

# Lecture 1

## Astrophysical Magnetic Fields

Magnetic Universe, Basics of MHD, Detecting magnetic fields and more

IIA Summer Programme 2022



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The Magnetic Universe

Detecting magnetic fields

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Mean-field dynamos

Why think about magnetic fields?

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# Agenda

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## Books and References

- Books and References on Magnetohydrodynamics
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  - Heald et. al., (2020) - *Magnetism Science with the Square Kilometre Array*, Galaxies, Vol. 8
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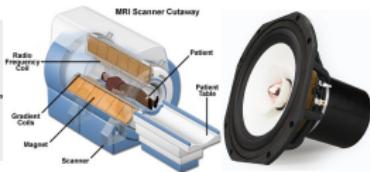
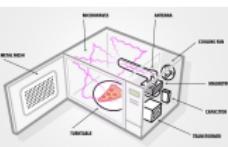
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Why think about magnetic fields?

# Introduction

- Magnets are ubiquitous in our lives



- From credit cards, microwave ovens, speakers to medical instruments
- Until recently, mariners relied on magnetic compass for navigation
- Some birds also seem to have a **Magnetic Sixth sense** - Magneto-reception
- In fact, magnetism seems to be everywhere in the **Cosmos!!**



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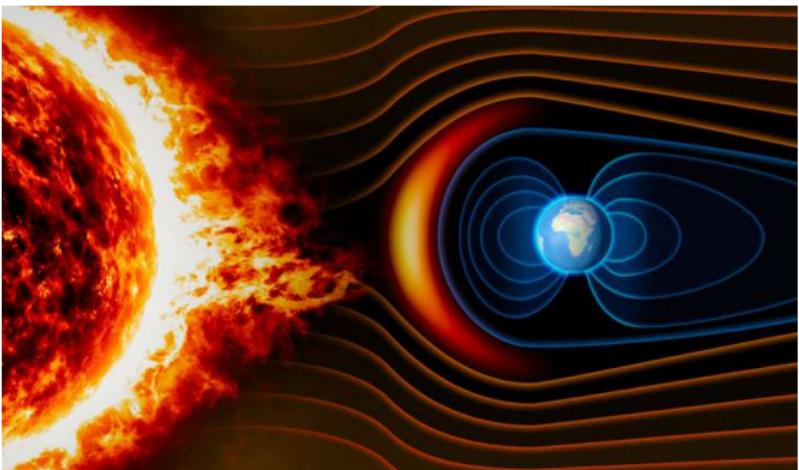
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#### Why think about magnetic fields?

- Earth has a Dipolar magnetic field structure



Credit : *The Week*, April 2019

- Field strengths  $\approx 1\text{ G}$ , with irregular reversals over a million years!
  - It is not just a handy navigation aid!
  - The Earth's magnetic field is **vital for the existence of life**
- Shields us from high energy particles arriving from the Sun
- Particles approach near the Earth's surface only at the poles



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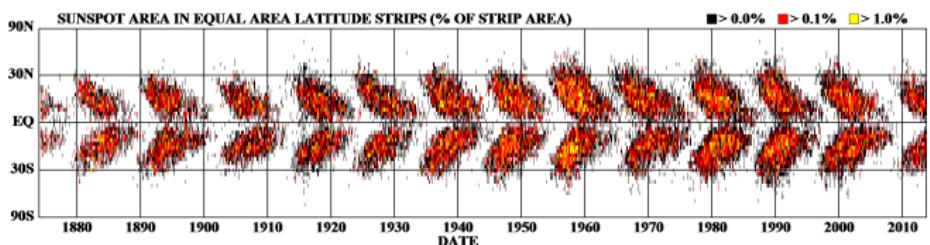
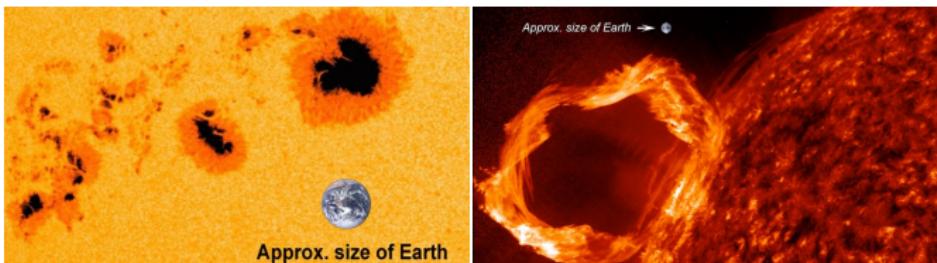
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# The Magnetic Universe

- Magnetism of the Sun responsible for a whole range of phenomena



- Strong fields  $\sim 3 \times 10^3$  G in sunspots
- Butterfly Diagram** : Variation of sunspot number
- Exhibits plethora of features - Solar prominences, Coronal mass ejections; test bed for MHD theories



