

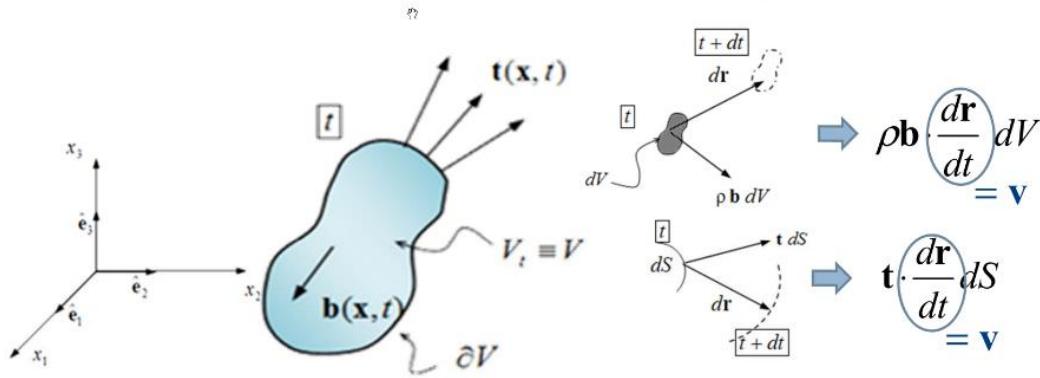
# A plastic-damage model for concrete

Online Course on Computational Solid Mechanics

J. Lubliner, J. Oliver, S. Oller, E. Oñate

Computational Solid Mechanics

$$P_e(t) = \int_V \rho \mathbf{b} \cdot \mathbf{v} dV + \int_{\partial V} \mathbf{t} \cdot \mathbf{v} dS$$



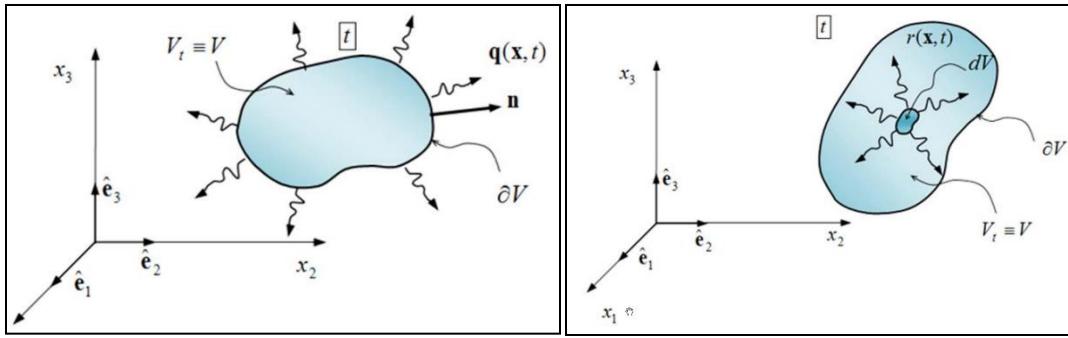
$$P_e(t) = \int_V \rho \mathbf{b} \cdot \mathbf{v} dV + \int_{\partial V} \mathbf{t} \cdot \mathbf{v} dS = \frac{d}{dt} \int_{V_t \equiv V} \frac{1}{2} \rho \mathbf{v}^2 dV + \int_V \sigma : \mathbf{d} dV$$

Kinetic energy  $\mathcal{K}$       Stress power  $\mathcal{P}_\sigma$

**external mechanical power**  
entering the medium

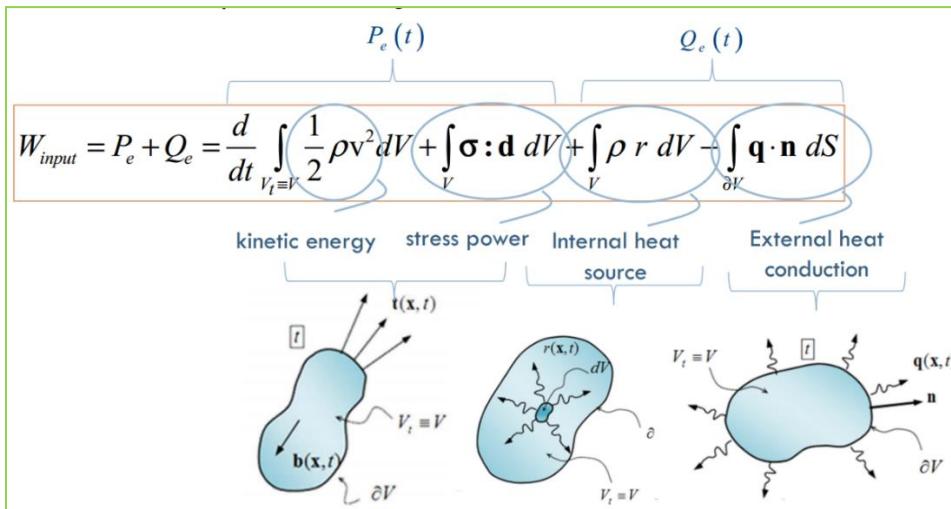
$\downarrow$

$$P_e(t) = \frac{d}{dt} \mathcal{K}(t) + \mathcal{P}_\sigma$$



$$\frac{\text{non convective (conduction) incoming heat}}{\text{unit of time}} = - \int_{\partial V} \mathbf{q} \cdot \mathbf{n} \, dS$$

$$\frac{\text{heat generated by internal sources}}{\text{unit of time}} = \int_V \rho r \, dV$$



$$P_{\text{stored}} = \frac{d}{dt} \underbrace{\int_{V_t \equiv V} \frac{1}{2} \rho v^2 dV}_{\text{Kinetic energy } \mathcal{K}} + \frac{d}{dt} \underbrace{\int_V \rho \psi \, dV}_{\text{Stored mechanical energy } \mathcal{V}} = \frac{d\mathcal{K}}{dt} + \frac{d\mathcal{V}}{dt}$$

$$\rho_0 \psi(\mathbf{x}, t) \rightarrow \begin{cases} \text{Density of free energy} & \rightarrow \frac{\text{Stored mechanical energy}}{\text{unit of volume}} \\ (\text{Helmholtz energy}) & \end{cases}$$

$$D_{\text{mech}} = \frac{P_e}{\frac{d\mathcal{K}}{dt} + \mathcal{P}_\sigma} - \frac{P_{\text{stored}}}{\frac{d\mathcal{K}}{dt} + \frac{d\mathcal{V}}{dt}} = \underbrace{\frac{\frac{d\mathcal{K}}{dt} + \int_V \boldsymbol{\sigma} : \mathbf{d} \, dV - \left( \frac{d\mathcal{K}}{dt} + \frac{d\mathcal{V}}{dt} \right)}{P_e}}_{\frac{d\mathcal{K}}{dt} + \frac{d\mathcal{V}}{dt}} = \int_V \boldsymbol{\sigma} : \mathbf{d} \, dV - \int_V \rho \dot{\psi} \, dV$$

□ Mechanical dissipation

$$P_e \rightarrow \frac{d\mathcal{K}}{dt} + \frac{d\mathcal{V}}{dt} + D_{\text{mech}}$$

$$D_{\text{mech}} = - \int_V \rho \dot{\psi} \, dV + \int_V \boldsymbol{\sigma} : \mathbf{d} \, dV$$

□ Thermal dissipation :

$$D_{\text{therm}} = - \int_V \rho s \dot{\theta} \, dV$$

$$D = D_{mech} + D_{therm} = - \int_V [\rho(\psi + s\dot{\theta}) + \sigma : \mathbf{d}] dV \geq 0 \quad \forall \Delta V \subset V$$

Global (integral) form of the second principle of thermodynamics

$$\mathcal{D}(\mathbf{x}, t) = -\rho(\psi + s\dot{\theta}) + \sigma : \mathbf{d} \geq 0 \quad \forall \mathbf{x} \quad \forall t$$

Local (differential) form of the second principle of the thermodynamics

$$u(\mathbf{x}, t) := \psi + s\theta \rightarrow \begin{cases} u(\mathbf{x}, t) \rightarrow \text{Total stored energy} \\ \psi(\mathbf{x}, t) \rightarrow \text{Mechanical stored energy} \\ s\theta(\mathbf{x}, t) \rightarrow \text{Thermal stored energy} \end{cases}$$

$$\mathcal{D} = -\rho(\dot{u} - \theta\dot{s}) + \sigma : \mathbf{d} \geq 0$$

Clausius-Planck Inequality  
in terms of the  
specific internal energy

### REMARK

For infinitesimal deformation,  $\mathbf{d} = \dot{\mathbf{\epsilon}}$ ,  
the Clausius-Planck inequality  
becomes:  $-\rho(\dot{\psi} + s\dot{\theta}) + \sigma : \dot{\mathbf{\epsilon}} \geq 0$

$$\boldsymbol{\sigma} = E \boldsymbol{\epsilon}$$

$$\sigma_{ij} = \mathbb{C}_{ijkl} \epsilon_{kl}$$

$$\mathbb{C}_{ijkl} = \lambda \mathbf{1} \otimes \mathbf{1} + 2\mu \mathbf{I}$$

$$[\mathbf{a} \otimes \mathbf{b}]_{i,j} = a_i b_j$$

$$[\mathbf{A} \otimes \mathbf{B}]_{ijkl} = A_{ij} B_{kl}$$

$$[\mathbf{1}]_{ij} = \delta_{ij}$$

$$[I]_{ijkl} = \frac{1}{2} [\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk}]$$

$$\begin{cases} \boldsymbol{\sigma} = \lambda \text{tr}(\boldsymbol{\epsilon}) + 2\mu \boldsymbol{\epsilon} \\ \sigma_{ij} = \lambda \epsilon_{kk} \delta_{ij} + 2\mu \epsilon_{ij} \quad i, j \in \{1, 2, 3\} \end{cases}$$

### Theorem

$$\mathcal{D}(x, y, \dot{x}, \dot{y}) = f(x, y)\dot{x} + g(x, y)\dot{y} \geq 0 \quad \forall \dot{x}, \dot{y} \Rightarrow \begin{cases} f(x, y) = 0 \\ g(x, y) = 0 \end{cases}$$

$$\mathcal{D} = f(\boldsymbol{\epsilon}, \boldsymbol{\epsilon}^i, \alpha) : \dot{\boldsymbol{\epsilon}} + g(\boldsymbol{\epsilon}, \boldsymbol{\epsilon}^i, \alpha) \dot{\boldsymbol{\epsilon}} \geq 0 \quad \forall \dot{\boldsymbol{\epsilon}}$$



## Variables

$$\mathbb{V} := \left\{ v_1, v_2, \dots, v_{n_v} \right\} \quad v_i(\mathbf{x}, t) \quad i \in \{1, 2, \dots, n_v\}$$

$$\mathbb{V} := \{\rho, \sigma, \epsilon, u, \psi, s, \theta, \alpha\}$$

1)  $\mathbb{F} := \{\lambda_1, \lambda_2, \dots, \lambda_{n_F}\} \quad \lambda_i(\mathbf{x}, t) \quad i \in \{1, 2, \dots, n_F\}$

$$\dot{\lambda}_i(\mathbf{x}, t) = \frac{\partial \lambda_i(\mathbf{x}, t)}{\partial t} \rightarrow \text{any}$$

$$\mathbb{F} := \{\rho, \epsilon\}$$

2)  $\mathbb{I} := \{\alpha_1, \alpha_2, \dots, \alpha_{n_I}\} \quad \alpha_i(\mathbf{x}, t) \quad i \in \{1, 2, \dots, n_I\}$

$$\dot{\alpha}_i = \frac{\partial \alpha_i(\mathbf{x}, t)}{\partial t} = \xi_i(\underbrace{\lambda(\mathbf{x}, t), \alpha(\mathbf{x}, t)}_{\substack{\text{instantaneous} \\ \text{values (at time } t\text{)}}}) \quad i \in \{1, 2, \dots, n_I\}$$

$$\mathbb{I} := \{\alpha\} \quad \dot{\alpha} = \gamma(\rho, \epsilon, \alpha)$$

3)  $\mathbb{D} := \{d_1, d_2, \dots, d_{n_D}\} \quad d_i(\mathbf{x}, t) \quad i \in \{1, 2, \dots, n_D\}$

$$d_i = \gamma_i(\lambda, \alpha) \rightarrow \dot{d}_i = \varphi_i(\lambda, \alpha, \dot{\lambda}) \quad i \in \{1, 2, \dots, n_D\}$$

$$\mathbb{D} := \{\rho, \sigma, \epsilon, u, \psi, s, \theta, \alpha\}$$

$$\psi = \psi(\rho, \epsilon, \alpha)$$

$$\psi = \gamma(\rho, \epsilon, \alpha, \dot{\rho}, \dot{\epsilon}, \dot{\alpha}(\rho, \epsilon, \alpha)) = \underbrace{\psi(\rho, \epsilon, \alpha, \dot{\rho}, \dot{\epsilon})}_{\substack{\text{not depending on} \\ \text{the internal variable} \\ \text{evolution, } \dot{\alpha}}}$$

$$\sigma = \sigma(\rho, \epsilon, \alpha)$$

$$\dot{\sigma} = \varphi(\rho, \epsilon, \alpha, \dot{\rho}, \dot{\epsilon}, \underbrace{\dot{\alpha}(\rho, \epsilon, \alpha)}_{\substack{\text{provided by the} \\ \text{evolution equation}}}) = \dot{\sigma}(\rho, \epsilon, \alpha, \dot{\rho}, \dot{\epsilon})$$

4) Postulate a specific form of the free energy:

## Elastic material

$$\mathcal{F} := \{\boldsymbol{\varepsilon}\}$$

$$\mathcal{I} := \{\emptyset\}$$

$$\mathcal{D} := \{\boldsymbol{\sigma}, \psi\}$$

$$\begin{cases} \boldsymbol{\sigma}(\boldsymbol{\varepsilon}) \rightarrow \dot{\boldsymbol{\sigma}} = \frac{\partial \boldsymbol{\sigma}(\boldsymbol{\varepsilon})}{\partial \boldsymbol{\varepsilon}} : \dot{\boldsymbol{\varepsilon}} \\ \rho_0 \psi(\boldsymbol{\varepsilon}) \rightarrow \rho_0 \dot{\psi} = \frac{\partial \rho_0 \psi(\boldsymbol{\varepsilon})}{\partial \boldsymbol{\varepsilon}} : \dot{\boldsymbol{\varepsilon}} \end{cases}$$

$$\rho_0 \psi(\boldsymbol{\varepsilon}) = \frac{1}{2} \boldsymbol{\varepsilon} : \mathbb{C} : \boldsymbol{\varepsilon}$$

