PHYS 243A Homework 1

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1.1

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If we wish to produce a thin-film magnetic storage device with 100 Gbits/in^2 , each bit is to be 5 times as big in one dimension as the other, the total amount of open area between bits is to be the same as the total area occupied by the bits, and the film thickness of the medium storing the information is to be 10 nm, how many atoms are involved in each bit? Assume for simplicity that the film is pure Co, with a density of $9.0 \times 1022 \text{ atoms/cm}^3$.

Let's first visualize our bit with a nice image.

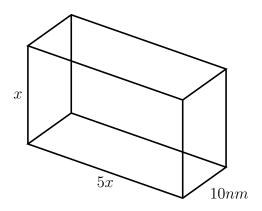


Figure 1: A single bit

We are told that we want 100 Gbits/in^2 and that half of the area is open, thus

$$A_{bits} = 0.5 \text{ in}^2 = 3.23 \times 10^{-4} \text{ m}^2.$$

Each square inch contains 100 Gbits or 10^11 bits. We can compare this to the knowledge that a single bit has an cross sectional area of $A = 5x^2$, so we set up a simple ratio to solve for x:

$$\frac{3.23 \times 10^{-4} \,\mathrm{m}^2}{10^{11} \,\mathrm{bits}} = \frac{5x^2 \,\mathrm{m}^2}{1 \mathrm{bit}}$$

$$x = \sqrt{\frac{3.23 \times 10^{-4}}{5 \times 10^{11}}} \,\mathrm{m}$$

$$x = 2.54 \times 10^{-8} \,\mathrm{m}$$

$$x = 25.4 \,\mathrm{nm}$$

We can now directly calculate the volume of the bit as pictured in Figure 1.

$$V_{bit_stack} = 5(25.4)^2 10$$

= 32258 nm³
= 3.23 × 10⁻¹⁷ cm³

Now we can use the density of Co atoms to find out how many atoms are in a single bit.

$$\frac{3.23 \times 10^{-17} \, \mathrm{cm}^3}{1 \, \mathrm{bit}} \cdot \frac{9 \times 10^{22} \, \mathrm{atoms}}{1 \, \mathrm{cm}^3} = \boxed{\frac{2.9 \times 10^6 \, \mathrm{atoms}}{1 \, \mathrm{bit}}}$$

1.2

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- (a) Begin with the Maxwell-Boltzmann distribution for molecular velocities in an ideal gas as expressed in x,y,z coordinates, and derive by integration the formula for the rate at which molecules strike a flat surface of unit area perpendicular to one of the axes. From this, determine the time necessary to form a monolayer of gas on a surface, assuming a general sticking probability of PS. Note that this is the same derivation that can be found in many physical chemistry texts to predict the rate at which a gas leaks out through an orifice.
- (b) Now make the assumption that only the bare surface area remaining on the surface after a given exposure time is active for a particular gas adsorption, and that PS = unity on the bare area, but zero elsewhere. Derive the general form of PS as a function of time for this case.
- (c) If a surface is exposed to 10^{-9} Torr of CO at ambient temperature, how long will it take to form the first monolayer:
- (i) If Ps = unity?
- (ii) If Ps follows the relationship of part (b)?

Part a:

We begin with the Maxwell-Boltzmann distribution for molecules in an ideal gas.

$$D(v) = \sqrt{\left(\frac{m}{2\pi kt}\right)^3} 4\pi v^2 \exp\left(-\frac{mv^2}{2kt}\right)$$

Where D is the probability of finding a molecule at speed v. The speed $v^2 = v_x^2 + v_y^2 + v_z^2$ is in three dimensions. We first want to find the average speed of the molecules, then the flux through a single side of a unit cube when the molecules have a given number density. To find an average speed we add all the speeds up weight by their probability.

$$\bar{v} = \sum v D(v) \Delta v,$$

where Δv represents a discrete spacing of speeds. In reality we can have continuum of speeds so to find the average speed we must sum over all possible speeds, which means an integral.

$$\bar{v} = \int_0^\infty v D(v) dv$$

$$= \int_0^\infty \sqrt{\left(\frac{m}{2\pi kt}\right)^3} 4\pi v^3 \exp\left(-\frac{mv^2}{2kt}\right)$$

$$= \left(\frac{m}{2\pi kt}\right)^{3/2} 4\pi \int_0^\infty v^3 \exp\left(-\frac{mv^2}{2kt}\right)$$

Now we make a substitution to make the integration easier:

$$x = mv^2/2kt$$

. Finding dx/dv and v^2 in terms of x, we make the substitutions and get,

$$\bar{v} = \left(\frac{m}{2\pi kt}\right)^{3/2} 4\pi \left(\frac{kt}{m}\right) \left(\frac{2kt}{m}\right) \int_0^\infty x \exp\left(-x\right)$$