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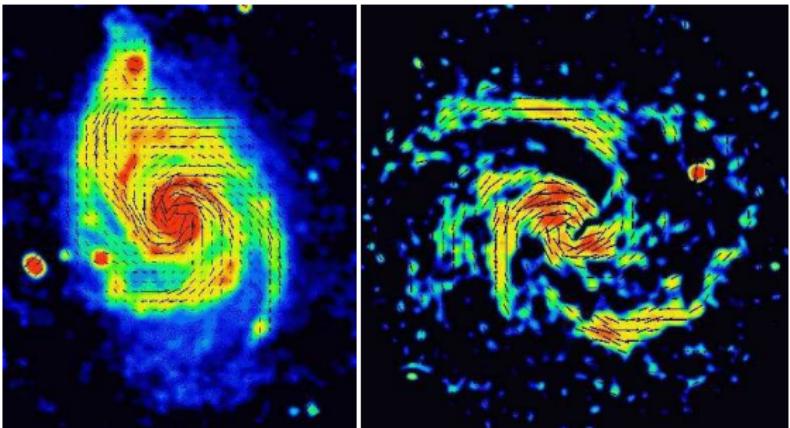
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- Magnetic fields in spiral galaxies (Left : M31, Right : NGC 6946)



- Spiral galaxies** : Thin rotating discs of  $\sim 10^{10}$  stars and interstellar gas, multiphase interstellar medium (ISM)
- Interstellar gas** :  $\langle n \rangle \sim 1 \text{ cm}^{-3}$ ,  $10^{-3} < n < 10^6 \text{ cm}^{-3}$ ,  
 $10 < T < 10^7 \text{ K}$
- $\langle \mathbf{B}_{\text{tot}} \rangle = 9 \mu\text{G}$  in a sample of 74 spirals; large-scale fields about half the value of the random field
  - Large-scale fields correlated on  $10 \text{ kpc}$  scales
  - Similar field strengths found in interacting galaxies



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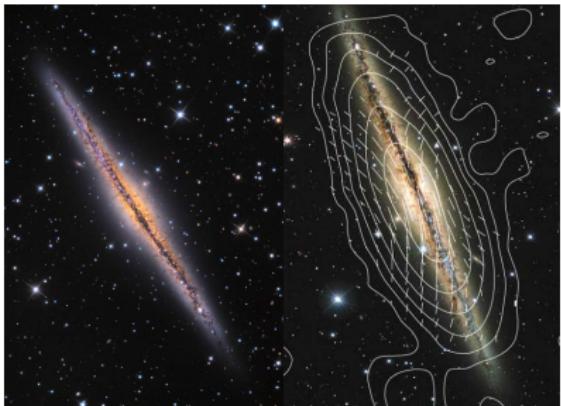
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- Magnetic fields in **gaseous halos of galaxies**
  - Hot, ionized, quasi-spherical envelopes of galactic disks



NGC 891, Credit : NASA & MPIfR, Bonn

- Field runs parallel to the plane near the disk, vertical components above and below the plane forms an **X shaped structure** in the halo
  - Field strength in the halos similar to those in the disks
  - **How to explain the occurrence of these fields?**



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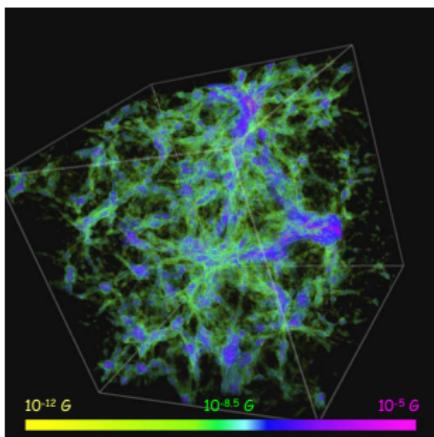
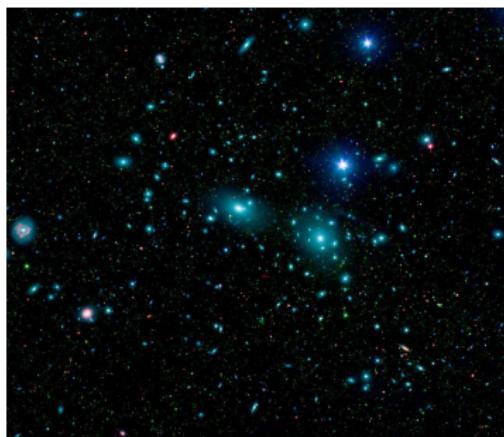
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- **Galaxy Clusters** : Largest gravitationally bound system in our Universe;  $M \sim 10^{14} - 10^{15} M_{\odot}$ , size of several Mpc
  - Baryonic matter contained in **hot X-ray emitting ICM** ( $T \sim 10^7 - 10^8$  K,  $n \sim 10^{-2} - 10^{-4}$  cm $^{-3}$ )
  - Field strengths  $\approx \mu\text{G}$  ordered on several kpc scales; fields detected with the help of **Faraday rotation measure**



- Left : Coma Cluster (SDSS/SST), Right : Ryu et al., 2008, Science
- No overall rotation; **What generates and maintains the magnetic field?**



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## Fundamental Questions

- How to **detect** the presence of magnetic fields?
- How did such magnetic fields **arise** and how are they **maintained**?
- How do we describe the motion of a conducting fluid?
  - MHD - study of the magnetic properties of electrically conducting fluids
- Why is it even necessary to **think about magnetic fields**?