Dissipation      Isothermal case

$$\mathcal{D} = -\rho_0 (\dot{\psi} + \cancel{s\dot{\theta}}) + \boldsymbol{\sigma} : \dot{\boldsymbol{\varepsilon}} \geq 0$$

$$\mathcal{D} = \underbrace{(\boldsymbol{\sigma} - \frac{\partial \rho_0 \psi(\boldsymbol{\varepsilon})}{\partial \boldsymbol{\varepsilon}})}_{f(\boldsymbol{\varepsilon})} : \dot{\boldsymbol{\varepsilon}} \geq 0 \quad \forall \dot{\boldsymbol{\varepsilon}} \Rightarrow f(\boldsymbol{\varepsilon}) = 0$$

## Inelastic material

$$\underline{\epsilon} = \underbrace{\underline{\epsilon}^e}_{\text{Elastic strain}} + \underbrace{\underline{\epsilon}^i}_{\text{Inelastic strain}}$$

$$\mathbb{F} := \{\underline{\epsilon}\}$$

$$\mathbb{I} := \{\underline{\epsilon}^i, \alpha\}$$

$$\begin{cases} \dot{\alpha} = \dot{\alpha}(\underline{\epsilon}, \dot{\underline{\epsilon}}, \alpha) \\ \dot{\underline{\epsilon}}^i = \dot{\underline{\epsilon}}^i(\underline{\epsilon}, \dot{\underline{\epsilon}}, \alpha) \end{cases}$$

$$\mathbb{D} := \{\sigma, \psi\}$$

$$\begin{aligned} \dot{\sigma} &= \frac{\partial \sigma(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \underline{\epsilon}} : \dot{\underline{\epsilon}} + \frac{\partial \sigma(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \underline{\epsilon}^i} : \dot{\underline{\epsilon}}^i + \frac{\partial \sigma(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \alpha} \dot{\alpha} \\ \dot{\psi} &= \frac{\partial \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \underline{\epsilon}} : \dot{\underline{\epsilon}} + \frac{\partial \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \underline{\epsilon}^i} : \dot{\underline{\epsilon}}^i + \frac{\partial \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \alpha} \dot{\alpha} \end{aligned}$$

$$\rho_0 \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha) = \underbrace{\frac{1}{2} \underline{\epsilon}^e : \mathbb{C} : \underline{\epsilon}^e}_{\text{Elastic potential}} + \underbrace{\mathcal{H}(\alpha)}_{\text{Hardening potential}} + \dot{\sigma} \dot{\theta} + \sigma : \dot{\underline{\epsilon}} \geq 0$$

□ Dissipation:  $\mathcal{D} = -\rho_0 (\dot{\psi} + \dot{\sigma} \dot{\theta}) + \sigma : \dot{\underline{\epsilon}} \geq 0$

$$\mathcal{D} = \underbrace{(\sigma - \frac{\partial(\rho_0 \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha))}{\partial \underline{\epsilon}} : \dot{\underline{\epsilon}})}_{f(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)} + \underbrace{\frac{\partial(\rho_0 \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha))}{\partial \underline{\epsilon}} : \dot{\underline{\epsilon}}^i}_{g_1(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)} - \underbrace{\frac{\partial(\rho_0 \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha))}{\partial \alpha} \dot{\alpha}}_{g_2(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)} \geq 0 \quad \forall \dot{\underline{\epsilon}}$$

Constitutive equation

$$\sigma = \underbrace{\frac{\partial \rho_0 \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \underline{\epsilon}}}_{\mathbb{C}} : \underline{\epsilon}^e = \mathbb{C} : (\underline{\epsilon} - \underline{\epsilon}^i)$$

$$\mathcal{D} = g(\underline{\epsilon}, \underline{\epsilon}^i, \alpha) = \underbrace{\frac{\partial \rho_0 \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \underline{\epsilon}} : \dot{\underline{\epsilon}}^i}_{\frac{\partial \rho_0 \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \underline{\epsilon}} = \sigma} - \underbrace{\frac{\partial \rho_0 \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \alpha} \dot{\alpha}}_{\mathcal{H}'(\alpha) = -q(\alpha)} \geq 0$$

Dissipation  $\mathcal{D}$

$$\mathcal{D} = \sigma : \dot{\underline{\epsilon}}^i - \underbrace{(\frac{\partial \rho_0 \psi(\underline{\epsilon}, \underline{\epsilon}^i, \alpha)}{\partial \alpha}) : \dot{\alpha}}_{\mathcal{H}'(\alpha) = -q(\alpha)} = \sigma : \dot{\underline{\epsilon}}^i + \underbrace{q(\alpha) \dot{\alpha}}_{\text{Hardening variable}} \geq 0$$



## Stress driven models

$$\mathcal{F} := \{\sigma\}$$

$$\mathcal{I} := \{q\} \quad \rightarrow \dot{q} = \zeta(\sigma, q)$$

$$\mathcal{D} := \{\varepsilon, G\}$$

Legendre transform:  $\varepsilon, \psi(\varepsilon) \rightarrow \sigma, G(\sigma)$

$$\rho_0 G(\sigma) = \sigma : \varepsilon - \rho_0 \psi(\varepsilon) \quad ; \quad \sigma = \frac{\partial \rho_0 \psi(\varepsilon)}{\partial \varepsilon}$$

$$\rho_0 \dot{G} = (\dot{\sigma} : \varepsilon + \sigma : \dot{\varepsilon}) - \rho_0 \dot{\psi}$$

$$\mathcal{D} = \sigma : \dot{\varepsilon} - \rho_0 \dot{\psi} \geq 0 \quad \rightarrow \quad \mathcal{D} = \rho_0 \dot{G} - \varepsilon : \dot{\sigma} \geq 0$$

$$\begin{cases} \psi(x) \leftrightarrow G(y) \quad ; \quad G(y) = xy - \psi(x) \\ x \leftrightarrow y ; \quad y \equiv \frac{d\psi(x)}{dx} ; \quad x \equiv \frac{dG(y)}{dy} \end{cases} \rightarrow \frac{\partial \rho_0 G(\sigma)}{\partial \sigma} = \varepsilon(\sigma)$$

## Elastic material

$$\mathcal{F} := \{\sigma\}$$

$$\mathcal{I} := \{\emptyset\}$$

$$\mathcal{D} := \{\varepsilon(\sigma), G(\sigma)\}$$

$$\begin{aligned}\rho_0 G(\sigma) &= \sigma : \varepsilon(\sigma) - \rho_0 \psi(\varepsilon(\sigma)) \quad ; \quad \sigma = \frac{\partial \rho_0 \psi(\varepsilon)}{\partial \varepsilon} \\ \rightarrow \frac{\partial \rho_0 G(\sigma)}{\partial \sigma} &= \varepsilon(\sigma)\end{aligned}$$

$$\dot{\varepsilon} = \frac{\partial \varepsilon(\sigma)}{\partial \sigma} : \dot{\sigma}$$

$$\begin{aligned}\mathcal{D} &= \rho_0 \dot{G} - \varepsilon : \dot{\sigma} \geq 0 \\ \mathcal{D} &= \underbrace{\left( \frac{\partial \rho_0 G(\sigma)}{\partial \sigma} - \varepsilon \right)}_{= \mathcal{E}} : \dot{\sigma} = 0 \quad \forall \dot{\sigma}\end{aligned}$$

Constitutive equation

$$\varepsilon = \overbrace{\frac{\partial \rho_0 G(\sigma)}{\partial \sigma}}$$

Dissipation

$$\mathcal{D} = 0$$

$$\begin{aligned}\text{Gibbs potential} \\ G(\sigma) = \frac{1}{2} \sigma : \mathbb{C}^{-1} : \sigma\end{aligned}$$

### Variable set definition

- Free variables:  $\mathcal{F} := \{\sigma\}$
- Internal variables:  $\mathcal{I} := \{q\} \rightarrow \dot{q} = \zeta(\dot{\sigma}, q)$
- Dependent:  $\mathcal{D} := \{\varepsilon(\sigma, q), G(\sigma, q)\} \rightarrow \dot{G} = \frac{\partial G(\sigma, q)}{\partial \sigma} : \dot{\sigma} + \frac{\partial G(\sigma, q)}{\partial q} \dot{q}$

### Potential

- Gibbs energy:  $\rho_0 G = \frac{1}{2} \sigma : \mathbb{C}^{-1} : \sigma + \eta(q)$

### Dissipation

$$\mathcal{D} = \rho_0 \dot{G} - \varepsilon : \dot{\sigma} \geq 0 \rightarrow \mathcal{D} = \underbrace{\left( \frac{\partial \rho_0 G(\sigma, q)}{\partial \sigma} - \varepsilon \right)}_{f(\sigma, q)} : \dot{\sigma} + \underbrace{\frac{\partial G(\sigma, q)}{\partial q} \dot{q}}_{g(\sigma, q)} \geq 0 \quad \forall \dot{\sigma}$$

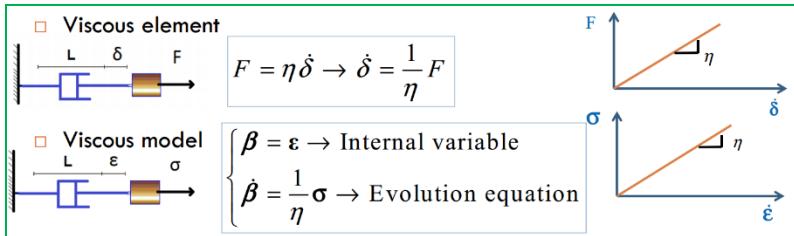
$$\begin{aligned}f(\sigma, q) &= \frac{\partial \rho_0 G(\sigma, q)}{\partial \sigma} - \varepsilon = 0 \\ g(\sigma, q) &= \frac{\partial \rho_0 G(\sigma, q)}{\partial q} \dot{q} \geq 0\end{aligned}$$

$$\begin{aligned}\varepsilon &= \frac{\partial \rho_0 G(\sigma, q)}{\partial \sigma} \\ \mathcal{D} &= \frac{\partial \rho_0 G(\sigma, q)}{\partial q} \dot{q} \geq 0\end{aligned}$$

$$\begin{aligned}\text{Constitutive equation} \\ \varepsilon = \overbrace{\frac{\partial \rho_0 G(\sigma, q)}{\partial \sigma}} = \mathbb{C}^{-1} : \sigma\end{aligned}$$

$$\begin{aligned}\text{Dissipation} \\ \mathcal{D} &= \overbrace{\frac{\partial \rho_0 G(\sigma, q)}{\partial q} \dot{q}}^{\alpha(q)} = \eta'(q) \dot{q} \geq 0 \\ \alpha(q) &= \eta'(q) = \frac{d\eta(q)}{dq} \rightarrow \begin{cases} \text{conjugate} \\ \text{internal variable} \end{cases}\end{aligned}$$





□ Stress driven model:  $\mathcal{F} := \{\sigma\}$     $\mathcal{I} := \{\beta\}$     $\mathcal{D} := \{\varepsilon, G\}$

$$\begin{cases} \rho_0 G(\sigma, \beta) = \sigma : \beta \\ \varepsilon = \frac{\partial \rho_0 G(\sigma, \beta)}{\partial \sigma} = \beta \end{cases} \Rightarrow \mathcal{D} = \underbrace{\frac{\partial \rho_0 G(\sigma, \beta)}{\partial \beta}}_{=\sigma} : \dot{\beta} = \sigma : \dot{\beta} = \frac{1}{\eta} \frac{\sigma : \sigma}{\eta} \geq 0 \Rightarrow \eta \geq 0$$

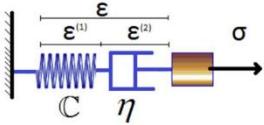
Free and internal variables  
must be different !!!!

