Digital Signal Processing

Class 16 03/20/2025

ENGR 71

- Class Overview
 - Frequency Analysis of Discrete Signals
- Assignments
 - Exam 1 due March 23
 - Lab 2 due March 28
 - Reading:

Chapter 5: Frequency-Domain Analysis of LTI Systems

- Key concept behind the action of LTI systems on signals:
 - Signals can be decomposed into superposition of frequency components
 - Basis function for this decomposition are sines and cosines (and complex exponential)
 - Frequency components of signal are unchanged when passed through Linear Time Invariant systems
 - Only amplitude and phase change

Changes magnitude and phase $H(s) \qquad \qquad H(s)$ $H(s) \qquad \qquad M \cdot A \cdot \sin(\omega t + \varphi + \theta)$

- Some interesting relationship for $H(\omega)$, $H^*(\omega)$, and $|H(\omega)|$
 - Section 5.2.1 of Proakis & Manolakis derives interesting relationships for the frequency response described by rational polynomials when the coefficients of the polynomials are real-valued numbers.
 - This is generally the case since discrete LTI systems are defined by difference equations like:

$$y(n) + \sum_{k=1}^{N} a_k y(n-k) = \sum_{k=0}^{M} b_k x(n-k)$$
 where the a's and b's are weights of delayed inputs and outputs

$$y(n) + \sum_{k=1}^{N} a_k y(n-k) = \sum_{k=0}^{M} b_k x(n-k) \text{ where the } a \text{'s and } b \text{'s are weights of delayed inputs and outputs.}$$

$$\text{Transfer function: } H(z) = \frac{B(z)}{A(z)} = \frac{\sum_{k=0}^{M} b_k z^{-k}}{1 + \sum_{k=1}^{N} a_k z^{-k}}$$

Frequency response:
$$H(\omega) = H(z)|_{z=e^{j\omega}} = \frac{B(\omega)}{A(\omega)} = \frac{\sum_{k=0}^{M} b_k e^{-j\omega k}}{1 + \sum_{k=1}^{N} a_k e^{-j\omega k}}$$

Writing this in terms of poles and zeros

$$H(\omega) = b_0 \frac{\prod_{k=1}^{M} (1 - z_k z^{-1})}{\prod_{k=1}^{N} (1 - p_k z^{-1})} = b_0 \frac{\prod_{k=1}^{M} (1 - z_k e^{-j\omega})}{\prod_{k=1}^{N} (1 - p_k e^{-j\omega})}$$

While the a's and b's are real-valued, the poles and zeros $(z_k$'s and p_k 's) can be real or complex, and, if complex, occur in complex conjugate pairs.

• What are the implications of this?

Skipping the details covered in the text:

$$H^*(\omega) = H(-\omega)$$

$$|H(\omega)|^2 = H(\omega)H(-\omega) = H(z)H(z^{-1})|_{z=e^{j\omega w}}$$

One of the properties of the z-transform is that $H(z)H(z^{-1})$ is the z-transform of the autocorrelation of the impulse response.

• What are the implications (continued)?

One of the properties of the z-transform is that $H(z)H(z^{-1})$ is the z-transform of the autocorrelation of the impulse response.

Finally, since $|H(\omega)|^2 = H(z)H(z^{-1})$ and using the fact that the energy spectral density of a signal is the Fourier transform of the autocorrelation. (Wiener-Khintchine theorem)

 $|H(\omega)|^2$ is the Fourier transform of the autocorrelation of the impulse response, $[r_{hh}(m)]$

• Other results as a consequence of transfer function being rational polynomial with coefficients of polynomial real numbers and using autocorrelation relationship to z-transform:

Can show that
$$|H(\omega)|^2 = \frac{d_0 + 2\sum_{k=1}^{M} d_k \cos(k\omega)}{c_0 + 2\sum_{k=1}^{M} c_k \cos(k\omega)}$$

where c's and d's are autocorrlation of a's and b's in polynomial

$$c_l = \sum_{k=0}^{N-|l|} a_k a_{k+l} - N \le l \le N, \quad c_l = c_{-l} \quad \text{since a's and b's are real}$$

$$d_l = \sum_{k=0}^{M-|l|} b_k b_{k+l} - M \le l \le M, \quad d_l = d_{-l} \quad \text{since a's and b's are real}$$