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## Detecting the presence of magnetic fields

- Magnetism is invisible!! Is there a way to detect their presence?
  - Recall the simple experiment with a bar magnet and iron filings



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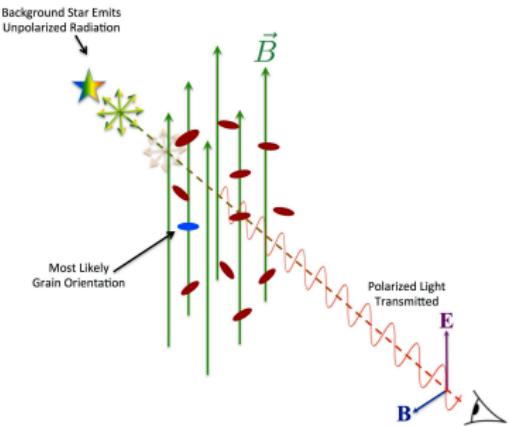
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## • How do we observe magnetic fields in the sky?

- Zeeman splitting of spectral lines - mostly used for the Sun
- Light polarization by interstellar dust
- Faraday rotation of polarised emission and synchrotron emission

## Detecting the presence of magnetic fields

- 1949 : John Hall & Albert Hiltner independently showed that **star light is polarised**



- Refers to the orientation of the oscillation of light waves
- Star light is expected to be **unpolarised!**
- **What then causes star light to become polarized?**

- Interstellar space is dusty; **Dust particles act like tiny compasses** in the presence of magnetic field
- Polarised starlight reveals the presence of magnetic fields



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# Detecting the presence of magnetic fields

- Faraday Rotation : Magneto - Optical Phenomenon

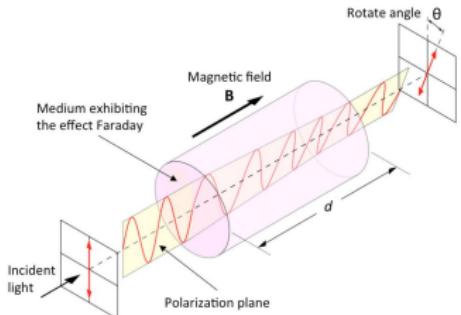


Image Credit : Wikipedia

- Discovered by Michael Faraday in 1845
- Angle of polarisation rotates as the light passes through a foreground magnetized region
  - Information about the line-of-sight component of the magnetic field and its direction
  - Stronger the field, the more rotation is produced

- Amount of rotation :  $\psi = \psi_0 + RM\lambda^2$ ;  $RM \propto \int_0^L n_e B_{||} ds$

## Faraday rotation measure in galaxies

$$RM = 0.81 \frac{\text{rad}}{\text{m}^2} \int_0^L \frac{n_e}{1 \text{ cm}^{-3}} \frac{B_{||}}{1 \mu\text{G}} \frac{dl}{1 \text{ pc}}$$

- Normalizations need to be appropriately adjusted for galaxy clusters



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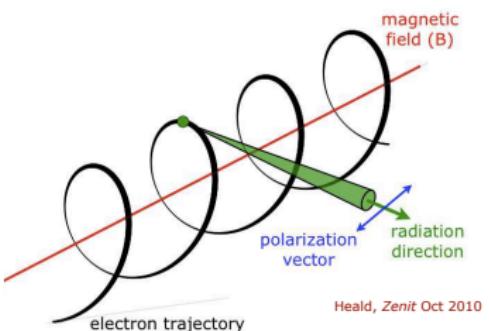
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# Detecting the presence of magnetic fields

- Polarized Synchrotron Emission



- Left : **Synchrotron emission**, Right : **100m Effelsburg radio telescope**
- Produced by relativistic electrons spiraling around magnetic field lines
  - Information about the total strength of the magnetic field
  - The Effelsburg radio telescope played a pioneering role in inferring about galactic magnetic fields way back in the 70-80's



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# Detecting the presence of magnetic fields

- Total Intensity of the synchrotron emission

$$I_\nu = \int_0^L \epsilon_\nu \, ds \propto \int_0^L n_{\text{cr}} B_\perp^2 \, ds,$$



- $\epsilon_\nu$  is the synchrotron emissivity and  $n_{\text{cr}}$  is the number density of cosmic ray electrons,  $B_\perp$  is the magnetic field in the sky plane
- Polarized intensity :  $PI_\nu$  and Fractional polarisation : ( $p_\nu$ )

$$PI_\nu = \sqrt{Q_\nu^2 + U_\nu^2}, \quad p_\nu = PI_\nu / I_\nu,$$

where  $Q_\nu, U_\nu$  are the Stokes parameters

- In galaxies :  $PI_\nu \propto \int_0^L n_{\text{cr}} \bar{B}_\perp^2 \, ds$ ;  $\mathbf{B} = \bar{\mathbf{B}} + \mathbf{b}$ ,  $\langle \mathbf{B} \rangle = \bar{\mathbf{B}}$ ,  $\langle \mathbf{b} \rangle = 0$ 
  - Traces the regular magnetic field whereas  $I - P$  traces the turbulent magnetic field
- Together with Faraday Rotation measure, synchrotron emission provides the most important observational methods for galactic and extragalactic magnetic fields

