also demonstrates is that our solar signal has a completely different radial dependence to these backgrounds. As a result, if one considers not just the energy of events but also their reconstructed radius, then it is possible to get an additional handle on the external backgrounds. The FV cut can then be pushed further out to larger radii, allowing one to gain more signal statistics.

Work by Tereza Kroupova [47] allows for additional means of distinguishing external backgrounds from the solar signal. The underlying assumption in the reconstruction of SNO+ events is that there was an electron at a single point, which is entirely valid for ^{8}B elastic scattering events. However, external backgrounds can fail this assumption in two ways. Firstly, these radioactive decays often generate γ radiation in addition to the main α/β particle, which creates a multi-site event. Because the `scintFit` position reconstruction algorithm is not prepared for a distribution of energy depositions in the scintillator, the t_{res} distribution will broaden. This allows an event classifier to be built that distinguishes between the t_{res} distributions of single-site events and externals, known as the “external background timing classifier”. Secondly, because external backgrounds that do reconstruct close to the centre of the detector typically have a γ that travelled a long distance towards the centre of the detector, we expect the earliest light that hits the PMTs to arrive most often along the direction of the reconstructed position vector. A distribution of PMT hits for a given event as a function of their angular distribution relative to the direction of position reconstruction can be built, and compared to the expected distributions for single-site and external background events. This is known as the “external background topological classifier”. Much like the classifier described in Section 6.1.2, the single-site events used for comparison were $0\nu\beta\beta$ events, but these have a similar single-site structure to the solar signal of interest in this analysis.

Fig. 6.9

Fig. 6.9 shows the correlation between the two classifier results for both a typical external background, and ${}^8\text{B} \nu_e$ events.

6.1.3 The Log-likelihood Test Statistic

At the highest level, this analysis involves taking the data observed in the scintillator-fill after applying a certain set of cuts, along with simulated PDFs for all processes believed to build up the observed data with those same cuts applied, and then attempting to fit the combined energy and radial distributions of the MC to that of the data. Given a set of PDFs, to try and match the distribution of observables in data we can modify a number of parameters. These consist of the normalisations of each PDF (i.e. the total number of events observed due to that process), and any systematic parameters that could modify the shapes of these distributions. For this analysis, the neutrino oscillation parameters act as *de facto* systematic parameters, as they modify the shape of the ${}^8\text{B}$ PDFs. Of course, unlike usual systematics the oscillation parameters are what we are actively trying to measure, instead of being a nuisance.

In order to perform a fit to data in this way, we must first answer a set of questions:

1. Which signal and background processes must we consider?
2. In addition to their normalisations, are there any further parameters necessary to specify the distributions of the PDFs for each of the processes? Systematics and oscillation parameters are good examples.
3. What is our test statistic?
4. What algorithm do we use to try and find the best-fit result?
5. How do we measure uncertainties on these best-fit values for each parameter?

¹ In section 6.1.2, question 1 was answered for this analysis. We now give the answer to
² question 3; all other questions on this list will be answered shortly.

³ The test statistic used for this analysis is the *binned extended log-likelihood*. Once
⁴ the data and MC PDFs have been binned in both the observables of interest, it is
⁵ assumed that the expected number of events in a given bin j is governed by a Poisson
⁶ distribution:

$$\text{⁷} P_j(n_j|\lambda_j) = \frac{\lambda_j^{n_j} e^{-\lambda_j}}{n_j!}, \quad (6.2)$$

⁸ where $P_j(n_j|\lambda_j)$ is the probability of observing n_j events in bin j , given an expectation
⁹ of λ_j events in total from signal and background processes in that bin. This λ_j can be
¹⁰ decomposed into each of the expected rates for each process, i :

$$\text{¹¹} \lambda_j = \sum_{i=1}^{N_{\text{PDFs}}} \mathcal{N}_i P_{ij}(\boldsymbol{\theta}), \quad (6.3)$$

¹² where N_{PDF} is the number of PDFs being considered in the analysis, \mathcal{N}_i is the nor-
¹³ malisation parameter of the i^{th} PDF, and $P_{ij}(\boldsymbol{\theta})$ is the probability of observing an
¹⁴ event of process type j in bin i , assuming a set of non-normalisation parameters $\boldsymbol{\theta}$. By
¹⁵ combining the probabilities of all the bins together, the total probability for a given
¹⁶ set of processes assuming these parameters to give rise to the data seen is:

$$\text{¹⁷} P(\mathbf{n}|\mathcal{N}, \boldsymbol{\theta}) = L(\mathcal{N}, \boldsymbol{\theta}|\mathbf{n}) = \prod_{j=1}^{N_{\text{bins}}} \frac{\left[\sum_{i=1}^{N_{\text{PDFs}}} \mathcal{N}_i P_{ij}(\boldsymbol{\theta}) \right]^{n_j} e^{-\sum_{i=1}^{N_{\text{PDFs}}} \mathcal{N}_i P_{ij}(\boldsymbol{\theta})}}{n_j!}, \quad (6.4)$$

¹⁹ where N_{bins} is the total number of bins being considered in the analysis. This probability
²⁰ can be re-framed as the likelihood of the vectors of parameters \mathcal{N} and $\boldsymbol{\theta}$ given
²¹ the vector of number of events in each bin, \mathbf{n} : $L(\mathcal{N}, \boldsymbol{\theta}|\mathbf{n})$. It is rare to see the
²² likelihood as-is, instead, for computational purposes the log-likelihood is used instead,
²³ $\mathcal{L}(\mathcal{N}, \boldsymbol{\theta}|\mathbf{n}) := \ln L(\mathcal{N}, \boldsymbol{\theta}|\mathbf{n})$. We can then get to the formula actually used for this

analysis:

$$\mathcal{L}(\mathcal{N}, \boldsymbol{\theta} | \mathbf{n}) = - \sum_{i=1}^{N_{\text{PDFs}}} \mathcal{N}_i + \sum_{j=1}^{N_{\text{bins}}} n_j \ln \left(\sum_{i=1}^{N_{\text{PDFs}}} \mathcal{N}_i P_{ij}(\boldsymbol{\theta}) \right). \quad (6.5)$$

6.1.4 The Bayesian Statistical Approach

There are two main schools of statistical inference, “Frequentist” and “Bayesian”. In the former, probabilities describe the fraction of times a situation can be found within the whole ensemble of possible worlds. For the latter, we care not about an ensemble of worlds but instead our degree of belief in this current one. We update our beliefs as we acquire knowledge of the world through Bayes’ Theorem:

$$P(\boldsymbol{\mu} | \mathbf{x}) = \frac{\mathcal{L}(\boldsymbol{\mu} | \mathbf{x}) P(\boldsymbol{\mu})}{P(\mathbf{x})}. \quad (6.6)$$

Here, $\boldsymbol{\mu}$ is the set of parameters that model our system, $P(\boldsymbol{\mu})$ is our *prior* (pre-existing) distribution for those model parameters, and \mathbf{x} is the data taken in our experiment. The updated, *posterior* distribution $P(\boldsymbol{\mu} | \mathbf{x})$ is then the prior multiplied by the likelihood of parameters $\boldsymbol{\mu}$ given observations \mathbf{x} , $L(\boldsymbol{\mu} | \mathbf{x})$, and divided by the total probability $P(\mathbf{x})$ of observing \mathbf{x} under any circumstance.

Both approaches to statistics are widely-used in statistical analysis, in both particle physics and beyond. The Bayesian approach was used for this analysis, as it was believed that this helps keep transparent what assumptions are being made in the analysis.

Now, if one is able to determine the overall posterior distribution, then it is possible to derive best-fit values with uncertainties for all parameters in the fit. This is done by “marginalising” the posterior distribution, i.e. integrating over all parameters other than the one of interest. A sensible best-fit value is then the modal marginalised posterior density, the highest value in the marginalised distribution. The uncertainty

on this value is derived from the spread of the marginalised posterior, by the calculation of the 1σ Credible Interval (CI): this is the set of values for a given parameter which has a total posterior probability of 68.3%, and contain the best-fit value. There are in fact an infinite number of CIs that satisfy this property; for this analysis, the values are chosen in decreasing order of marginalised posterior probability density.

6.1.5 Markov Chain Monte Carlo

Of course, all of this assumes that one can accurately determine the posterior density distribution. Whilst the likelihood and prior distribution are straightforward enough to calculate, often-times $P(\mathbf{x})$ (which acts as a normalisation) is very challenging to determine. This is because calculating this normalisation involves integrating the likelihood over all the parameter space, and if there are a large number of parameters this can become enormously numerically complex. An alternative approach is needed!

That alternative comes in the form of *Markov Chain Monte Carlo*, MCMC. A Markov Chain is any mathematical system for which the next state of the system is dependent only on its current state; the system is in some sense “memoryless”. For a large class of Markov Chains — those that are “ergodic” and “aperiodic” — one can prove that regardless of the initial position on the chain, the probability distribution converges to the same distribution π . MCMC uses such a Markov Chain which attempts to converge towards the posterior density distribution in particular. In MCMC, after choosing the initial position in the parameter space, successive states are chosen at random with a probability dependent only on the properties of the current position in parameter space and the proposed position. The convergence property of Markov Chains means that the set of steps made in the parameter space after some initial number of steps will have a distribution that converges to that of the posterior density distribution.

There are a number of MCMC algorithms, and the particular one used in this analysis is that of the *Random-Walk Metropolis-Hastings Algorithm*. In this algorithm, after the initial position in the parameter space $\boldsymbol{\mu}$, a new step is proposed some distance from the current one, $\boldsymbol{\mu}'$. This step is chosen at random from a multivariate Gaussian distribution centred on the current position, with widths in each dimension of the parameter space chosen beforehand as constants for tuning the MCMC process. This choosing of a new proposed step at random is what gives the algorithm its Monte Carlo and Random Walk titles. Once a new step is proposed, it is accepted as the new step with a probability $S(\boldsymbol{\mu}'|\boldsymbol{\mu})$ according to the condition of *detailed balance*:

$$\begin{aligned} S(\boldsymbol{\mu}'|\boldsymbol{\mu}) &= \min \left(1, \frac{P(\boldsymbol{\mu}'|\mathbf{x})}{P(\boldsymbol{\mu}|\mathbf{x})} \frac{R(\boldsymbol{\mu}|\boldsymbol{\mu}')}{R(\boldsymbol{\mu}'|\boldsymbol{\mu})} \right) = \min \left(1, \frac{L(\boldsymbol{\mu}'|\mathbf{x}) P(\boldsymbol{\mu}')}{L(\boldsymbol{\mu}|\mathbf{x}) P(\boldsymbol{\mu})} \frac{R(\boldsymbol{\mu}|\boldsymbol{\mu}')}{R(\boldsymbol{\mu}'|\boldsymbol{\mu})} \right) \\ &= \min \left(1, \frac{R(\boldsymbol{\mu}|\boldsymbol{\mu}')}{R(\boldsymbol{\mu}'|\boldsymbol{\mu})} \exp \left[\mathcal{L}(\boldsymbol{\mu}'|\mathbf{x}) - \mathcal{L}(\boldsymbol{\mu}|\mathbf{x}) + \ln \frac{P(\boldsymbol{\mu}')}{P(\boldsymbol{\mu})} \right] \right). \end{aligned} \quad (6.7)$$

$R(\boldsymbol{\mu}'|\boldsymbol{\mu})$ is the probability density that position $\boldsymbol{\mu}'$ is proposed as a step from position $\boldsymbol{\mu}$, and vice versa for $R(\boldsymbol{\mu}|\boldsymbol{\mu}')$. In most cases, because of the use of the same multivariate Gaussian in choosing proposals, $\frac{R(\boldsymbol{\mu}|\boldsymbol{\mu}')}{R(\boldsymbol{\mu}'|\boldsymbol{\mu})} = 1$ simply. This component only becomes important at the edges of the parameter space, preventing the sampling probability being incorrectly impacted if a proposed step goes outside the allowed parameter space.

It is the detailed balance condition that ensures convergence of the MCMC algorithm to specifically the posterior density distribution. Crucially, because it is only dependent on the ratio of posterior densities, the hard-to-calculate normalisation $P(\mathbf{x})$ in both posterior density terms cancels out, meaning one only needs to calculate the likelihood and priors for each step, as well as $\frac{R(\boldsymbol{\mu}|\boldsymbol{\mu}')}{R(\boldsymbol{\mu}'|\boldsymbol{\mu})}$.