From trigonometric identity
$$\cos(k\omega) = \sum_{m}^{k} \beta_{m} (\cos \omega)^{m}$$

• The upshot of all this:

The magnitude squared of the frequency response can always be written as a ratio of series expression of powers of cosines

$$|H(\omega)|^2 = \frac{\beta_0 + \beta_1 \cos \omega + \beta_2 \cos^2 \omega + \cdots}{\alpha_0 + \alpha_1 \cos \omega + \alpha_2 \cos^2 \omega + \cdots}$$

• Why is this worth knowing:

Has significance for design of filters (and determining their stability).

Also, computations involving cosines are numerically stable compared to complex exponential calculations

- Another thing to keep in mind about the relationship between H(z) and $|H(\omega)|^2$
 - You can get $|H(\omega)|^2$ from H(z), but you can't get H(z) from $|H(\omega)|^2$
 - You loose the phase information.

- Analysis of frequency response in terms of zero & pole locations
 - Writing the frequency response in terms of zeros and poles

$$H(\omega) = b_0 \frac{\prod_{k=1}^{M} \left(1 - z_k e^{-j\omega}\right)}{\prod_{k=1}^{N} \left(1 - p_k e^{-j\omega}\right)} = b_0 e^{j\omega(N-M)} \frac{\prod_{k=1}^{M} \left(e^{j\omega} - z_k\right)}{\prod_{k=1}^{N} \left(e^{j\omega} - p_k\right)}$$

Each factor can be written as a magnitude and phase:

$$\begin{split} &\left(e^{j\omega}-z_{k}\right)=V_{k}\left(\omega\right)e^{j\Theta_{k}\left(\omega\right)} \quad \text{where } V_{k}\left(\omega\right)=\left|e^{j\omega}-z_{k}\right| \\ &\left(e^{j\omega}-p_{k}\right)=U_{k}\left(\omega\right)e^{j\Phi_{k}\left(\omega\right)} \quad \text{where } U_{k}\left(\omega\right)=\left|e^{j\omega}-p_{k}\right| \\ &\left|H\left(\omega\right)\right|=\left|b_{0}\right|\frac{V_{1}\left(\omega\right)V_{2}\left(\omega\right)\cdots V_{M}\left(\omega\right)}{U_{1}\left(\omega\right)U_{2}\left(\omega\right)\cdots U_{n}\left(\omega\right)}, \\ & \measuredangle H\left(\omega\right)=\left.\omega\left(N-M\right)+\left[\Theta_{1}\left(\omega\right)+\Theta_{2}\left(\omega\right)+\cdots+\Theta_{M}\left(\omega\right)\right]-\left[\Phi_{1}\left(\omega\right)+\Phi_{2}\left(\omega\right)+\cdots+\Phi_{N}\left(\omega\right)\right] \end{split}$$

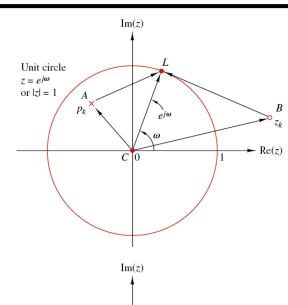
 $e^{j\omega}$ is a point on the unit circle

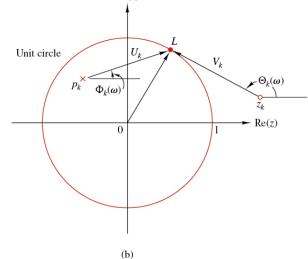
Each $(e^{j\omega} - p_k)$ and $(e^{j\omega} - z_k)$ is a vector from

the pole or zero to the $e^{j\omega}$ point on unit circle.

