

Quantum Mechanics : : Homework 04

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Problem 01

a

I do the matrix multiplication by hand.

$$\begin{aligned}\hat{\sigma}_1 \hat{\sigma}_2 &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} = i\hat{\sigma}_3 \\ \hat{\sigma}_2 \hat{\sigma}_1 &= \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} -i & 0 \\ 0 & i \end{pmatrix} = -i\hat{\sigma}_3 \\ \hat{\sigma}_2 \hat{\sigma}_3 &= \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} = i\hat{\sigma}_1 \\ \hat{\sigma}_3 \hat{\sigma}_2 &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix} = -i\hat{\sigma}_1 \\ \hat{\sigma}_3 \hat{\sigma}_1 &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = i\hat{\sigma}_2 \\ \hat{\sigma}_1 \hat{\sigma}_3 &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = -i\hat{\sigma}_2\end{aligned}$$

The cross product

$$\begin{aligned}\hat{\vec{\sigma}} \times \hat{\vec{\sigma}} &= \begin{vmatrix} \vec{n}_x & \vec{n}_y & \vec{n}_z \\ \sigma_1 & \sigma_2 & \sigma_3 \\ \sigma_1 & \sigma_2 & \sigma_3 \end{vmatrix} \\ &= \vec{n}_x (\sigma_2 \sigma_3 - \sigma_3 \sigma_2) + \vec{n}_y (\sigma_3 \sigma_1 - \sigma_1 \sigma_3) + \vec{n}_z (\sigma_1 \sigma_2 - \sigma_2 \sigma_1) \\ &= \vec{n}_x (2i\sigma_1) + \vec{n}_y (2i\sigma_2) + \vec{n}_z (2i\sigma_3) \\ &= 2i (\sigma_1 \vec{n}_x + \sigma_2 \vec{n}_y + \sigma_3 \vec{n}_z) \\ &= 2i\hat{\vec{\sigma}}\end{aligned}$$

Note this also is alluring to the Levi-Civita Symbol because

$$\text{cyclic } i \rightarrow j \rightarrow k \rightarrow i \rightarrow j \rightarrow k \implies \sigma_i \sigma_j = i\epsilon_{ij} \sigma_k \quad (\text{where } \{i, j, k\} \in \{1, 2, 3\})$$

b

$$(\vec{U} \cdot \hat{\vec{\sigma}}) (\vec{V} \cdot \hat{\vec{\sigma}}) = (U_1 \sigma_1 + U_2 \sigma_2 + U_3 \sigma_3) (V_1 \sigma_1 + V_2 \sigma_2 + V_3 \sigma_3)$$

$$\begin{aligned}
(\vec{U} \cdot \hat{\vec{\sigma}}) (\vec{V} \cdot \hat{\vec{\sigma}}) &= \left(\sum_{i=1}^3 U_i \sigma_i \right) \left(\sum_{j=1}^3 V_j \sigma_j \right) \\
&= \sum_{i,j=1}^3 U_i V_j \sigma_i \sigma_j \\
&= \sum_{i,j=1, i \neq j}^3 U_i V_j \sigma_i \sigma_j + \sum_{n=1}^3 U_n V_n \sigma_n \sigma_n \\
&= \sum_{i,j=1, i \neq j}^3 i U_i V_j \epsilon_{ij} \sigma_k + \sum_{n=1}^3 U_n V_n \sigma_n \sigma_n \\
&= [i(U_1 V_2 - U_2 V_1) \sigma_3 + i(U_2 V_3 - U_3 V_2) \sigma_1 + i(U_3 V_1 - U_1 V_3) \sigma_2] + (U_1 V_1 + U_2 V_2 + U_3 V_3) \hat{I} \\
&= [\vec{U} \times \vec{V}]_3 i \sigma_3 + [\vec{U} \times \vec{V}]_1 i \sigma_1 + [\vec{U} \times \vec{V}]_2 i \sigma_2 + (\vec{U} \cdot \vec{V}) \hat{I} \\
&= i [\vec{U} \times \vec{V}] \cdot \hat{\vec{\sigma}} + (\vec{U} \cdot \vec{V}) \hat{I}
\end{aligned}$$