The specific implementation of MCMC used for this analysis is that of `OXO`, a C++ analysis framework first developed by Jack Dunger [50]. `OXO` is able to run the Metropolis-Hastings algorithm on multidimensional binned data, using the log-

¹ likelihood defined in 6.1.3. This framework also allows one to include systematic
² parameters that can float within the fit, and define non-uniform priors for normalisations
³ and systematics that have constraints: the details of this will be discussed shortly.

⁴ 6.1.6 Choosing Priors

⁵ For this analysis, the suggestions made by Biller & Oser in [55] about choosing prior
⁶ distributions are followed: for parameters that do not have some pre-existing constraint,
⁷ a flat prior is used. A nice consequence of this choice is that $\ln \frac{P(\mu')}{P(\mu)} = 0$, so the actual
⁸ value of the prior for these variables never needs to be calculated when running the
⁹ MCMC algorithm. For the bulk of this analysis, uniform priors are assumed on the
¹⁰ neutrino oscillation parameters Δm_{21}^2 and θ_{12} , as the magnitudes of these parameters
¹¹ are now well-established.

¹² For parameters with existing asymmetric constraints $\beta_{-\sigma_-}^{+\sigma_+}$, this analysis uses an
¹³ asymmetric Gaussian prior, equivalent to the logarithm of the prior being an asymmetric
¹⁴ quadratic:

$$\ln P(\mu) = \mathcal{A} - \begin{cases} \frac{(\mu-\beta)^2}{2\sigma_+^2} & \text{if } \mu \geq \beta, \\ \frac{(\mu-\beta)^2}{2\sigma_-^2} & \text{if } \mu < \beta. \end{cases} \quad (6.8)$$

¹⁶ Here, \mathcal{A} is the logarithm of the prior's normalisation constant, and cancels out in the
¹⁷ detailed balance condition. For parameters with symmetric constraints, $\sigma_+ = \sigma_-$, then
¹⁸ $\ln P(\mu)$ reduces to a quadratic with maximum at $\mu = \beta$.

¹⁹ 6.1.7 Including Systematics in the Fit

²⁰ One important implementation detail is how systematics are applied within the MCMC
²¹ fitting process. Once systematics are added to the fit, at every step the binned PDFs
²² for all the processes considered in the fit must get modified appropriately, which
²³ can become extremely computationally-intensive if not approached carefully. The

Fig. 6.10

strategy used in the **OXO** framework starts by thinking of the contents of a binned PDF as a vector of bin probabilities, $\mathbf{p} = (p_1, p_2, \dots, p_{N_{\text{bins}}})^T$. Then, we can think of a systematic acting on the PDF as a linear transformation, and hence a matrix S acting on this vector: $\mathbf{p}' = S\mathbf{p}$. We only need to calculate this matrix once for a given set of systematic parameter values, and can then use the same matrix on all the PDFs in the fit. Furthermore, when multiple systematics are applied, the matrix for each systematic can then be combined via matrix multiplication into one single “detector response” matrix. **OXO** uses the **Armadillo** [] linear algebra package for efficient matrix manipulation.

There is a problem that can arise when considering the impact of systematics near the edge of the analysis ROI. Many systematics such as shifts, scalings, and convolutions use information about the contents of nearby bins to determine the contents of a particular bin. However, for bins near the edge some of that information does not exist — it has been lost to the cuts that define the ROI. This can lead to a bias in the generation of the modified PDFs, and therefore also the posterior distribution.

As an example, consider the impact of an energy scale systematic on the energy distribution of ^{234m}Pa events in the detector, shown in Fig. 6.10. Because the events seen for this process in the ROI are merely the high-energy tail, any systematic energy scaling $E'_{\text{reco}} := \beta E_{\text{reco}}$ should have a large impact on the number of events observed at the low end of the ROI. However, given that the information about data below E_{min} is lost to the ROI cuts, any energy scaling of $\beta > 1$ will not be applied correctly at all.

The solution to this problem is defining a “buffer region” of bins on either side of the ROI, which allow for tracking of events in and out of the ROI due to systematics, but aren’t considered when calculating the likelihood. This is also shown in Fig. 6.10. After the scaling systematic is applied, although incorrect bin values are found in the

¹ buffer region, this is fine because we are no longer calculating the likelihood with those
² bins. Note that because of this modification, the normalisation parameters we put into
³ the model no longer represent the expected number of events in the ROI. Instead, they
⁴ represent the number of events expected in both the ROI and buffer region, before any
⁵ systematics have been applied.

⁶ 6.1.8 Including Oscillations in the Fit

⁷ Within the analysis MCMC code, the process of neutrino oscillations are thought of
⁸ as a *de facto* systematic that acts only on the ${}^8\text{B}$ ν_e and ν_x signal spectra. Three
⁹ parameters relevant to the signal are floated within the MCMC fit: Δm_{21}^2 , θ_{12} , and
¹⁰ $\Phi_{{}^8\text{B}}$, the unoscillated ${}^8\text{B}$ flux relative to the expected rate. For the two signal PDFs, a
¹¹ third “bookkeeping” dimension is added on top of reconstructed energy and radius:
¹² the true neutrino energy, E_ν . This is necessary for correctly applying oscillations, as
¹³ the oscillation probability is a function of the neutrino’s energy, not the scattered
¹⁴ electron’s. Before the fit, these 3D PDFs are given normalisations corresponding to
¹⁵ the expectation of the number of events for each type, ν_e and ν_x , after cuts but before
¹⁶ oscillations have been applied. Strictly speaking there should be zero ν_x events before
¹⁷ neutrino oscillations: the pre-oscillation rate used here is the post-cut number of events
¹⁸ expected if 100% of the neutrinos oscillated to the ν_x type.

¹⁹ During the MCMC fit, for a given set of parameters $\boldsymbol{\theta} = (\Delta m_{21}^2, \theta_{12}, \Phi_{{}^8\text{B}})$ the
²⁰ following is performed to oscillate the signal PDFs. Firstly, the normalisations are scaled
²¹ by the factor $\Phi_{{}^8\text{B}}$. Then, for each E_ν bin the survival probability $P_{ee}(E_\nu, \Delta m_{21}^2, \theta_{12})$ is
²² calculated. Each bin then has their probability scaled by either P_{ee} or $1 - P_{ee}$, for ν_e and
²³ ν_x respectively. Of course, within the structure of the OXO framework these bin-by-bin
²⁴ scaling aren’t immediately applied, but instead a matrix describing the impact of
²⁵ oscillations on each of the PDFs is made. Because the oscillation transformation is

purely a bin-by-bin scaling, the resulting matrices are diagonal, with diagonal elements $\Phi_{\text{S}_B} \cdot P_{ee}(E_\nu, \Delta m_{21}^2, \theta_{12})$ or $\Phi_{\text{S}_B} \cdot (1 - P_{ee}(E_\nu, \Delta m_{21}^2, \theta_{12}))$ for ν_e and ν_x respectively.

After the oscillation matrix along with all other systematic matrices are applied to the signal PDFs, the PDFs are then marginalised over the E_ν dimension so that the signal PDFs match the dimensionality of all other processes.

Calculations of the survival probability are handled with **PSelmaa**, an algorithm written by Nuno Barros for the SNO 3-phase Analysis []. This considers not only the neutrino oscillations through the vacuum of space between the Sun and Earth, but also the impact of matter effects in both the Sun and Earth. This can usually be a very computationally-intensive process, but **PSelmaa** takes advantage of the assumption that the solar oscillation parameters are in the so-called “Large Mixing Angle” regime, making the calculation much faster. As seen in Section 1.2.1, previous solar oscillation experiments demonstrate that this assumption is reasonable. For this analysis, the standard MSW effect is assumed with neutrinos obeying the Normal Hierarchy, with the Sun following the **B16_GS98** metallicity model [] and the **PREM** model being used for the Earth [].

One final thing **PSelmaa** needs to know to calculate survival probabilities is the distribution of solar zenith angles during the data-taking. The solar zenith θ_z is the angle between the two following vectors: one going from the centre of the Earth through the centre of the SNO+ detector, and another starting from the detector’s centre and pointing towards the Sun. As an example, if the Sun were ever to be directly overhead the detector, both vectors would be along direction $\hat{\mathbf{z}}$ in detector coordinates, leading to a solar zenith angle of $\theta_z = 0$. The position of the SNO+ detector on Earth, as well as the times at which the detector was live, determine the solar zenith angle distribution. If not accounted for, this can lead to a bias in the result of the analysis, as a preponderance of livetime taken at night (say) would lead to a larger fraction of

¹ solar neutrinos having to pass through the bulk of the Earth to get to the detector,
² and hence the impact of the MSW effect would be greater.

³ Even after using the Large Mixing Angle approximation, having to call `PSelmaa`
⁴ numerous times for every step in the MCMC algorithm would lead to exorbitant run
⁵ times for the fitting. Therefore, a further approximation is made. Before running the
⁶ MCMC fit, `PSelmaa` is used to calculate P_{ee} over the necessary 3D space of parameters.
⁷ To get a fine scan of this space, 101 E_ν values from 1 MeV to 20 MeV, 101 Δm_{21}^2 values
⁸ from 3×10^{-6} eV² to 1×10^{-3} eV², and 151 values for θ_{12} from 5° to 65° were looked
⁹ over. This 3D grid of $101 \cdot 101 \cdot 151$ P_{ee} values is then written to disk, and loaded
¹⁰ into memory for use during the fit as a lookup table. At run-time, as the Metropolis-
¹¹ Hastings algorithm samples this 3D space the survival probability is estimated through
¹² a trilinear interpolation of the 3D grid loaded in: a version of linear interpolation for
¹³ three dimensions [].

¹⁴ 6.2 Analysis on Scintillator-Phase data

¹⁵ 6.2.1 Dataset and Livetime

¹⁶ The data used in this analysis was chosen to be scintillator phase data after the end of
¹⁷ the PPO top-up campaign that completed in April 2022. Following an initial validation
¹⁸ of the analysis tools was performed on data between 29th April and 10th May 2022 [],
¹⁹ data for the full analysis was chosen from runs taken after these dates. Not all data
²⁰ taken during this time was considered usable for this analysis, however: this is the role
²¹ of *run selection*. The ‘Preliminary Scintillator Gold’ run selection list was used as the
²² basis for this analysis. This run list requires:

- ²³ • The run type must be in ‘Physics’, as opposed to running in calibration or
²⁴ maintenance;

- The run must last at least 30 minutes; 1
- Detector electronics must be working in a stable manner without any alarms, 2
and with all crates online; 3
- There are no abnormal rates of tagged muons, and the OWL PMTs are correctly 4
functioning; 5
- There are no unusual conditions from e.g. earthquakes, blasting activity in the 6
mine, or loss of power. 7

In the end, data for this analysis used ‘Gold’ runs selected between 17th May–30th 8
November 2022, run numbers 300733–306498. The raw livetime associated with this 9
dataset was calculated by looking at the start and end times of each run using the 10
detector’s 10 MHz clock: 84.977 days. 11

Importantly, the livetime of the data actually used in the analysis ends up being 12
somewhat less than this raw value. This is because two of the cuts used in the analysis 13
(described in Section 6.2.2), the muon tagger and the high-**nhit** event tagger, also veto 14
events for a time period following a tagged event. This means that any events in such a 15
time window are automatically thrown out of consideration for analysis. An algorithm 16
was written to determine the loss of livetime from both of these tagging processes for 17
all the runs selected in this analysis, allowing for the net livetime to be calculated. To 18
ensure accuracy in the value of this lost livetime, the algorithm took care to handle any 19
overlaps in the time veto windows. This ends up being a quite common occurrence, as 20
tagged muon events and their followers often have a very high **nhit**. The net livetime 21
was calculated to be 80.615 days. Data processing and simulations used RAT versions 22
7.0.8 and 7.0.9. 23

As discussed in Section 6.1.8, in order to account for the impact of the MSW effect 24
through the Earth on P_{ee} , the solar zenith angular distribution of the dataset is needed. 25