Magnitudes are given by U_k or V_k and angle determined from Φ_k or Θ_k

In terms of the V_k and U_k terms and associated phases:





• From the diagram you can see that the V_k and U_k correspond to distances from the zeros and poles (respectively) to the point on the unit circle associated with frequency ω

$$V_{k}(\omega) = \left| e^{j\omega} - z_{k} \right|$$

$$U_{k}(\omega) = \left| e^{j\omega} - p_{k} \right|$$

$$\left| H(\omega) \right| = \left| b_{0} \right| \frac{V_{1}(\omega)V_{2}(\omega)\cdots V_{M}(\omega)}{U_{1}(\omega)U_{2}(\omega)\cdots U_{n}(\omega)},$$

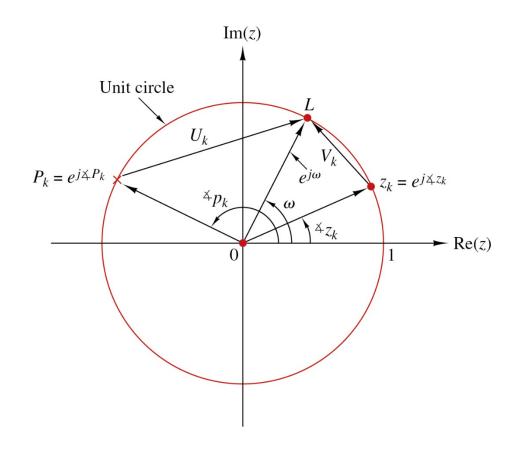
$$\angle H(\omega) = \omega(N-M) + \left[\Theta_{1}(\omega) + \Theta_{2}(\omega) + \cdots + \Theta_{M}(\omega) \right]$$

$$- \left[\Phi_{1}(\omega) + \Phi_{2}(\omega) + \cdots + \Phi_{N}(\omega) \right]$$

$$|H(\omega)| = |b_0| \frac{V_1(\omega)V_2(\omega)\cdots V_M(\omega)}{U_1(\omega)U_2(\omega)\cdots U_n(\omega)}$$

A zero close to a point at frequency ω on the unit circle causes the $|H(\omega)|$ to be small.

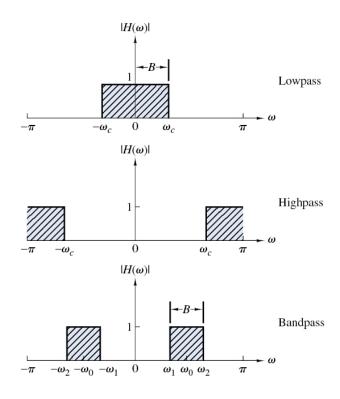
A pole close to a point at frequency ω on the unit circle causes $|H(\omega)|$ to be big.

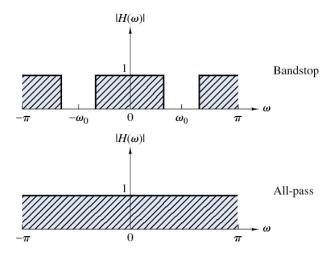


$$Y(\omega) = H(\omega)X(\omega)$$

- The system function H acts on the input X to produce an output Y
 - Its effect on the input depends on the functional dependence of H on ω
- -Any LTI system can be thought of as a filter
 - It may: reduce high frequencies in the input, reduce low frequencies, enhance high frequencies, enhance low frequencies, ... mess up all frequencies
 - Essentially, $H(\omega)$ can be designed to do almost anything you want to the frequency composition of the input to produce an output.

- -We have talked about low-pass, high-pass, band-pass, etc. filters
- -Ideal ones (unrealizable in practice) look like these:





Notice that Bandwidth (*B*) only considers the width in the positive frequency half of the diagrams.

- -Phase of filters
 - For an ideal filter:

$$Y(\omega) = H(\omega)X(\omega)$$
 $\omega_1 < \omega < \omega_2$
 $Y(\omega) = Ce^{-j\omega n_0}X(\omega)$

- It scales the magnitude of the input by C shifts the phase linearly with ω
- Linear phase filters are "good,"because they only introduce a time delay in the input signal
 - » Time shift property of Fourier transform: $y(n) = Cx(n-n_0)$
- What would be "bad" would be if phase of the input changed as a function of frequency, i.e., different frequency components would be delayed by different amounts.