Problem 02

a

Considering the simplest basis in column vector forms

$$|1\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad |2\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad |3\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

We want to compute the following operator \hat{H}

$$\hat{H} = E_0 (|1\rangle\langle 1| + |2\rangle\langle 2| + |3\rangle\langle 3|) - J (|1\rangle\langle 2| + |2\rangle\langle 1| + |2\rangle\langle 3| + |3\rangle\langle 2| + |3\rangle\langle 1| + |1\rangle\langle 3|)$$

Matrices for the first term

$$\begin{aligned} |1\rangle\langle 1| + |2\rangle\langle 2| + |3\rangle\langle 3| &= \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix} \begin{bmatrix} 0 & 1 & 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \begin{bmatrix} 0 & 0 & 1 \end{bmatrix} \\ &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} = \hat{I} \end{aligned}$$

Matrices for the second term

$$\begin{aligned} &|1\rangle\langle 2| + |2\rangle\langle 1| + |2\rangle\langle 3| + |3\rangle\langle 2| + |3\rangle\langle 1| + |1\rangle\langle 3| \\ &= \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \begin{bmatrix} 0 & 1 & 0 \end{bmatrix} + \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \begin{bmatrix} 1 & 0 & 0 \end{bmatrix} + \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \begin{bmatrix} 0 & 0 & 1 \end{bmatrix} + \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \begin{bmatrix} 0 & 1 & 0 \end{bmatrix} + \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \begin{bmatrix} 1 & 0 & 0 \end{bmatrix} + \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \begin{bmatrix} 0 & 0 & 1 \end{bmatrix} \\ &= \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 1 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \\ &= \begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix} \end{aligned}$$

Put them all together

$$\begin{aligned} \hat{H} &= E_0 (|1\rangle\langle 1| + |2\rangle\langle 2| + |3\rangle\langle 3|) - J (|1\rangle\langle 2| + |2\rangle\langle 1| + |2\rangle\langle 3| + |3\rangle\langle 2| + |3\rangle\langle 1| + |1\rangle\langle 3|) \\ &= \begin{bmatrix} E_0 & -J & -J \\ -J & E_0 & -J \\ -J & -J & E_0 \end{bmatrix} \end{aligned}$$

b (i)

The $|E_1\rangle$ in column vector representation

$$|E_1\rangle = \frac{1}{\sqrt{3}} (|1\rangle + |2\rangle + |3\rangle) = \frac{1}{\sqrt{3}} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

Computing $\hat{H}|E_1\rangle$

$$\begin{aligned}
\hat{H}|E_1\rangle &= \begin{bmatrix} E_0 & -J & -J \\ -J & E_0 & -J \\ -J & -J & E_0 \end{bmatrix} \left(\frac{1}{\sqrt{3}} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} \right) \\
&= \frac{1}{\sqrt{3}} \begin{bmatrix} E_0 - J - J \\ -J + E_0 - J \\ -J - J + E_0 \end{bmatrix} \\
&= \frac{1}{\sqrt{3}} \begin{bmatrix} E_0 - 2J \\ E_0 - 2J \\ E_0 - 2J \end{bmatrix} \\
&= \frac{E_0 - 2J}{\sqrt{3}} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} \\
&= (E_0 - 2J)|E_1\rangle
\end{aligned}$$

We see the eigen-equation for the Hamiltonian

$$\hat{H}|E_1\rangle = E_1|E_1\rangle \implies E_1 = E_0 - 2J$$

The special property can be realized by seeing that,

$$\begin{bmatrix} E_0 & -J & -J \\ -J & E_0 & -J \\ -J & -J & E_0 \end{bmatrix} = (E_0 + J) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} + (-J) \begin{bmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{bmatrix}$$

