

Matrix hydrodynamics and 2-D Euler

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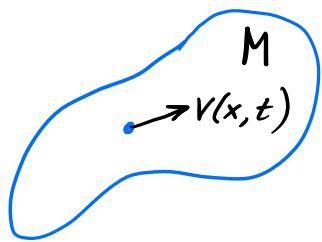


Dedicated to Vladimir Zeitlin

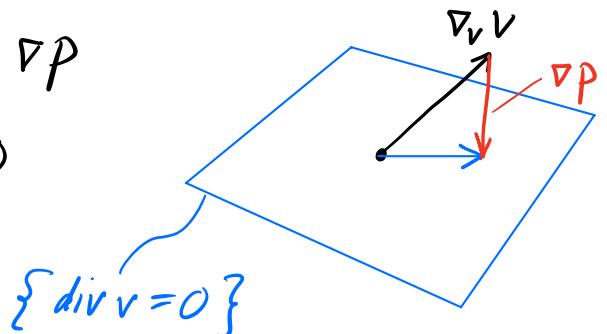


The Father of Matrix Hydrodynamics

Incompressible inviscid hydrodynamics (Euler 1757)



$$\left\{ \begin{array}{l} \dot{v} + \nabla_v v = -\nabla p \\ \operatorname{div} v = 0 \end{array} \right.$$



Arnold's viewpoint

Integrate: $\dot{x}(t) = v(x(t), t)$

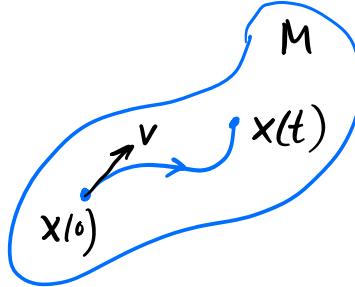
Flow map (all $x \in M$ at once):

$$\dot{\varphi}_t = v_t \circ \varphi_t, \quad \varphi_0(x) = x$$

tracks fluid particles

NOTICE:

$$\varphi_t \in \text{Diff}_{\mu}(M) = \{\eta \in \text{Diff}(M) \mid \eta_* \mu = \mu\}$$



Newton's equation (differential)

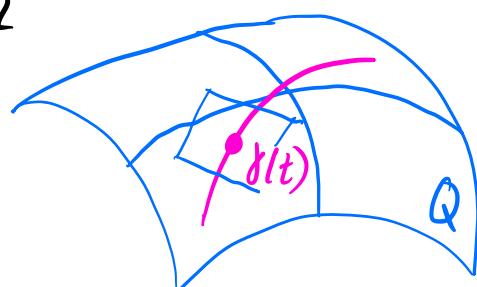
$$\begin{aligned} \ddot{\varphi} &= \dot{v} \circ \varphi + \underbrace{\nabla v \circ \varphi}_{v \circ \dot{\varphi}} \cdot \dot{\varphi} = (\dot{v} + \nabla_v v) \circ \varphi \\ &= (-\nabla_v v - \nabla p + \nabla V) \circ \varphi \end{aligned}$$

$$\Rightarrow \ddot{\varphi} = -\nabla p \circ \varphi \quad (*)$$

What is a geodesic curve?

Idea: take $Q = \text{Diff}_{\mu}(M)$

int. dim
manifold

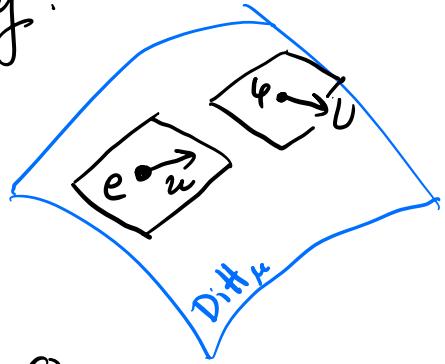


$$\gamma \in (T_q Q)^\perp$$

Is $\nabla p \circ \varphi \perp$ to $T_{\varphi} \text{Diff}_{\mu} \subset T_{\varphi} \text{Diff}$?