-Ideal filter:
$$H(\omega) = |H(\omega)| e^{j\Theta(\omega)} = Ce^{-j\omega n_0}$$
 $\omega_1 < \omega < \omega_2$
 $|H(\omega)| = C$
 $\Theta(\omega) = -\omega n_0$

- On previous slide:
 - Delay is given by: $y(n) = Cy(n n_0)$ $\tau_g = -\frac{d\Theta(\omega)}{d\omega} = -\frac{d(-\omega n_0)}{d\omega} = n_0$
- Generalize definition of "group delay" (or "envelope delay") for arbitrary phase:

Group delay:
$$\tau_g = -\frac{d\Theta(\omega)}{d\omega}$$

- -Consider effect of phase on a sinusoidal input: $x(n) = \sin(\omega n)$
 - Filter shifts phase by $\Theta(\omega)$: $y(n) = C \sin[\omega n + \Theta(\omega)]$
 - Ideal filter:

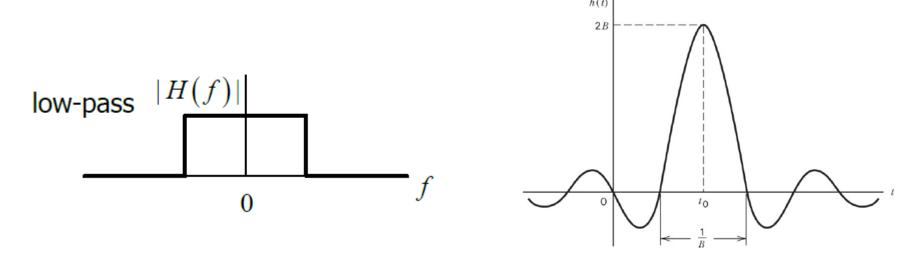
$$y(n) = C\sin\left[\omega n + \Theta(\omega)\right] = C\sin\left[\omega\left(n - n_0\right)\right] = C\sin\left[\omega\left(n - \frac{-\Theta(\omega)}{\omega}\right)\right], \text{ for } \Theta(\omega) = -\omega n_0$$

Write this in terms of a delay term, $\tau_{pd} = -\frac{\Theta(\omega)}{\omega}$

- For ideal filter: $y(n) = C \sin \left[\omega \left(n \tau_{pd} \right) \right]$
- Generalize definition of "phase delay" for filter phase:

Phase delay: $\tau_{pd} = -\frac{\Theta(\omega)}{\omega}$

- Why aren't ideal filters realizable?
 - Ideal low-pass filter Filter is a rectangular pulse in the frequency domain:



In the time domain, there is some response from the filter before t = 0, so the ideal filter is non-casual.

Remember condition for casual system is that there is no response before time = 0

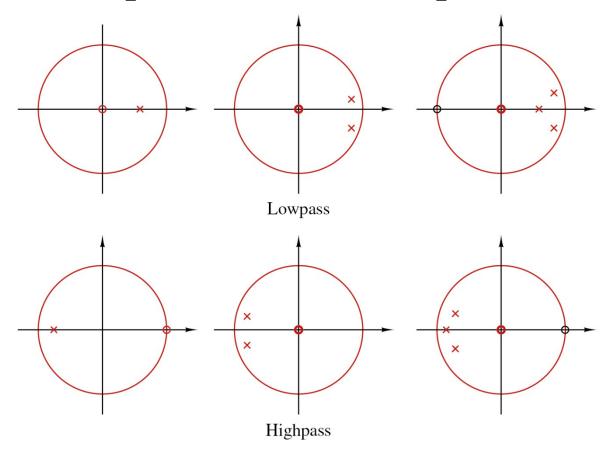
• Simple filter design by pole-zero placement

- Place poles and zeros to affect the frequency response in the desired way
 - For stability, poles must be inside unit circle
 - Zeros can be anywhere
 - Complex poles and zeros must include complex conjugate pairs
 - Usually normalize frequency response so that the gain is 1 at specified frequency
 - Do this by setting b_0 so that $|H(\omega_0)| = 1$