It's very obvious that $|E_1\rangle$ is an eigenvector for both of the upper matrices. And if $\hat{H} = \hat{A} + \hat{B}$ then if $|E_1\rangle$ is eigenvector for both \hat{A}, \hat{B} , then

$$\hat{H}|E_1\rangle = \hat{A}|E_1\rangle + \hat{B}|E_1\rangle = (a + b)|E_1\rangle = E_1|E_1\rangle$$

Just from looking at the matrices it's obvious that $a = E_0 + J$ and $b = -3J$ so $E_1 = E_0 - 2J$

b (ii) and (iii)

We algebraically initialize the elements of the matrix to minimize the computational effort by hand. Keep in mind that $\hat{H} - E\hat{I} = \hat{0} = \frac{1}{J}\hat{H} - \frac{1}{J}E\hat{I}$ hence

$$\begin{aligned}
\frac{1}{J}\hat{H} &= \begin{bmatrix} \frac{E_0}{J} & -1 & -1 \\ -1 & \frac{E_0}{J} & -1 \\ -1 & -1 & \frac{E_0}{J} \end{bmatrix} \\
\left(\frac{1}{J}\hat{H} \right) - \frac{E}{J}\hat{I} &= \begin{bmatrix} \frac{E_0-E}{J} & -1 & -1 \\ -1 & \frac{E_0-E}{J} & -1 \\ -1 & -1 & \frac{E_0-E}{J} \end{bmatrix} \\
&= (-1) \begin{bmatrix} \lambda & 1 & 1 \\ 1 & \lambda & 1 \\ 1 & 1 & \lambda \end{bmatrix} = 0 \quad (\lambda = (E - E_0)/J) \\
\implies \det \begin{bmatrix} \lambda & 1 & 1 \\ 1 & \lambda & 1 \\ 1 & 1 & \lambda \end{bmatrix} &= 0 \\
\implies \lambda^3 + 3\lambda + 2 &= 0 \\
\implies (\lambda - 1)^2(\lambda + 2) &= 0 \\
\implies (\lambda^2 - 2\lambda + 1)(\lambda + 2) &= 0 \\
\implies \lambda = 1, 1, -2 &
\end{aligned}$$

As stated in the problem statement, $0 = (\lambda - \lambda_1)(\lambda^2 + b\lambda + c) \implies 0 = (\lambda - (-2))(\lambda^2 + (-2)\lambda + 1)$ giving us $(b, c) = (-2, 1)$

For $\lambda = -2$ the eigen-energy is $\frac{E - E_0}{J} = -2 \implies E = E_0 - 2J$

For $\lambda = 1$ the eigen-energy is $\frac{E - E_0}{J} = 1 \implies E = E_0 + J$

The eigenvalues are

$$\lambda_1 = -2, \quad \lambda_2 = 1, \quad \lambda_3 = 1$$

And corresponding eigen-energy

$$E_1 = E_0 - 2J, \quad E_2 = E_0 + J, \quad E_3 = E_0 + J$$

Problem 03

a

$$\begin{aligned}
 \hat{T}|n\rangle = |n+1\rangle &\implies \langle n|T^\dagger = \langle n+1| \\
 &\text{now, } \langle n|T^\dagger|n+1\rangle = \langle n+1||n+1\rangle = 1 \\
 &\text{or, } \langle n|T^\dagger T|n\rangle = 1 \\
 &\text{and, } \langle n|T^\dagger|k+1\rangle = \langle n+1||k+1\rangle = 0 \quad (n \neq k) \\
 &\implies T^\dagger T = \hat{I}
 \end{aligned}$$