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## MHD : Introduction and Objective

- MHD = Equations of fluid dynamics + Maxwell's equations
- Plasma treated as a continuous medium, described by a single temperature, density and bulk velocity
- For a pure hydrodynamical fluid, description completely specified by
  - $\rho \rightarrow$  Mass Density,  $\mathbf{v} \rightarrow$  Flow velocity,  $p \rightarrow$  Pressure
  - Governing equations derived from conservation laws
- However, description of a conducting fluid requires additional variables
  - $\rho_e \rightarrow$  Charge density,  $\mathbf{J} \rightarrow$  Current density
  - $\mathbf{E} \rightarrow$  Electric field,  $\mathbf{B} \rightarrow$  Magnetic field
- Objective : To find a set of closed equations describing the evolution of these variables



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# Conservation Laws

- Conservation laws can be used to derive evolution equations
- Conservation of Mass  $\Rightarrow$  **Continuity Equation**

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

- Need an evolution equation for the fluid velocity
- Conservation of Momentum  $\Rightarrow$  **Navier-Stokes Equation**

$$\rho \left[ \frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} \right] = -\nabla p + \rho \nu \nabla^2 \mathbf{v} + \rho \mathbf{F} + \text{Extra terms}$$

where  $\nu$  is the viscosity of the fluid

- Reynolds number :  $\text{Re} = |(\mathbf{v} \cdot \nabla) \mathbf{v}| / |\nu \nabla^2 \mathbf{v}| = v L / \nu$
- Additional terms when the fluid is conducting
- **Check** : Fluid Mechanics by Landau & Lifshitz



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## Maxwell's Equations (Gaussian CGS units)

$$\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}, \quad \nabla \cdot \mathbf{B} = 0$$

$$\frac{1}{c} \frac{\partial \mathbf{E}}{\partial t} = \nabla \times \mathbf{B} - \frac{4\pi}{c} \mathbf{J}, \quad \nabla \cdot \mathbf{E} = 4\pi\rho_e$$



- Need a relation between the current density and the fields
- If the fluid is moving, what fields should we use?
  - Fields in the fluid's local rest frame :  $\{\mathbf{J}', \mathbf{E}', \mathbf{B}'\}$
  - Fields in the laboratory frame :  $\{\mathbf{J}, \mathbf{E}, \mathbf{B}\}$
- Ohm's Law in the fluid's local rest frame  $\mathbf{J}' = \sigma \mathbf{E}'$ , where  $\sigma$  is the conductivity
- Relate  $\mathbf{E}'$  and  $\mathbf{B}'$  to  $\mathbf{E}$  and  $\mathbf{B}$  in the lab frame
  - Carry out Lorentz transformation between frames

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## Lorentz transformation between frames

$$E'_{\parallel} = E_{\parallel}, \quad B'_{\parallel} = B_{\parallel}$$

$$\mathbf{E}'_{\perp} = \gamma \left( \mathbf{E}_{\perp} + \frac{\mathbf{v}}{c} \times \mathbf{B}_{\perp} \right), \quad \mathbf{B}'_{\perp} = \gamma \left( \mathbf{B}_{\perp} - \frac{\mathbf{v}}{c} \times \mathbf{E}_{\perp} \right),$$

where  $\gamma = 1/\sqrt{1 - v^2/c^2}$  is the Lorentz factor

- Can be simplified further if velocities are assumed to be **non-relativistic**,  $\Rightarrow$  neglect terms of order  $v^2/c^2$ 
  - Lorentz factor  $\gamma \approx 1$
  - The electric and magnetic fields are related by

$$\mathbf{E}' = \mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B}, \quad \mathbf{B}' = \mathbf{B},$$

- **Exercise :** Show that to order  $|v|/c$ ,  $\mathbf{J}' = \mathbf{J}$
- $\mathbf{J}/\sigma = \lambda \mathbf{J} = \mathbf{E}' = \mathbf{E} + \mathbf{v} \times \mathbf{B}/c$ , where  $\lambda = \sigma^{-1}$  is the resistivity of the fluid



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- Ohm's Law for a conducting fluid

$$\mathbf{J} = \sigma \left( \mathbf{E} + \frac{\mathbf{v} \times \mathbf{B}}{c} \right), \quad |E| \approx \frac{|v|}{c} |B|$$

- Solve for  $\mathbf{E}$ , substitute in the Faraday equation and neglect displacement current  $\Rightarrow$  Induction Equation

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B} - \eta \nabla \times \mathbf{B}),$$

where,  $\eta = c^2 / 4\pi\sigma$  is the magnetic diffusivity

## Simple consequences

- $\mathbf{v} = 0 \Rightarrow$  Pure diffusion and decay
- $\eta \rightarrow 0$ , the flux  $\Phi = \int_S \mathbf{B} \cdot d\mathbf{S}$  is frozen,  $d\Phi/dt \rightarrow 0$
- Magnetic Reynolds number :  $Rm = v L / \eta$ , for astrophysical systems  $Rm \gg 1$