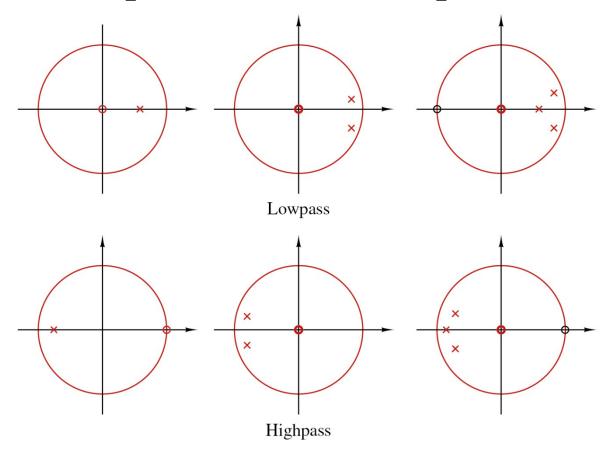
$$H(z) = \frac{\sum_{k=0}^{M} b_k z^{-k}}{1 + \sum_{k=0}^{N} a_k z^{-k}} = b_0 \frac{\prod_{k=1}^{M} (1 - z_k e^{-j\omega})}{\prod_{k=1}^{N} (1 - p_k e^{-j\omega})}$$

Usually have more poles than zeros

• Picture of pole-zero placement for lowpass and highpass filters



• Picture of pole-zero placement for lowpass and highpass filters



- Single pole lowpass filter
 - -Put pole close to unit circle at $\omega = 0$ (1 on real axis) to emphasize low frequencies
 - -Put zero near $\omega = \pi$ (highest frequency, Nyquist) to suppress high frequencies
 - -Normalize so H(0) = 1

$$H_{LP}(z) = \frac{1-a}{2} \frac{1+z^{-1}}{1-az^{-1}}$$

$$H_{LP}(0) = \frac{1-a}{2} \frac{1+e^{-j0}}{1-ae^{-j0}} = \frac{1-a}{2} \frac{2}{1-a} = 1$$

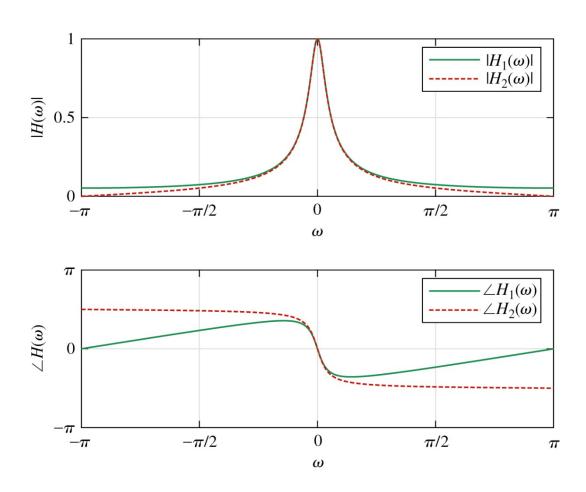
$$H_{LP}(\omega) = \frac{1-a}{2} \frac{1+e^{-j\omega}}{1-ae^{-j\omega}}$$

$$H_{LP}(\pi) = \frac{1-a}{2} \frac{1+e^{-j\pi}}{1-ae^{-j\pi}} = \frac{1-a}{2} \frac{1-1}{1+a} = 0$$

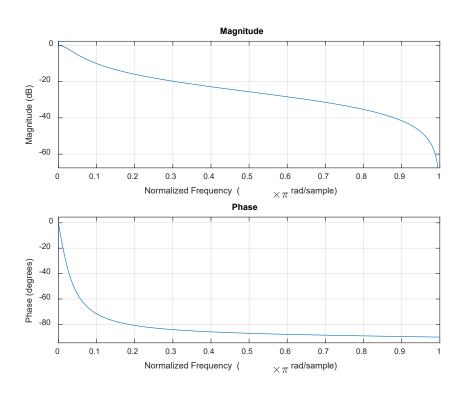
$$a = 0.9 \text{ is close to } 1$$

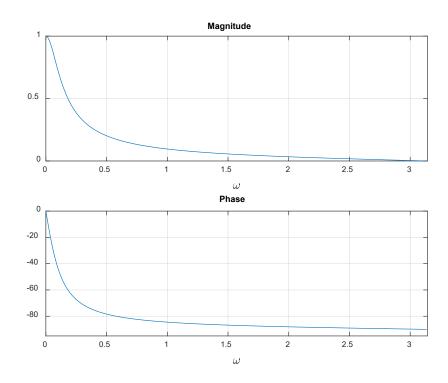
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Book's picture, red line is this filter