b

To compute $[\hat{H}, \hat{T}]$ we compute the two operator terms individually

$$\begin{aligned}
 \hat{H}\hat{T} &= \left(\sum_{n=-\infty}^{\infty} [E_0|n\rangle\langle n| + J(|n+1\rangle\langle n| + |n\rangle\langle n+1|)] \right) \hat{T} \\
 \hat{H}\hat{T}|n\rangle &= \left(\sum_{n=-\infty}^{\infty} [E_0|n\rangle\langle n| + J(|n+1\rangle\langle n| + |n\rangle\langle n+1|)] \right) |n+1\rangle \\
 &= [E_0|n\rangle\langle n| + J(|n+1\rangle\langle n| + |n\rangle\langle n+1|)] |n+1\rangle + [E_0|n+1\rangle\langle n+1| + J(|n+2\rangle\langle n+1| + |n+1\rangle\langle n+2|)] |n+1\rangle \\
 &= J|n\rangle + E_0|n+1\rangle + J(|n+2\rangle) \\
 \hat{T}\hat{H}|n\rangle &= \hat{T} \left([E_0|n\rangle\langle n| + J(|n+1\rangle\langle n| + |n\rangle\langle n+1|)] |n\rangle + [E_0|n-1\rangle\langle n-1| + J(|n\rangle\langle n-1| + |n-1\rangle\langle n|)] |n\rangle \right) \\
 &= \hat{T}(E_0|n\rangle + J|n+1\rangle + J|n-1\rangle) \\
 &= E_0|n+1\rangle + J|n+2\rangle + J|n\rangle = \hat{H}\hat{T}|n\rangle
 \end{aligned}$$

$$\therefore (\hat{H}\hat{T} - \hat{T}\hat{H})|n\rangle = |0\rangle \implies [\hat{H}, \hat{T}] = 0$$

C

Using the form given for the energy, and the eigen-equation for \hat{T} , we determine the general formula for $\psi_{E,n}$ in terms of $\psi_{E,0}$

$$\begin{aligned}
|E\rangle &= \sum_{n=-\infty}^{\infty} |n\rangle \psi_{E,n} \\
T|E\rangle &= \sum_{n=-\infty}^{\infty} T|n\rangle \psi_{E,n} \\
e^{-i\phi}|E\rangle &= \sum_{n=-\infty}^{\infty} |n+1\rangle \psi_{E,n} \\
\sum_{n=-\infty}^{\infty} e^{-i\phi}|n\rangle \psi_{E,n} &= \sum_{n=-\infty}^{\infty} |n+1\rangle \psi_{E,n} \\
\Rightarrow e^{-i\phi}|n+1\rangle \psi_{E,n+1} &= |n+1\rangle \psi_{E,n} \\
\therefore e^{-i\phi} &= \frac{\psi_{E,n}}{\psi_{E,n+1}}
\end{aligned}$$

The inductive relation between two coefficient is

$$\begin{aligned}
\psi_{E,n+1} = e^{i\phi} \psi_{E,n} &\Rightarrow \psi_{E,1} = e^{i\phi} \psi_{E,0} \Rightarrow \psi_{E,2} = e^{i\phi} \psi_{E,1} = e^{2i\phi} \psi_{E,0} \\
&\Rightarrow \psi_{E,n} = e^{in\phi} \psi_{E,0}
\end{aligned}$$