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## MHD Basics

- Instructive to clarify the role of  $\nabla \times (\mathbf{v} \times \mathbf{B})$  term

$$\nabla \times (\mathbf{v} \times \mathbf{B}) = -(\mathbf{v} \cdot \nabla) \mathbf{B} + (\mathbf{B} \cdot \nabla) \mathbf{v} - \mathbf{B}(\nabla \cdot \mathbf{v})$$

- Advection (1st term), Stretching (2nd term) and Compression (3rd term)
- Exercise :** Find the solution of the induction equation when  $\mathbf{v} = (0, Sx, 0)$ ,  $\mathbf{B}_0 = (B_0, 0, 0)$  for  $\eta = 0$ .
- Important parameters in different astrophysical settings

	T [K]	$\rho$ [g cm $^{-3}$ ]	$P_m$	$u_{rms}$ [cm s $^{-1}$ ]	L [cm]	$R_m$
Solar CZ (upper part)	$10^4$	$10^{-6}$	$10^{-7}$	$10^6$	$10^8$	$10^6$
Solar CZ (lower part)	$10^6$	$10^{-1}$	$10^{-4}$	$10^4$	$10^{10}$	$10^9$
Protostellar discs	$10^3$	$10^{-10}$	$10^{-8}$	$10^5$	$10^{12}$	10
CV discs and similar	$10^4$	$10^{-7}$	$10^{-6}$	$10^5$	$10^7$	$10^4$
AGN discs	$10^7$	$10^{-5}$	$10^4$	$10^5$	$10^9$	$10^{11}$
Galaxy	$10^4$	$10^{-24}$	( $10^{11}$ )	$10^6$	$10^{20}$	( $10^{18}$ )
Galaxy clusters	$10^8$	$10^{-26}$	( $10^{29}$ )	$10^8$	$10^{23}$	( $10^{29}$ )

- Magnetic Prandtl number :  $Pr_M = Rm/Re = \nu/\eta$
- Galaxies and clusters have very large  $Pr_M$  due to very low densities and much higher temperatures
- Reference : Brandenburg & Subramanian, Physics Reports, 2005



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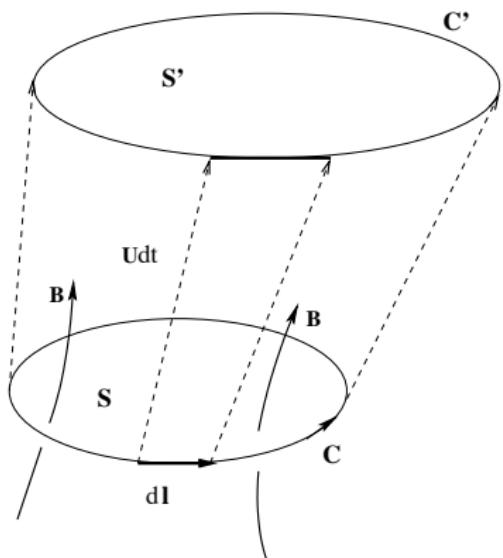
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# Flux Freezing



- Magnetic flux

$$\Phi = \int_S \mathbf{B} \cdot d\mathbf{S}$$

- Interested to know the time rate of change of  $\Phi$

- Change in the flux :

$$d\Phi = \int_{S'} \mathbf{B}(t + dt) \cdot d\mathbf{S} - \int_S \mathbf{B}(t) \cdot d\mathbf{S}$$



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# Flux Freezing

- Using the fact that  $\nabla \cdot \mathbf{B} = 0$  at time  $t + dt \Rightarrow$

$$\int_{S'} \mathbf{B}(t + dt) \cdot d\mathbf{S} = \int_S \mathbf{B}(t + dt) \cdot d\mathbf{S} - \oint_C \mathbf{B}(t + dt) \cdot (d\mathbf{l} \times \mathbf{U} dt),$$

- Therefore,

$$d\Phi = \int_S [\mathbf{B}(t + dt) - \mathbf{B}(t)] \cdot d\mathbf{S} - \oint_C \mathbf{B}(t + dt) \cdot (d\mathbf{l} \times \mathbf{U} dt).$$

$$\frac{d\Phi}{dt} = \int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} - \oint_C (\mathbf{U} \times \mathbf{B}) \cdot d\mathbf{l}$$

- Using  $\oint_C (\mathbf{U} \times \mathbf{B}) \cdot d\mathbf{l} = \int_S \nabla \times (\mathbf{U} \times \mathbf{B}) \cdot d\mathbf{S}$

$$\frac{d\Phi}{dt} = \int_S \left[ \frac{\partial \mathbf{B}}{\partial t} - \nabla \times (\mathbf{U} \times \mathbf{B}) \right] \cdot d\mathbf{S} = \eta \int_S (\nabla^2 \mathbf{B}) \cdot d\mathbf{S}$$

- As  $\eta \rightarrow 0$ ,  $d\Phi/dt \rightarrow 0$ . Therefore  $\Phi$  is a constant
- Flows in a conducting fluid can amplify magnetic fields



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# The Lorentz force - Influence of magnetic field on velocity

- Magnetic field influences the velocity through the **Lorentz force**

$$\mathbf{F}_L = q \left[ \mathbf{E} + \frac{\mathbf{V} \times \mathbf{B}}{c} \right]$$