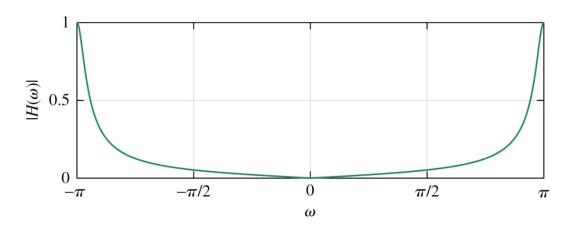
Matlab (freqz db, freqz results linear)

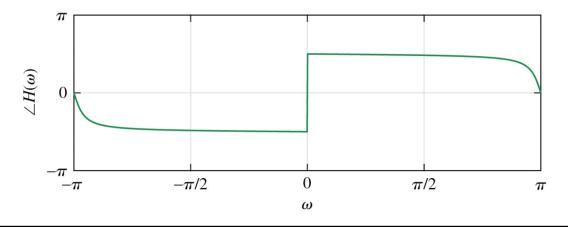




- Highpass filter by "folding" pole-zero locations:
 - -Normalize so H(0) = 1

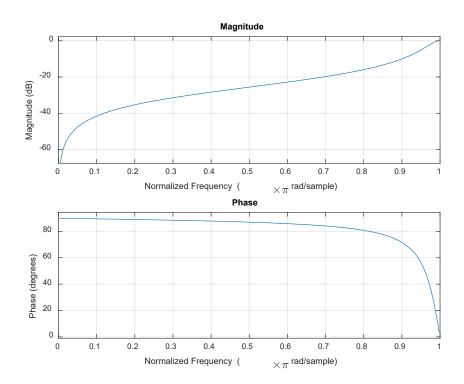
$$H_{HP}(z) = \frac{1-a}{2} \frac{1-z^{-1}}{1+az^{-1}}$$

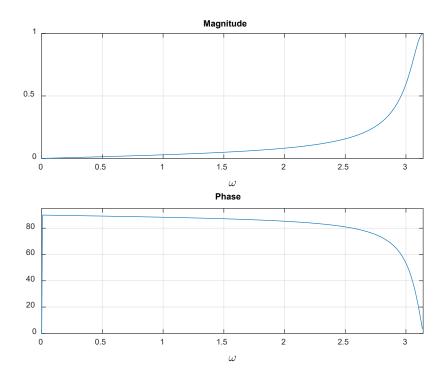




- Highpass filter by "folding" pole-zero locations:
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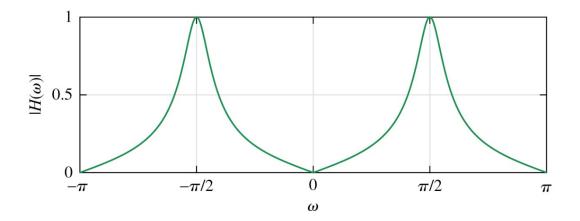
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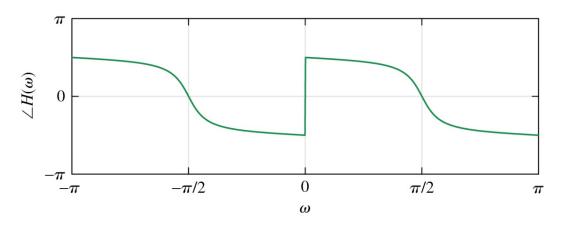




- See examples in book for
 - -Two-pole bandpass filter centered at $\omega = \pi/2$ $H(\pi/2) = 1$,

$$H(0) = H(\pi) = 0$$
, $H(4\pi/9) = 1/\sqrt{2}$





- See examples in book for
 - -Two-pole lowpass filter
 - Two-pole bandpass filter with pass band centered at
 - -p/2, zero at 0 and pi, and

centered at $\omega = \pi/2$ $H(\pi/2) = 1$, $H(0) = H(\pi) = 0$, $H(4\pi/9) = 1/\sqrt{2}$