□ Strain driven model  $\mathcal{F} := \{\varepsilon\}$     $\mathcal{I} := \underbrace{\{\beta\}}_{=\varepsilon} = \mathcal{F}$     $\mathcal{D} := \{\sigma, \psi\}$

$$\begin{cases} \rho_0 \psi = \rho_0 G - \sigma : \varepsilon = \sigma : \dot{\beta} - \sigma : \varepsilon = 0 \rightarrow \text{Free energy does not exist} \\ \sigma = \frac{\partial \rho_0 \psi(\varepsilon, \beta)}{\partial \varepsilon} \end{cases}$$

 A strain driven version of the model has no sense

□ Viscous Model



Let's define

$$\begin{cases} \boldsymbol{\varepsilon} = \boldsymbol{\varepsilon}^e + \boldsymbol{\varepsilon}^i \\ \boldsymbol{\varepsilon}^e = \boldsymbol{\beta} \\ \boldsymbol{\varepsilon}^i = \boldsymbol{\beta} \end{cases} \quad \text{being} \quad \begin{cases} \boldsymbol{\varepsilon}^e = \mathbb{C}^{-1} : \boldsymbol{\sigma} \\ \dot{\boldsymbol{\varepsilon}}^i = \frac{1}{\eta} \boldsymbol{\sigma} \end{cases}$$

$$\mathbb{F} := \{\boldsymbol{\varepsilon}\}, \quad \mathbb{I} := \{\boldsymbol{\beta}\} \quad \text{and} \quad \mathbb{D} := \{\boldsymbol{\sigma}, \psi\} \quad \dot{\boldsymbol{\beta}} = \frac{1}{\eta} \boldsymbol{\sigma}$$

$$\rho_0 \psi(\boldsymbol{\varepsilon}) = \rho_0 \psi^{(1)}(\boldsymbol{\varepsilon}) + \rho_0 \psi^{(2)}(\boldsymbol{\varepsilon})$$

$$\text{Elastic} \quad \rho_0 \psi^{(1)}(\boldsymbol{\varepsilon}, \boldsymbol{\beta}) = \frac{1}{2} \boldsymbol{\varepsilon}^{(1)} : \mathbb{C} : \boldsymbol{\varepsilon}^{(1)} = \frac{1}{2} (\boldsymbol{\varepsilon} - \boldsymbol{\beta}) : \mathbb{C} : (\boldsymbol{\varepsilon} - \boldsymbol{\beta})$$

$$\text{Viscous} \quad \rho_0 \psi^{(2)}(\boldsymbol{\varepsilon}) = 0$$

$$\rho_0 \hat{\psi}(\boldsymbol{\varepsilon}) = \rho_0 \psi^{(1)}(\boldsymbol{\varepsilon}, \boldsymbol{\beta}) + \rho_0 \psi^{(2)}(\boldsymbol{\varepsilon}) = \frac{1}{2} (\boldsymbol{\varepsilon} - \boldsymbol{\beta}) : \mathbb{C} : (\boldsymbol{\varepsilon} - \boldsymbol{\beta})$$

Constitutive equation  
 $\dot{\boldsymbol{\varepsilon}} = \mathbb{C}^{-1} : \dot{\boldsymbol{\sigma}} + \frac{1}{\eta} \boldsymbol{\sigma}$

$$\mathcal{D} = \boldsymbol{\sigma} : \dot{\boldsymbol{\beta}} = \eta \underbrace{\dot{\boldsymbol{\beta}} : \dot{\boldsymbol{\beta}}}_{\geq 0} \geq 0 \quad \Rightarrow \quad \eta \geq 0$$

$$\mathbb{F} := \{\boldsymbol{\sigma}\}, \quad \mathbb{I} := \{\boldsymbol{\beta}\}, \quad \mathbb{D} := \{\boldsymbol{\varepsilon}, G\}, \quad \dot{\boldsymbol{\beta}} = \frac{1}{\eta} \boldsymbol{\sigma}$$

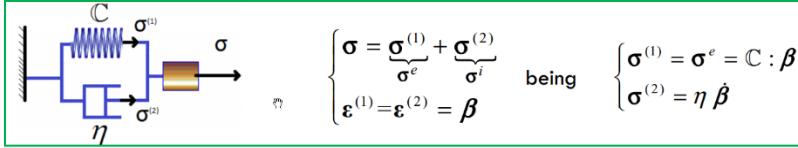
$$\text{Elastic} \quad \rho_0 G^{(1)}(\boldsymbol{\sigma}, \boldsymbol{\beta}) = \frac{1}{2} \boldsymbol{\sigma} : \mathbb{C}^{-1} : \boldsymbol{\sigma}$$

$$\text{Viscous} \quad \rho_0 G^{(2)}(\boldsymbol{\sigma}, \boldsymbol{\beta}) = \boldsymbol{\sigma} : \boldsymbol{\beta}$$

$$\Rightarrow \quad \rho_0 G(\boldsymbol{\sigma}, \boldsymbol{\beta}) = \frac{1}{2} \boldsymbol{\sigma} : \mathbb{C}^{-1} : \boldsymbol{\sigma} + \boldsymbol{\sigma} : \boldsymbol{\beta}$$

Constitutive equation  
 $\dot{\boldsymbol{\varepsilon}} = \mathbb{C}^{-1} : \dot{\boldsymbol{\sigma}} + \frac{1}{\eta} \boldsymbol{\sigma}$

$$\eta \geq 0$$



□ Stress driven model

□ Variable set definition

$$\mathbb{F} := \{\sigma\} \quad \text{and} \quad \mathbb{I} := \{\beta\} \quad \dot{\beta} = \frac{1}{\eta} \sigma^{(2)} = \frac{1}{\eta} (\sigma - \sigma^{(1)}) = \frac{1}{\eta} (\sigma - \mathbb{C} : \beta)$$

□ Potential  $\rho_0 G(\sigma, \beta) = \rho_0 G^{(1)}(\sigma, \beta) + \rho_0 G^{(2)}(\sigma, \beta)$

$$= \frac{1}{2} \sigma^{(1)} : \mathbb{C}^{-1} : \sigma^{(1)} = \frac{1}{2} \underbrace{(\beta : \mathbb{C})}_{\sigma^{(1)}} : \mathbb{C}^{-1} : \underbrace{(\mathbb{C} : \beta)}_{\sigma^{(1)}} = \frac{1}{2} \beta : \mathbb{C} : \beta$$

Viscous  $\rho_0 G^{(2)}(\sigma, \beta) = \sigma^{(2)} : \beta = (\sigma - \sigma^{(1)}) : \beta = \sigma : \beta - \underbrace{\beta : \mathbb{C} : \beta}_{\sigma^{(1)}}$

$$\rho_0 G(\sigma, \beta) = \rho_0 G^{(1)} + \rho_0 G^{(2)} = \sigma : \beta - \frac{1}{2} \beta : \mathbb{C} : \beta$$

Constitutive equation

$$\begin{cases} \dot{\beta} = \frac{1}{\eta} \sigma^{(2)} = \frac{1}{\eta} (\sigma - \underbrace{\mathbb{C} : \beta}_{\sigma^{(1)}}) \rightarrow \beta(\sigma) \\ \varepsilon = \beta(\sigma) \end{cases}$$

$\eta \geq 0$

□ Strain driven model

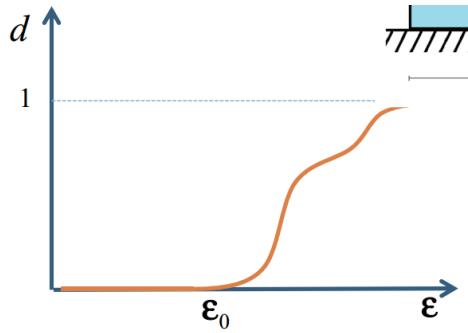
□ Variables set definition

$$\mathbb{F} := \{\varepsilon\}, \quad \mathbb{I} := \{\beta\} \quad \mathcal{D} := \{\sigma, \psi\} \quad \dot{\beta} = \frac{1}{\eta} \sigma^{(2)}$$

NOT PHYSICAL

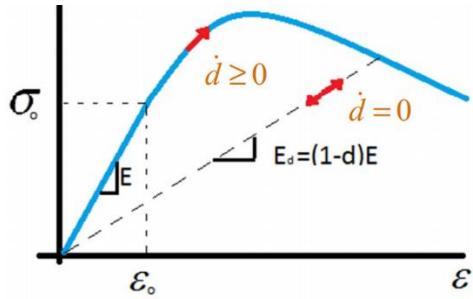
Free and internal variables  
must be different !!!!

## Damage



- Damage or (degradation) is initiated when the strain (or stress) exceeds the initial damage threshold ( $\varepsilon_0, \sigma_0$ )

$$d = 0 \quad \text{if} \quad \varepsilon \leq \varepsilon_0 \quad \text{or} \quad \sigma \leq \sigma_0$$



$$\begin{cases} \sigma = \underbrace{(1-d)\mathbb{C}}_{\mathbb{C}^d} : \varepsilon = \mathbb{C}^d : \varepsilon = (1-d)\bar{\sigma} \\ \bar{\sigma} = \mathbb{C} : \varepsilon \rightarrow \text{effective stress} \end{cases}$$

$$\mathbb{C}^d = (1-d)\mathbb{C} = (1-d)[\lambda \mathbf{1} \otimes \mathbf{1} + 2\mu \mathbb{I}]$$

$$\mathcal{F} := \{\varepsilon\}$$

$$\mathcal{I} := \{r\} \quad \dot{r} = \lambda(\varepsilon, r)$$

$$\mathcal{D} := \{\sigma(\varepsilon, r), \psi(\varepsilon, r), d(\varepsilon, r) \dots\}$$

$$\psi(\varepsilon, d(r)) = (1 - d(r))\psi_0(\varepsilon) = (1 - d(r)) \frac{1}{2}(\varepsilon : \mathbb{C} : \varepsilon)$$

$$\dot{\psi} = \frac{\partial \psi(\boldsymbol{\varepsilon}, d(r))}{\partial \boldsymbol{\varepsilon}} : \dot{\boldsymbol{\varepsilon}} + \frac{\partial \psi(\boldsymbol{\varepsilon}, d(r))}{\partial d} \dot{d}$$

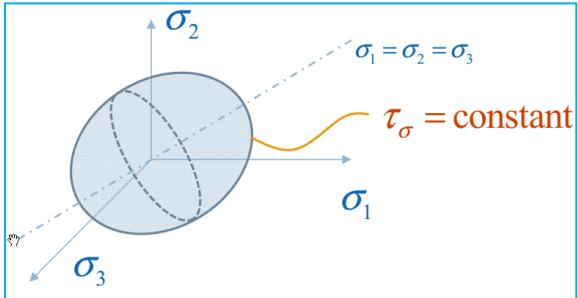
$$\mathcal{D} = (\boldsymbol{\sigma} - \frac{\partial \psi(\boldsymbol{\varepsilon}, d(r))}{\partial \boldsymbol{\varepsilon}}) : \dot{\boldsymbol{\varepsilon}} - \frac{\partial \psi(\boldsymbol{\varepsilon}, \dot{d}(r))}{\partial d} \dot{d}(r, \dot{r}) \geq 0 \quad \forall \dot{\boldsymbol{\varepsilon}}$$

$$\psi(\boldsymbol{\varepsilon}, d(r)) = (1 - d(r))\psi_0(\boldsymbol{\varepsilon}) = (1 - d(r)) \underbrace{\frac{1}{2}(\boldsymbol{\varepsilon} : \mathbb{C} : \boldsymbol{\varepsilon})}_{\psi_0(\boldsymbol{\varepsilon}) \geq 0}$$

$$\begin{aligned}\boldsymbol{\sigma} &= \frac{\partial \psi(\boldsymbol{\varepsilon}, d(r))}{\partial \boldsymbol{\varepsilon}} = (1 - d(r)) \frac{\partial \psi_0(\boldsymbol{\varepsilon}, d(r))}{\partial \boldsymbol{\varepsilon}} = (1 - d(r)) \mathbb{C} : \boldsymbol{\varepsilon} \\ \mathcal{D} &= - \frac{\partial \psi(\boldsymbol{\varepsilon}, d(r))}{\partial d} \dot{d} = \underbrace{\psi_0}_{\geq 0} \dot{d} \geq 0 \quad \xrightarrow{\hspace{1cm}} \quad \boxed{\dot{d} \geq 0}\end{aligned}$$

$$\tau_\sigma = \sqrt{\sigma : M : \sigma} = \|\sigma\|_M$$

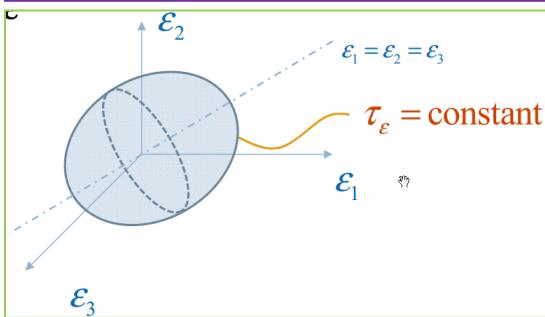
$$M = \alpha \mathbf{1} \otimes \mathbf{1} + \beta \mathbb{I} \text{ with } \begin{cases} \alpha + \beta \geq 0 \\ \beta \geq 0 \end{cases}$$



□ STRAIN NORM

$$\tau_\varepsilon = \|\varepsilon\|_M = \sqrt{\varepsilon : M : \varepsilon} = (\varepsilon : M : \underbrace{\mathbb{C}_\varepsilon}_{\mathbb{C}^*})^{\frac{1}{2}} = (\varepsilon : \underbrace{M : \mathbb{C}}_{\mathbb{C}^*} : \varepsilon)^{\frac{1}{2}} = \sqrt{\varepsilon : C^* : \varepsilon}$$

$$\tau_\varepsilon = \|\varepsilon\|_{C^*} = \sqrt{\varepsilon : C^* : \varepsilon}$$



$$\tau_\sigma = (1-d)\tau_\varepsilon$$

□ Damage function (in stress space)

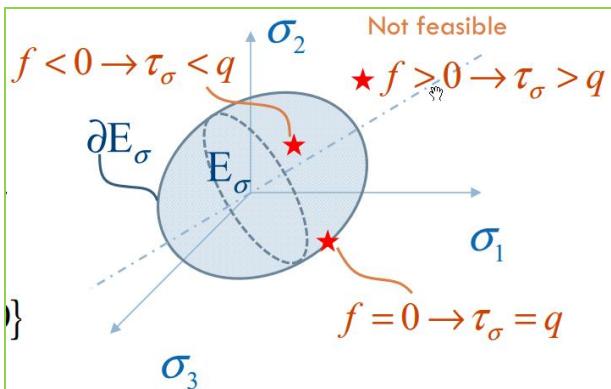
$$f(\sigma, r) = \tau_\sigma - q(r)$$

Elastic domain (in stress space)

$$E_\sigma := \{\sigma \in \mathbb{S} \mid f(\sigma, r) = \tau_\sigma - q(r) < 0\}$$

Damage surface (in stress space)

$$\partial E_\sigma := \{\sigma \in \mathbb{S} \mid f(\sigma, r) = \tau_\sigma - q(r) = 0\}$$



Damage function (in strain space)

$$g(\varepsilon, r) = \tau_\varepsilon - r$$

Elastic domain (in strain space)

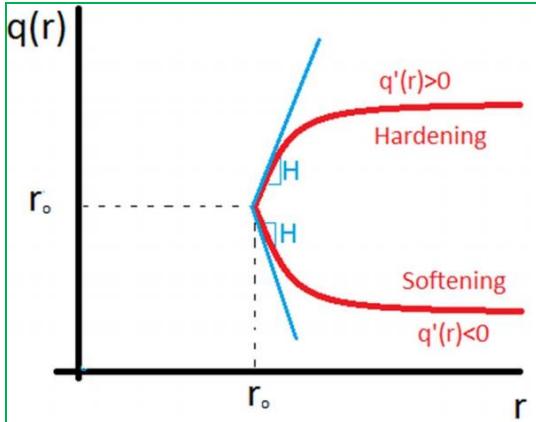
$$E_\varepsilon := \{\varepsilon \in \mathbb{S} \mid g(\varepsilon, r) = \tau_\varepsilon - r < 0\}$$