The general form of energy is then

$$|E\rangle = \psi_{E,0} \sum_{n=-\infty}^{\infty} e^{in\phi} |n\rangle$$

$$\begin{aligned}
\hat{H}|E_n\rangle &= \psi_{E,n} \hat{H}|n\rangle = \psi_{E,n} (E_0|n\rangle + J|n+1\rangle + J|n-1\rangle) \\
\hat{H}|E\rangle &= \hat{H} \left(\sum_{n=-\infty}^{\infty} \psi_{E,n} |n\rangle \right) = \sum_{n=-\infty}^{\infty} \psi_{E,n} (E_0|n\rangle + J|n+1\rangle + J|n-1\rangle) \\
\langle n'|\hat{H}|E\rangle &= \sum_{n=-\infty}^{\infty} \psi_{E,n} (E_0\langle n'|n\rangle + J\langle n'|n+1\rangle + J\langle n'|n-1\rangle) \\
&= E_0\psi_{E,n'} + J\psi_{E,n'-1} + J\psi_{E,n'+1} \\
&= E_0\psi_{E,n'} + J\frac{\psi_{E,n'}}{e^{i\phi}} + J e^{i\phi} \psi_{E,n'} \\
&= \psi_{E,n'} (E_0 + J e^{-i\phi} + J e^{i\phi}) \\
&= \psi_{E,n'} (E_0 + 2J \cos(\phi))
\end{aligned}$$

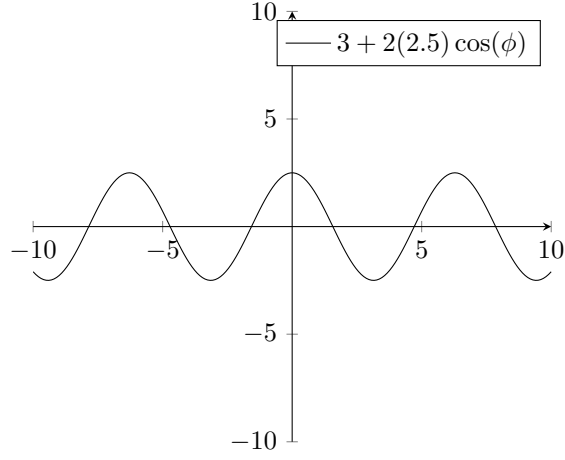


Figure 1: Simple pgfplot for aesthetic purposes $E = 3, J = 2.5$

$$[E_0 + 2J \cos(\phi)]_{\max} = E_0 + 2J$$

$$[E_0 + 2J \cos(\phi)]_{\min} = E_0 - 2J$$

Problem 04

$$\begin{aligned}
\langle 1 | \hat{X} | 1 \rangle &= \langle 1 | \hat{I} \hat{X} \hat{I} | 1 \rangle \\
&= \langle 1 | \left(\int_0^L dx |x\rangle \langle x| \right) \hat{X} \left(\int_0^L dx' |x'\rangle \langle x'| \right) | 1 \rangle \\
&= \left(\int_0^L dx \langle 1 | x \rangle \langle x | \right) \hat{X} \left(\int_0^L dx' |x'\rangle \langle x' | 1 \rangle \right) \\
&= \int_0^L dx \int_0^L dx' \langle 1 | x \rangle \langle x' | 1 \rangle \langle x | \hat{X} | x' \rangle \\
&= \int_0^L dx \int_0^L dx' \langle 1 | x \rangle \langle x' | 1 \rangle x' \delta(x - x') \\
&= \int_0^L dx \int_0^L dx' \left(\sqrt{\frac{2}{L}} \sin\left(\frac{\pi}{L}x\right) \right) \left(\sqrt{\frac{2}{L}} \sin\left(\frac{\pi}{L}x'\right) \right) x' \delta(x - x') \\
&= \int_0^L dx' \left(\sqrt{\frac{2}{L}} \sin\left(\frac{\pi}{L}x'\right) \right) \left(\sqrt{\frac{2}{L}} \sin\left(\frac{\pi}{L}x'\right) \right) x' \delta(x - x') \\
&= \frac{2}{L} \int_0^L dx' x' \sin^2\left(\frac{\pi}{L}x'\right) \\
&= \frac{2}{L} \left(\frac{L}{2} \right)^2 \\
&= \frac{L}{2}
\end{aligned}$$

$$\begin{aligned}
\langle 1|\hat{X}|2\rangle &= \int_0^L dx \int_0^L dx' \langle 1|x\rangle \langle x'|2\rangle \langle x|\hat{X}|x'\rangle \\
&= \int_0^L dx \int_0^L dx' \left(\sqrt{\frac{2}{L}} \sin\left(\frac{\pi}{L}x\right) \right) \left(\sqrt{\frac{2}{L}} \sin\left(\frac{2\pi}{L}x'\right) \right) x' \delta(x-x') \\
&= \frac{2}{L} \int_0^L dx' x' \sin\left(\frac{\pi}{L}x'\right) \sin\left(\frac{2\pi}{L}x'\right) \\
&= \frac{2}{L} \left(-\frac{8L^2}{9\pi^2} \right) \\
&= -\frac{16L}{9\pi^2}
\end{aligned}$$

