Electrodynamics and Introductory Problems A \mathbf{S}^1

I begin ascertaining that the motivation to write such a booklet has been got from Prof. Jaan Kalda. He has made many more great booklets on topics in Physics that concentrate on Competitive Physics problems.

The pedagogy of Electrodynamics has two informal parts, one is using the Coulomb's Law and the other is using Vector Calculus. Coulomb's law shall easily give the magnitude and the direction of the force between two charges, and the Vector calculus approach can solve the thing that would be tiresome to be done with coulomb's. Problems can be as easy as finding the *Force* in between two charged spheres; or as deep as finding the charge density of two spherical conductor in the aforementioned configuration. One set of problem we use the Coulombs law, basically the fact of using Vectors. Actually the development in Electrodynamics has been majorly contributed by its similarity with fluid, recalling the way we learn Gauss's Law at first. The analogy is taken to the next level, we think as if the dipole is a source and sink pair that conducts water, curiously, the water flow, the streamline in the real world takes the path same as the Electric field lines when it flows above a surface. The Momentum of the fluid particles through some wide river and the Field equation of a Homogenous² Electric field both obey non vortex flow, mathematically,

$$\nabla \times \mathbf{P} = 0$$
$$\nabla \times \mathbf{E} = 0$$

The $\nabla \times$ above with a field equation tells the magnitude of rotation or vortex that a field might have. Like, the Magnetic field is always circular around the current carrying wire, so we can say that the magnetic field that case has non zero value as in the equation above. The $\nabla \times$ is called the curl. We will discover the idea's on these later.

Doing the maths, we find that one point there is nothing left but the equations like the Gauss's law ones. When it is required to practically solve maths, the application of things like the ∇ becomes important. And the math's name is $Vector\ Calculus$. No worries, vector calculus is just taking derivatives with equation with a coefficient $\hat{x}, \hat{y}, \hat{z}$. The booklet is intended to make an introduction of the skills that would be useful in doing such sophisticated (but interesting) Physics problems, and also a directly use-able application of the vector calculus. The reader as assumed to have familiarity with basic calculus, its ideas, and simple Halliday Resnick or University Physics type Electrodynamics. Griffith and $Purcell\ and\ Morin\ will\ be\ used\ as\ reference.$

Our target is to make an arsenal of basic ideas through

with we can have intuition behind every problem that we see in Electrodynamics.

If you need more basics, then try out a few simple problems and exercises from the Sears and Zemansky's University Physics

1 General Electrics

We begin with a set of simple Irodov problems. It is assumed that the reader knows Calculus (basics), has idea on the general school textbook Electric Physics, and looking forward to have insight on the topic through problems. The document is not self contained and frequently refers to *University Physics*, by Sears and Zemasky and Physics - Halliday, Resnick, Krane. As an addition, we follow Electricity and Magnetism, by Purcell³ and Morin. The first part of the booklet tries to make the reader comfortable with the usual facts, in general, the first part 1 is not meant to be esoteric or technical enough, but gives a bit of the Taste.

Note that most of the medium and hard electrodynamics problems require the delicate handling of integration and usage of $\sin\theta,\cos\theta$. So during time pressure, it is important to do the math once and for all. Because reviewing the math again is going to consume enough amount of time, so it is wise to take time DUR-ING the math, not AFTER it. Refer to the appendix adjoined if needed.

Facts below inside the boxes are bookish theories. Ideas are creative tactics and some sophisticated facts, or even advice that prove helpful to use to solve a problem.

For completeness, we note the Coulomb's Law.

Fact 1. The Coulomb's Law and the Electric Field The force is a vector, making the separation of the *vector components*.

$$F = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} \hat{r}$$

$$F_x = F \cos \theta = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} \cos \theta$$

$$F_y = F \sin \theta = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} \sin \theta$$

But it is much good to take the $\frac{1}{4\pi\epsilon_0}$ as a mere k. That is much better and easy to deal with.

$$\mathbf{F} = k \frac{q_1 q_2}{r^2} \hat{r}$$

$$F_x = F \cos \theta = k \frac{q_1 q_2}{r^2} \cos \theta$$

$$F_y = F \sin \theta = k \frac{q_1 q_2}{r^2} \sin \theta$$

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²"Homogenous" is same as saying "Independent of Position"

³Nobel Laureate

We define the field as $\it The force per unit charge; one coulomb.$

 $oldsymbol{E} = rac{oldsymbol{F}}{q}$

Thus it stands up that,

$$E = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$$

$$E_x = E \cos \theta = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \cos \theta$$

$$E_y = E \sin \theta = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \sin \theta$$

The differential form is worth mentioning,

$$d\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} \hat{r}$$

Fact 2. The Newton Materials If the momentum is P, mass is m then

$$\boldsymbol{F} = m\boldsymbol{a} \tag{1}$$

$$F = \frac{dP}{dt} \tag{2}$$

Fact 3. The idea of Potential When there is energy associated with any electrical system, then the *Energy* per unit charge is the potential.

$$V = \frac{E}{q} \tag{3}$$

If work is done, by the intuition of the Energy and Work relation,

$$V = \frac{W}{q}$$
$$V = \frac{dW}{da}$$

Fact 4. The Gauss's Law Also the First *Maxwell's Equation*, describes the flux through a surface of the electric field. Flux is the streamlines going through a surface.

$$\int_{S} \mathbf{E} \cdot d\mathbf{a} = \frac{q}{\epsilon_0} \tag{4}$$

The technique is usually taking a small surface with some symmetry in it, then using the Gauss's law, without carrying any integration that would be required for the Coulomb one. Good to know that for a sphere the area,

$$A = 4\pi r^2$$

For a cylinder, area of the bases and the curved surface.

$$A = 2(\pi r^2) + (2\pi r)h$$

As to ensure the fluency through out, some simpler problems are given.