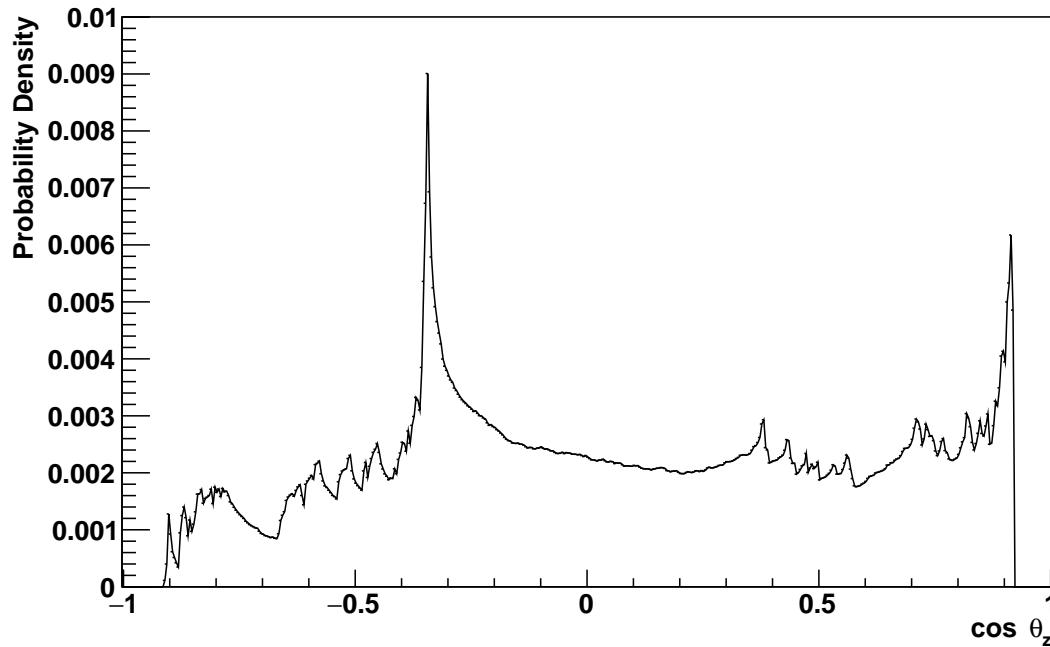


Fig. 6.11: Distribution of the cosine of the solar zenith angle, $\cos \theta_z$, for the events within this dataset. If an event has $\cos \theta_z > 0$, then at the time of that event the position of the Sun is above the detector's equator.

This was achieved by taking the recorded trigger time of each event within the dataset as given by the GPS-calibrated 10 MHz clock, determining the position of the Sun at that time, and then deriving the value of $\cos \theta_z$ given that solar position. The resulting distribution of $\cos \theta_z$ is shown in Fig. 6.11. The structure of this distribution are a result of the times of day and year of the runs selected for this analysis.

6.2.2 Event Selection

Once the runs of processed data and matching MC have been selected, the next step was to perform cuts on both in order to obtain our analysis ROI. The full list of cuts used are shown in Table 6.1, most of which have already been explained in Section 6.1.2. Most of the other cuts used in this analysis that haven't yet been discussed come from handling differences between data and MC. In MC, it is possible to simulate an

Cut Description	Cut
Event triggers detector	<code>evIndex</code> ≥ 0
Pass data cleaning	<code>ANALYSIS_MASK</code> = <code>0x2100000042c2</code>
High- <code>nhit</code> veto	<code>correctedNhits</code> ≥ 5000 ; veto for 20 s
<code>ScintFitter</code> used	<code>scintFit</code> = true
Results of <code>ScintFitter</code> valid	<code>fitValid</code> = true
In energy Region of Interest	$2.5 \text{ MeV} < E < 14.0 \text{ MeV}$
Fiducial Volume cut	$R < 5.0 \text{ m}$
Remove BiPo Out-of-window tags	See Section 6.1.2
Remove BiPo In-window tags	Tagged if <code>alphaBeta212</code> > 0 and <code>alphaBeta214</code> > 0
Remove External tags	Tagged if $E < 5.0 \text{ MeV}$, <code>ext0NuTimeTl208AVNaive</code> > -0.007 , and <code>ext0NuAngleTl208AV</code> > -4.7

Table 6.1: Cuts used in this analysis.

event but not have it trigger the detector. In this case, the stored index of that event (evIndex) will be less than zero. We never see these ‘events’ in data of course, but they do need to be removed from simulation. On the flip side, there are types of events that we observe in the detector but do not simulate. Many of these events are caught during ‘data cleaning’, which looks for a wide variety of problems in the data when the data is being processed. Any subset of the checks made during data cleaning can be chosen to consider in a given analysis by defining a specific ‘analysis mask’; a number that specifies for each of its binary digits whether to consider the data cleaning process associated with that bit. The mask chosen used in this analysis is the Collaboration’s current standard for the scintillator phase, `0x2100000042c2`, which corresponds to:

- **00 cut:** Tags every event which has a GTID ending in 00 in hexadecimal, which is associated with a long-standing issue in the detector triggering system.
- **Junk cut:** Tags events that have channels with multiple recorded hits for that event. These are associated with so-called ‘orphaned’ hits that the event building process doesn’t know what to do with.

- ¹ • **OWL cut:** Tags events that have at least three OWL PMTs hit, a sign that the
- ² event came from a cosmic-ray muon.
- ³ • **Polling cut:** Tags all events that occur during ‘polling’, a detector monitoring
- ⁴ process.
- ⁵ • **CAEN cut:** Tags all events that have missing CAEN trace data.
- ⁶ • **Muon cut:** Tags muon-like events.
- ⁷ • **Muon follower cut:** Tags all events coming up to 20 s after a tagged muon
- ⁸ event.

⁹ Beyond the formal data cleaning cuts, a `high-nhit` veto was used to tag and remove
¹⁰ instrumental backgrounds known as ‘flashers’, in which a PMT will generate large
¹¹ amounts of light in the detector. After any event with a sufficiently large number of
¹² hits, both that event and any following it for 20 s are removed. This method is likely
¹³ to be fairly conservative, and there exist different methods for cutting flashers within
¹⁴ data cleaning. However, these were originally coordinated for the water phase, and
¹⁵ hadn’t been updated for the scintillator phase at the time of writing.

¹⁶ The results of all these cuts can be seen for the data in Table 6.2. 654 out of
¹⁷ the 18,112,495,770 events within the dataset pass all cuts. These same cuts were also
¹⁸ run over the MC, with results summarised in Table 6.3. Important to note was the
¹⁹ difference between the number of physics events simulated for a given process and
²⁰ the number of triggered events. The latter ends up being larger than the former for
²¹ all processes discussed here, because a given physics event in the detector will often
²² generate a retrigger event following the primary triggered event. Also shown in this
²³ table are the combined cut efficiencies for each process. This efficiency is defined as
²⁴ the number of triggered events that pass all cuts, divided by the number of simulated
²⁵ physics events for that process.

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Cut	# Events Remaining	Cut Efficiency (%)
Before cuts	18,112,495,770	100.00
Event triggers detector	18,112,495,770	100.00
Data Cleaning	17,663,711,108	97.52
high-nhit veto	17,039,514,332	96.47
ScintFitter used	10,433,875,757	61.23
ScintFitter results valid	3,431,328,125	32.89
Energy cut	132,673	0.00387
FV cut	2093	1.58
BiPo OOW cut	817	39.03
BiPo IW cut	719	88.00
Externals cut	652	90.68

Table 6.2: Impact of each cut on the quantity of events in data.

MC process	# Physics Events Simulated	# Triggered Events Simulated	# Remaining Events after cuts	Overall Cut Efficiency (%)
$^8\text{B } \nu_e$	2,830,425	5,511,568	929,551	32.8
$^8\text{B } \nu_{\mu,\tau}$	1,898,088	3,675,261	572,949	30.2
AV ^{214}Bi	27,138,894	38,267,431	127	0.000468
Ropes ^{214}Bi	67,541,698	84,734,663	221	0.000327
AV ^{208}Tl	4,071,850	6,689,331	1471	0.0361
Ropes ^{208}Tl	6,105,956	8,758,457	2407	0.0394
External Water ^{214}Bi	81,300,835	85,056,233	62	0.0000763
External Water ^{208}Tl	36,602,319	38,963,663	2786	0.00761
PMT $\beta - \gamma$	4,069,055	6,615,601	13,570	0.333
Internal $^{212}\text{BiPo}$	9,761,554	24,331,685	8276	0.0848
Internal ^{208}Tl	816,233	1,626,536	409,629	50.2
Internal $^{214}\text{BiPo}$	4,883,417	19,293,285	1030	0.0211
Internal ^{210}Tl	815,256	1,625,012	342,415	42.0
Surface (α, n) (combined)	22,627,103	48,834,413	2238	0.00458
Internal (α, n)	1,219,279	4,740,313	655	0.0537
Internal ^{228}Ac	9,764,480	19,294,643	0	0
Internal ^{234m}Pa	16,959,995	30,806,833	1	0.0000059

Table 6.3: Combined impact of cuts on each MC process. Overall cut efficiency is defined as the number of remaining triggered events after all cuts have been applied, divided by the number of physics events simulated for that process. Processes have been split into three general categories: signal, externals, and internals. A final category corresponds to various processes that have been considered for this analysis, but either all events are cut out, or Section 6.2.3 will show a negligible number of these events are expected in our dataset.

1 Variable bin widths in energy were chosen for this analysis: 0.1 MeV bins between
 2 2.5 MeV–5.0 MeV, 0.25 MeV bins between 5.0 MeV–13.0 MeV, and a single energy bin
 3 between 13.0 MeV–14.0 MeV. Four equally-spaced bins were used in the parameter
 4 $r_3 = (r_{\text{reco}}/R_{\text{AV}})^3$ in the range $0 \text{ m} \leq r_{\text{reco}} < 5.0 \text{ m}$. The r_3 parameter was used to
 5 allow for equal volume weighting for each radial bin.

6 6.2.3 Expected Rates and their Constraints

7 Because this analysis is dependent on deriving the shape and size of the ${}^8\text{B}$ spectrum
 8 from the data observed, it is important to have expectations about the number of events
 9 expected to be seen within the dataset for each process, both signal and background.
 10 This is done in two stages: firstly, the expected rates for each process before any cuts
 11 applied is determined, along with any constraints. Then, by using the cut efficiencies
 12 calculated in the previous section, we can derive estimates for the expected number of
 13 events after cuts. When combined with a constraining uncertainty, this can be used
 14 as a prior within the MCMC fit. The resulting expected rates before and after cuts,
 15 along with any constraints being used, afer shown in Table 6.5.

16 ${}^8\text{B}$ Signal

17 Starting with the ${}^8\text{B}$ signal, numerous theoretical predictions and experimental mea-
 18 surements have been made of the ${}^8\text{B}$ flux. For this analysis, two different val-
 19 ues have been used. The SSM predicts a relatively loose constraint of $\Phi_{{}^8\text{B}} =$
 20 $(5.46 \pm 12\%) \times 10^6 \text{ cm}^{-2} \text{ s}^{-1}$ [], whereas a recent global fit of neutrino oscillation
 21 experiments leads to a much stronger constraint of $\Phi_{{}^8\text{B}} = (5.16^{+2.5\%}_{-1.7\%}) \times 10^6 \text{ cm}^{-2} \text{ s}^{-1}$ [].
 22 We shall use both for calculating the expected rate of the signal events in the detector,
 23 and their fractional uncertainties to constrain this rate.

6.2 Analysis on Scintillator-Phase data

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To calculate the number of electrons within the liquid scintillator, we follow the method used in [56], which uses the formula:

$$n_e = \frac{(f_{\text{LAB}} n_{\text{LAB}} + f_{\text{PPO}} n_{\text{PPO}}) N_A M_{\text{LAB}}}{m}. \quad (6.9)$$

Here, f_{LAB} and f_{PPO} are the fraction by weight of the LAB and PPO within the scintillator cocktail, respectively. Because this analysis is done over data taken during the scintillator phase with a PPO concentration of 2.2 g for every litre of LAB, these take values 99.744% and 0.256%, respectively. n_{LAB} and n_{PPO} are the mean number of electrons per molecule of LAB and PPO, respectively. PPO has the chemical formula $\text{C}_{15}\text{H}_{11}\text{NO}$, leading to $n_{\text{PPO}} = 116$. The LAB used in SNO+ has varying alkyl chain lengths, leading to a varying number of electrons per molecule. This distribution is known to have changed between batches of LAB made by the manufacturer, and is also impacted by the distillation process used during the purification of the LAB before it was put into the AV. At the time of writing, no final molecular breakdown has been made for the LAB within the detector; for now we shall use the breakdown provided here [], from a representative tanker truck of LAB: there $n_{\text{LAB}} = 131.68$. Finally, $N_A = 6.02 \times 10^{23} \text{ mol}^{-1}$ is Avogadro's Constant, $M_{\text{LAB}} = 780.2 \text{ t}$ is the total mass of scintillator within a sphere of radius the size of the AV, and $m = 235 \text{ g mol}^{-1}$ is the molecular weight of the scintillator. This leads to a value for the number of electron targets:

$$n_e = 2.63 \times 10^{32} \text{ electrons.}$$

We can modify Eq. 6.1 to get an equation for the total rate of solar neutrino events by flavour, before oscillations or analysis cuts are considered:

$$R_i = \Phi_{\text{sB}} n_e \int S_\nu(E_\nu) \sigma_{\nu_i}(E_\nu) dE_\nu. \quad (6.10)$$

¹ Note that for $i = \mu, \tau$, this rate corresponds to the interaction rate if 100% of the
² neutrinos have flavour $\nu_{\mu,\tau}$.