- Consider a conducting fluid with  $n_i$  ions and  $n_e$  electrons per unit volume
- The Lorentz force density is given by,

$$\begin{aligned}\mathbf{f}_L &= +en_i \left[ \mathbf{E} + \frac{\mathbf{u}_i \times \mathbf{B}}{c} \right] - en_e \left[ \mathbf{E} + \frac{\mathbf{u}_e \times \mathbf{B}}{c} \right] \\ &= \rho_e \mathbf{E} + [+en_i \mathbf{u}_i - en_e \mathbf{u}_e] \times \mathbf{B}/c \sim \frac{\mathbf{J} \times \mathbf{B}}{c}\end{aligned}$$

- $|\rho_e \mathbf{E}| / |(\mathbf{J} \times \mathbf{B})/c| \sim v^2/c^2 \ll 1, \Rightarrow \mathbf{F}_L$  due to  $\mathbf{E}$  negligible
- Using  $\nabla \times \mathbf{B} = (4\pi/c)\mathbf{J}$

$$\mathbf{F}_L = \frac{(\nabla \times \mathbf{B}) \times \mathbf{B}}{4\pi} = -\nabla \left( \frac{\mathbf{B}^2}{8\pi} \right) + \frac{(\mathbf{B} \cdot \nabla) \mathbf{B}}{4\pi}$$

- For straight field lines,  $(\mathbf{B} \cdot \nabla) \mathbf{B} = 0$
- Magnetic pressure can still be non zero



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# The Lorentz force - Influence of magnetic field on velocity

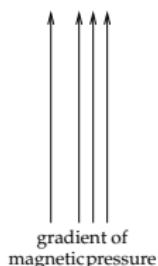
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# MHD equations

## Full set of MHD equations

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) &= 0, \\ \rho \left[ \frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} \right] &= -\nabla p - \rho \nabla \phi + \rho \nu \nabla^2 \mathbf{v} + \frac{\mathbf{J} \times \mathbf{B}}{c}, \\ \frac{\partial \mathbf{B}}{\partial t} &= \nabla \times (\mathbf{v} \times \mathbf{B} - \eta \nabla \times \mathbf{B}), \\ \frac{\partial \rho E}{\partial t} + \nabla \cdot [\mathbf{v}(\rho E + p_*) - \mathbf{B}(\mathbf{v} \cdot \mathbf{B})] &= \rho \mathbf{g} \cdot \mathbf{v} + \nabla \cdot (\mathbf{v} \cdot \tau + \sigma \nabla T) \\ &+ \nabla \cdot [\mathbf{B} \times (\eta(\nabla \times \mathbf{B})], \\ \nabla^2 \phi &= 4\pi G \rho. \end{aligned}$$

- $p_* = p + \mathbf{B}^2/8\pi$
- $E = \mathbf{v}^2/2 + \epsilon + \mathbf{B}^2/8\pi\rho$



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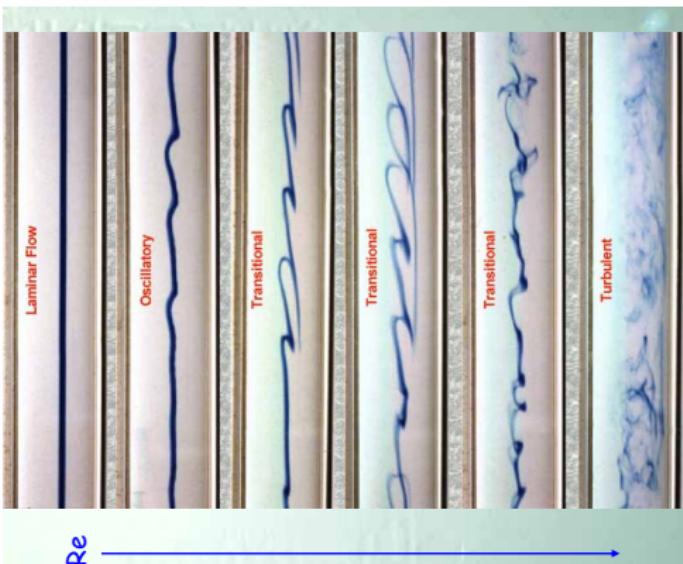
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# Turbulence

- What kind of velocities are we talking about?
  - In astrophysical systems, velocities are turbulent
- Turbulence - a flow regime characterized by random variations of pressure and velocity in space and time
  - Onset of turbulence is determined by the Reynolds number with  $Re > 1000 \Rightarrow$  turbulent regime



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# Turbulence in Astrophysics

- Turbulence requires a continuous supply of energy

## Sources

- Instabilities in a flow - Shear instability
- Solar convection
- Cosmological structure formation shocks, merger events
- Supernovae explosion in the ISM
- From subsonic (in cluster cores) to supersonic (in the ISM)

## Significance

- Energy transfer from large scales of motion
- Jupiter's great Red spot
- Augments molecular transport - causing mixing of the fluid
- Large/Small -scale field generation via turbulent dynamo