Pr 1. If a proton and an electron are released with a distance of $d = 2 \times 10^{-10} m$ apart (typical atomic distance), then find their initial acceleration.

Pr 2. Two charges are placed at a distance of d = 5.25m. The ratio of the charges are $\mu = \frac{5^2}{9^2}$. Where is the place between the line joining the two charges, if some other small amount of charge kept would feel no force?

Pr 3. So we have found the place where the force on small charge is zero as in the above problem. Now can you use the analogy of the impossibility of balancing a sharp pencil on its tip, **to defend the statement with physical logic**, that the charge in the "Zero Force place" cannot stay there forever? No need to use maths or something subtle. Refer to fact 11.

Pr 4. NASA planned to put some charge on the head of a rocket and same amount of charge in the launch pad. Given the rocket is Atlas V, whose height is d = ?? and the mass m = ??, what is the amount of charge Q to launch the rocket so that it flies off to the infinite space and never come back?

Pr 5. Four identical charges Q has been placed at the corners of a square of side L.

- a. In a free body diagram, show all of the forces that act on one of the charges. The diagram should be scaled for sake of visualization and the force vector lengths should be proportionate either.
- b. Find the magnitude and direction of the total force exerted on one of the charges by the other three.
- c. Should there be any acceleration or any motion?

We move on to more sophisticated ones.

Pr 6. Two small identical balls carrying the charges of the same sign are suspended from the same point by insulating threads of equal length. When the surrounding space was filled of Kerosene, the divergence angle between the threads remain constant. **What is the density of the material of which the balls are made?**

Idea 1. A whole lot of uninteresting thing can be excluded only by using *Ratios*. Like the initial value be $V_0 = \beta \times \varphi$ and the final value be $V = \beta \times \varphi'$. If we know that $\varphi = \frac{1}{x}$ and $\varphi' = \frac{1}{x'}$, then it's obviously,

$$\frac{V}{V_0} = \frac{\varphi}{\varphi'} = \frac{x'}{x}$$

And sometimes, some of the ratio can turn into trigonometric ratios. Like the ratio of the vertical y and horizontal x is $\tan \theta$.

Idea 2. If something stays constant, then the properties causing it shall keep it stay constant by equalizing (or balancing). And the derivative would be a zero.

Note that this idea gains importance if the "Property" is one in number. As a hint, the only thing causing the threads divergence angle to stay constant is the "Tension" of the string. So our idea will intuitively mean the property of the constant divergence angle α is the Tension force. The Tension of the string is caused by the Weight of the ball and the Electrical force. In kerosene, the tension force arises from the (Weight of the ball - Buoyant force from kerosene) and the Electric force inside the kerosene. And both shall cause the tension the equal. And finding the way to exploit the equal tension idea is left as the mystery of the problem. All is to make use of idea 2 and to use the trigonometric way of idea 1.

In vacuum, the constant $k = \frac{1}{4\pi\epsilon_0}$. In kerosene it will take some other $\epsilon_{kerosene}$ value that is the *Electric permittivity of Kerosene*. It's value would be in the data sheet. And to find the density, the target, use the simple reason why buoyancy should arise.

Pr 7. Two small equally charged spheres, each with a mass m, are suspended from the same point by silk threads of length l. The separation of the two sphere is x, and x << l. Find the rate of charge leak $\frac{dq}{dt}$ with which the charge leaks off sphere if their approach velocity varies by $v = \frac{a}{\sqrt{x}}$, while a is a constant.

To solve for the $\frac{dq}{dt}$, the idea lies in the tricky application of the way calculus can be used, it is that if it appears the we have an equation say somewhat like $q=\alpha x$, then little increment in both side is $\mathrm{d}q=\alpha\mathrm{d}x$ is we assume that α is a mere constant. Now dividing the both side by a $\mathrm{d}t$ yields us $\frac{dq}{dt}=\alpha\frac{dx}{dt}$. That gives us a velocity, and it's mathematically legal.

And the other thing is knowing the use of the approximation $\sin x \approx \tan x \approx x$. Then, if there is some function f(x), it's derivative can be written as $f'(x) = \varsigma f(x)$ in order to reduce something uninteresting.

Pr 8. A thin wire has electric charge q and and radius r. What would be the increment of the tension in the wire (the force of stretching) if a q_0 is introduced in the center of the wire?

The technique to the problem 8 is really easy. Draw the two tension forces placed at two points nearby on the string. The points makes small θ with the center. Now the vertical components of the tension should cancel out the electric forces. As $\theta << 1$, $\sin\theta \tilde{=}\theta$. That's how the equation should be found. Notice that the tension of the charged ring don't really effect the new tension increment. So you can work out the math totally forgetting about that. The ability of the problem weakens drastically if you have seen the way how in Halliday, Resnick, Krane the equation of the velocity of the wave on a string was derived.