³ Using the cross-section formula from Eq. 1.1, the ^8B spectral shape from [51], and
⁴ $\Phi_{^8\text{B}} = 5.46 \times 10^6 \text{ cm}^{-2} \text{ s}^{-1}$ ($5.16 \times 10^6 \text{ cm}^{-2} \text{ s}^{-1}$), we expect a rate before considering
⁵ oscillations or cuts of:

$$\begin{aligned} \text{for } i = e, \\ R_i = \begin{cases} 2743.2(2592.5) \text{ events/yr} & \text{for } i = e, \\ 489.7(462.8) \text{ events/yr} & \text{for } i = \mu, \tau. \end{cases} \end{aligned}$$

⁷ In the above rates, we have included an additional volume correction factor of 1.0139,
⁸ because in MC events are simulated within the neck of the AV in addition to the main
⁹ spherical bulk [].

¹⁰ Internal Uranium- and Thorium-Chain Backgrounds

¹¹ As mentioned in Section 6.1.2, $^{214}\text{BiPo}$ and $^{212}\text{BiPo}$ decays, coming from the ^{238}U
¹² and ^{232}Th decay chains respectively, are capable of generating distinctive delayed
¹³ coincidence events. Rafael Hunt-Stokes was able to use a series of cuts to isolate both
¹⁴ types of coincidence signals [], looking at coincidence signals within the same 5.0 m FV
¹⁵ as my analysis. After correcting for his cut efficiencies, he obtains a rate in the whole
¹⁶ detector of []:

$$\text{BiPo rate} = \begin{cases} 0.94 \text{ events/hour} & \text{for } ^{212}\text{BiPo}, \\ 6.06 \text{ events/hour} & \text{for } ^{214}\text{BiPo}. \end{cases}$$

¹⁸ These rates assume a uniform concentration of ^{222}Rn and ^{220}Rn throughout the detector:
¹⁹ this is how internal backgrounds are simulated by default in RAT. However, it has
²⁰ been shown that this is not at all the case within SNO+; radon ingress from the AV
²¹ neck propagates through the bulk of the detector via convection currents induced

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by the thermal gradient present throughout the detector []. This radon ingress also
 1
 breaks the secular equilibrium of the two decay chains, so that one can no longer
 2
 make straightforward predictions about rates for decays in the chains above radon.
 3
 Fortunately for this analysis, Table 6.3 showed that the cut efficiencies for ^{228}Ac and
 4
 ^{234m}Pa decays, both in their respective decay chain above radon, are negligible. This
 5
 means that even if their decay rates were much greater than what would be predicted
 6
 from the tagged BiPo rates, a negligible number of events are expected for both in the
 7
 livetime of the dataset considered in this analysis.
 8

The combination of broad coincidence tagging as well as using an in-window BiPo
 9
 classifier cut ensures that the expected rate of BiPo events that survive all the cuts
 10
 is 1.5 and 2.5 for $^{212}\text{BiPo}$ and $^{214}\text{BiPo}$, respectively. Because the branching ratio of
 11
 $^{212}\text{Bi} \rightarrow ^{208}\text{Tl}$ is 36%, and the branching ratio of $^{214}\text{Bi} \rightarrow ^{210}\text{Tl}$ is 0.021% [], after
 12
 considering the cut efficiencies in this analysis we expect 512.4 internal ^{208}Tl and 1.1
 13
 internal ^{210}Tl events.
 14

In theory, the non-uniformity of all the radon-induced backgrounds could impact
 15
 the results of the oscillation analysis. However, the rates of each after cuts have
 16
 been applied mean that only the internal ^{208}Tl could plausibly have any effect. For
 17
 this particular background, each slice of the binned PDF in r_3 was allowed to float
 18
 independently, with no constraints. The other three related backgrounds considered
 19
 in the fit (^{210}Tl , $^{212}\text{BiPo}$, and $^{214}\text{BiPo}$) did not have this splitting applied to them. A
 20
 loose 25% constraint on each of these three processes' rates was applied, to account for
 21
 the current uncertainty in the calculation of R. Hunt-Stokes' tagging efficiency.
 22

 (α, n) Reactions

23

As discussed in Section 6.1.2, (α, n) reactions are induced by α -particles generated in
 24
 the detector. In the SNO+ scintillator phase, the dominant source of α -particles are
 25

¹ decays of ^{210}Po located both within the scintillator and on the AV surfaces. Because
² of substantial quenching by the scintillator, the reconstructed energy of the $^{210}\text{Po} \alpha$
³ events fall substantially below E_{\min} . The rate of the internal and surface ^{210}Po events
⁴ have been tracked throughout the scintillator phase by Serena Ricetto and Shengzhao
⁵ Yu, respectively.

⁶ Within the 5 m FV, S. Ricetto measured a total of 1.60×10^8 ^{210}Po events over
⁷ the runs considered for this analysis []. Assuming the internal rate of these events are
⁸ uniform throughout the scintillator, this leads to a total rate of 2.78×10^8 events for
⁹ all the scintillator. After using a conversion rate of 6.11×10^{-8} neutrons generated per
¹⁰ α on ^{13}C in liquid scintillator [], we can derive an expected number of internal (α, n)
¹¹ events before cuts to be 17. However, because a coincidence cut is used in this analysis,
¹² the vast majority of these should be removed, leading to merely 0.009 events expected
¹³ after cuts.

¹⁴ Similarly, S. Yu measured the surface ^{210}Po rate to vary between $2\text{--}5 \text{ Bq m}^{-2}$ over
¹⁵ the time period of this analysis as well as a function of height in the detector []. Using
¹⁶ the midpoint value of 3.5 Bq m^{-2} , one can derive an expected rate of 554.4 surface (α, n)
¹⁷ events occurring over the dataset's livetime, before cuts. Once again, the coincidence
¹⁸ cut as well as the FV cut removes the vast majority of these events, so that only 0.07
¹⁹ events are expected after all cuts have been applied. Because both classes of (α, n)
²⁰ events have negligible expected rates in our dataset, they were not included within the
²¹ MCMC fit.

²² External Backgrounds

²³ During the water phase of SNO+, an analysis of the rates of the external backgrounds
²⁴ was performed by Tony Zumbo []. The results of this analysis are shown in Table 6.4,
²⁵ giving the measured rates as a fraction of the nominal values given in []. Although the

Background Type	Rate (Fraction of Nominal)
AV & Ropes	$0.21 \pm 0.009^{+0.64}_{-0.21}$
External Water	$0.44 \pm 0.003^{+0.32}_{-0.27}$
PMTs	$1.48 \pm 0.002^{+1.65}_{-0.60}$

Table 6.4: Measured rates of the external backgrounds during the water phase of SNO+, by Tony Zumbo [].

statistical uncertainty for these measurements were quite small, the total systematic uncertainty was quite substantial because T. Zummo's analysis involved looking at events in the tail of energy distributions, so that any uncertainty in the energy scale systematic had an outsized impact on the number of events observed within that energy window. These measured rates and their systematic uncertainties were used to predict the expected rate and constrain the external backgrounds in this scintillator phase solar analysis. For the AV, ropes, and PMTs, this seems reasonable as we do not expect there to be any substantial change in these backgrounds between the phases. The exception to this is the external water, where various aspects of the water purification process over the years since T. Zummo's analysis dataset was taken []. Because of this, the water phase measured rate was used as a starting point for the solar analysis, but no constraint was applied.

Before the expected number of triggered events can be calculated for each external background process, two subtleties must be dealt with. Firstly, MC production of the externals did not include the hold-up ropes at the time of performing the analysis. This is not a major problem, because the hold-up ropes are of substantially lower mass than the hold-down ropes, and also their average radius from the centre of the detector is much larger. This means that the expected rate of events from the hold-up ropes that manage to reconstruct inside the FV should be sub-dominant to their hold-down counterparts. For the purposes of this thesis, the combined rate for both kinds of ropes

¹ were calculated, with the overall cut efficiency and derived PDFs coming from just the
² hold-down ropes.

³ Secondly, because the attenuation length of a 2.6 MeV γ particle is 23 cm in water [],
⁴ the vast majority of external backgrounds which start from a substantial distance away
⁵ from the AV do not generate γ s that make it into the scintillator. As a result, an
⁶ enormous amount of computational resources could be wasted on simulating events
⁷ that never get into the ROI. To work around this, PMT $\beta - \gamma$ events are modelled
⁸ as 2.6 MeV γ s generated at a radius of 6.2 m (just outside the AV), pointing radially
⁹ inwards. Making some basic assumptions about these events, one can derive the
¹⁰ expected survival rate of these γ particles as a function of radius [], which leads to a
¹¹ correction factor of 1.17×10^{-6} [].

¹² Similarly, simulations of background events in the external water and ropes are
¹³ restricted from starting beyond certain maximum radii depending on the specific
¹⁴ process, such that there remains a negligible impact on the fraction of simulated events
¹⁵ that actually deposit any energy into the scintillator []. This leads to rate correction
¹⁶ factors of 0.35 and 0.50 for external water ^{214}Bi and ^{208}Tl events, as well as 0.50 and
¹⁷ 0.35 for hold-down and hold-up ropes [].

¹⁸ 6.2.4 Systematics

¹⁹ There are a number of systematic effects in the analysis that could possibly have
²⁰ some impact on the resulting posterior densities of θ_{12} and Δm_{21}^2 . Consideration of
²¹ these effects is important to ensure that the uncertainty in the measurement is not
²² underestimated. However, we shall find in Section 6.3 that for this dataset, the analysis
²³ is statistically-limited, and so it is not necessary to perform exhaustive measurements
²⁴ to determine the contribution of all possible systematic contributions. Instead, the

MC process	Expected # Events Pre-Cuts	Overall Cut Efficiency (%)	Expected # Events Post-Cuts	Con- straint (%)
${}^8\text{B} \nu_e$	572.58	32.8	187.94	+2.5% -1.7%
${}^8\text{B} \nu_{\mu,\tau}$	102.22	30.2	30.86	+2.5% -1.7%
AV ${}^{214}\text{Bi}$	594,000	0.000468	2.78	+304.8% -100%
Ropes ${}^{214}\text{Bi}$	140,000	0.000327	0.46	+304.8% -100%
AV ${}^{208}\text{Tl}$	69,600	0.0361	25.14	+304.8% -100%
Ropes ${}^{208}\text{Tl}$	79,900	0.0394	31.48	+304.8% -100%
External Water ${}^{214}\text{Bi}$	4,490,000	0.0000763	3.42	—
External Water ${}^{208}\text{Tl}$	190,000	0.00761	14.46	—
PMT $\beta - \gamma$	16,800	0.333	56.12	+111.5% -40.5%
Internal ${}^{212}\text{BiPo}$	1814	0.0848	1.54	$\pm 25\%$
Internal ${}^{208}\text{Tl}$	1021	50.2	512.39	—
Internal ${}^{214}\text{BiPo}$	11,728	0.0211	2.47	$\pm 25\%$
Internal ${}^{210}\text{Tl}$	2.5	42.0	1.05	$\pm 25\%$
Surface (α, n) (combined)	554.4	0.00458	0.07	—
Internal (α, n)	17.0	0.0537	0.009	—
Internal ${}^{228}\text{Ac}$	2835	0	0	—
Internal ${}^{234m}\text{Pa}$	11,726	0.0000059	0.0007	—

Table 6.5: Number of events expected both before and after cuts have been applied, along with any constraints. As in Table 6.3, processes have been split into broad categories, with the last one being processes which have negligible rates after cuts. These processes are not included within the analysis fit.