- Check out : Kolmogorov's hypothesis



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# The Bicycle Dynamo

- Consider the simple example of a **Bicycle Dynamo**



- Mechanical energy transformed to electrical energy
  - Peddling action rotates the magnet; changing magnetic field induces electric currents ⇒ **illuminates the light bulb**
- In astrophysical objects **there are no external magnets, wires, frames etc.**
  - Kinetic energy in fluid motions tapped to amplify magnetic energy



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# Generation of Magnetic fields

- Magnetic fields evolve as

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B},$$

where  $\mathbf{v}$  is a solution of the momentum equation

$$\rho \left[ \frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} \right] = -\nabla p + \rho \nu \nabla^2 \mathbf{v} + \frac{\mathbf{J} \times \mathbf{B}}{c}$$

- Dynamos** : class of velocity fields that allow a weak seed magnetic field to grow
- What if we start with a very weak magnetic field?**
  - Ratio of  $|(\mathbf{J} \times \mathbf{B})/c|/|\rho(\mathbf{v} \cdot \nabla) \mathbf{v}| \approx (B^2/8\pi)/(\rho v^2/2) \ll 1$
  - Either the field decays or it grows such that  $(B^2/8\pi)/(\rho v^2/2) \sim 1$
- Galaxies and Galaxy clusters are **turbulent systems**
  - If  $\langle B^2 \rangle$  grow? How? → **Turbulent Dynamos**
  - Fluctuation dynamos & Mean-field dynamos**
  - When and how do dynamos saturate - active area of research



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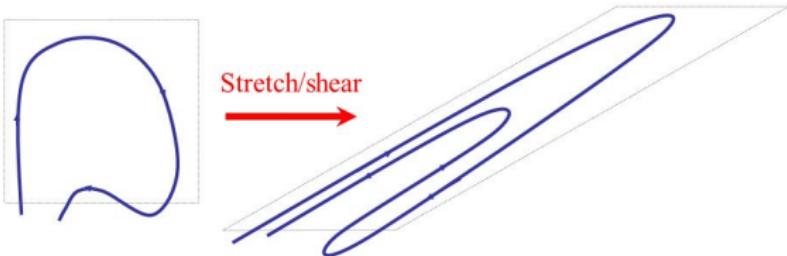
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## Fluctuation dynamos

- Ideally suited for amplifying fields in the ICM, may also operate in galaxies
  - Growth by random stretching by turbulent eddies
  - Field grows exponentially at first and then saturates
  - Saturation achieved **on a scale-by-scale basis** with smaller scales saturating the field at that scale first



*Batchelor 1950, Ruzmaikin & Zeldovich 1990, Childress & Gilbert 1995*

- Average magnetic energy evolution governed by **stretching**, **compression** and **dissipation** terms;  $\nabla \cdot \mathbf{u}$  term is negligible in subsonic flows

$$\frac{\delta}{\delta t} \langle \mathbf{b}^2 / 2 \rangle = \langle \mathbf{s}_{ij} \mathbf{b}_i \mathbf{b}_j \rangle - \left\langle \frac{1}{2} |\mathbf{b}|^2 (\nabla \cdot \mathbf{u}) \right\rangle - \langle \eta |\mathbf{j}|^2 \rangle$$



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## Characteristics of Fluctuation dynamos

- Generic in any random/turbulent flow for  $Rm > Rm_{cr} \sim 200$
- Field is amplified on the eddy-turnover time-scale;  $\tau_I \sim l/v_I \propto l^{2/3}$
- Growth time  $\sim 0.3$  Gyr in galaxy clusters
- Fields correlated on scales at most on the scale of turbulence
- Field structure appears to be highly intermittent; long tails in the PDF or increased values of the kurtosis



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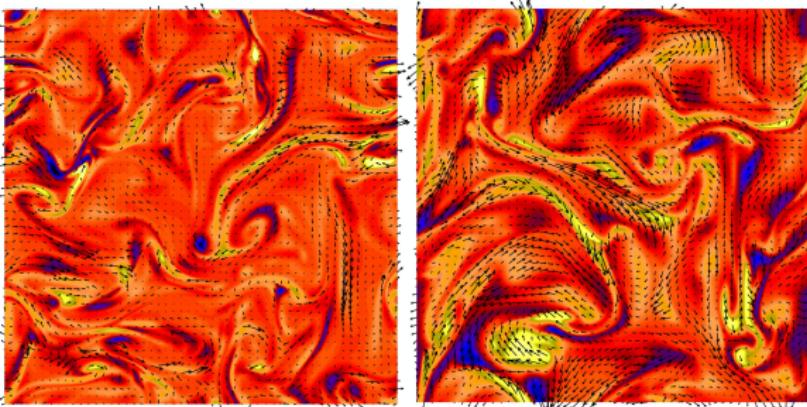
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- 2D snapshots of  $B_z/B_{\text{rms}}$  in the **kinematic** (left) and **saturated** phase (right)