Pr 9. A system of consists of a thin charged wire of radius R and there is a very long charged thread that is aligned along the axis of the ring. The charge consisted by the thread is q and the charge density of the thread

is λ . Find the *Forces of interaction* in the system to the thread.

Fact 5. Reference to the calculus way of finding the E field over a charged string.

Suppose we are r distance above an infinite string. If we need to carry out any integration so that the full string is covered, we rather than taking the integration limits from $-\infty$ to $+\infty$ over the x axis of the string, we turn the equations so that there is a trigonometric quantity, that is $r \tan \theta = x$, the θ made with the r line perpendicular to the string. Now, carrying out integration about $\theta = -\frac{\pi}{2}$ to $\theta = \frac{\pi}{2}$ is same as taking the infinite over x.

Using this the maths can be made easier to do.

The problem 9 is the technique of carrying out a Genius integration, you might see the appendix for the math.

Pr 10. Find the force on a charge Q that is h/2 height above an infinite sheet conductor.

Definitely the preceding ideas and facts have to be used.

Fact 6. Image Charge If any charge is above an infinite conductor by y height, then the charge causes opposite charges to come and accumulate over the surface. Thus a force towards the surface on the charge develops. The magnitude of the force is exactly same of the configuration of an opposite charge staying 2y distance away. As if, the conductor was a mirror, and the place where the charge's image is the position of another opposite charge attracting. An in depth analysis in appendix 2.6.

Idea 3. Mirrors in Electrostatic Force Laws of reflection definitely works for the solution of conductor-charge pair problems.

Fact 7. If the conductors of the "Image Charge" system is grounded, then during the vicinity (means "near") of the real charge, the same polarity of charge of the real charge goes away in the ground, same way as in the "Electroscope" charging process.

Pr 12. APhO-EuPHO BGdTST2019 There is a plastic solid sphere and it has a spherical cavity inside it a distance b from the center of the ball and the radius of the cavity is a that follows that a < b. The uniform charge density of the ball is ρ . Then **Find the Electric Field inside the cavity**.

Whenever there is any absurd hole in a system, then one way of attacking the system is to assume *there was no hole* and find the equation, then replacing some other parameter to return back to the initial system and solve for the target variable.

In this case in problem 12, the cavity (hole) can be assumed a superimposition of a positive area (like the solid sphere without any cavity) and a negative area. So the net charge is same as a void in that area, so meets the condition

Remark: When I was in the APhO-EuPhO Selection test, I tried to do the problem by using the Coulomb type formula and integrate. But simply it was impossibly difficult. And the cause of downfall then. This motivates the next idea

Idea 4. Avoid the calculation using extremely difficult equations and too long mathematics, there is definitely a simpler way, confirm if it is a competition.

Fact 8. Gradient For any function, the vector that we use to show the *Tangent* of any function is the Gradient (rough statement). Or simply, the Tangent vector itself is the Gradient of any function. Note that this function is a function of coordinates (x, y, z).

Fact 9. Electric Field as the "Gradient" of Potential Electric Field is the *Tangent to the rate of change of Potential with respect of Position*, that is, in Vector Calculus language, Electric Field is the *Gradient* of the *Potential*. A derivation has been shown in the Appendix.

$$\mathbf{E} = -\left(\frac{dV}{dx}\hat{x} + \frac{dV}{dy}\hat{y} + \frac{dV}{dz}\hat{z}\right) \tag{5}$$

In the Gradient form,

$$\mathbf{E} = -\nabla V \tag{6}$$

$$\mathbf{E} = -\operatorname{grad} V \tag{7}$$

Fact 10. Total work done on a unit charge to bring it close to any other charge Q to a distance r from infinity ∞ is the Potential Energy at the point r. The same is taking the charge from r to infinity ∞ , but it takes the opposite sign.

Fact 11. Something in equilibrium means that the total force on the particle is 0 and it tends to stay in that condition. Which means if a very small displacement δx is caused, then there should be some small δF force that brings the particle in its initial position. That is to zero out the displacement. IF that small force is not present, or it doesn't bring the particle into the initial position, then it should known that the Equilibrium is not stable. For the calculation, approximations as the Taylor Series is used.

Pr 13. A finite system of point charges has been made stable, they are in a stationary equilibrium because of there own electric fields. What is the electrostatic interaction energy of the system?

Use the following facts. Specially the 10 fact.

Pr 15. Find the electric field strength vector if the potential of the field is dependent on the x, y coordinates by

$$\Psi = a(x^2 + y^2)$$

$$\Psi = axy$$

Refer to the appendix for the method.

If we have a charge, and we keep it in an electric field, definitely we will have the displacement of the charge in the direction where the field acts in the position of the charge, given by $\ddot{x} = a = \mathbf{E}q/m$. It thus changes the potential, and we can call the points equipotential, if for points P_1 and P_2 , we have, $V(P_1) = V(P_2)$, where V us potential dependent on the position of Point. An electric field is always perpendicular to the Equipotential lines. So, if one takes a charge in a trajectory around an Equipotential, total work is zero.