¹ focus shall be on the subset of the systematics that can plausibly have the most impact
² on the final measurement.

³ In this analysis, the measurement of the oscillation parameters is based on the fitted
⁴ shape and normalisation of the ${}^8\text{B}$ signal energy spectrum. Therefore, for a systematic
⁵ effect to have an impact on the final measurement it must first impact the fit of the
⁶ signal PDF. One straightforward way this could occur is for some background events
⁷ to be misattributed as signal events within the fit (or vice versa). If the rates of signal
⁸ and background processes are strongly constrained to incorrect values, then this could
⁹ easily happen. In this analysis, the strongest constraint on the rate of events comes
¹⁰ from the global fit constraint of $\Phi_{\text{S}B}$. Because of this, the fit will also be run with the
¹¹ looser SSM constraint on the signal flux.

¹² Another class of systematic effects that can plausibly have a substantial impact
¹³ come from the mismodelling of the detector response, particularly on the reconstructed
¹⁴ energy and radius, as they are the observables used within the fit. The most important
¹⁵ of these for this analysis is a global mismodelling of the energy calibration by some
¹⁶ linear factor known as the energy scale, α : $E_{\text{reco}} \rightarrow \alpha E_{\text{reco}}$.

¹⁷ Because of this importance, it is worthwhile constraining this α parameter. This can
¹⁸ be done by comparing the reconstructed energy distributions of ${}^{214}\text{BiPo}$ events tagged
¹⁹ by R. Hunt-Stokes [] in the dataset to the equivalent production MC, in the same 5 m
²⁰ FV. For a given α , the energies of the tagged ${}^{214}\text{Bi}$ in MC were scaled by α , and then
²¹ compared to the equivalent (unscaled) distribution in data, via a log-likelihood ratio

²² \mathcal{L}_{Bi} :

$$\mathcal{L}_{\text{Bi}} = \sum_{i=1}^{N_{\text{bins}}} n_i^{\text{data}} \ln \frac{n_i^{\text{MC}}}{N_{\text{MC}}}. \quad (6.11)$$

²⁴ Here, n_i^{data} and n_i^{MC} are the number of events for data and (scaled) MC in energy bin
²⁵ i , with N_{bins} being the total number of bins and N_{MC} being the total number of tagged

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events in MC. To account for the edge effects discussed in Section 6.1.7, the first and last bins in energy were used as buffers and not considered in calculation of \mathcal{L}_{Bi} .

Values of α were scanned over, generating a log-likelihood distribution. In the usual way, a constraint on α was obtained by obtaining the energy scaling factor which maximises the log-likelihood, followed by looking for the α values with values of \mathcal{L}_{Bi} less than the maximum by 1/2. The constraint was found to be $\alpha = 0.9969 \pm 0.0022$, with the data and scaled MC at the best fit value of α shown in Fig. 6.12.

To handle energy scaling as a systematic in this analysis, we allow the energy scale to be floated within the MCMC fit, with the constraint given above. The first and last energy bin of the MCMC fit were used as buffer bins to let this floating occur; this allows for the energy scale to float in the range $\frac{13}{14} \leq \alpha \leq \frac{2.6}{2.5}$.

In addition to a global energy scaling, it is possible that the resolution of the energy reconstruction could be systematically off. Because of the effects discussed in Section 6.1.1, changes to the oscillation parameters only impact the observed signal distribution over a broad range of energies. This means that mismodelling the energy resolution of the detector is only liable to have an impact on the measurement if that mismodelling is substantial. However, by looking again at Fig. 6.12 we see that there appears to be a good qualitative match of the energy resolution between data and MC for the tagged ${}^{214}\text{Bi}$ events. As a result, the systematic effect of additional energy smearing will not be considered in this thesis.

Similar to the systematics in reconstructed energy, there are also systematics possible in reconstructed position. These can be decomposed into three kinds of systematic: radial scaling, radial resolution, and position bias along each coordinate axis. The dominant effect that these systematics will have on the PDFs is to change their normalisations. Because we are floating all normalisations within the MCMC fit, we expect the impact of these systematics to be sub-dominant.

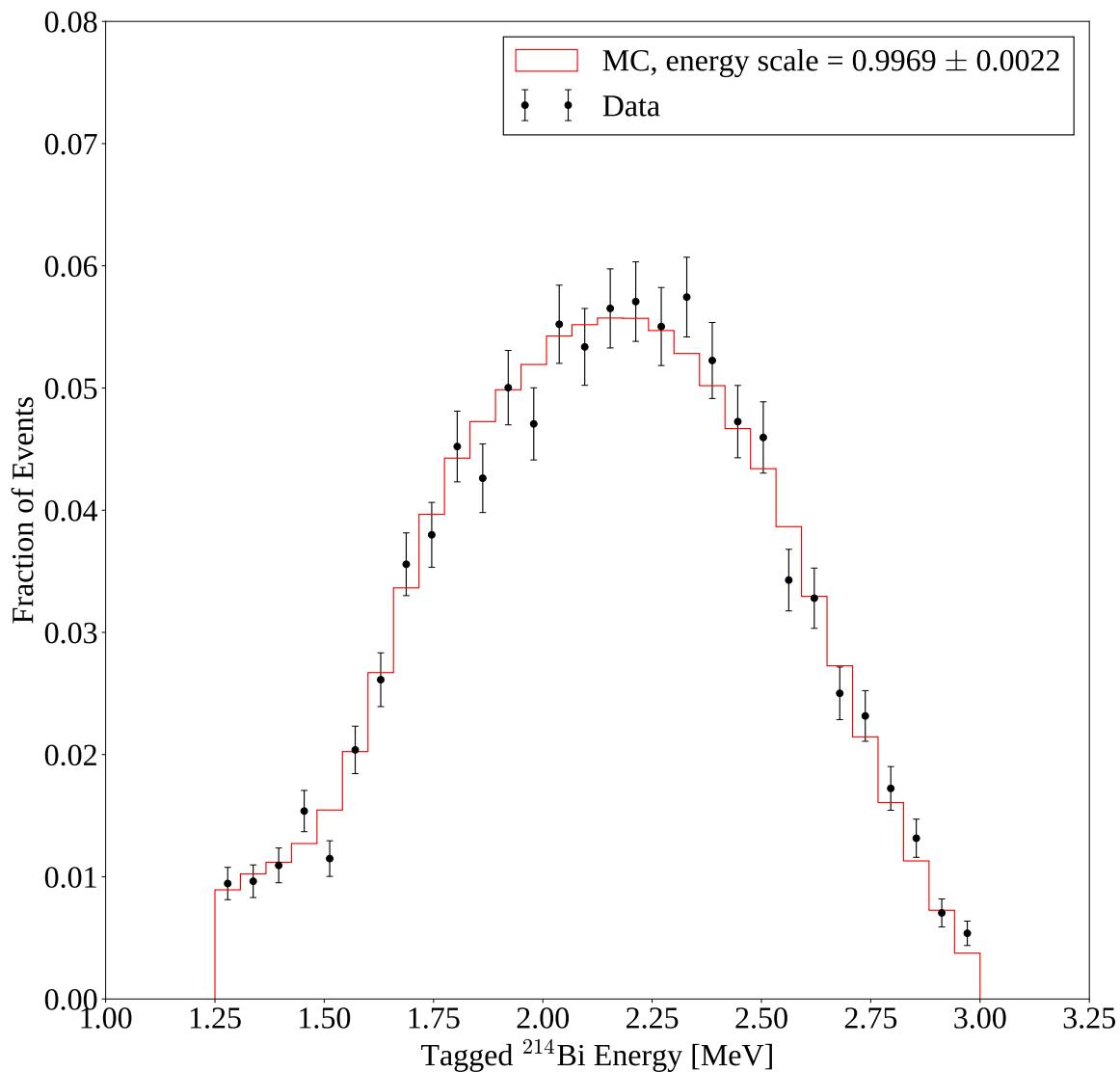


Fig. 6.12: Comparison of tagged ^{214}Bi reconstructed energy distributions between data and MC, after MC has been scaled by the best-fit factor $\alpha = 0.9969$.

Finally, there can also be mismodelling of the radial distributions for given processes in MC compared to data. For example, it is known that the internal ^{208}Tl events are non-uniform as a function of radius; this is why the normalisation of each radial slice for this background is floated independently. The external backgrounds can also plausibly have PDF shape systematics, especially if the assumptions made to model them (such as the use of simulation shells) are not fully accurate. These shape systematics can also arise from the lack of statistics in the PDFs used to model them: as an example, Table 6.3 shows that only 62 external water ^{214}Bi events are used in the creation of the associated PDF. Handling this type of modelling uncertainties for the external backgrounds is not considered in this thesis.

6.2.5 Results

Fit Validation

In theory, the convergence of the sampled distribution within an MCMC fit to the posterior density distribution is guaranteed to eventually occur, regardless of the specific shape of the multivariate Gaussian proposal distribution. However, the speed of this convergence is highly dependent on the widths of the proposal distribution. This is very important as in practice we can only run the MCMC fit for a finite amount of time. For a given parameter within the fit, if the associated width (also known as the step size) in the proposal distribution is too large, then the vast majority of proposals will be to points in the parameter space where the log-likelihood is less than before, and so these proposals will be consistently rejected. Alternatively, if the width is chosen too small, the acceptance rate of the chain will be very high, but too much time will be spent for the chain to go from sampling one point in the space to another.

By choosing sensible width parameters for the proposal distribution, as well as running the chain for as long as possible, one can attempt to maximise the “effective sample size”

1 of the MCMC fit. It has been shown that, for a wide of situations, an appropriate size
2 for the width of a given parameter is the standard deviation of the true posterior density
3 distribution []. Because this posterior density distribution is precisely the aim of the
4 whole MCMC procedure, this can make choosing these widths challenging in practice.
5 In this analysis, an iterative approach was used: an MCMC fit was first run using
6 width parameters that were guessed. Once this fit completed, the resulting sampled
7 distribution marginalised onto each fit parameter were looked at to find approximately
8 the magnitude of their standard deviations. This then informed the widths chosen for
9 the next MCMC fit. This process was repeated until the autocorrelation of the sampled
10 distribution for each fit parameter appeared to be minimal; this is an indication that
11 the effective sample size has been maximised. A chain length of 1,000,000 steps was
12 found to be sufficient to obtain large enough effective sample sizes for this analysis.

13 In order to guard against the possibility of a given chain missing some important
14 part of the parameter space because of its starting position, 100 chains were run
15 simultaneously. The initial positions of each in the parameter space were chosen
16 randomly according to the prior distribution. Fig. 6.13 shows the sampled values for
17 four different parameters in a given chain, after the tuning of the width parameters
18 had been completed.