Looking up this solution in a table (I used Schaums) we find that the integral evaluates to unity and we end up with.

$$\bar{v} = \left(\frac{m}{2\pi kt}\right)^{3/2} 4\pi \left(\frac{kt}{m}\right) \left(\frac{2kt}{m}\right)$$
$$= \sqrt{\frac{8kT}{\pi m}}$$

The flux of particles on the surface of the entire chamber is

$$J_{cube} = n\bar{v},$$

where n is the number density of particles (or molecules) and \bar{v} is the average speed of the particles. The particles are traveling in all directions with equal probability and at the same speed. We are only interested in the flux through a single face so we can integrate over all the particles that pass through a plane of area A in the chamber,

$$J_{plane} = \int_{\theta=0}^{2\pi} \int_{\phi=0}^{\pi/2} \vec{J} \cdot d\vec{A} \sin(\theta) d\theta d\phi$$
$$= JA2\pi \int_{\phi=0}^{\pi/2} \cos(\phi) \sin(\phi) d\phi$$
$$= JA2\pi \left[\frac{\sin(\phi)^2}{2} \right]_0^{\pi/2}$$
$$= \pi JA$$

With this we can say that the number of molecules that will stick on a typical surface of area A per second is,

$$\alpha = \pi J P_s A t$$

where P_s is some general sticking probability for the impinging molecules/atoms/particles. If we divide out the area of the surface then we get the surface density of the layer on the surface after a time t.

$$\rho_s = \pi J P_s t$$

The time it takes to form a monolayer is given by,

$$t = \frac{\rho_s}{\pi J P_s}$$

$$= \frac{\rho_s}{\pi n \bar{v} P_s}$$

$$= \frac{\rho_s}{\pi n P_s} \sqrt{\frac{\pi m}{8k_b T}}$$

$$= \frac{\rho_s [pts/m^2]}{n[pts/m^3] P_s} \sqrt{\frac{m[kg]}{8\pi k_b [J/K]T[K]}}$$

We make a few substitutions to make things a bit easier,

$$n = \frac{N}{V} = \frac{P[Pa]}{k_b[J/K]T[K]} = 7.5 \times 10^{-3} \frac{P[Torr]}{k_b[J/K]T[K]}$$
$$m[kg] = 10^{-3} \frac{M}{N_A}$$

Using these substitutions, we find that the time to form a monolayer, in seconds, is,

$$t = \frac{\rho_s k_b T}{7.5 \times 10^{-3} P P_s} \sqrt{\frac{10^{-3} M}{8\pi k_b T N_A}}$$

$$= \frac{0.84 \rho_s}{P P_s} \sqrt{\frac{M k_b T}{N_A}}$$

$$= \frac{0.4 \times 10^{-23} \rho_s [pts/m^2]}{P [Torr] P_s} \sqrt{M [AMU] T [K]}$$

Part b:

Assuming the initial sticking probability is unity then the fraction of the covered surface on the substrate is going to increase very quickly and eventually converage to unity at time inifinty. Let's call this fraction β . We can then write the sticking probability as a function of β .

$$P_s(t) \propto 1 - \beta(t)$$

$$\propto 1 - \left(1 - e^{-c_0 t}\right)$$

$$\propto e^{-c_0 t}$$

$$= c_1 e^{-c_0 t}$$

At t = 0 the sticking probability should be unity thus $c_1 = 1$. We do not have a way of obtaining the constant c_0 so we leave the sticking probability in this general form.

$$P_s(t) = e^{-c_0 t}$$

Part c:

We will assume that the chamber is filled with nitrogen gas so M=28, at ambient temperature T=298K, and a typical surface density is $\rho_s=10^19$ atoms /m². If $P_s=1$ and $P=10^{-9}$ Torr then the time to form a monolayer is,

$$t(10^{-9}) = \frac{0.4 \times 10^{-23} \cdot 10^{19}}{10^{-9}} \sqrt{28 \cdot 298}$$
$$= 3.654 \times 10^{6} \text{ seconds}$$

If the conditions from part b describe the sticking probability then a complete monolayer will probably never form since the sticking probability gets arbitrarily close to zero which means the time to form a monolayer will go to infinity.

1.3

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- (a) What would be the minimum energy required to take a cube of Pt metal 1.0 cm on a side at room temperature and disperse it into tiny cubic "nanoparticles" of 10-6 cm on a side? Assume that this is done in a perfect ultrahigh vacuum environment, with the surfaces in equilibrium with the very low vapor pressure of Pt, and make use of the argument in Zangwill, p. 12, but with a key equation corrected to
- (b) Estimate the fraction of the atoms that are on the surface of the 1 cm cube. Of the 10-6 cm cube. Assume that the density of Pt atoms is 6.62×1022 cm-3.