Pr 17. There is a dipole, whose potential is given by,

$$\Phi(r,\theta) = kp \frac{\cos \theta}{r^2}$$

Where r is distance from the center of the dipole, assuming the dipole line makes along the \hat{y} axis, θ is the angle made by r with the vertical. Draw the equipotential line where r = 5m distance away is a point in the equipotential in x axis ($\theta = \pi/2$).

Idea 5. For having a coordinate equation of an equipotential, we should have some f(x) = y or $r(\theta)$ function, that defines positions as functions.

We know that potential of dipole is $\Phi = \Phi(r, \cos \theta)$. At some defined r_0, θ_0 , potential has value $\Phi(r_0, \cos \theta_0)$. Then this can be written as a function $r(\theta)$ that yields an equipotential function. Usage of this is the requirement of above problem. This, in fact, can be used for other type of potentials too.

Pr 18. A point dipole has a dipole moment p. It is oriented in the positive direction of the z axis, that is located at the origin of the coordinates. Find the projecion of the E_z and E_{\perp} of the electric field strength vector (on the plane perpendicular to the z axis at the point S). The point S makes an angle θ and is r distance from the point source.

Idea 6. Sometimes, the usage of Vectors and exploiting the definition of vectors are preferable.

⁴Adding more coordinates e_1, e_2, \cdots is the motivation of "Tensor".

Pr 19. Three small positively charged pearls lie one at each vertex of a triangle, their masses are m_1, m_2, m_3 . And respectively, their charges are Q_1, Q_2, Q_3 . When the pearls are released from rest, they move and follow a different straight line each. Neglect gravity. Their charge to mass ratio is given by

$$\frac{Q_1}{m_1}: \frac{Q_2}{m_2}: \frac{Q_3}{m_3} = 1:2:3$$

What special conditions should be imposed so that they follow their straight line? And what should be the angles of the Triangle?

Being an extremely useful problem, it requires bit more algebra, as a corresponding hint,

Fact 12. Vector Multiplication The cross product of two parallel vector is always 0. And the dot product is always the maximum.

Cross product is,

$$a \times b = r$$

The magnitude of the vector \mathbf{r} is,

$$r = a b \sin \theta$$

Where the θ is the angle between the vectors. And the direction of r vector points along the thumb direction by assuming curling our right hand finger from a to b. Very necessarily, we are interested on the perpendicular projection of one of the vectors, and its multiplication. Reference to Halliday, Resnick's Physics.

Dot product is,

$$\mathbf{a} \cdot \mathbf{b} = r$$
 scalar "r"

The magnitude is same,

$$ab \cos \theta = r$$

Now we have a scalar r and the projection of one vector on the other is now our interest. Cross product can be used to perpendicular things or constraining the directions or things.

Above in the problem, the cross product of the acceleration vector \boldsymbol{a} with the small displacement vector $\boldsymbol{\Delta}\boldsymbol{x}$ should be 0. And, because there is no internal force, the next idea can be implied.

Idea 7. Total force in a systems particle is $F_{\Sigma} = F_{external} + F_{internal}$. When the external force is absent, then by Newton's third law, the internal force, whatever the direction, shall conserve the point where the center of mass is. Hence if r_i is the position vector of an i-th particle with respect to the center of mass,

$$\sum m_i \boldsymbol{r_i} = 0$$

Same said that, for three bodies,

$$m_1 \mathbf{r_1} + m_2 \mathbf{r_2} + m_3 \mathbf{r_3} = 0$$

By taking the derivative, the idea can be extended further.

$$m_1 \mathbf{r_1} + m_2 \mathbf{r_2} + m_3 \mathbf{r_3} = 0$$

$$m_1 \mathbf{v_1} + m_2 \mathbf{v_2} + m_3 \mathbf{v_3} = 0$$

$$m_1 a_1 + m_2 a_2 + m_3 a_3 = 0$$

For the problem, one can find a few equations, using the r_i vector positions respect to the Center of Mass. As we need to concentrate on the Geometry of the triangle, the sides d_1, d_2, d_3 of the triangle are our main interest. This is visible that for any two particle, say particle 2 and 3, that $d_3 = r_3 - r_2$. And the forces (hence acceleration) of the particle are dependent on the sides of the triangle d_1, d_2, d_3 . From idea 7, a few equations are found out and then using the straight line conditions, a curl product with a required vector should be taken to reduce on side of the equation to be 0 and solve for sought answer.

Here we apply an important mathematical idea that is used in all of Physics.

Idea 8. It is customary to use equation and totally cancel out an unimportant variable, it makes the math simple to handle.

Idea 9. "As many variables, that many equations". That is, if we have three variables, but two equations, we cannot find the three variables each independently. It can only be done if we have three equations, not to.

A brief explanation of idea 9 has been made in appendix 2.5.