19 Given some initial start point in the parameter space, a given chain will typically
20 take some time in finding where the region of greatest likelihood is. This leads to an
21 initial set of steps where the chain moves a large distance in the parameter space in the
22 same direction, so that the autocorrelation of the chain in this period is substantial.
23 These initial samples are not representative of the actual posterior density distribution:
24 as a result, in this analysis the first 100,000 steps of each chain were declared as a
25 “burn-in” period. This number of steps for the burn-in was motivated by looking
26 at the auto-correlation plots for the likelihood evaluation for the MCMC chains, as

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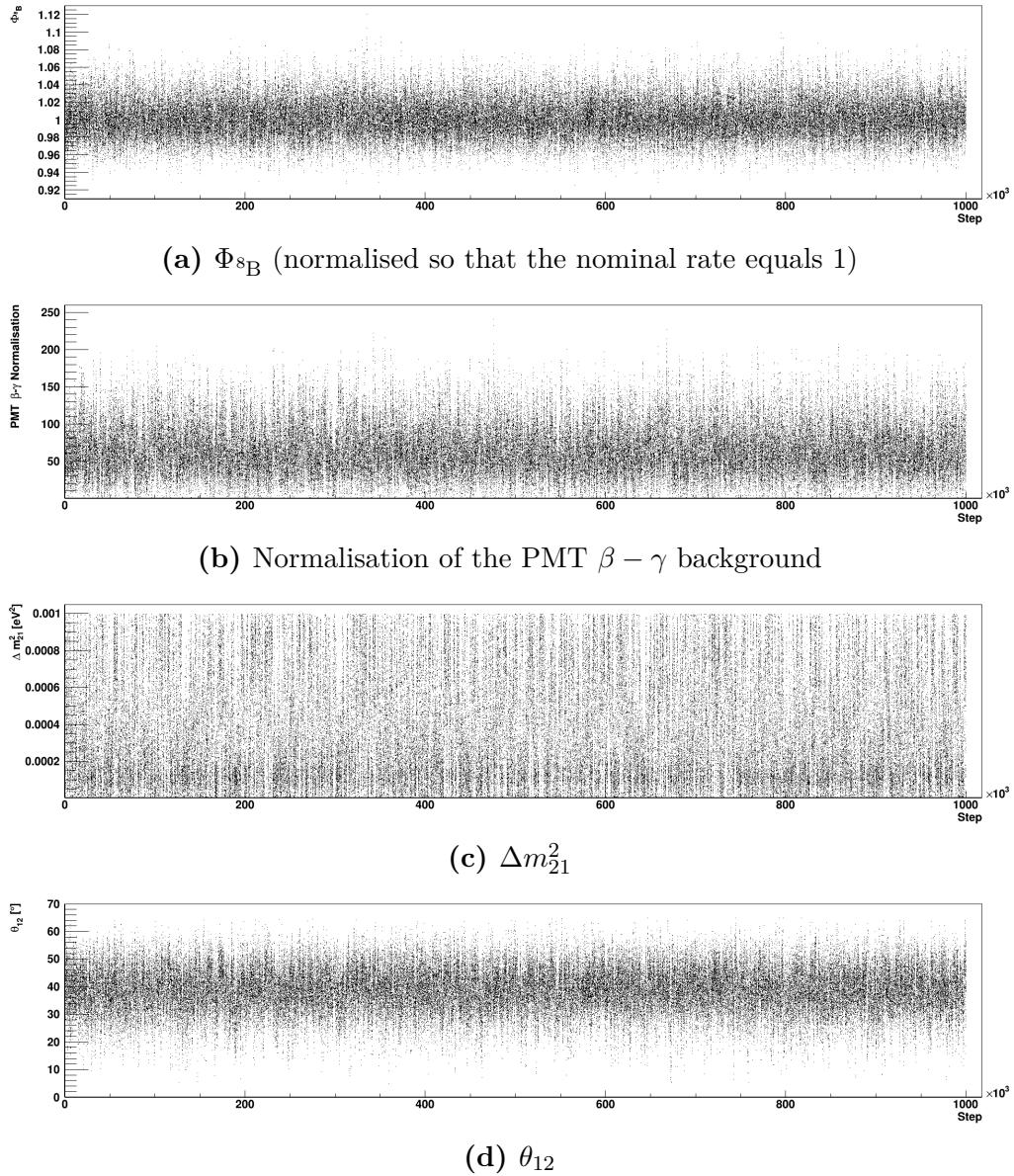


Fig. 6.13: Examples of the sampled values for parameters within a given MCMC chain as a function of the step number, after tuning the proposal distribution width parameters.

Fig. 6.14

seen in Fig. 6.14. When the number of steps differed by 100,000, the autocorrelation of each chain fell to zero, indicating the likelihood evaluations in the chains were independent when they differed by at least 100,000 steps. The overall set of samples used to approximate the posterior density distribution corresponds to collection of all sampled points on all the chains, after burn-in.

6 Oscillation Fit Results

Now that there is confidence in the convergence of the MCMC fit to the posterior density distribution, we can look at this distribution to obtain the results of this analysis. Fig. 6.15 shows the posterior density distribution, marginalised onto each parameter of the fit other than the two oscillation parameters. We can see that each marginalised distribution appears smooth, another qualitative indication of a sufficient effective sample size in the MCMC fit. For each 1D distribution, the bin centre associated with the Highest Posterior Density (HPD) has been used as the best-fit value for that parameter, and the vertical band indicates the Bayesian 1σ CI for that parameter. Also shown in these plots is the prior distribution for each of the parameters, for comparison.

For a number of the parameters, the posterior distributions are more strongly peaked than their associated priors. This indicates that there was enough information in the dataset further constrain those parameters. One such parameter is the rate of the external water ^{214}Bi . The resulting measurement of this is $34.1^{+18.5}_{-18.9}$ events in the dataset. This is 10 times larger than the expected rate shown in Table 6.5, providing good evidence for a change in the level of this particular background relative to the water phase. In contrast, all other external backgrounds have HPDs consistent with their prior expected values.

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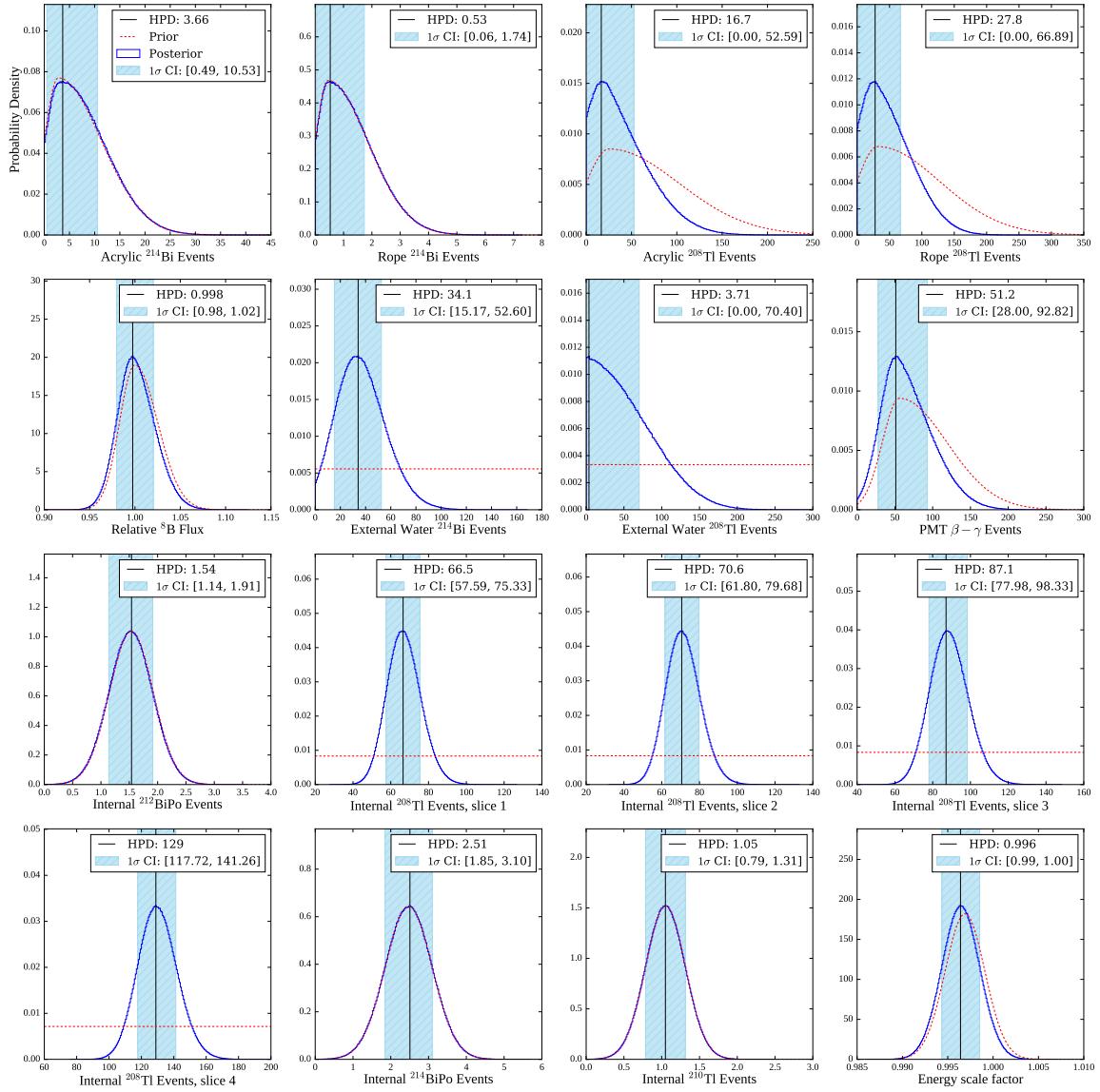


Fig. 6.15: Marginalised 1D posterior density distributions for all non-oscillation parameters, shown in blue. The HPD and 1σ Bayesian CI are shown for each parameter. Also shown in dashed red are the associated prior distributions.

Another set of parameters that have been well-constrained in the fit are the normalisations of each of the internal ^{208}Tl slices. As expected, slices corresponding to larger radii have a greater fitted rate, with the number of events measured in the outermost radial slice to be twice that of the innermost one. However, the posteriors of the other internal ^{238}U - and ^{232}Th -chain are the same as their priors, indicating there was no evidence from data within the fit to constrain those processes any further. Similarly, the posteriors of the energy scale parameter and Φ_{sB} only show very mild differences with their strong priors.

If instead of looking only at the 1D marginalised distributions of each parameter we look at the equivalent 2D distributions for every combination of parameters, we can generate a table of correlation coefficients. This is shown in Fig. 6.16. From this, we can see that there is some clear anti-correlation between a number of the external backgrounds. This seems intuitive: the PDFs for many of these external backgrounds are similar, and so the fit struggles to distinguish between them.

The only other pair of parameters which have a substantial correlation between one another are the two oscillation parameters. Fig. 6.17 shows their 2D marginalised posterior distribution, along with contours for a number of Bayesian CIs. Because the priors on both oscillation parameters are uniform, this can be directly converted into the log-likelihood ratio $\ln L_R$ for a given bin, where $L_R = p/p_{max}$: p is the posterior density of a given bin in the 2D space, and p_{max} is the posterior density of the HPD bin, equivalent to the bin with the maximum likelihood. This allows for Frequentist CI to be derived as well. For this plot, the HPD is at $\theta_{12} = 34.1^\circ$, $\Delta m_{21}^2 = 1.80 \times 10^{-5} \text{ eV}^2$.

As can be seen from this plot, the fit has been able to successfully constrain θ_{12} to some success. However, beyond a value of $5 \times 10^{-4} \text{ eV}^2$, the posterior density appears flat as a function of Δm_{21}^2 . This implies that this dataset provides little ability to reject the possibility of large values of Δm_{21}^2 .

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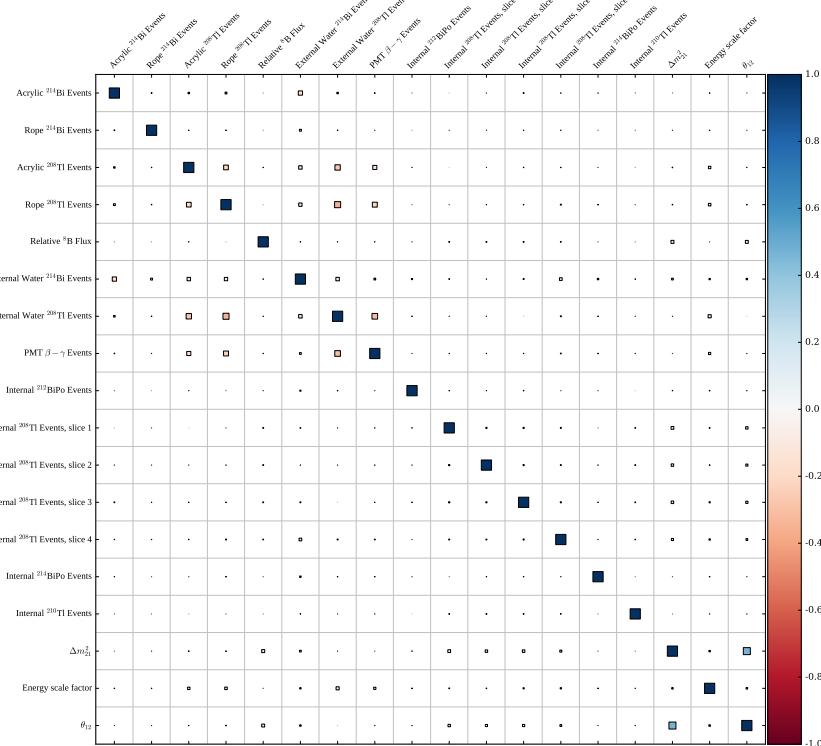


Fig. 6.16: Correlation matrix between all parameters in the MCMC fit.