Damage surface (in strain space)

$$\partial E_\varepsilon := \{\varepsilon \in \mathbb{S} \mid g(\varepsilon, r) = \tau_\varepsilon - r = 0\}$$

**THEOREM :** The elastic domains in stress and strain spaces are equivalent, i.e.:

$$\text{If } \sigma \in E_\sigma \Leftrightarrow \varepsilon \in E_\varepsilon$$



$$\dot{r} = \lambda(\varepsilon, r) \geq 0, \quad r|_{t=0} = r_0 > 0, \quad r \in [r_0, \infty]$$

### Damage variable

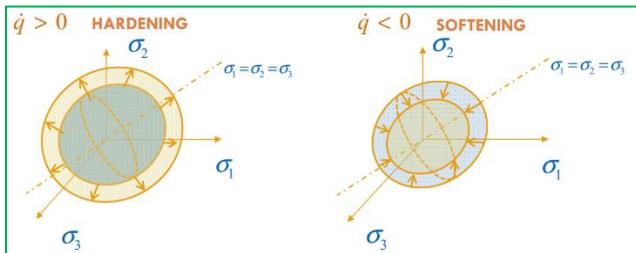
$$d := 1 - \frac{q(r)}{r} \quad 0 \leq d \leq 1$$

$$q(r_0) = r_0 \quad \rightarrow \quad d(r_0) = 0$$

□ Hardening parameter

$$H := \frac{dq(r)}{dr} = q'(r)$$

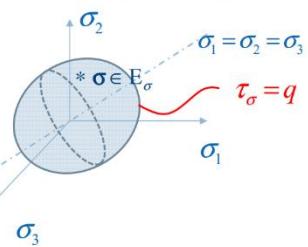
$H > 0$	Hardening
$H < 0$	Softening



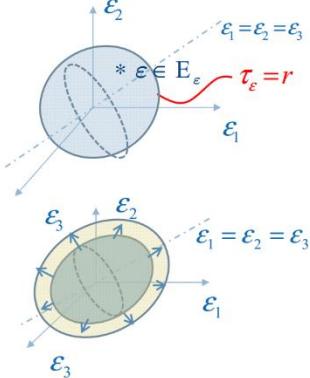
$$H \leq \frac{q(r)}{r}$$

□ Elastic domain (in stress and strain spaces)

Stress space  $\rightarrow E_\sigma$



Strain space  $\rightarrow E_\varepsilon$



Remark

$\dot{r} \geq 0 \Rightarrow E_\varepsilon$  always grows, regardless hardening or softening take place

## Constitutive equation-additional elements: Karush/Kuhn/Tucker conditions

### MATHEMATICAL EXPRESSIONS OF THE K-K-T CONDITIONS

$$\begin{aligned} \lambda \geq 0 ; g \leq 0 ; \lambda g = 0 &\rightarrow \text{Loading/unloading conditions} \\ \text{for } g = 0 \quad \lambda \dot{g} = 0 &\rightarrow \text{Persistency/consistency conditions} \end{aligned}$$

$$\psi(\boldsymbol{\varepsilon}, d(r)) = (1 - d(r))\psi_0(\boldsymbol{\varepsilon}) = (1 - d(r))\frac{1}{2}(\boldsymbol{\varepsilon} : \mathbb{C} : \boldsymbol{\varepsilon})$$

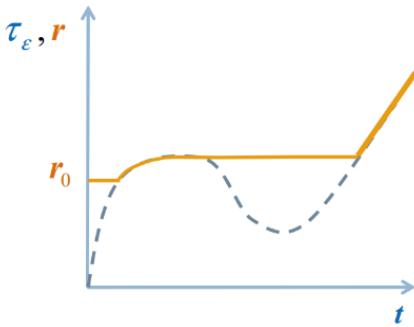
$$d := 1 - \frac{q(r)}{r} \quad 0 \leq d \leq 1 \quad q(r_0) = r_0 \quad \text{for } d = 0$$

$$\sigma = \frac{\partial \psi(\boldsymbol{\varepsilon}, d(r))}{\partial \boldsymbol{\varepsilon}} = (1 - d(r)) \frac{\partial \psi_0(\boldsymbol{\varepsilon}, d(r))}{\partial \boldsymbol{\varepsilon}} = (1 - d(r)) \mathbb{C} : \boldsymbol{\varepsilon}$$

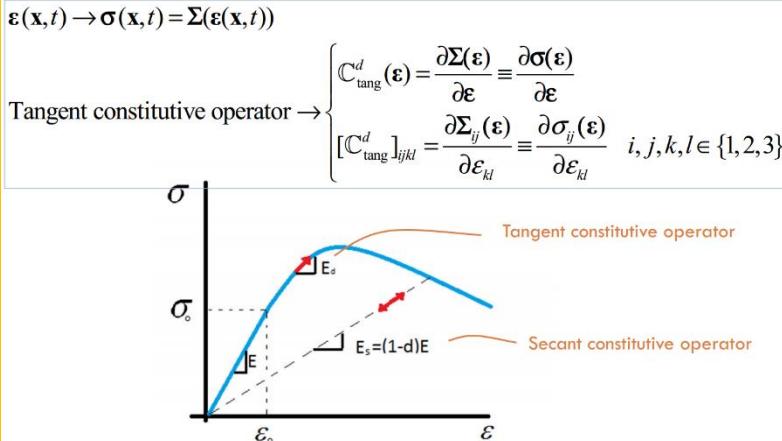
$$r(t) = \max_{\substack{\text{historical maximum of} \\ [r_0, \tau_\varepsilon(s)]}} (r_0, \tau_\varepsilon(s)) \quad s \in [0, t]$$

$$\begin{aligned} \boldsymbol{\varepsilon}_s \rightarrow \tau_\varepsilon(s) \rightarrow r_t \rightarrow q_t = q(r_t) \rightarrow \\ d_t = 1 - \frac{q_t}{r_t} \rightarrow \sigma_t = (1 - d_t) \mathbb{C} : \boldsymbol{\varepsilon}_t \end{aligned}$$

The integration is exact (closed form).  
No dependence on  $\Delta t$



SUMMARY :	<ol style="list-style-type: none"> <li>1) the initial value of <math>r</math> is <math>r_0 &gt; 0</math></li> <li>2) the initial value of <math>\tau_\varepsilon</math> is <math>\tau_\varepsilon(0) = 0</math></li> <li>3) <math>r</math> never decreases (<math>\dot{r} \geq 0</math>)</li> <li>4) <math>\tau_\varepsilon</math> is always smaller or equal than <math>r</math> (<math>g \equiv \tau_\varepsilon - r \leq 0</math>)</li> <li>5) When <math>\tau_\varepsilon</math> equals <math>r</math> then           <div style="display: inline-block; vertical-align: middle;"> <math>\begin{cases} \text{if } \tau_\varepsilon \text{ grows} \rightarrow \dot{\tau}_\varepsilon = \dot{r} \geq 0 \\ \text{if } \tau_\varepsilon \text{ decreases} \rightarrow \dot{\tau}_\varepsilon &lt; 0 ; \dot{r} = 0 \end{cases}</math> </div> </li> </ol>
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$$\dot{\boldsymbol{\sigma}} = (1-d)\mathbb{C} : \dot{\boldsymbol{\varepsilon}} - d\mathbb{C} : \boldsymbol{\varepsilon} = \underbrace{\mathbb{C}_{\text{tang}}^d}_{\mathbb{C}_{\text{tang}}} : \dot{\boldsymbol{\varepsilon}}$$

$$\dot{r} = 0 \rightarrow \dot{d} = 0 \rightarrow \dot{\boldsymbol{\sigma}} = \underbrace{(1-d)\mathbb{C}}_{\mathbb{C}_{\text{tang}}^d} : \dot{\boldsymbol{\varepsilon}}$$

$$\mathbb{C}_{\text{tang}}^d = (1-d)\mathbb{C} = \mathbb{C}_{\text{sec}}^d \quad (\text{secant constitutive operator})$$

$$\dot{d}(t) = \frac{q - Hr}{r^3} \dot{r} \geq 0$$

$$\dot{r} = \frac{1}{r} \bar{\boldsymbol{\sigma}} : \mathbb{A} : \dot{\boldsymbol{\varepsilon}}$$

$$\dot{d} = \frac{q - Hr}{r^3} \bar{\boldsymbol{\sigma}} : \mathbb{A} : \dot{\boldsymbol{\varepsilon}}$$

$$\dot{\boldsymbol{\sigma}} = (1-d)\mathbb{C} : \dot{\boldsymbol{\varepsilon}} - \dot{d} \underbrace{\mathbb{C} : \boldsymbol{\varepsilon}}_{\bar{\boldsymbol{\sigma}}} = (1-d)\mathbb{C} : \dot{\boldsymbol{\varepsilon}} - \dot{d} \underbrace{\bar{\boldsymbol{\sigma}}}_{\bar{\boldsymbol{\sigma}} \otimes \dot{d}}$$

$$\mathbb{C}_{\text{tang}}^d(\boldsymbol{\varepsilon}) = (1-d)\mathbb{C} - \frac{q - Hr}{r^3} (\bar{\boldsymbol{\sigma}} \otimes [\bar{\boldsymbol{\sigma}} : \mathbb{A}])$$

Symmetric model  
(tension/compression)  $\rightarrow \begin{cases} \mathbb{M} = \mathbb{C}^{-1} \\ \mathbb{A} = \mathbb{M} : \mathbb{C} \end{cases} \rightarrow \begin{cases} \mathbb{A} = \mathbb{C}^{-1} : \mathbb{C} = \mathbb{I} \\ \bar{\boldsymbol{\sigma}} : \mathbb{A} = \bar{\boldsymbol{\sigma}} : \mathbb{I} = \bar{\boldsymbol{\sigma}} \end{cases}$

$$\mathbb{C}_{\text{tang}}^d(\boldsymbol{\varepsilon}) = \begin{cases} (1-d)\mathbb{C} \rightarrow (\text{elastic/unloading}) \\ (1-d)\mathbb{C} - \frac{q - Hr}{r^3} (\bar{\boldsymbol{\sigma}} \otimes \bar{\boldsymbol{\sigma}}) \rightarrow (\text{loading}) \end{cases}$$

Step 1 → Compute

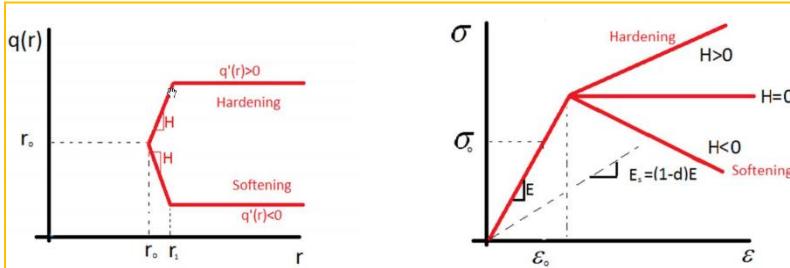
$$\begin{cases} \bar{\sigma}_{t+\Delta t} = \mathbb{C} : \boldsymbol{\varepsilon}_{t+\Delta t} \\ \tau_{\varepsilon_{t+\Delta t}} = \sqrt{\boldsymbol{\varepsilon}_{t+\Delta t} : \mathbb{C} : \boldsymbol{\varepsilon}_{t+\Delta t}} = \sqrt{\boldsymbol{\varepsilon}_{t+\Delta t} : \bar{\sigma}_{t+\Delta t}} \end{cases}$$

Step 2 → If  $\tau_{\varepsilon_{t+\Delta t}} \leq r_t \rightarrow$

$r_{t+\Delta t} = r_t$ $d_{t+\Delta t} = d_t = 1 - \frac{q(r_{t+\Delta t})}{r_{t+\Delta t}}$ $\sigma_{t+\Delta t} = (1 - d_{t+\Delta t})\bar{\sigma}_{t+\Delta t}$ $(\mathbb{C}_{\text{tang}}^d)_{t+\Delta t} = (1 - d_{t+\Delta t})\mathbb{C}$	<b>Elastic</b> <b>Unloading</b> <b>Neutral loading</b>
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Step 3 → If  $\tau_{\varepsilon_{t+\Delta t}} > r_t \rightarrow$  (Loading)

$r_{t+\Delta t} = \tau_{\varepsilon_{t+\Delta t}}$ $d_{t+\Delta t} = 1 - \frac{q(r_{t+\Delta t})}{r_{t+\Delta t}}$ $\sigma_{t+\Delta t} = (1 - d_{t+\Delta t})\bar{\sigma}_{t+\Delta t}$ $(\mathbb{C}_{\text{tang}}^d)_{t+\Delta t} = (1 - d_{t+\Delta t})\mathbb{C} - \frac{q(r_{t+\Delta t}) - H_{t+\Delta t}r_{t+\Delta t}}{(r_{t+\Delta t})^3}(\bar{\sigma}_{t+\Delta t} \otimes \bar{\sigma}_{t+\Delta t})$
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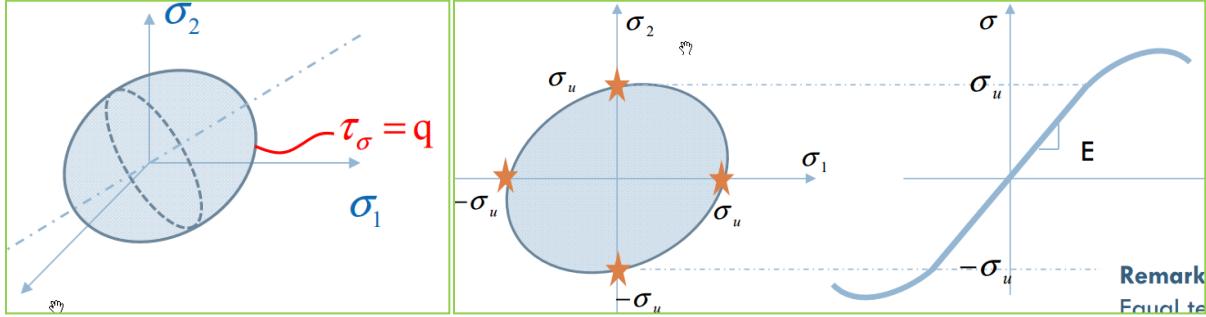
$$q(r) = q_\infty - (q_\infty - r_0)e^{-A(1-\frac{r}{r_0})} \quad (A > 0)$$

$$H(r) = \frac{dq(r)}{dr} = A \frac{(q_\infty - r_0)}{r_0} e^{-A(1-\frac{r}{r_0})}$$