Simple method to convert lowpass to highpass

$$H_{HP}(\omega)=H_{LP}(\omega-\pi)$$

Lowpass:

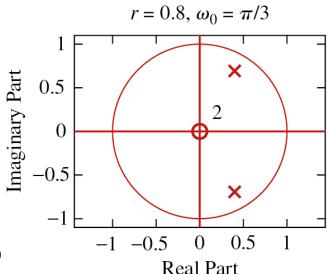
Highpass:

$$y(n) = -\sum_{k=1}^{N} a_k y(n-k) + \sum_{k=0}^{M} b_k x(n-k)$$

$$y(n) = -\sum_{k=1}^{N} (-1)^k a_k y(n-k) + \sum_{k=0}^{M} (-1)^k b_k x(n-k)$$
Frequency domain:
$$H_{LP}(\omega) = \frac{\sum_{k=0}^{M} b_k e^{-j\omega k}}{1 + \sum_{k=0}^{N} a_k e^{-j\omega k}}$$

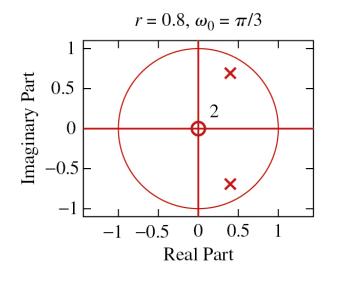
$$H_{HP}(\omega) = \frac{\sum_{k=0}^{M} (-1)^k b_k e^{-j\omega k}}{1 + \sum_{k=0}^{N} (-1)^k a_k e^{-j\omega k}}$$

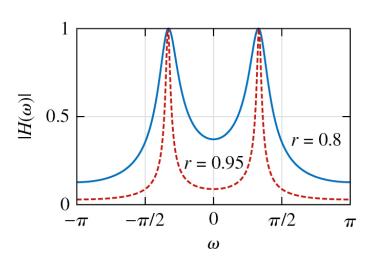
- Proakis & Manolakis give several examples of simple filters designed by pole placement:
 - -Resonators:
 - Pole placement of complex conjugate pair near unit circle
 - Enhances frequency corresponding to location of poles so "resonates" at that frequency



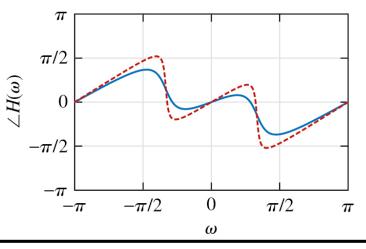
$$p_{1,2} = re^{\pm j\omega_0}$$

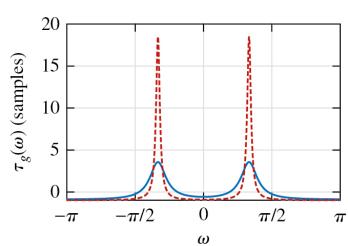
• Resonators:



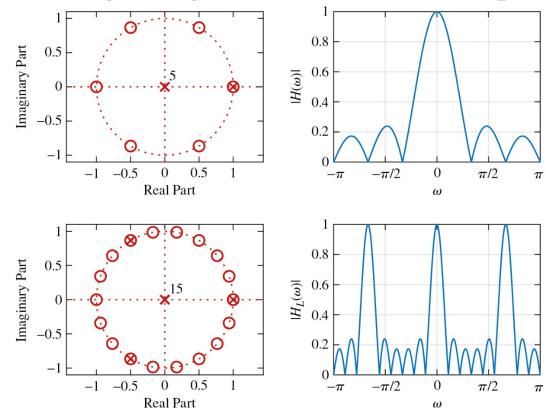


Notice group delay is negative of the derivative of the phase.