Sur 2019, Sur, Basu & Subramanian 2021



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## Mean-field dynamos

- Theoretical ansatz :  $\mathbf{U} = \bar{\mathbf{U}} + \mathbf{u}$ ,  $\mathbf{B} = \bar{\mathbf{B}} + \mathbf{b}$ ,  $\bar{\mathbf{u}} = 0$ ,  $\bar{\mathbf{b}} = 0$
- The mean-field satisfies the dynamo equation

$$\frac{\partial \bar{\mathbf{B}}}{\partial t} = \nabla \times (\bar{\mathbf{U}} \times \bar{\mathbf{B}} + \bar{\mathcal{E}}) + \eta \nabla^2 \bar{\mathbf{B}}$$

- $\bar{\mathcal{E}} = \bar{\mathbf{u}} \times \bar{\mathbf{b}} \approx \alpha \bar{\mathbf{B}} - \eta_t \bar{\mathbf{J}}$
- $\alpha = (-\tau/3)\langle \omega \cdot \mathbf{u} \rangle$  is the mean helicity of turbulence; also known as the  $\alpha$  – effect
- $\eta_t = (\tau/3)\langle \mathbf{u}^2 \rangle$  is the turbulent magnetic diffusivity
- Large-scale magnetic field generation in galaxies : Interstellar medium in spiral galaxies are :
  - Rotating, stratified, contains electrically conducting fluid
  - Randomly stirred by supernovae and stellar winds
- Perfect conditions for the Galactic Dynamo to operate and generate large-scale fields



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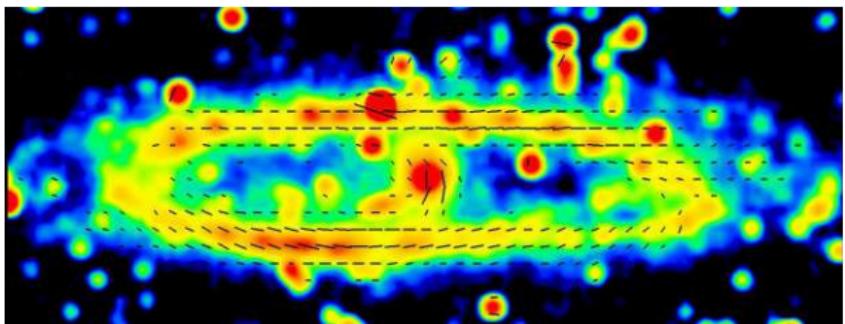
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- Large-scale magnetic fields in M31



Credit : MPIfR, Bonn

- Galactic shear generates  $\bar{B}_\phi$  from  $\bar{B}_r$
- Supernovae drive **helical turbulence** in the disk
- Helical motions generate  $\bar{B}_r$  from  $\bar{B}_\phi$  through the  $\alpha$ -effect
- Growth time of the magnetic field  $\sim 10^9$  yr



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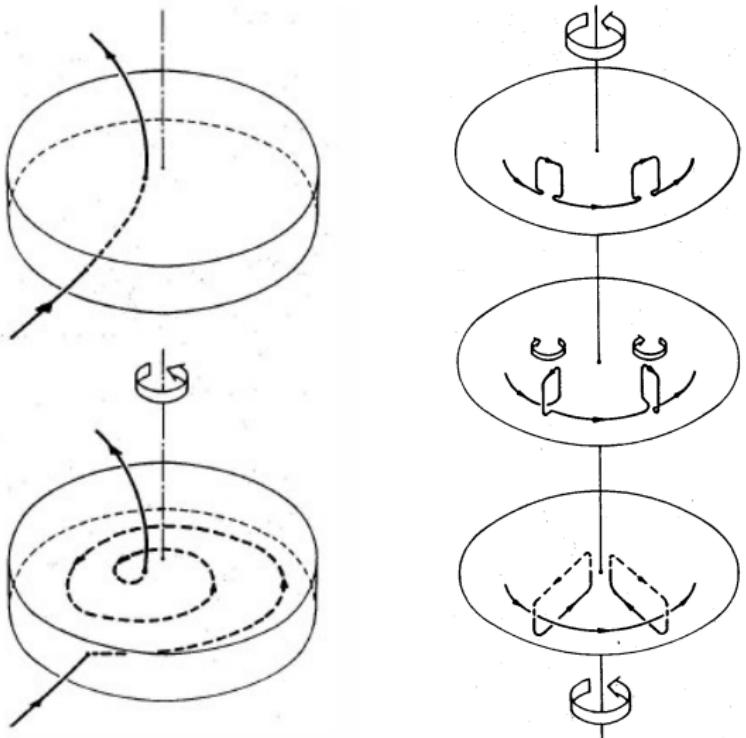
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# Mean-field dynamos

- Magnetic field generation in disk galaxies

Ruzmaikin, Sokoloff & Shukurov, 1988



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- Pursuit of one's curiosity! ⇒ Hallmark of a civilized society
- May lead to technology spin-offs; Potential to improve the material quality
- Implications for **Space Weather** - A variety of technologies rely heavily on near-Earth space conditions
- **Astrophysical context :**
  - Effects on Star-formation - collapse and fragmentation of clouds
  - Can affect mixing properties of fluids
  - Fields generated in the **First Stars** provides seed fields for the **First Galaxies**
- A new **Magnetic Era** to usher in with the **Square Kilometre Array (SKA)**



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### Quote from astronomer Lo Woltjer

The larger the one's ignorance, the stronger the magnetic field!

### Contact Details

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