### 1. Symmetric (tension/compression) model

$$\mathbb{M} = \mathbb{C}^{-1}$$

$$\begin{aligned}\tau_\sigma &= \sqrt{\sigma : \mathbb{C}^{-1} : \sigma} = (1-d)\tau_\varepsilon = (1-d)\sqrt{\varepsilon : \mathbb{C} : \varepsilon} \\ \tau_\varepsilon &= \sqrt{\varepsilon : \mathbb{C} : \varepsilon}\end{aligned}$$



$$[\sigma] \rightarrow \text{diagonalization} \rightarrow [\sigma]_{diag} = \begin{bmatrix} \sigma_1 & 0 & 0 \\ 0 & \sigma_2 & 0 \\ 0 & 0 & \sigma_3 \end{bmatrix}$$

$$[\sigma]_{diag}^+ = \begin{bmatrix} \langle \sigma_1 \rangle & 0 & 0 \\ 0 & \langle \sigma_2 \rangle & 0 \\ 0 & 0 & \langle \sigma_3 \rangle \end{bmatrix}$$

return to original system of coordinates  $\rightarrow [\sigma]^+$   
Positive counterpart of  $[\sigma]$

*eigenvector "i"*

$$\sigma = \sum_{i=1}^{i=3} \underbrace{\sigma_i}_{\text{eigenvalue "i" }} \hat{\mathbf{p}}_i \otimes \hat{\mathbf{p}}_i$$

$$\sigma^+ = \sum_{i=1}^{i=3} \langle \sigma_i \rangle \hat{\mathbf{p}}_i \otimes \hat{\mathbf{p}}_i$$

$$\sigma = \underbrace{(1-d)}_{\geq 0} \bar{\sigma} \Rightarrow \sigma^+ = (1-d) \bar{\sigma}^+$$

$\sigma$  and  $\sigma^+$  have the same eigenvectors  
 $\sigma^+$  shares the positive eigenvalues of  $\sigma$   
 $\sigma^+$  has those negative eigenvalues of  $\sigma$  null

□ Strain norm redefinition

$$\tau_\sigma^+ \equiv \sqrt{\sigma^+ : \mathbb{C}^{-1} : \sigma} = \sqrt{(1-d)^2 \bar{\sigma}^+ : \mathbb{C}^{-1} : \bar{\sigma}} = (1-d) \sqrt{\bar{\sigma}^+ : \mathbb{C}^{-1} : \bar{\sigma}} = (1-d) \tau_\varepsilon^+$$

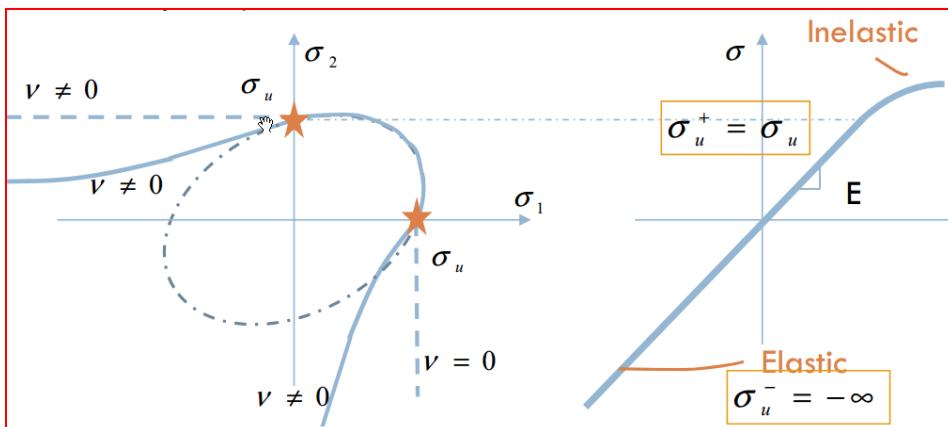
$$\tau_\varepsilon^+ = (\bar{\sigma}^+ : \mathbb{C}^{-1} : \bar{\sigma})^{\frac{1}{2}} = \sqrt{\bar{\sigma}^+ : \varepsilon}$$

Pure tensile state  $\rightarrow \sigma = \sigma^+$

$$\rightarrow \tau_\sigma^+ \equiv \sqrt{\sigma^+ : \mathbb{C}^{-1} : \sigma} = \tau_\sigma$$

Pure compressive state

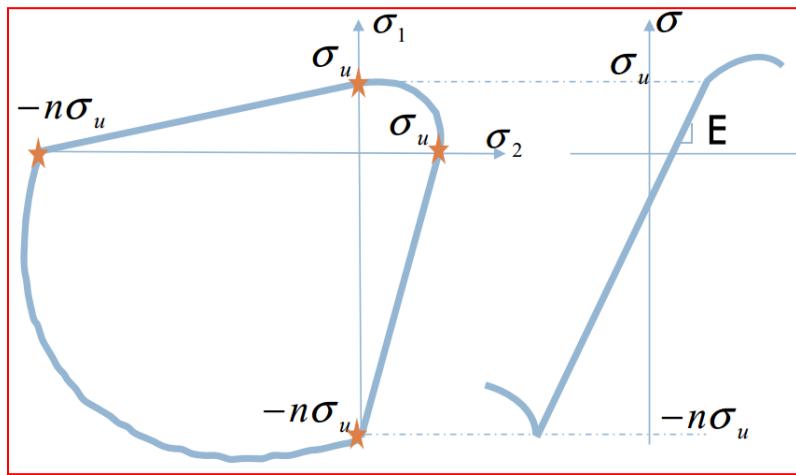
$$\sigma^+ = 0 \rightarrow \tau_\sigma^+ \equiv \sqrt{\sigma^+ : \mathbb{C}^{-1} : \sigma} = 0 \rightarrow f = \tau_\sigma^+ - \sigma_u \stackrel{=0}{\rightarrow} \text{The state is always elastic}$$



$$\tau_{\sigma} = \left[ \theta + \frac{1-\theta}{n} \right] \sqrt{\sigma : \mathbb{C}^{-1} : \sigma}$$

$$\theta = \frac{\sum_1^3 \langle \sigma_i \rangle}{\sum_1^3 |\sigma_i|} = \frac{\sum_1^3 \langle \bar{\sigma}_i \rangle}{\sum_1^3 |\bar{\sigma}_i|}$$

$$\begin{cases} \sigma_1, \sigma_2, \sigma_3 > 0 \rightarrow \theta = 1 \rightarrow \tau_{\sigma} = \sqrt{\sigma : \mathbb{C}^{-1} : \sigma} \\ \sigma_1, \sigma_2, \sigma_3 < 0 \rightarrow \theta = 0 \rightarrow \tau_{\sigma} = \frac{1}{n} \sqrt{\sigma : \mathbb{C}^{-1} : \sigma} \end{cases}$$



<sup>8</sup>Karush-Kuhn-Tucker and persistency conditions

$$\dot{r} = \lambda \geq 0, \quad g \leq 0 \quad \rightarrow \quad \lambda g \leq 0$$

$$\text{if } g = 0 \quad \rightarrow \quad \lambda \dot{g} \leq 0$$

Visco regularization (Perzyna's Regularization)

Obtained by replacing  $\lambda \rightarrow \frac{1}{\eta} \langle g \rangle$  in the inviscid model

Evolution equation

$$\dot{r} = \lambda(\boldsymbol{\varepsilon}, r) = \frac{1}{\eta} \langle g \rangle$$

where  $\eta \geq 0$  is the viscosity  
The only change is the evolution law  
 $\langle \rangle$  is the McAuley bracket

Constitutive equation

Damage function

$$g(\boldsymbol{\varepsilon}, r) \equiv \tau_{\varepsilon} - r$$

Karush – Kuhn –Tucker and persistency condition

Not necessary

□ Remark 1

$$\frac{1}{\eta} \geq 0; \quad \langle g \rangle \geq 0 \quad \rightarrow \quad \lambda = \frac{1}{\eta} \langle g \rangle \geq 0$$

$$\lambda g = \frac{1}{\eta} \underbrace{\langle g \rangle g}_{\langle g \rangle^2} = \frac{1}{\eta} \langle \frac{g}{\lambda} \rangle^2 = \frac{1}{\eta} (\lambda \eta)^2 = \eta \underbrace{\lambda^2}_{=\dot{r}} = \eta \dot{r}^2 \quad \text{As } \eta \rightarrow 0 \quad \lambda g = 0$$

□ Remark 2

$$\langle g \rangle = \lambda \eta \geq 0$$

$$\text{As } \eta \rightarrow 0 \quad \langle g \rangle = 0 \quad \rightarrow \quad g \leq 0$$

□ Remark 3

$$\lambda g = 0 \quad \dot{\lambda} g + \lambda \dot{g} = 0$$

$$\text{If } g = 0 \quad \rightarrow \quad \lambda \dot{g} = 0$$

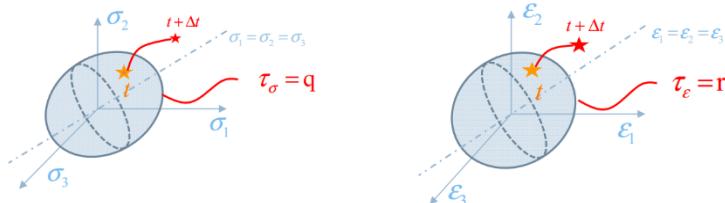
$$\eta \rightarrow 0 \quad \begin{cases} \lambda \geq 0 \quad g \leq 0 \quad \lambda g = 0 \rightarrow \text{Loading/unloading conditions} \\ \text{If } g = 0 \quad \lambda \dot{g} = 0 \rightarrow \text{Persistency conditions} \end{cases}$$

As  $\eta \rightarrow 0$  (inviscid case ) the rate independent continuum damage model is recovered

**Remark 3**  
 For  $\eta \neq 0$  and  $\lambda \neq 0$  then  $\langle g \rangle = \lambda \eta \neq 0 \rightarrow \begin{cases} g \equiv \tau_e - r > 0 \rightarrow \tau_e > r \\ f \equiv \tau_\sigma - q > 0 \rightarrow \tau_\sigma > q \end{cases}$

$$\begin{cases} f \leq 0 \Leftrightarrow g \leq 0 \rightarrow \dot{r} = \frac{1}{\eta} \langle g \rangle = 0 \rightarrow \text{No evolution (unloading / neutral loading)} \\ f > 0 \Leftrightarrow g > 0 \rightarrow \dot{r} = \frac{1}{\eta} \langle g \rangle > 0 \rightarrow \text{Evolution (loading)} \end{cases}$$

The stress/strain state can lay outside the elastic domain



$$\psi(\boldsymbol{\varepsilon}, r) = (1-d)\psi_0(\boldsymbol{\varepsilon}) = (1-d)\frac{1}{2}(\boldsymbol{\varepsilon} : \mathbb{C} : \boldsymbol{\varepsilon})$$

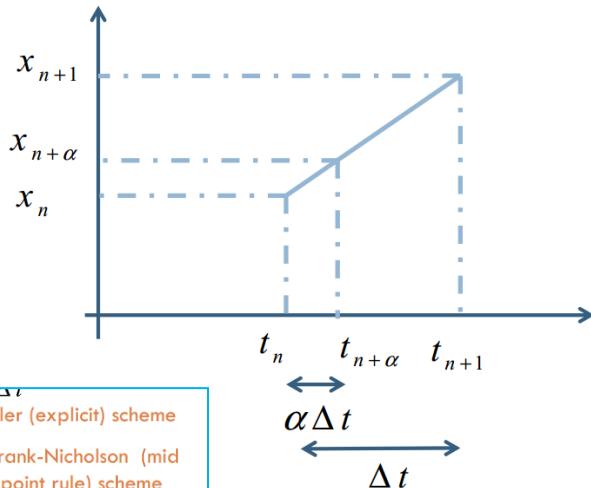
$$\dot{d} \geq 0$$

$$\begin{aligned} \mathbb{C}_{\text{tang}}^{\text{vd}} &= \frac{\partial \boldsymbol{\sigma}}{\partial \boldsymbol{\varepsilon}} = (1-d)\mathbb{C} \\ \frac{\partial \boldsymbol{\sigma}}{\partial t} &= \frac{1}{\eta} d'(r) \langle g(\boldsymbol{\varepsilon}, r) \rangle \mathbb{C} : \boldsymbol{\varepsilon} \end{aligned}$$

$$\dot{r} = \frac{1}{\eta} \langle \tau_\varepsilon(\boldsymbol{\varepsilon}(t)) - r(t) \rangle \rightarrow \begin{cases} \text{Data: } \boldsymbol{\varepsilon}(s) \\ s \in [0, t] \\ \text{Result: } r(t) \end{cases}$$

$$\begin{aligned} x_n &:= x(t_n) \\ x_{n+1} &:= x(t_{n+1}) \\ x_{n+\alpha} &:= x(t_n + \alpha \Delta t) \end{aligned}$$

$$x_{n+\alpha} \simeq (1-\alpha)x_n + \alpha x_{n+1}$$



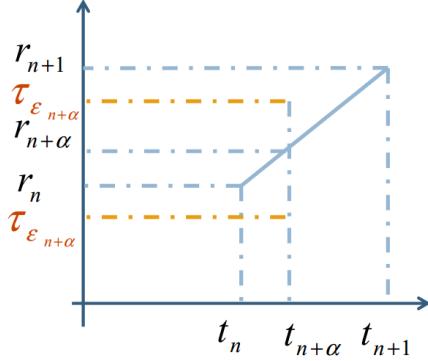
**Integration**

$$\dot{x}_{n+1} = \frac{x_{n+1} - x_n}{\Delta t} = f(x_{n+\alpha}, t_{n+\alpha})$$

$\alpha = 0 \quad x_{n+\alpha} = x_n$ $\alpha = 0.5 \quad x_{n+\alpha} = \frac{x_{n+1} + x_n}{2}$ $\alpha = 1 \quad x_{n+\alpha} = x_{n+1}$	$\Delta t$ $\text{Forward Euler (explicit) scheme}$ $\text{Crank-Nicholson (mid point rule) scheme}$ $\text{Backward Euler (implicit) scheme}$
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$$\dot{r} = \frac{1}{\eta} \langle \tau_{\varepsilon}(\varepsilon(t)) - r(t) \rangle$$