- Comb filter: Multiple, narrow passbands
 - -Good for good for getting rid of noise with repeating harmonics

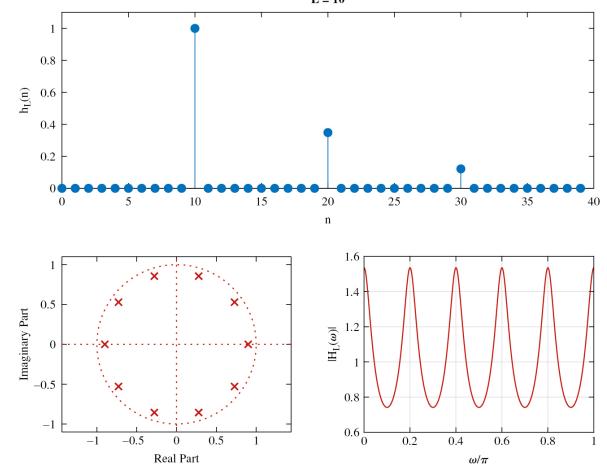


- Reverberator: Multiple, narrow passbands
 - Example: Music in a live performance bounces off walls of auditorium music in studio sounds "dry" so add some reverb

$$y(n) = ax(n-L) + a^2x(n-2L) + a^3x(n-3L) + \cdots$$

0 < a < 1 is strength of bounce

• Reverberator: Multiple narrow passbands



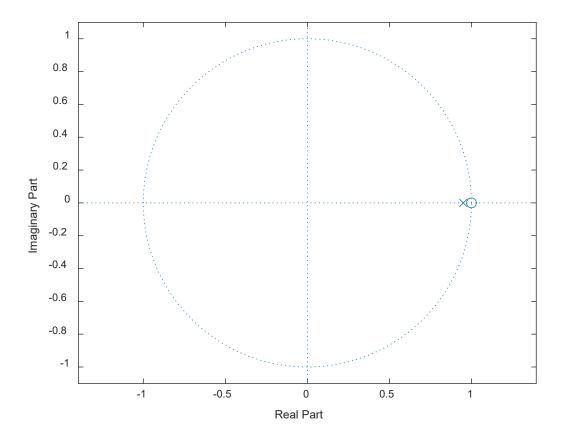
• Notch filters: (This one I'll go through in more detail) DC notch filter:

$$H(z) = \frac{z-1}{z-\lambda}$$

Zero in the numerator makes magnitude of impulse response zero at z=1 If you put a pole at λ very close to 1, it almost cancels out the zero in the numerator. You need to normalize the filter so you have unit gain.

$$H(z) = \left(\frac{1+\lambda}{2}\right) \frac{z-1}{z-\lambda}$$

Try this for different values of λ



Show demo for DC notch filter: notch_signal.m

Notch filters at other frequencies

$$H(e^{j\omega}) = \frac{\left(z - e^{j\omega_0}\right)\left(z - e^{-j\omega_0}\right)}{\left(z - \alpha e^{j\omega_0}\right)\left(z - \alpha e^{-j\omega_0}\right)} = \frac{\left(z^2 - 2\cos\omega_0 + 1\right)}{\left(z^2 - 2\alpha\cos\omega_0 + \alpha^2\right)}$$

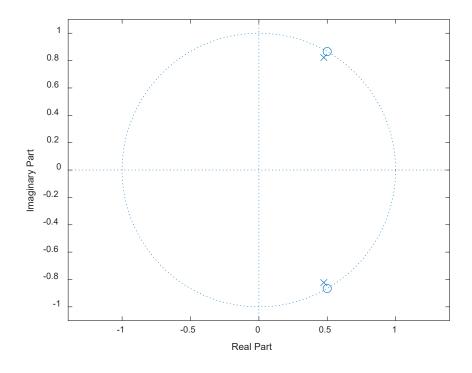
Zero in the numerator makes magnitude of impulse response zero at $z = e^{\pm j\omega_0}$

If you put a pole very close to zero, it almost cancels it out.

You need to normalize the filter so you have unit gain at DC.

$$H(z) = \left(\frac{1+\alpha^2 - 2\alpha\cos\omega_0}{2 - 2\cos\omega_0}\right) \frac{\left(z^2 - 2\cos\omega_0 + 1\right)}{\left(z^2 - 2\alpha\cos\omega_0 + \alpha^2\right)}$$

• Notch filter at 60 Hz



$$\omega_0 = 2\pi f/F_s$$

for 60 Hz with $F_s = 360$ Hz
 $\omega_0 = \pi/3$

Show demo for DC notch filter: filter_60_120Hz.m