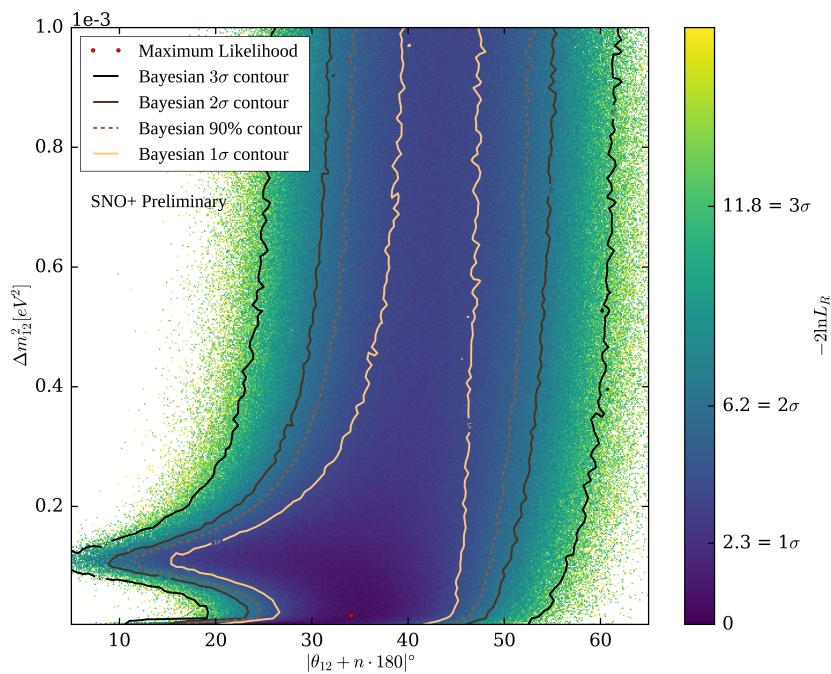


Fig. 6.17: 2D posterior density distribution marginalised onto Δm_{21}^2 and θ_{12} . The colour axis shows $-2 \ln L_R$, proportional to the natural logarithm of the posterior density, with labels corresponding to the values associated with the $N\sigma$ Frequentist CIs. Also shown, in shades of brown, are the Bayesian CIs.

6.2 Analysis on Scintillator-Phase data

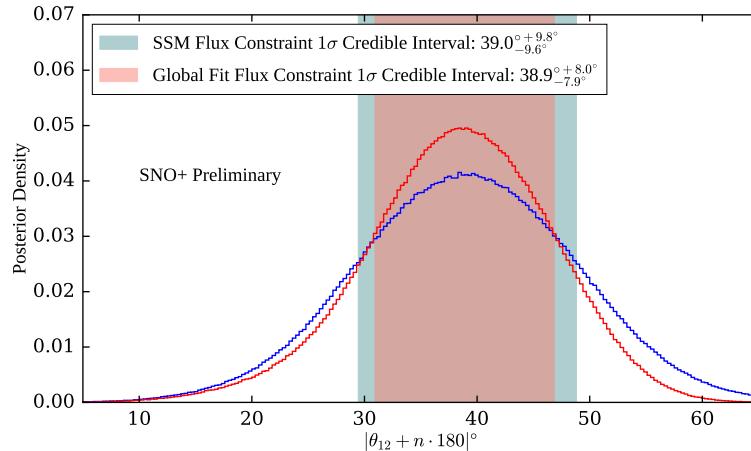


Fig. 6.18: Comparison of 1D posterior density distributions marginalised onto θ_{12} , using both forms of ${}^8\text{B}$ flux constraint. The 1σ Bayesian CI are shown for both.

Given this lack of constraining power on Δm_{21}^2 , it is reasonable to only obtain a measurement for θ_{12} . The result of marginalising over Δm_{21}^2 to obtain the 1D posterior distribution in θ_{12} is shown in Fig. 6.18. From this, we can obtain a measurement of:

$$|\theta_{12} + n \cdot 180|^\circ = 38.9^\circ_{-7.9^\circ}.$$

Here, the point of HPD has been estimated by fitting a quadratic function to the peak of the posterior distribution, and the uncertainties by obtaining the Bayesian 1σ CI. For comparison, the NuFit 5.1 global fit results have []:

$$|\theta_{12} + n \cdot 180|^\circ = 33.44^\circ_{-0.74^\circ}.$$

As can be seen, the measurement made in this thesis is consistent with the global fit value, although with a substantially larger uncertainty. This is not surprising, given the limited statistics of the 80.6 days of data used.

The fact that the measured value of θ_{12} is somewhat above the global fit value also has a reasonable explanation. Looking back at Fig. 6.17, we find that the preferred

values of θ_{12} for large Δm_{21}^2 are somewhat greater on average than for Δm_{21}^2 values close to the global fit value of 7.4×10^{-5} eV². As a result, the final value quoted for θ_{12} is dependent on our choice of the prior distribution for Δm_{21}^2 : this is an inevitable result of using a Bayesian framework for this analysis. We can see the impact of choosing a less conservative prior for Δm_{21}^2 by ignoring all sampled points with $\Delta m_{21}^2 \geq 5 \times 10^{-4}$ eV², say. The resulting measurement of θ_{12} becomes:

$$|\theta_{12} + n \cdot 180|^\circ = 36.4^\circ_{-7.8^\circ}.$$

Although there is negligible change to the uncertainty of the measurement, this change in prior shifts the HPD value down by 2.5°, in the direction of the global fit value.

Also shown in Fig. 6.17 is the posterior density for θ_{12} if the looser SSM flux constraint is used. Under these conditions, the measurement becomes:

$$|\theta_{12} + n \cdot 180|^\circ = 39.0^\circ_{-9.6^\circ}.$$

Using this constraint has little impact on the HPD value, but does increase the uncertainty by $\sim 1.8^\circ$. This gives an indication of the extent to which the strong global flux constraint allows θ_{12} to be constrained on its own, without any information from the dataset.

One could naïvely expect that the point in parameter space where each fit parameter corresponds to the HPD value obtained in the above 1D marginalisation would correspond to the maximum likelihood. This is not the case, because of the correlations between parameters in the fit. Because an MCMC fit does not actually attempt to directly find this maximum likelihood point, deciding on the overall ‘best-fit’ point in parameter space is a little ambiguous. For this analysis, the sampled point in all of the chains which had the greatest log-likelihood was declared as the best-fit

6.3 Sensitivity Projections

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Fit Parameter	1D HPD value	Maximum likelihood value
Relative ${}^8\text{B}$ flux	0.9975	1.006
$\Delta m_{21}^2 [\times 10^{-5} \text{ eV}^2]$	11.02	5.50
$\theta_{12} [^\circ]$	39.05	33.73
AV ${}^{214}\text{Bi}$ Events	3.66	8.49
Ropes ${}^{214}\text{Bi}$ Events	0.53	0.52
AV ${}^{208}\text{Tl}$ Events	16.71	19.08
Ropes ${}^{208}\text{Tl}$ Events	27.81	31.20
External Water ${}^{214}\text{Bi}$ Events	34.09	45.93
External Water ${}^{208}\text{Tl}$ Events	3.71	88.03
PMT $\beta - \gamma$ Events	51.22	46.88
Internal ${}^{212}\text{BiPo}$ Events	1.54	1.45
Internal ${}^{208}\text{Tl}$ Events, Slice 1	66.52	68.78
Internal ${}^{208}\text{Tl}$ Events, Slice 2	70.60	66.62
Internal ${}^{208}\text{Tl}$ Events, Slice 3	87.09	87.33
Internal ${}^{208}\text{Tl}$ Events, Slice 4	128.92	125.25
Internal ${}^{214}\text{BiPo}$ Events	2.51	2.56
Internal ${}^{210}\text{Tl}$ Events	1.05	1.00
Energy Scale Factor, α	0.9964	0.9970

Table 6.6: Comparison of the fit parameter values obtained when getting the HPD values after marginalising onto each parameter, versus looking at the sampled point with the maximal likelihood found.

point. A comparison between the parameter values obtained at the HPD of the 1D marginalisations to the maximum likelihood point is shown in Table 6.6.

Using the maximal likelihood fit parameters, the comparison of data to MC is shown in Fig. 6.19. We see that the ${}^8\text{B}$ signal distribution only dominates above 5 MeV, as well as at low energies for the first two radial slices. This is because for larger radii, external backgrounds become the dominant processes in the 2.5–2.8 MeV region. Between 3–5 MeV, at all radii the internal ${}^{208}\text{Tl}$ shows a clear peak well above the ${}^8\text{B}$ signal. Qualitatively, the MC appears to give a good fit to the data.

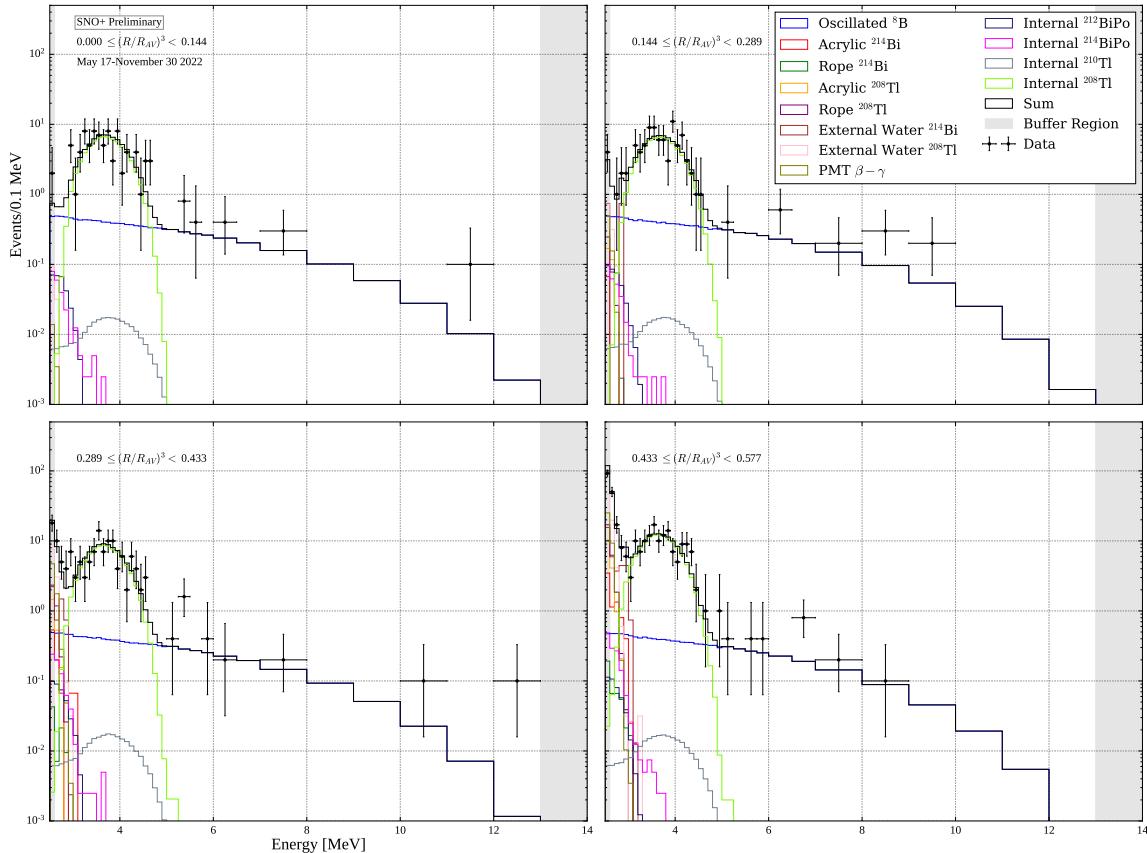


Fig. 6.19: Comparison of data to MC for each radial slice, using the best-fit parameter values derived from the MCMC results. The buffer bins used for varying the energy scale parameter within the fit are shown in grey.

6.3 Sensitivity Projections

Given the limited statistics used in the dataset analysed in Section 6.2, the question of how much better we could measure the solar oscillation parameters with more data naturally arises. In particular, it is worthwhile knowing the expected sensitivity of this analysis as a function of livetime, and hence whether SNO+ could eventually make a world-leading measurement of θ_{12} via ^8B neutrinos. Furthermore, it is useful to find out to what extent improvements such as additional background reduction could help the measurement.