- Data  $\begin{cases} r|_{t=t_0} = r_0 \\ r_n, \varepsilon_n, \varepsilon_{n+1} \Rightarrow \tau_{\varepsilon_n}(\varepsilon_n), \tau_{\varepsilon_{n+1}}(\varepsilon_{n+1}) \end{cases}$
- Unknowns  $r_{n+1} \rightarrow q_{n+1} \rightarrow d_{n+1} \rightarrow \sigma_{n+1}$
- Integration  $\begin{cases} \dot{r}_{n+1} = \frac{r_{n+1} - r_n}{\Delta t} = \frac{1}{\eta} \langle \tau_{\varepsilon_{n+\alpha}} - r_{n+\alpha} \rangle \\ \tau_{\varepsilon_{n+\alpha}} = (1-\alpha)\tau_{\varepsilon_n} + \alpha\tau_{\varepsilon_{n+1}} \rightarrow Data \\ r_{n+\alpha} = (1-\alpha)r_n + \alpha r_{n+1} \end{cases}$



**THEOREM:**  $\boxed{\tau_{\varepsilon_{n+\alpha}} \leq r_{n+\alpha} \Leftrightarrow \tau_{\varepsilon_{n+\alpha}} \leq r_n} \rightarrow (equivalence)$

**COROLLARY:**  $\boxed{\tau_{\varepsilon_{n+\alpha}} > r_{n+\alpha} \Leftrightarrow \tau_{\varepsilon_{n+\alpha}} > r_n} \rightarrow (equivalence)$

## 1 ) Unloading / Neutral loading

$$\Leftrightarrow \boxed{\tau_{\varepsilon_{n+a}} \leq r_n}$$

$$\boxed{r_{n+1} = r_n}$$

## 2 ) Loading

$$\boxed{\tau_{\varepsilon_{n+a}} > r_n}$$

$$\boxed{r_{n+1} = \frac{[\eta - \Delta t(1 - \alpha)]r_n + \Delta t\tau_{\varepsilon_{n+\alpha}}}{\eta + \alpha\Delta t}}$$

**Inviscid case and implicit integration**

$$\boxed{r_{n+1} = \tau_{\varepsilon_{n+1}}}$$

Numerical integration inherits the properties of the model and recovers the solution of the rate independent problem for the **inviscid case and implicit integration**

## Numerical integration: stability analysis

Crank-Nicholson (mid-point rule)

$$\frac{1}{2} \leq \alpha \leq 1$$

Backward-Euler

$$\begin{cases} f(x(t), \dot{x}(t)) = 0 \\ x|_{t=0} = x_0 \end{cases} \rightarrow x_{n+1} = ax_n + b \quad \boxed{|a| \leq 1} \rightarrow \begin{cases} \text{errors do not propagate} \\ (\text{stable integration}) \end{cases}$$

## Numerical integration: accuracy analysis

### □ Accuracy

$$\dot{r}(t) = \frac{1}{\eta}(\tau_\varepsilon(t) - r(t))$$

**First order accuracy/consistency:**

Numerical solution fulfills this equation as  $\Delta t \rightarrow 0$

$$\ddot{r}(t) = \frac{1}{\eta}(\dot{\tau}_\varepsilon(t) - \dot{r}(t))$$

**Second order accuracy:**

Numerical solution fulfills this equation as  $\Delta t \rightarrow 0$

### □ Alpha method

$$r_{n+\alpha} = (1-\alpha)r_n + \alpha r_{n+1} \quad \Rightarrow \quad r_{n+\alpha}(t_{n+1}) = (1-\alpha)r_n + \alpha r_{n+1}(t_{n+1})$$

$$\dot{r}_{n+\alpha} = \frac{\partial r_{n+\alpha}}{\partial t_{n+1}} = \alpha \frac{\partial r_{n+1}}{\partial t_{n+1}} = \alpha \dot{r}_{n+1} ; \quad \ddot{r}_{n+\alpha} = \frac{\partial \dot{r}_{n+\alpha}}{\partial t_{n+1}} = \alpha \frac{\partial \dot{r}_{n+1}}{\partial t_{n+1}} = \alpha \ddot{r}_{n+1}$$

$$\tau_{n+\alpha} = (1-\alpha)\tau_n + \alpha \tau_{n+1} \quad \Rightarrow \quad$$

$$\dot{\tau}_{n+\alpha} = \frac{\partial \tau_{n+\alpha}}{\partial t_{n+1}} = \alpha \dot{\tau}_{n+1} ; \quad \ddot{\tau}_{n+\alpha} = \frac{\partial \dot{\tau}_{n+\alpha}}{\partial t_{n+1}} = \alpha \frac{\partial \dot{\tau}_{n+1}}{\partial t_{n+1}} = \alpha \ddot{\tau}_{n+1}$$

$$\boxed{\dot{r}(t) = \frac{1}{\eta}(\tau_\varepsilon(t) - r(t)) \quad ; \quad \ddot{r}(t) = \frac{1}{\eta}(\dot{\tau}_\varepsilon(t) - \dot{r}(t))}$$

$$\text{Stable: } \alpha = \left[ \frac{1}{2}, 1 \right] \quad \text{First order accurate: } \alpha = [0, 1] \quad \text{Second order accurate: } \alpha = \frac{1}{2}$$

$$\begin{cases} \mathbb{C}_{\text{tang}}^{vd} = \frac{\partial \sigma(\boldsymbol{\varepsilon})}{\partial \boldsymbol{\varepsilon}} = (1-d)\mathbb{C} \rightarrow \text{Analytic tangent operator} \\ \mathbb{C}_{\text{alg},n+1}^{vd} := \frac{\partial \sigma_{n+1}(\boldsymbol{\varepsilon}_{n+1})}{\partial \boldsymbol{\varepsilon}_{n+1}} \rightarrow \text{Algorithmic tangent operator} \end{cases} \quad \begin{cases} \mathbb{C}_{\text{alg},n+1}^{vd} \neq \mathbb{C}_{\text{tang}}^{vd}(\boldsymbol{\varepsilon}_{n+1}) \\ \lim_{\Delta t \rightarrow 0} \mathbb{C}_{\text{alg},n+1}^{vd} = \mathbb{C}_{\text{tang}}^{vd}(\boldsymbol{\varepsilon}_{n+1}) \end{cases}$$

$$\begin{cases} \boldsymbol{\sigma}_{n+1} = [1 - d_{n+1}] \mathbb{C} : \boldsymbol{\varepsilon}_{n+1} \\ d_{n+1}(r_{n+1}) = 1 - \frac{q(r_{n+1})}{r_{n+1}} \\ r_{n+1} = \begin{cases} r_n \text{ (elastic/unloading)} \\ \frac{[\eta - \Delta t(1-\alpha)]}{\eta + \alpha \Delta t} r_n + \frac{\Delta t}{\eta + \alpha \Delta t} \tau_{\varepsilon_{n+\alpha}}(\boldsymbol{\varepsilon}_{n+1}) \text{ (loading)} \end{cases} \end{cases} \rightarrow \begin{cases} \boldsymbol{\sigma}_{n+1} = \Sigma_{n+1}(\boldsymbol{\varepsilon}_{n+1}) \\ \delta \boldsymbol{\sigma}_{n+1} = \mathbb{C}_{\text{alg},n+1}^{vd}(\boldsymbol{\varepsilon}_{n+1}) : \delta \boldsymbol{\varepsilon}_{n+1} \end{cases}$$

$$\delta \boldsymbol{\sigma}_{n+1} = (1 - d_{n+1}) \mathbb{C} : \delta \boldsymbol{\varepsilon}_{n+1} - d'(r_{n+1}) \delta r_{n+1} \otimes \underbrace{\mathbb{C} : \boldsymbol{\varepsilon}_{n+1}}_{\bar{\sigma}_{n+1}} ; \quad \delta r_{n+1} = \begin{cases} 0 \text{ (elastic/unloading)} \\ \frac{\alpha \Delta t}{\eta + \alpha \Delta t} \delta \tau_{\varepsilon_{n+1}} \text{ (loading)} \end{cases}$$

$$\tau_{\varepsilon_{n+1}} = \sqrt{\boldsymbol{\varepsilon}_{n+1} : \mathbb{C} : \boldsymbol{\varepsilon}_{n+1}} \quad \Rightarrow \quad \delta\tau_{\varepsilon_{n+1}} = \frac{1}{\tau_{\varepsilon_{n+1}}} \underbrace{\boldsymbol{\varepsilon}_{n+1} : \mathbb{C} : \delta\boldsymbol{\varepsilon}_{n+1}}_{\bar{\boldsymbol{\sigma}}_{n+1}} = \frac{1}{\tau_{\varepsilon_{n+1}}} \bar{\boldsymbol{\sigma}}_{n+1} : \delta\boldsymbol{\varepsilon}_{n+1}$$

$$\delta\boldsymbol{\sigma}_{n+1} = \begin{cases} (1-d_n)\mathbb{C} : \delta\boldsymbol{\varepsilon}_{n+1} & \text{(elastic/unloading)} \\ [(1-d_{n+1})\mathbb{C} - \frac{\alpha\Delta t}{\eta + \alpha\Delta t} d'_{n+1} \bar{\boldsymbol{\sigma}}_{n+1} \otimes \bar{\boldsymbol{\sigma}}_{n+1}] : \delta\boldsymbol{\varepsilon}_{n+1} & \text{(loading)} \end{cases} = \mathbb{C}_{\text{alg},n+1}^{vd} : \delta\boldsymbol{\varepsilon}_{n+1}$$

$$\mathbb{C}_{\text{alg},n+1}^{vd} = \begin{cases} (1-d_n)\mathbb{C} & \text{(elastic/unloading)} \\ \underbrace{(1-d_{n+1})\mathbb{C} - \frac{\alpha\Delta t}{\eta + \alpha\Delta t} d'_{n+1} \bar{\boldsymbol{\sigma}}_{n+1} \otimes \bar{\boldsymbol{\sigma}}_{n+1}}_{\mathbb{C}_{\text{tang},n+1}^{vd}} & \text{(loading)} \end{cases} \quad (d'_{n+1} = -\frac{H_{n+1}r_{n+1} - q(r_{n+1})}{(r_{n+1})^2})$$

$$\mathbb{C}_{\text{alg},n+1}^{vd} = \begin{cases} (1-d_n)\mathbb{C} & \text{(elastic/unloading)} \\ \underbrace{(1-d_{n+1})\mathbb{C} - \frac{\alpha\Delta t}{\eta + \alpha\Delta t} d'_{n+1} \bar{\boldsymbol{\sigma}}_{n+1} \otimes \bar{\boldsymbol{\sigma}}_{n+1}}_{\mathbb{C}_{\text{tang},n+1}^{vd}} & \text{(loading)} \end{cases} \quad \text{Additional term } \rightarrow \mathcal{O}(\Delta t)$$

- **Remark 1**
  - If:  $\alpha = 0$   $\mathbb{C}_{\text{alg},n+1}^{vd} = \mathbb{C}_{\text{tang},n+1}^{vd}$
  - If:  $\Delta t = 0$   $\mathbb{C}_{\text{alg},n+1}^{vd} = \mathbb{C}_{n+1}^d$

Algorithmic and analytical tangent operators match
- **Remark 2**
  - If:  $\eta \rightarrow 0$

$\mathbb{C}_{\text{alg},n+1}^{vd} = (1-d_{n+1})\mathbb{C} - \frac{1}{\tau_{\varepsilon_{n+1}}} d'_{n+1} \bar{\boldsymbol{\sigma}}_{n+1} \otimes \bar{\boldsymbol{\sigma}}_{n+1}$

Algorithmic viscous tangent operator match (in the inviscid limit) the one for the rate independent damage model

INPUT DATA  $[t_n, t_n + \Delta t = t_{n+1}] \rightarrow \boldsymbol{\varepsilon}_n, r_n, \boldsymbol{\varepsilon}_{t_n+1}$

$$\text{Step 1} \rightarrow \text{Compute } \begin{cases} \bar{\boldsymbol{\sigma}}_{n+1} = \mathbb{C} : \boldsymbol{\varepsilon}_{n+1} \rightarrow \begin{cases} \tau_{\varepsilon_n} = \sqrt{\boldsymbol{\varepsilon}_n : \mathbb{C} : \boldsymbol{\varepsilon}_n} \\ \tau_{\varepsilon_{n+1}} = \sqrt{\boldsymbol{\varepsilon}_{n+1} : \mathbb{C} : \boldsymbol{\varepsilon}_{n+1}} = \sqrt{\bar{\boldsymbol{\sigma}}_{n+1} : \boldsymbol{\varepsilon}_{n+1}} \end{cases} \\ \tau_{\varepsilon_{n+\alpha}} = (1 - \alpha)\tau_{\varepsilon_n} + \alpha\tau_{\varepsilon_{n+1}} \end{cases}$$

$$\text{Step 2} \rightarrow \text{If } \tau_{\varepsilon_{n+\alpha}} \leq r_n \rightarrow \begin{cases} \text{Elastic} \\ \text{Unloading} \end{cases}$$

$$\rightarrow \begin{cases} r_{n+1} = r_n ; d_{n+1} = d_n = 1 - \frac{q(r_{n+1})}{r_{n+1}} ; \boldsymbol{\sigma}_{n+1} = (1 - d_{n+1})\bar{\boldsymbol{\sigma}}_{n+1} \\ \mathbb{C}_{\text{alg}, n+1}^{vd} = (1 - d_{n+1})\mathbb{C} \end{cases}$$

$$\text{Step 3} \rightarrow \text{If } \tau_{\varepsilon_{n+\alpha}} > r_n \rightarrow \text{(Loading)}$$

$$\rightarrow \begin{cases} r_{n+1} = \frac{[\eta - \Delta t(1 - \alpha)]}{\eta + \alpha\Delta t} r_n + \frac{\Delta t}{\eta + \alpha\Delta t} \tau_{\varepsilon_{n+\alpha}} ; d_{n+1} = 1 - \frac{q(r_{n+1})}{r_{n+1}} \\ \boldsymbol{\sigma}_{n+1} = (1 - d_{n+1})\bar{\boldsymbol{\sigma}}_{n+1} \\ \mathbb{C}_{\text{alg}, n+1}^{vd} = (1 - d_{n+1})\mathbb{C} + \\ \quad + \frac{\alpha\Delta t}{\eta + \alpha\Delta t} \frac{1}{\tau_{\varepsilon_{n+1}}} \frac{H_{n+1}r_{n+1} - q(r_{n+1})}{(r_{n+1})^2} (\bar{\boldsymbol{\sigma}}_{n+1} \otimes \bar{\boldsymbol{\sigma}}_{n+1}) \end{cases}$$