2 Appendix

2.1 The derivation for the Gradient of Potential to be the Electric Field

We can show that by the principle of work done on a test charge moving towards the vicinity of a bigger charge (both same type of charge; both positive or both negative) is,

$$W = F \times (-r)$$

Negative as we have to reduce the mutual distance in between the charges, and the force vector points opposite the way we are displacing. By a small increment,

$$dW = -F dr$$

Integrating from initial position r_1 to final position r_2 ,

$$W = \int dW = -\int_{r_1}^{r_2} F \, dr$$

Now, we know that V = W/q, so putting that, means dividing both side with q,

$$\frac{1}{q}W = \frac{1}{q} \int dW = -\frac{1}{q} \int_{r_1}^{r_2} F \, dr$$

This is thus trivially,

$$V = -\int_{r_1}^{r_2} E \, dr$$

More specifically, it can be easily shown that

$$V = -\int \boldsymbol{E} \cdot d\boldsymbol{r}$$

This is also useful in the both side differentiated form, that reaches to vector calculus introduction,

$$E_x = -\frac{dV}{dx}$$

Note this will only work if the E has no components except the x axis, otherwise the formula for the system is,

$$\boldsymbol{E} = -\left(\frac{dV}{dx}\hat{x} + \frac{dV}{dy}\hat{y} + \frac{dV}{dz}\hat{z}\right)$$

And Physicists isolated the common above to turn it to something simple and less time consuming, can't we just write the $\frac{d}{dr}\hat{r}$ type of things in the equation once and for all?

We had that,

$$\left(\frac{dV}{dx}\hat{x} + \frac{dV}{dy}\hat{y} + \frac{dV}{dz}\hat{z}\right)$$

If we keep reducing and introduce a symbol,

$$= \left(\frac{d}{dx}\hat{x} + \frac{d}{dy}\hat{y} + \frac{d}{dz}\hat{z}\right)V$$

Our new symbol should stand for the bracketed part above, let it be ∇ ,

$$\left(\frac{d}{dx}\hat{x} + \frac{d}{dy}\hat{y} + \frac{d}{dz}\hat{z}\right) = \nabla$$

Thus with our new symbol,

$$\left(\frac{d}{dx}\hat{x} + \frac{d}{dy}\hat{y} + \frac{d}{dz}\hat{z}\right)V = \nabla V$$

Our Electric Equation,

$$E = -\nabla V$$

Some also like to say it in a grand style, this has some similarities with a *Gradient*, so some also say it,

$$\mathbf{E} = -\operatorname{grad} V$$

This is really easy and very fun, we take three derivatives respect to x, y, z and have answers.

2.2 The Gradient of a Scalar Function

The math is more easy to understand through an example. 5

Pr 20. If a function $\Phi = axy^3 + xyz + x^2z$, then find the $\nabla \Phi$

Sol 1. I will show the method directly over here, same problems also included after it.

Let's do one by one, when differentiating about x, then y and z acts as a constant

$$\frac{d}{dx}\Phi = \frac{d}{dx}(axy^3 + xyz + x^2z) = ay^3 + yz + 2xz$$

$$\frac{d}{dy}\Phi = \frac{d}{dy}(axy^3 + xyz + x^2z) = 3axy^2 + xz + x^2z$$

$$\frac{d}{dz}\Phi = \frac{d}{dz}(axy^3 + xyz + x^2z) = axy^3 + xy + x^2$$

So, all together, as we know,

$$\left(\frac{d\Phi}{dx}\hat{x} + \frac{d\Phi}{dy}\hat{y} + \frac{d\Phi}{dz}\hat{z}\right) = \nabla\Phi$$

Our sought answer is,

$$\nabla \Phi$$
 =

$$(ay^3+yz+2xz)\hat{x} + (3axy^2+xz+x^2z)\hat{y} + (axy^3+xy+x^2)\hat{z}$$

Pr 21. Keep finding the gradient, or $\nabla \Phi$

- 1. $\Phi = x^4$
- 2. $\Phi = xyz^2$
- 3. $\Phi = x + y + z$
- 4. $\Phi = ax^2 + by^4$
- 5. $\Phi = x \frac{y}{x^3}$
- 6. $\Phi = \frac{xy}{z^3} + \frac{abx^3}{y-z}$
- 7. $\Phi = a \frac{45x}{6z^3}$
- 8. $\Phi = e^x + ye^z$

2.3 Approximation method

Let a function be $f(x)=x^2$. Then, putting any small number δ in the function should give nearly zero. As an example, the function is to take a square, so if δ is assumed to equal 0.001, its square, $\delta^2=0.000001$. Because, $\delta^2=\frac{1}{1000^2}=\frac{1}{1000000}=0.000001$. The small the δ , the small its exponent is gonna be. So if we put a small "displacement" δ along x in the function, what we will have is,

$$f(x) = x^2$$

$$f(x+\delta) = x^2 + 2x\delta + \delta^2$$

From our previous learning, $\delta << 1$, δ^2 is too small to use. We ignore the δ^2 in the equation, so,

$$f(x+\delta) = x^2 + 2x\delta$$

To understand how much efficient this is, let x = 5 and $\delta = 0.01$ (not too small!), so directly,

$$(5+0.01)^2 = 25.1001$$

Using the approximated system,

$$(5+0.01)^2 = 5^2 + (2 \times 5 \times 0.01) = 25.1$$

Of course, rounding off the value, both are same, 25.1001=25.1. So, the approximation works.