Riemannian metric: kinetic energy:

$$\langle U, U \rangle_{\varphi} = \int_M |U(x)|^2 d\mu(x)$$



$$U \in T_\varphi \text{Diff}_\mu(M) \iff \text{div} \left(\underbrace{U \circ \varphi^{-1}}_u \right) = 0$$

$$U: M \rightarrow TM$$

$$x \mapsto U(x) \in T_{\varphi(x)}M$$

Back to Euler equations: $\ddot{\varphi} = -\nabla p \cdot \dot{\varphi}$

Thun (Arnold 1966)

$\varphi(t) \in \text{Diff}_\mu(M)$ is geodesic curve

prost:

$$\begin{aligned} \langle U, \nabla p \circ \varphi \rangle_\varphi &= \int_M U \cdot \nabla p \circ \varphi \, d\mu = \int_M (u \cdot \nabla p) \circ \varphi \, d\mu \\ &= \int_M u \cdot \nabla p \underbrace{\varphi_* \, d\mu}_{d\mu} = \int_M u \cdot \nabla p \, d\mu = 0 \quad \square \end{aligned}$$

$$\langle u, \nabla p \rangle_e = \langle u, \nabla p \rangle_{L^2}$$

NOTE: Metric is right-invariant: $\langle U \circ \eta, V \circ \eta \rangle_{\psi \circ \eta} = \langle U, V \rangle_p$

Hamiltonian description: Lie-Poisson dynamics

G : Lie group

Hamiltonian $\bar{H}: T^*G \rightarrow \mathbb{R}$ is right-invariant:

$$\bar{H}((q,p) \cdot g) = \bar{H}(q,p)$$

Symmetry group G : dynamics reduce to $T^*G/G \cong \mathfrak{g}^*$

dual of Lie algebra

$$[(q,p)] \longleftrightarrow (q,p) \cdot q^{-1} = (e, \underbrace{p \cdot q^{-1}}_m) \cong m \in T_e^*G = \mathfrak{g}^*$$

$$\bar{H}(q,p) = H(\underbrace{p \cdot q^{-1}}_m), \quad H: \mathfrak{g}^* \rightarrow \mathbb{R}$$

Hamilton's eq:

$$\begin{cases} \dot{q} = \frac{\partial H}{\partial p} \\ \dot{p} = -\frac{\partial H}{\partial q} \end{cases} \quad m \equiv p \cdot q^{-1} \implies$$

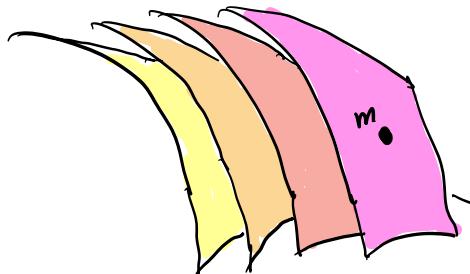
Lie-Poisson system

$$\dot{m} - ad_{dH(m)}^* m = 0$$

Preserves Poisson structure on \mathfrak{g}^*

$$\{dF, dG\}(m) = \langle m, [dF(m), dG(m)] \rangle$$

\mathfrak{g}^* foliated in co-adjoint orbits (symplectic leafs)



$$O_m = \{Ad_g^*(m) \mid g \in G\}$$

Special case M Riemannian (compact)

$G = \text{Diff}_\mu(M)$ (Frechet Lie group)

$$\mathfrak{g} = \mathcal{X}_\mu(M) = \{v \in X(M) \mid \text{div } v = 0\}$$

$$\mathfrak{g}^* \cong \Omega^1(M)/d\Omega^0(M) \cong d\Omega^1(M) \quad (\text{exact 2-forms})$$