OUTPUT DATA  $[t_n, t_n + \Delta t = t_{n+1}] \rightarrow r_{n+1}, \boldsymbol{\sigma}_{n+1}, \mathbb{C}_{\text{alg}, n+1}^{vd}$

For  $\eta = 0$  and  $\alpha = 1$  the inviscid model is recovered !!!!

$$\dot{\boldsymbol{\varepsilon}}_v^{pl} = \frac{1}{\mu}(\boldsymbol{\varepsilon}^{pl} - \boldsymbol{\varepsilon}_v^{pl})$$

- Taking into account the symmetry of the stress and strain tensors, they can be written in vector form:

$$\boldsymbol{\varepsilon} = \begin{bmatrix} \varepsilon_x & \varepsilon_{xy} & \varepsilon_{xz} \\ \varepsilon_{xy} & \varepsilon_y & \varepsilon_{yz} \\ \varepsilon_{xz} & \varepsilon_{yz} & \varepsilon_z \end{bmatrix} \stackrel{\text{not.}}{=} \begin{bmatrix} \varepsilon_x & \frac{1}{2}\gamma_{xy} & \frac{1}{2}\gamma_{xz} \\ \frac{1}{2}\gamma_{xy} & \varepsilon_y & \frac{1}{2}\gamma_{yz} \\ \frac{1}{2}\gamma_{xz} & \frac{1}{2}\gamma_{yz} & \varepsilon_z \end{bmatrix} \xrightarrow{\text{VOIGT'S NOTATION}} \{\boldsymbol{\varepsilon}\} \stackrel{\text{def}}{=} \begin{bmatrix} \varepsilon_x \\ \varepsilon_y \\ \varepsilon_z \\ \gamma_{xy} \\ \gamma_{xz} \\ \gamma_{yz} \end{bmatrix} \in \mathbb{R}^6$$

**REMARK**  
The double contraction  $(\boldsymbol{\sigma} : \boldsymbol{\varepsilon})$  is equivalent to the scalar (dot) product  $(\{\boldsymbol{\sigma}\} \cdot \{\boldsymbol{\varepsilon}\})$ :

$\overset{\text{vectors}}{\textcircled{S} : \boldsymbol{\varepsilon} = \{\boldsymbol{\sigma}\} \cdot \{\boldsymbol{\varepsilon}\}}$   $\Leftrightarrow \overset{\text{2nd order tensors}}{\sigma_{ij}\varepsilon_{ij} = \sigma_i\varepsilon_i}$

$$\boldsymbol{\sigma} \equiv \begin{bmatrix} \sigma_x & \tau_{xy} & \tau_{xz} \\ \tau_{xy} & \sigma_y & \tau_{yz} \\ \tau_{xz} & \tau_{yz} & \sigma_z \end{bmatrix} \xrightarrow{\text{VOIGT'S NOTATION}} \{\boldsymbol{\sigma}\} \stackrel{\text{def}}{=} \begin{bmatrix} \sigma_x \\ \sigma_y \\ \sigma_z \\ \tau_{xy} \\ \tau_{xz} \\ \tau_{yz} \end{bmatrix} \in \mathbb{R}^6$$

- Hooke's law in terms of the stress and strain vectors:

$$\begin{cases} \boldsymbol{\sigma} = \mathbb{C} : \boldsymbol{\varepsilon} \\ \mathbb{C} = \lambda \mathbf{1} \otimes \mathbf{1} + 2\mu \mathbb{I} \end{cases} \xrightarrow{\quad} \begin{cases} \{\boldsymbol{\sigma}\} = \mathbf{D} \cdot \{\boldsymbol{\varepsilon}\} \\ \sigma_i = D_{ij}\varepsilon_j \quad i \in \{1, \dots, 6\} \end{cases}$$

Where  $\mathbf{D}$  is the matrix of elastic constants:

$$\mathbf{D} = \frac{E(1-\nu)}{(1+\nu)(1-2\nu)} \begin{bmatrix} 1 & \frac{\nu}{1-\nu} & \frac{\nu}{1-\nu} & 0 & 0 & 0 \\ \frac{\nu}{1-\nu} & 1 & \frac{\nu}{1-\nu} & 0 & 0 & 0 \\ \frac{\nu}{1-\nu} & \frac{\nu}{1-\nu} & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1-2\nu}{2(1-\nu)} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1-2\nu}{2(1-\nu)} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1-2\nu}{2(1-\nu)} \end{bmatrix}$$

$$\begin{cases} \boldsymbol{\varepsilon} = \mathbb{C}^{-1} : \boldsymbol{\sigma} \\ \mathbb{C} = -\frac{\nu}{E} \mathbf{1} \otimes \mathbf{1} + \frac{1+\nu}{2E} \mathbb{I} \end{cases} \xrightarrow{\quad} \begin{cases} \{\boldsymbol{\varepsilon}\} = \mathbf{D}^{-1} \cdot \{\boldsymbol{\sigma}\} \\ \varepsilon_j = (\mathbf{D}^{-1})_{ij} \sigma_i \quad i \in \{1, \dots, 6\} \end{cases}$$

Where  $\mathbf{D}^{-1}$  is the elastic compliance matrix:

$$\mathbf{D}^{-1} = \begin{bmatrix} 1 & -\nu & -\nu & 0 & 0 & 0 \\ -\nu & 1 & -\nu & 0 & 0 & 0 \\ -\nu & -\nu & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{G} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{G} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{G} \end{bmatrix}; \quad G = \frac{E}{2(1+\nu)}$$

- The expressions used in the algorithms of the model read:

$$\tau_\varepsilon = \sqrt{\boldsymbol{\varepsilon} : \mathbb{C} : \boldsymbol{\varepsilon}} = \sqrt{\boldsymbol{\sigma} : \boldsymbol{\varepsilon}} = \sqrt{\{\boldsymbol{\sigma}\} \cdot \{\boldsymbol{\varepsilon}\}} = \sqrt{\{\boldsymbol{\varepsilon}\} \cdot \{\mathbf{D}\} \cdot \{\boldsymbol{\varepsilon}\}} = \sqrt{\{\boldsymbol{\sigma}\} \cdot \{\mathbf{D}^{-1}\} \cdot \{\boldsymbol{\sigma}\}}$$

$$\{\boldsymbol{\sigma}\}_{n+1} = (1 - d(r_{n+1})) \{\bar{\boldsymbol{\sigma}}\}$$

$$\{\bar{\boldsymbol{\sigma}}\}_{n+1} = \mathbf{D} \cdot \{\boldsymbol{\varepsilon}\}_{n+1}$$

$$\tau_{\varepsilon_{n+1}} = \sqrt{\{\boldsymbol{\varepsilon}\}_{n+1} \cdot \mathbf{D} \cdot \{\boldsymbol{\varepsilon}\}_{n+1}} = \sqrt{\{\bar{\boldsymbol{\sigma}}\}_{n+1} \cdot \mathbf{D}^{-1} \cdot \{\bar{\boldsymbol{\sigma}}\}_{n+1}} = \sqrt{\{\bar{\boldsymbol{\sigma}}\}_{n+1} \cdot \{\boldsymbol{\varepsilon}\}_{n+1}}$$

$$\underbrace{\mathbf{D}_{\text{alg},n+1}^{vd}}_{6 \times 6} = \underbrace{\frac{\partial \{\boldsymbol{\sigma}\}_{n+1}}{\partial \{\boldsymbol{\varepsilon}\}_{n+1}}}_{6 \times 6} = (1 - d_{n+1}) \underbrace{\mathbf{D}}_{6 \times 6} -$$

$$-\frac{\alpha \Delta t}{\eta + \alpha \Delta t} \frac{1}{\tau_{\varepsilon_{n+1}}} \frac{H_{n+1} r_{n+1} - q(r_{n+1})}{(r_{n+1})^2} \left( \underbrace{\{\bar{\boldsymbol{\sigma}}\}_{n+1}}_{6 \times 1} \otimes \underbrace{\{\bar{\boldsymbol{\sigma}}\}_{n+1}}_{1 \times 6} \right)$$



*Plasticity model based on effective stress*

$$\begin{aligned}f_p(\bar{\boldsymbol{\sigma}}, \kappa_p) &= \tilde{\sigma}(\bar{\boldsymbol{\sigma}}) - \sigma_Y(\kappa_p) \\ \dot{\varepsilon}_p &= \dot{\lambda} \frac{\partial g_p}{\partial \bar{\boldsymbol{\sigma}}}(\bar{\boldsymbol{\sigma}}, \kappa_p) \\ \dot{\kappa}_p &= \dot{\lambda} k_p(\bar{\boldsymbol{\sigma}}, \kappa_p) \\ f_p &\leqslant 0, \quad \dot{\lambda} \geqslant 0, \quad \dot{\lambda} f_p = 0\end{aligned}$$

$$\begin{aligned}f_d(\boldsymbol{\varepsilon}, \boldsymbol{\varepsilon}_p, \kappa_d) &= \tilde{\varepsilon}(\boldsymbol{\varepsilon}, \boldsymbol{\varepsilon}_p) - \kappa_d \\ \omega &= g_d(\kappa_d) \\ f_d &\leqslant 0, \quad \dot{\kappa}_d \geqslant 0, \quad \dot{\kappa}_d f_d = 0\end{aligned}$$

$$\dot{f}_p \leqslant 0, \quad \dot{\lambda} \geqslant 0, \quad \dot{\lambda} \dot{f}_p = 0$$

$$\dot{f}_p = \frac{\partial f_p}{\partial \bar{\boldsymbol{\sigma}}} : \dot{\bar{\boldsymbol{\sigma}}} + \frac{\partial f_p}{\partial \kappa_p} \dot{\kappa}_p = \frac{\partial f_p}{\partial \bar{\boldsymbol{\sigma}}} : \mathbf{D}_e : \dot{\boldsymbol{\varepsilon}} - \dot{\lambda} \left( \frac{\partial f_p}{\partial \bar{\boldsymbol{\sigma}}} : \mathbf{D}_e : \frac{\partial g_p}{\partial \bar{\boldsymbol{\sigma}}} + H_p k_p \right)$$

$$H_p = - \frac{\partial f_p}{\partial \kappa_p} = \frac{d\sigma_Y}{d\kappa_p}$$

$$\frac{\partial f_p}{\partial \bar{\boldsymbol{\sigma}}} : \mathbf{D}_e : \frac{\partial g_p}{\partial \bar{\boldsymbol{\sigma}}} + H_p k_p > 0$$



$$\dot{\mathbf{f}} \leqslant \mathbf{0}, \quad \dot{\boldsymbol{\lambda}} \geqslant \mathbf{0}, \quad \dot{\boldsymbol{\lambda}}^{\mathrm{T}} \dot{\mathbf{f}} = 0$$

$$\dot{\mathbf{f}}=\left\{\begin{array}{c}\dot{f}_{\text { p }} \\ \dot{f}_{\text { d }}\end{array}\right\} \dot{\boldsymbol{\lambda}}=\left\{\begin{array}{c}\dot{\lambda} \\ \dot{\kappa}_{\text { d }}\end{array}\right\}$$

$$\dot{f}_{\text { d }}=\frac{\partial f_{\text { d }}}{\partial \boldsymbol{\varepsilon}}:\dot{\boldsymbol{\varepsilon}}+\frac{\partial f_{\text { d }}}{\partial \boldsymbol{\varepsilon}_{\text { p }}}:\dot{\boldsymbol{\varepsilon}}_{\text { p }}+\frac{\partial f_{\text { d }}}{\partial \kappa_{\text { d }}} \dot{\kappa}_{\text { d }}=\frac{\partial f_{\text { d }}}{\partial \boldsymbol{\varepsilon}}:\dot{\boldsymbol{\varepsilon}}+\dot{\lambda} \frac{\partial f_{\text { d }}}{\partial \boldsymbol{\varepsilon}_{\text { p }}}:\frac{\partial g_{\text { p }}}{\partial \bar{\boldsymbol{\sigma}}}+\frac{\partial f_{\text { d }}}{\partial \kappa_{\text { d }}} \dot{\kappa}_{\text { d }}$$

$$\mathbf{A}=\begin{pmatrix} \frac{\partial f_{\text { p }}}{\partial \bar{\boldsymbol{\sigma}}}:\mathbf{D}_{\text { e }}:\frac{\partial g_{\text { p }}}{\partial \bar{\boldsymbol{\sigma}}}+H_{\text { p }} k_{\text { p }} & 0 \\ -\frac{\partial f_{\text { d }}}{\partial \boldsymbol{\varepsilon}_{\text { p }}}:\frac{\partial g_{\text { p }}}{\partial \bar{\boldsymbol{\sigma}}} & -\frac{\partial f_{\text { d }}}{\partial \kappa_{\text { d }}} \end{pmatrix} \quad \mathbf{b}=\left\{\begin{array}{c}\frac{\partial f_{\text { p }}}{\partial \bar{\boldsymbol{\sigma}}}:\mathbf{D}_{\text { e }}:\dot{\boldsymbol{\varepsilon}} \\ \frac{\partial f_{\text { d }}}{\partial \boldsymbol{\varepsilon}}:\dot{\boldsymbol{\varepsilon}}\end{array}\right\}$$

$$\mathbf{A} \dot{\boldsymbol{\lambda}}-\mathbf{b} \geqslant \mathbf{0}, \quad \dot{\boldsymbol{\lambda}} \geqslant \mathbf{0}, \quad \dot{\boldsymbol{\lambda}}^{\mathrm{T}}(\mathbf{A} \dot{\boldsymbol{\lambda}}-\mathbf{b})=0$$