For all kinds of general case, let $f(x) = x^n$. For defining any small δ displacement, by the fundamental idea of calculus (derivatives),

$$m = \frac{\Delta y}{\Delta x} = \frac{\Delta f(x)}{\Delta x}$$

Through ideas of limits and other things, in Leibnitz way, taking $m = \frac{df(x)}{dx}$,

$$\frac{df(x)}{dx} = \frac{f(x+h) - f(x)}{(x+\delta) - x} = \frac{f(x+h) - f(x)}{\delta}$$

The δ be taken left side and so as the f(x) part, yields,

$$f(x+\delta) = f(x) + \delta \frac{df(x)}{dx}$$
 (8)

As a proof, for the previous numerical example, $f(x) = x^2$, so $\frac{df(x)}{dx} = 2x$. Plug it in the equation 8,

$$f(x+\delta) = (x+\delta)^2 = x^2 + (2x)\delta$$

Like so, any function having small displacement can be approximated.

$$f(x+\delta) = f(x) + \delta \frac{df(x)}{dx} \tag{9}$$

⁵"The best way to *learn* Mathematics is *doing* it

2.4 Taylor Series

The idea of the Taylor series is to make a function into some power series, **not too necessarily for Approximation**. The series is

$$f(x) = f(a) + f'(a)(x - a) + \frac{1}{2!}f''(a)(x - a)^{2} + \frac{1}{3!}f'''(a)(x - a)^{3} + \cdots$$
 (10)

NOT TO BE CONFUSED WITH THE APPROXIMATION SERIES. The a is any number, most likely 1 or 0. $f''(a) = \frac{d^2 f(a)}{dx}$ in short and so on. As a part of example, why we need it, let's think what can it do to $\sin x$.

$$a = 0$$

Now,

$$f(a) = \sin 0 = 0$$

$$f'(a) = \cos 0 = 1$$

$$f''(a) = -\sin 0 = 0$$

$$f'''(a) = -\cos 0 = -1$$

$$f^{(4)}(a) = \sin 0 = 0$$

$$f^{(5)}(a) = \cos 0 = 1$$
(11)

So put these where they belong and yields,

$$f(x) = \sin x = x - \frac{x^3}{3!} + \frac{x^5}{5!} + \cdots$$
 (12)

See that if the x << 1 then the Taylor series can also be used to approximate.

We can keep showing,

$$\cos x = 1 - \frac{x^2}{2!} + \frac{x^4}{4!} \tag{13}$$

$$\sqrt{1+x} = 1 + \frac{x}{2} - \frac{x^2}{8} + \dots \tag{14}$$

$$\frac{1}{\sqrt{1+x}} = 1 - \frac{x}{2} + \frac{3x^2}{8} + \dots \tag{15}$$

Pr 22. Keep finding the Taylor series,

- $\bullet \quad \frac{1}{1-x}$
- $\bullet \quad \frac{1}{\sqrt{1+x^2}}$
- $\bullet \quad \frac{1}{(1-x)^2}$
- ln(1-x)
- e^x

Pr 23. What would be the approximations if x is small in a way that till x^3 cannot be *ignored* but later than x^3 can be?⁶

- \bullet $\frac{1}{1-\tau}$
- $\bullet \quad \frac{1}{\sqrt{1+x^2}}$
- $\bullet \quad \frac{1}{(1-x)^2}$
- ln(1-x)
- e^x

2.5 "As many variables, that many equations"

Suppose we want to solve x, y, z from the set of equations, so number of variable $n_{variable} = 3$ but number of equations $n_{equation} = 2$.

$$x + y + z = 6$$
$$x - y - z = -4$$

Now can it be solved? No. It can be tried. We can write 1st equation putting x in one side and put this in 2nd equation and get y + z = 5 and x = 1. But the thing remains unfinished, as y and z can be (1, 2, 3, 4) and (4, 3, 2, 1).

If we supply another equation, making $n_{equation} = 2 + 1 = 3$, all together,

$$x + y + z = 6$$
$$x - y - z = -4$$
$$x + y + 2z = 9$$

Now the solution is well guessed. (x, y, z) = (1, 2, 3).

This means that, for any equation system to be solved,

$$n_{equation} = n_{variable}$$
 (16)

2.6 Image Charges

NB: NOT FINAL AND DRAFT, REQUIRES TECHNICAL REVISION

What would be the Magnitude and Direction of the Force on a Positive Electric Charge "+Q" that is "h" distance above a sheet of an infinite sized conductor?

In this case, remember that the charge we have above the conductor is Positive +Q. Thus a lot of electrons move toward part of the conductor that is near the charge. A negative charge made field would be made on the conductor sheet that pulls the Positive charge above towards it, so the direction of force would be towards the conductor. This was easy, but the fun is in finding out the magnitude. A net positive charge will be away the part of the vicinity (vicinity means nearby) of charge +Q. Making the calculation of the net accumulated charge (electron) near the vicinity of the +Q above the conductor is going to be a very difficult task. Because, if we assume some -q charge in the form of circular disk is made on vicinity of the conductor, this

⁶Hint: x^4 is too much small, but x^3 isn't.

-q charge will resist any further charge on the conductor to be near, but the charge above +Q stronger, more charges in the conductor will join the circular disk. Carrying out the math of the circular disk would be interrupted as we hardly can find how many charge has been gathered, also given the positive set of charge in conductor far away the vicinity will push the +Q a bit. A big mathematical mess.