↑
smooth
dual

↑ trivial
1-cohomology

Hamiltonians on $\mathfrak{g}^* \cong d\Omega^1(M)$

$$H(\hat{\omega}) = \frac{1}{2} \int_M i_v d\bar{d}\hat{\omega}^\mu, \quad dV^b = \hat{\omega}, \quad \text{div } v = 0$$

exact
2-form

2-D case: $\hat{\omega} = \omega^\mu$ symplectic structure on M

$$v = X_\psi = \nabla^\perp \psi$$

Hamiltonian on M

$$H(\omega) = \frac{1}{2} \int_M \psi \omega^\mu, \quad -\Delta \psi = \omega$$

Poisson bracket on M

$$\langle \text{ad}_{X_\psi}^* \hat{\omega}, X_\varphi \rangle = \langle \omega, \{\psi, \varphi\} \rangle_{L^2} = \langle -\{\psi, \omega\}, \varphi \rangle_{L^2}$$

$$\dot{\hat{\omega}} - \text{ad}_{X_\psi}^* \hat{\omega} = 0 \iff \boxed{\dot{\hat{\omega}} + \{\psi, \hat{\omega}\} = 0 \quad (*)}$$

$$\iff \dot{\hat{\omega}} + \mathcal{L}_{X_\psi} \hat{\omega} = 0 \iff \dot{\hat{\omega}} + \text{div}(\hat{\omega} \nabla^\perp \psi) = 0$$

Casimirs: $C_f(\omega) = \int_M f \circ \omega \mu$ preserved
for each $f: \mathbb{R} \rightarrow \mathbb{R}$

Proof: C_f constant on co-adjoint orbits

Zeithui's idea (1991)

Use quantization theory to replace $(*)$ by
finite dim. lie-Poisson system

Quantization in a nutshell (for $M = S^2$)

Vorticity state $\omega: S^2 \rightarrow \mathbb{R}$

Aim: mapping $T_N: \omega \mapsto W \in U(N)$ such that

$$T_N(\{\omega_1, \omega_2\}) \approx \frac{1}{\hbar} [T_N \omega_1, T_N \omega_2] \quad \text{← matrix commutator}$$