To perform these sensitivity projections, the same analysis method was employed as with the real dataset. However, for each projection scenario a fake “Asimov” dataset was generated. This was a 2D histogram binned in the same way as the MC PDFs, corresponding to the total expected rate of all signal and background processes. Because the same PDFs were used to build both the MC and fake dataset, the maximum likelihood should occur when the fit parameters are identical to those which generated the fake dataset. Although this method of performing sensitivity projections leads to the slightly odd situation of having non-integer numbers of events in the bins of the fake dataset, it has been shown that this approach allows us to estimate the median sensitivity of a scenario [57].

The assumptions used to make the fake datasets are as follows. To begin with, no cuts were changed between the main analysis and this projection, and the same PDFs were used in both. Furthermore, the ^8B rate was determined by the same global fit flux value and neutrino oscillation parameters used in the rest of this analysis. Because some of the dataset used in the main analysis was taken during a time of known elevated ^{238}U - and ^{232}Th -chain backgrounds, R. Hunt-Stokes provided an additional estimate of the $^{214}\text{BiPo}$ and $^{212}\text{BiPo}$ rates in the 5.0 m FV for only the period of the dataset where the internal background levels had stabilised. Those rates went from 6.06

¹ events per hour before excluding the higher background period down to 4.87 events
² per hour for the $^{214}\text{BiPo}$, and from 0.94 events per hour down to 0.89 events per hour.
³ These relative changes for the two background chains were used to scale the rates from
⁴ their expectation in the original dataset. For the internal ^{208}Tl , the rate of each slice
⁵ relative to one another was taken from their fitted HPD values in the original analysis,
⁶ with the absolute value of those rates scaled by 0.89/0.94.

⁷ External backgrounds were in general set to the rates expected from T. Zummo, as
⁸ described in Section 6.2.3. The exceptions were the backgrounds from the external
⁹ water: for the ^{208}Tl component, the nominal rate without T. Zummo’s correction
¹⁰ factor was used. For the ^{214}Bi component, the higher rate as given by the fitted HPD
¹¹ value was used instead. The same constraints on all the rates were used as in the
¹² main analysis. For the energy scale parameter, because no additional systematics were
¹³ applied to the Asmiov dataset the true and expected value of α was set to 1, with the
¹⁴ same uncertainty of ± 0.0022 as before.

¹⁵ Using this baseline set of assumptions, fake datasets were generated and fit over a
¹⁶ series of livetime scenarios between 150 days and 5 years. The results of these MCMC
¹⁷ fits, in terms of the 2D and 1D posterior densities for the oscillation parameters are
¹⁸ shown in Figs. 6.20 and 6.21.

¹⁹ These choices of assumptions for this baseline set of projections were deliberately
²⁰ chosen to be fairly conservative. However, they do not account for the loading of
²¹ BisMSB and TeLS into the detector that is likely to occur in the coming years. For
²² BisMSB, the main impact is expected to be a dramatic increase in the observed
²³ light yield in the detector, leading to a smaller energy resolution and hence improved
²⁴ measurement []. Once the TeLS is also loaded, the observed light yield will likely
²⁵ decrease back down to levels similar to that of the current LABPPO. The biggest
²⁶ expected impact of the Te phase on this analysis will instead be to the background

6.3 Sensitivity Projections

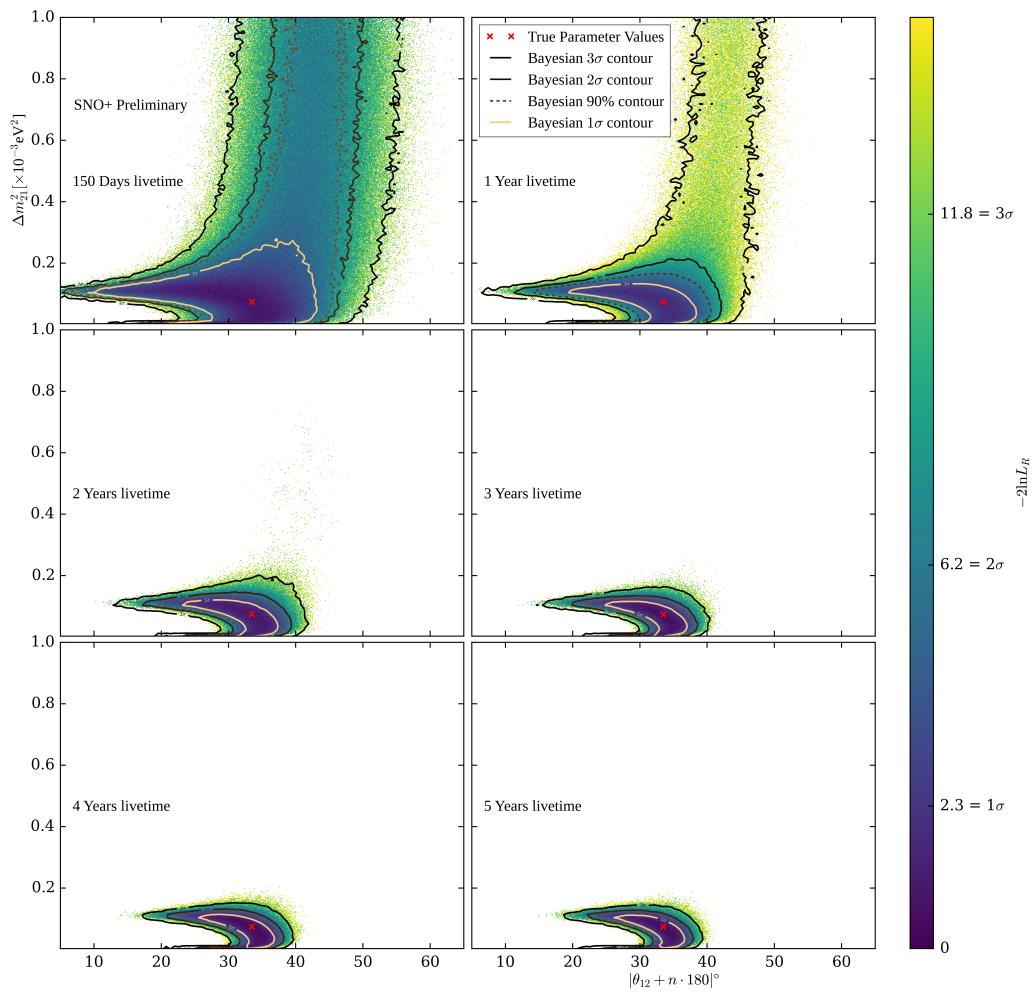


Fig. 6.20: Posterior density posterior distributions marginalised onto the two solar oscillation parameters, for each livetime scenario.

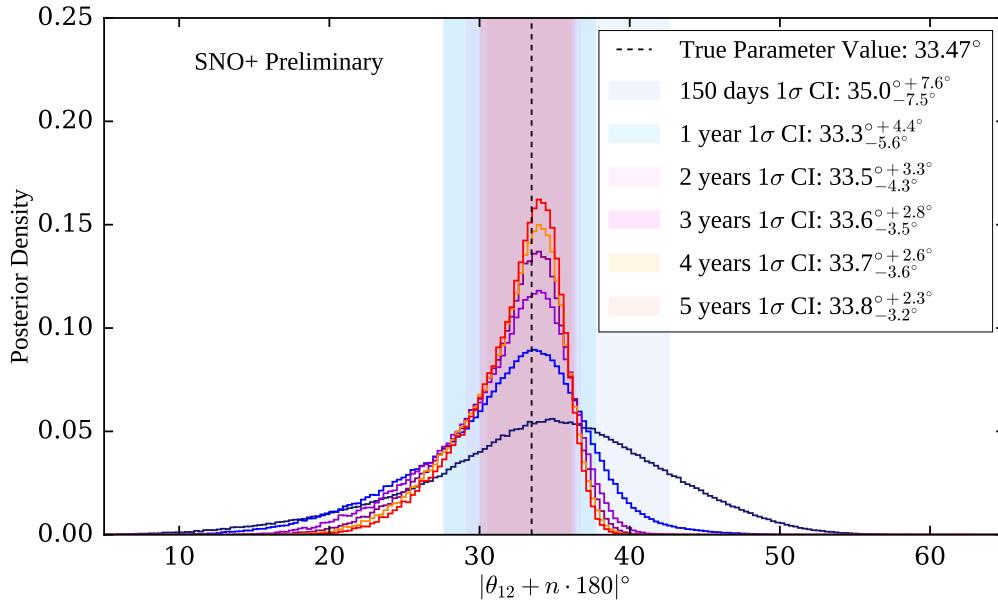


Fig. 6.21: Posterior density posterior distributions marginalised onto θ_{12} , for each livetime scenario.

¹ levels within the detector: the target rate of ^{238}U -chain backgrounds in the Te phase is
² over 100 times that of the scintillator phase [54].

- ³ • Using the same production MC, generate PDFs with the expected normalisations
- ⁴ for longer periods of livetime: 1, 3, and 5 years. This still assumes a scintillator-fill
- ⁵ with identical detector conditions on average, so not considering the impact of
- ⁶ BisMSB loading at any point.
- ⁷ • Describe any further constraints to be assumed on top of the existing analysis
- ⁸ of data: expected improved constraints on various backgrounds, as well as a
- ⁹ possible energy-scale calibration constraint from the AmBe source or an internally-
- ¹⁰ deployed source, for example.
- ¹¹ • Run MCMC fits to these Asimov PDFs for each livetime scenario. Describe
- ¹² results in terms of the improvement to the sensitivity to θ_{12} as a function of
- ¹³ livetime.

6.3 Sensitivity Projections**133**

- Also consider scenario of lower backgrounds! How does that impact the result? 1
- Could consider scenario of BisMSB deployment, leading to substantially greater 2
light yield and hence energy resolution. Given time it would take for Production 3
MC to be generated for this scenario, I would likely have to come up with 4
something clever and quick to actually do this. There quite possibly won't be 5
enough time. 6

[8 pages]

7

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Chapter 7

Conclusions and Suggestions for Future Work

- Say what has been achieved in this thesis! In particular:
 - Substantial improvement to the SMELLIE generator in terms of speed and dynamic range
 - A much stronger understanding of the discrepancies between data and MC in SMELLIE
 - The creation of two analyses of SMELLIE data, designed explicitly around being robust to these systematics
 - A measurement of the extinction length of scintillator *in-situ* with SMELLIE at 375 nm, monitored over time
 - A first measurement of the scattering length of the scintillator *in-situ*, monitored over time
 - The creation of an analysis of ${}^8\text{B}$ solar neutrinos in the scintillator phase to measure the solar neutrino oscillation parameters

- ¹ – The first measurement of θ_{12} using ${}^8\text{B}$ neutrinos in SNO+
- ² – Projections of this solar analysis' precision at longer livetimes
- ³ • Give suggestions for further work that could be done on both SMELLIE and the
⁴ solar oscillation analysis:
 - ⁵ – Inclusion of LABPPO's polarisability anisotropy in the detector's optical
⁶ scattering simulation, and determination of its impact on both Physics and
⁷ SMELLIE measurements
 - ⁸ – Further investigation of SMELLIE's wavelength-dependence of the beam
⁹ profiles, looking both at possible origins for the phenomenon and correcting
¹⁰ for this in simulation
 - ¹¹ – Various computational improvements to allow for faster MCMC run-times,
¹² including the calculation of the systematic matrices, and possible use of
¹³ GPUs to parallelise some parts of the computation.
 - ¹⁴ – Inclusion of additional solar neutrino components at lower energy into the
¹⁵ fit, e.g. Be7. Maybe the addition of lower energies also helps to naturally
¹⁶ constrain backgrounds within the fit?
 - ¹⁷ – Looking at the impact of various advanced background-rejection procedures,
¹⁸ such as event directionality or topology. How much do they help with the
¹⁹ sensitivity?
 - ²⁰ – Looking at the impact of splitting data into day and night parts, to try
²¹ and provide further constraints on any matter effects. Not expected to be
²² significant, so was ignored for this analysis so far.
 - ²³ – Performing a combined fit with the reactor anti-neutrino oscillation analysis,
²⁴ which allows for the handling of correlated uncertainties, such as detector
²⁵ response systematics.

[3 PAGES TOTAL]

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