But none of this assumption has to be done if we leave this way of facing the problem. We know that if a conductor is near a charge, the field lines entering or exiting the conductor has to be perpendicular near the surface, this leaves us with a system that looks like the figure. Remember charge is induced on the surface of the conductor. And tactfully applying Gauss law, what comes is that the system is totally similar that a -Q charge 2h distance away is attracting the +Q charge. The electric field lines of the +Q enters perpendicularly into the conductor, and to obey the Gauss Law, that the net charge in a small part of the conductor should be zero. Thus somewhat a new field line emerge that seems to exit the conductor through the other side. That new alien charge that we call is Image Charge.

The force has to be

$$F = \frac{1}{4\pi\epsilon_0} \frac{Q^2}{(2h)^2} \tag{17}$$

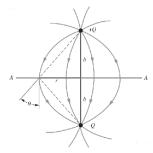


Figure 1: The image charge is the Q below, and it is nonexistent. Note that the field line through conductor enters and gets out perpendicularly.

In a simple summary, my point is that for any system, where there is some charge near by a conductor, the force it faces is same as an opposite charge is attracting it, from a position where the virtual image of the charge would have formed if the conductor was a Mirror. Applying Fermat's Least Time Principle, the math can be taken to the next level.

2.7 Introducing "Curl" and Analogy of Electric Field as "Fluid Flow"

NB: NOT FINAL AND DRAFT, REQUIRES TECHNICAL REVISION

The Maxwell's Equations are a wonderful set of formulas that describes Electromagnetism with excellent amount of brilliance. The Faradays law of the Maxwell's Equation in differential form is written as,

$$\operatorname{curl} \boldsymbol{E} = -\frac{\partial \boldsymbol{B}}{\partial t} \qquad \nabla \times \boldsymbol{E} = -\frac{\partial \boldsymbol{B}}{\partial t}$$
 (18)

Where the triangle is called the nabla and is extremely interesting to work with. It would be a whole lot of writing worth a book to explain its meaning. But if interested, give a read to the great Vector Calculus book "Div Grad Curl and All that" by Schey. Nice book on the topic.

Now I don't want to write an article where the reader needs to know vector calculus beforehand. Let's see the thing with models and diagrams in a more basic manner, making an understanding what the math really tries to make the sense of. Imagine a water flow in a river and the water is composed of particles (of course the particle here is a molecule). The flow is constant, the molecules have mass, doesn't awkwardly change and the water is incompressible. We imagine that we know the velocity of every particle, and know which direction it points to; simply that we know about the velocity vector of every particle. Multiply the mass and we know the momentum of all particle. Now this momentum vectors of the particles won't be chaotic or unpredictable, because it is in a flow, the momentum vectors will usually point along the flow. We can describe a function that can tell the momentum of the particle in a given location. The momentum vector of particles can be similarly understood using a general formula. Simply, we input that the position is 2 m from the bank of the river and can get the momentum 6kmkg/h and 0.3m from the bank 2kmkg/h (just an example). The viscous drag or the friction of the particles with the wall (bank of the river) reduces the velocity. Let we call the special function that gives the momentum of a particle at a given point as V, that mathematically is a field function.

There is no vortex or whirlpool (that tornado like spiral flow) in the flow. That roughly means that the rotation of the flow at any point is zero. The momentum vectors don't point as if they are circling. This all in mathematical formalism is said by

$$curl \mathbf{V} = 0 \tag{19}$$

Given that for any point in the flow. This is thus can a mean any rotation in a field. The formulas to do any math related to curl as above is a part of the vector calculus. The curl value is positive when the rotation is same as the right hand rule.

Now let us see another interesting case. What if we take the rate of change of V along the coordinates? Let's assume for now that see the river from above and only are interested how the upper layer particles move. The rate of change of V along the projection on the coordi-

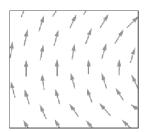


Figure 2: The curl is non zero at a little rightwards point outside the image. It is where the center of the rotation is. Curl of any point outside the center is zero.

nates (I mean position) x and y is

$$\frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} \tag{20}$$

Now what if the value of this above equation solves value more than 0? That means there is some increase of the momentum. How will the momentum increase? If we put a water hose and start it so that it adds water in the flow, hence increasing it in mass a certain point (where the mouth of the hose is). So if it happens that if the value is negative then what can it mean? It means that someone is taking away water from that point. Mathematically, the thing is called the Divergence of the flow, simply

$$div \mathbf{V} = \nabla \cdot \mathbf{V} = \frac{\partial \mathbf{V_x}}{\partial x} + \frac{\partial \mathbf{V_y}}{\partial y}$$
 (21)

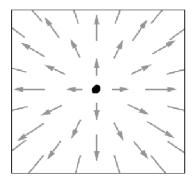


Figure 3: The center black part has the attribute of source, that's where the div is positive. Because it feels as if the flow emerging from the black dot at the center.