$$A_{11}>0,\quad A_{22}>0,\quad A_{11}A_{22}-A_{21}A_{12}>0$$

$$\frac{\partial f_{\text { p }}}{\partial \bar{\boldsymbol{\sigma}}}:\mathbf{D}_{\text { e }}:\frac{\partial g_{\text { p }}}{\partial \bar{\boldsymbol{\sigma}}}+H_{\text { p }} k_{\text { p }}>0 \\ -\frac{\partial f_{\text { d }}}{\partial \kappa_{\text { d }}}>0$$



$$f_{\mathrm{p}}(\boldsymbol{\sigma}, \kappa_{\mathrm{p}}) = \tilde{\sigma}(\boldsymbol{\sigma}) - \sigma_{\mathrm{Y}}(\kappa_{\mathrm{p}})$$

$$\dot{\boldsymbol{\varepsilon}}_{\mathrm{p}} = \dot{\lambda}\frac{\partial g_{\mathrm{p}}}{\partial \boldsymbol{\sigma}}$$

$$\dot{\kappa}_{\mathrm{p}} = \dot{\lambda} k_{\mathrm{p}}(\boldsymbol{\sigma}, \kappa_{\mathrm{p}})$$

$$\dot{f}_{\mathrm{p}} = \frac{\partial f_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} : \dot{\boldsymbol{\sigma}} + \frac{\partial f_{\mathrm{p}}}{\partial \kappa_{\mathrm{p}}} \dot{\kappa}_{\mathrm{p}} = (1-\omega)\frac{\partial f_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} : \mathbf{D}_{\mathrm{e}} : \dot{\boldsymbol{\varepsilon}} - \dot{\lambda} \bigg[ (1-\omega)\frac{\partial f_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} : \mathbf{D}_{\mathrm{e}} : \frac{\partial g_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} + H_{\mathrm{p}} k_{\mathrm{p}} \bigg] - \dot{\kappa}_{\mathrm{d}} \frac{\mathrm{d} g_{\mathrm{d}}}{\mathrm{d} \kappa_{\mathrm{d}}} \frac{\partial f_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} : \bar{\boldsymbol{\sigma}}$$

$$\mathbf{A}=\begin{pmatrix}(1-\omega)\frac{\partial f_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} : \mathbf{D}_{\mathrm{e}} : \frac{\partial g_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} + H_{\mathrm{p}} k_{\mathrm{p}} & \frac{\mathrm{d} g_{\mathrm{d}}}{\mathrm{d} \kappa_{\mathrm{d}}} \frac{\partial f_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} : \bar{\boldsymbol{\sigma}} \\ -\frac{\partial f_{\mathrm{d}}}{\partial \boldsymbol{\varepsilon}_{\mathrm{p}}} : \frac{\partial g_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} & 1\end{pmatrix}$$

$$(1-\omega)\frac{\partial f_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} : \mathbf{D}_{\mathrm{e}} : \frac{\partial g_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} + H_{\mathrm{p}} k_{\mathrm{p}} > 0$$

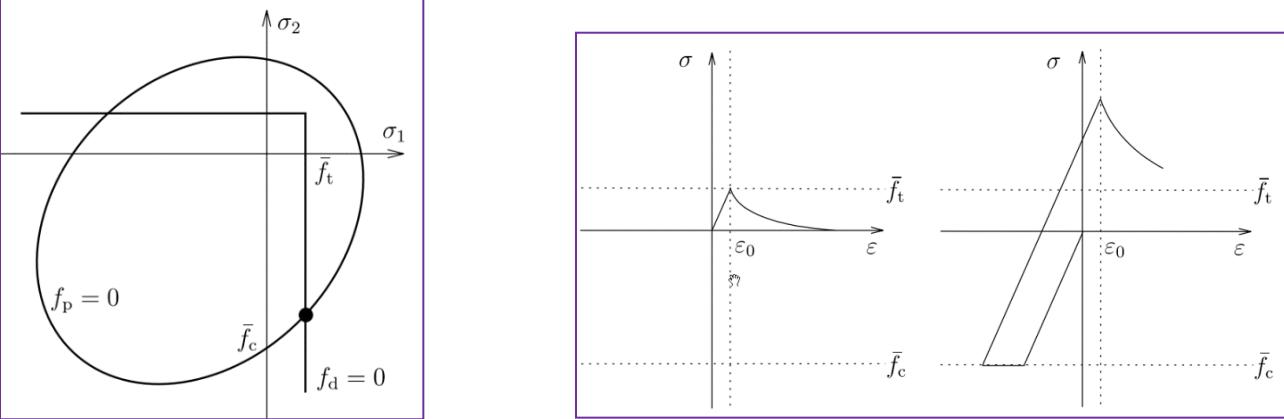
$$1>0$$

$$(1-\omega)\frac{\partial f_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} : \mathbf{D}_{\mathrm{e}} : \frac{\partial g_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} + H_{\mathrm{p}} k_{\mathrm{p}} > -\frac{\partial f_{\mathrm{d}}}{\partial \boldsymbol{\varepsilon}_{\mathrm{p}}} : \frac{\partial g_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} \left( \frac{\mathrm{d} g_{\mathrm{d}}}{\mathrm{d} \kappa_{\mathrm{d}}} \frac{\partial f_{\mathrm{p}}}{\partial \boldsymbol{\sigma}} : \bar{\boldsymbol{\sigma}} \right)$$

$$H_{\mathrm{p}} k_{\mathrm{p}} > 0$$

$$f_{\rm p}(\boldsymbol{\sigma}, \kappa_{\rm p}) = c_\phi I_1(\boldsymbol{\sigma}) + \sqrt{J_2(\boldsymbol{\sigma})} - \sigma_{\rm Y}(\kappa_{\rm p})$$

$$g_{\rm p}(\boldsymbol{\sigma}) = c_\psi I_1(\boldsymbol{\sigma}) + \sqrt{J_2(\boldsymbol{\sigma})}$$



$$\widetilde{\varepsilon}(\boldsymbol{\varepsilon},\boldsymbol{\varepsilon}_{\mathrm{p}})=\frac{1}{E}\max_{I=1,2,3}\langle\bar{\sigma}_I(\boldsymbol{\varepsilon},\boldsymbol{\varepsilon}_{\mathrm{p}})\rangle$$

$$g_{\rm d}(\kappa_{\rm d})=\left\{\begin{array}{ll} 0 & \text{if } \kappa_{\rm d}\leqslant \varepsilon_0 \\ 1-\dfrac{\varepsilon_0}{\kappa_{\rm d}}\exp\left(-\dfrac{\kappa_{\rm d}-\varepsilon_0}{\varepsilon_{\rm f}-\varepsilon_0}\right) & \text{if } \kappa_{\rm d}\geqslant \varepsilon_0 \end{array}\right.$$

$$\dot{\hat{f}}_{\rm d}=\frac{\partial f_{\rm d}}{\partial \pmb{\varepsilon}_{\rm e}}:\dot{\pmb{\varepsilon}}_{\rm e}+\frac{\partial f_{\rm d}}{\partial \kappa_{\rm d}}\dot{\kappa}_{\rm d}=\frac{\partial \widetilde{\varepsilon}}{\partial \pmb{\varepsilon}_{\rm e}}:\left(\dot{\pmb{\varepsilon}}-\dot{\lambda}\frac{\partial g_{\rm p}}{\partial \pmb{\sigma}}\right)-\dot{\kappa}_{\rm d}$$

$$(1-\omega)\frac{\partial f_{\rm p}}{\partial \pmb{\sigma}}:\mathbf{D}_{\rm e}:\frac{\partial g_{\rm p}}{\partial \pmb{\sigma}}+H_{\rm p}k_{\rm p}>\frac{\partial \widetilde{\varepsilon}}{\partial \pmb{\varepsilon}_{\rm e}}:\frac{\partial g_{\rm p}}{\partial \pmb{\sigma}}\left(\frac{{\rm d}g_{\rm d}}{{\rm d}\kappa_{\rm d}}\frac{\partial f_{\rm p}}{\partial \pmb{\sigma}}:\bar{\pmb{\sigma}}\right)$$

$$(1-\omega)E\biggl(\frac{3c_{\phi}c_{\psi}}{1-2\nu}+\frac{1}{2(1+\nu)}\biggr)+H_{\rm p}k_{\rm p}>\biggl(\frac{1}{\kappa_{\rm d}}+\frac{1}{\varepsilon_{\rm f}-\varepsilon_0}\biggr)\biggl(\frac{c_{\psi}}{(1-2\nu)}+\frac{s_1}{2(1+\nu)\sqrt{J_2}}\biggr)\sigma_{\rm Y}(\kappa_{\rm p})$$

$$H_{\rm p}k_{\rm p}>\biggl(\frac{1.8257}{1-\varepsilon_0/\varepsilon_{\rm f}}-0.4342\biggr)E>1.3915E$$



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$$\rho\psi(\boldsymbol{\varepsilon}, \boldsymbol{\varepsilon}_p, \omega) = \frac{1}{2}(1-\omega) \times (\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) : \mathbf{D}_e : (\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p)$$

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$$\dot{\mathbf{D}} = \boldsymbol{\sigma} : \dot{\boldsymbol{\varepsilon}} - \rho\dot{\psi} = \boldsymbol{\sigma} : \dot{\boldsymbol{\varepsilon}} - \left( \frac{\partial\rho\psi}{\partial\boldsymbol{\varepsilon}} \dot{\boldsymbol{\varepsilon}} + \frac{\partial\rho\psi}{\partial\boldsymbol{\varepsilon}_p} \dot{\boldsymbol{\varepsilon}}_p + \frac{\partial\rho\psi}{\partial\omega} \dot{\omega} \right)$$

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$$\frac{\partial\rho\psi}{\partial\boldsymbol{\varepsilon}} \dot{\boldsymbol{\varepsilon}} = (1-\omega)(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) : \mathbf{D}_e : \dot{\boldsymbol{\varepsilon}} = \boldsymbol{\sigma} : \dot{\boldsymbol{\varepsilon}}$$

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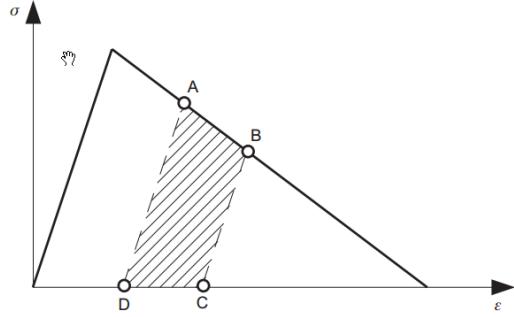
$$\dot{\mathbf{D}} = \dot{\mathbf{D}}^p + \dot{\mathbf{D}}^e = \frac{\partial\rho\psi}{\partial\boldsymbol{\varepsilon}_p} \dot{\boldsymbol{\varepsilon}}_p + \frac{\partial\rho\psi}{\partial\omega} \dot{\omega}$$

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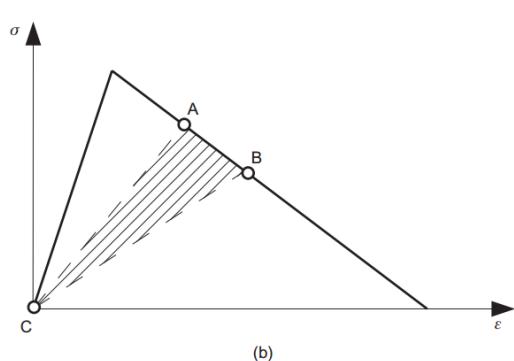
$$\dot{\mathbf{D}}^p = \frac{\partial\rho\psi}{\partial\boldsymbol{\varepsilon}_p} \dot{\boldsymbol{\varepsilon}}_p = (1-\omega)(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) : \mathbf{D}_e : \dot{\boldsymbol{\varepsilon}}_p$$

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$$\dot{\mathbf{D}}^d = \frac{\partial\rho\psi}{\partial\omega} \dot{\omega} = \frac{1}{2}(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) : \mathbf{D}_e : (\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) \dot{\omega}$$



(a)



(b)

$$\rho\psi(\boldsymbol{\varepsilon}, \boldsymbol{\varepsilon}_p, \omega) = \frac{1}{2}(1-\omega)(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) : \mathbf{D}_e : (\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p)$$

$$\mathcal{D} = \boldsymbol{\sigma} : \dot{\boldsymbol{\varepsilon}} - \rho\dot{\psi} = \left( \boldsymbol{\sigma} - \rho \frac{\partial\psi}{\partial\boldsymbol{\varepsilon}} \right) : \dot{\boldsymbol{\varepsilon}} - \rho \frac{\partial\psi}{\partial\boldsymbol{\varepsilon}_p} : \dot{\boldsymbol{\varepsilon}}_p - \rho \frac{\partial\psi}{\partial\omega} \dot{\omega}$$

$$\boldsymbol{\sigma} = \rho \frac{\partial\psi}{\partial\boldsymbol{\varepsilon}} = (1-\omega)\mathbf{D}_e : (\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p)$$

$$\mathcal{D} = -\rho \frac{\partial\psi}{\partial\boldsymbol{\varepsilon}_p} : \dot{\boldsymbol{\varepsilon}}_p - \rho \frac{\partial\psi}{\partial\omega} \dot{\omega} = (1-\omega)(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) : \mathbf{D}_e : \dot{\boldsymbol{\varepsilon}}_p + \frac{1}{2}(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) : \mathbf{D}_e : (\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) \dot{\omega} = \boldsymbol{\sigma} : \dot{\boldsymbol{\varepsilon}}_p + Y\dot{\omega} \geq 0$$

$$Y = -\rho \frac{\partial\psi}{\partial\omega} = \frac{1}{2}(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) : \mathbf{D}_e : (\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p)$$

$$\mathcal{D}_p = \boldsymbol{\sigma} : \dot{\boldsymbol{\varepsilon}}_p = \lambda \boldsymbol{\sigma} : \frac{\partial g_p}{\partial \bar{\boldsymbol{\sigma}}} = \lambda(1-\omega) \bar{\boldsymbol{\sigma}} : \frac{\partial g_p}{\partial \bar{\boldsymbol{\sigma}}}$$

$$\bar{\boldsymbol{\sigma}} = \mathbf{D}_e : (\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p)$$