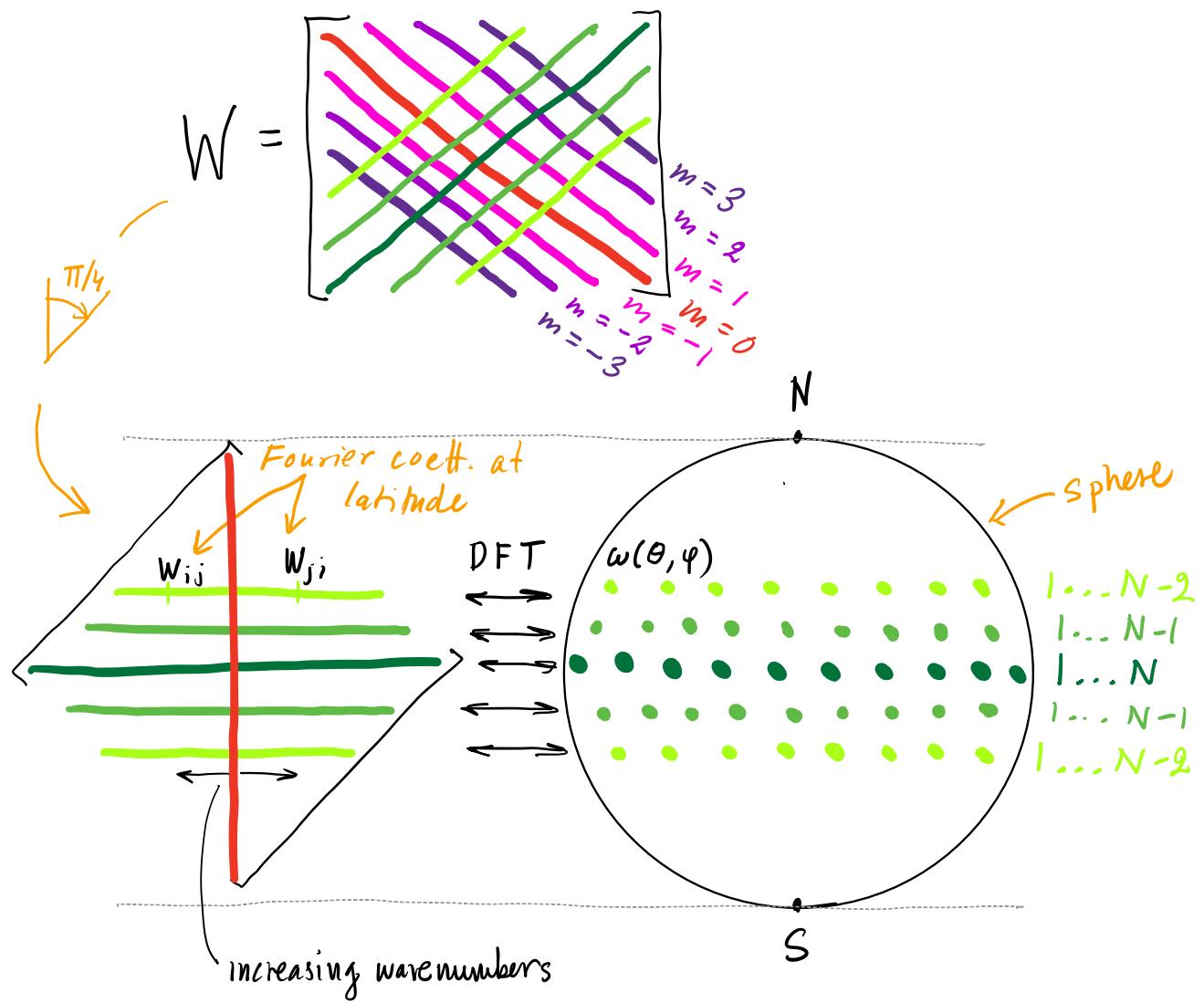
$$\hbar_N = \frac{1}{\sqrt{N^2 - 1}}$$

Hoppe (1989): explicit quantization on T^2 and S^2

$$\text{Route: } \omega \mapsto \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \omega_{\ell m} Y_{\ell m} \mapsto \sum_{\ell=0}^N \sum_{m=-\ell}^{\ell} \omega_{\ell m} T_{\ell m}^N \quad \mathfrak{U}(N)$$

\uparrow
spherical
harmonics

What is the Matrix? (Neo, 1999)



Euler-Zeitlin equation

$$\begin{aligned} \omega \in C^\infty(\mathbb{S}^2) &\longleftrightarrow W \in \mathcal{U}(N) \\ \Delta \psi = \omega &\longleftrightarrow \Delta_N P = W \end{aligned}$$

Hoppe-Yan Laplacian (1998)

$$\dot{\omega} = \{ \psi, \omega \} \longleftrightarrow \dot{W} = \frac{i}{\hbar} [P, W]$$

Stream matrix Vorticity matrix

Dictionary Hydrodynamics — Matrix theory

"Classical" hydrodynamics

vorticity $\omega \in L^\infty(\mathbb{S}^2)$

Casimir $I_f(\omega) = \int_{\mathbb{S}^2} f(\omega(x)) d\mu(x)$

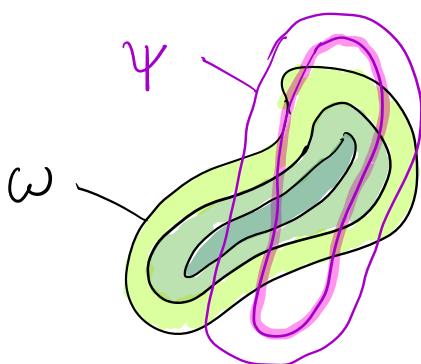
Hamiltonian $H(\omega) = \frac{1}{2} \int_{\mathbb{S}^2} \psi \omega d\mu$

values of ω

levelsets of ω

$\|\omega\|_{L^\infty}$

average ω along levelsets of ψ



matrix hydrodynamics

$W \in U(N)$

$I_f^N(W) = \text{tr}(f(W))$

$H^N(W) = \frac{1}{2} \text{tr}(PW)$

eigenvalues of W

eigenvectors of W

spectral norm

projections of W onto

$\text{stab}_P = \{W_s \in U(N) \mid [P, W_s] = 0\}$