That might be any value. A source will yield a positive div, a sink will do negative, and just the flow shall do 0. Intuitively, the div of the flow of water from a tap is zero, except the mouth of the tap from where the water is coming. There the div is more than zero. The more strongly the water comes the tap, more value the div V for that position. It makes sense, more the velocity of the particles then the more momentum, so

as if the mouth is wide, then more mass of water shall come. The sink is where the water reduces, hence the div is negative. And yes, curl is mostly seen as the water draining away in a vortex as you flush a commode, worth trying! All of it is a formal way to calculate fields that take similar to flow.

For those who look forward out of curiosity what nabla means, give a look to the tiny equation below. Phi is scalar that depends on the coordinate. And the V is some vector field, same as we described for water flow.

$$\nabla = \frac{d}{dx}\hat{x} + \frac{d}{dy}\hat{y} + \frac{d}{dz}\hat{z} \tag{22}$$

$$div V = \left(\frac{d}{dx}\hat{x} + \frac{d}{dy}\hat{y} + \frac{d}{dz}\hat{z}\right) \cdot (V_x\hat{x} + V_y\hat{y} + V_z\hat{z}) = \frac{dV}{dx}\hat{x} + \frac{dV}{dy}\hat{y} + \frac{dV}{dz}\hat{z}$$
(23)

Okay, so we have an idea what a curl and a divergence can do. Curl makes a numerical value of rotation of the field and the Divergence makes a numerical value on the magnitude of how much extra flow is contributed by a source or the opposite sink. Have you guessed why we had entered hydrodynamics instead of doing electromagnetism? Yes, to understand Maxwell's Equation, but it seems that the Maxwell's Equation in differential form all take the Hydrodynamic identity, see the four Maxwell's Equations,

$$div \mathbf{E} = \frac{\rho}{\epsilon_0} \tag{24}$$

$$div \mathbf{B} = 0 \tag{25}$$

$$curl \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$
 (26)

$$curl \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$
 (27)

The Rho is the charge density, J is the current density vector (current per unit cross section area), Mu-not is the permeability constant of free space and c is well known Leightgeschwindigkeit (speed of light, as in German by Einstein). Now doesn't the Electric Vector Field and the Magnetic Vector Field seems interesting? If not, then see that the source of Electric field is dependent on the charge density. What is that? That can nicely be seen when you imagine drawing the field lines of two opposite charges, from positive it emerge and get inside negative charge. This is same as water flow, there is a big plate, and in point, there is a water pipe vertically putting water in the plate. And near is a hole from where the water should sink, and the water flow takes the shape of the electric field lines when we draw a negative charge near a positive.

So this extremely interesting that mathematically, Electric Fields are same a water flow! They can vortex out like water, they can flow like water. And some

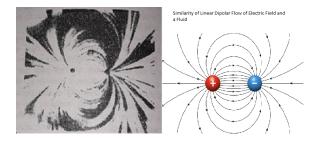


Figure 4: The similarity of the two kind of "flow" and the rough black and white Image is from the HRK which is the image of a Linear Dipolar Water flow.

math proves that Light, being an Electromagnetic Wave is actually like a Water Wave in the sea of Electric Field! Almost same to the Magnetic Field, the div of it is zero, because there is no isolated Magnetic Pole, a north pole always has a south pole, different from the charge. So net div of north pole is zeroed by the south pole. But there is no limit in rotating, so that rotation of the magnetic field depends on the current density (remember that magnetic field is always accompanied with a Current, by Ampere's Law) and the rate of change of Electric Field.

This parallelism of the fields is really astonishing because the general idea of the electromagnetic fields don't assume any Hydrodynamic things. It only makes an idea of Flux and from it comes the analogy of water flow with that. But will it help in solving severe problems in the competition? Definitely, Yes!

See the last year Physics Cup problem 2019, Prof. Jaan Kalda had given hint to it. The problem is given.

When a body moves in liquid, the motion of the body puts the liquid into motion, too. The motion of liquid contributes to the total kinetic energy of the system, and hence, leads to an increased effective mass of the body. The difference of the effective mass and the actual mass of the body is referred to as the added mass. The added mass depends on the size and shape of the body. Consider a certain metallic body of volume V and polarizability a along its symmetry axis x (i.e. homogeneous externally applied electric field E induces total dipole moment p = aE on this body). Additionally, the body shape is such that if it were made from a homogeneous dielectric material and put into homogeneous electric field, the electric field inside the body would be also homogeneous. Find the added mass of this body when it starts moving transnationally, parallel to the x-axis, in an incompressible initially motionless liquid of density rho. (Best solutions can be found in the Physics Cup 2019 website)

The technique is knowing the boundary conditions and writing the Maxwell's Equations in the hydrodynamic analogy of the problem text. And doing some very basis vector calculus finishes the job. A way of finding equation is that the vortex absence makes the curl of the fluid field be zero. It is quite exaggeration for people like us, but this is a great problem that ac-

tually makes the sense of how awesome Physics can be against Biology and Chemistry.

Electrodynamics thus is two form like, one place is of the realm of the equations of Coulombs Law and simple University Physics (Sears and Zemansky) topic; the other is place of the mathematics that takes Electrodynamics same as a fluid. That's where the Laplacian and Poisson's Equation takes over.

$$\nabla^2 V = 0 \qquad \text{(Laplacian Equation)} \tag{28}$$

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