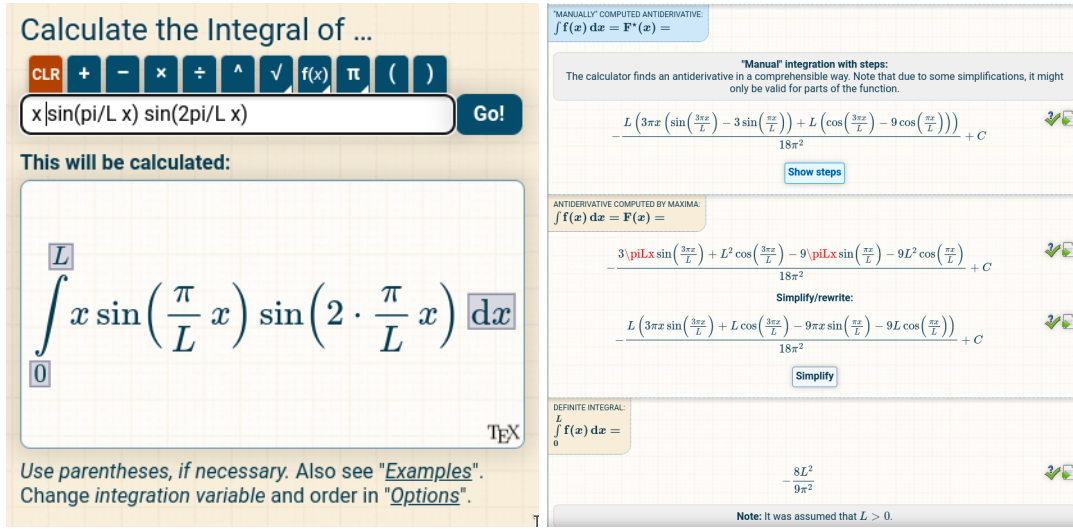


Figure 2: ss/inti02.png

$$\begin{aligned}
\langle 1|\hat{K}|1\rangle &= \int_0^L dx' \langle 1|\hat{K}|x'\rangle \langle x'|1\rangle \\
&= \int_0^L dx' \int_0^L dx \langle 1|x\rangle \langle x|\hat{K}|x'\rangle \langle x'|1\rangle \\
&= \int_0^L dx' \int_0^L dx \langle 1|x\rangle \langle x|\hat{K}|x'\rangle \langle x'|1\rangle \\
&= \frac{2}{L} \int_0^L dx' \int_0^L dx \sin\left(\frac{\pi}{L}x\right) \langle x|\hat{K}|x'\rangle \sin\left(\frac{\pi}{L}x'\right) \\
&= \frac{2}{L} \int_0^L dx' \int_0^L dx \sin\left(\frac{\pi}{L}x\right) \delta(x-x') \left(-i \frac{d}{dx'}\right) \sin\left(\frac{\pi}{L}x'\right) \\
&= \frac{2}{L} \int_0^L dx' \int_0^L dx \sin\left(\frac{\pi}{L}x\right) \delta(x-x') (-i) \left(\frac{\pi}{L}\right) \cos\left(\frac{\pi}{L}x'\right) \\
&= -i \frac{2\pi}{L^2} \int_0^L dx' \sin\left(\frac{\pi}{L}x'\right) \cos\left(\frac{\pi}{L}x'\right) \\
&= 0
\end{aligned}$$

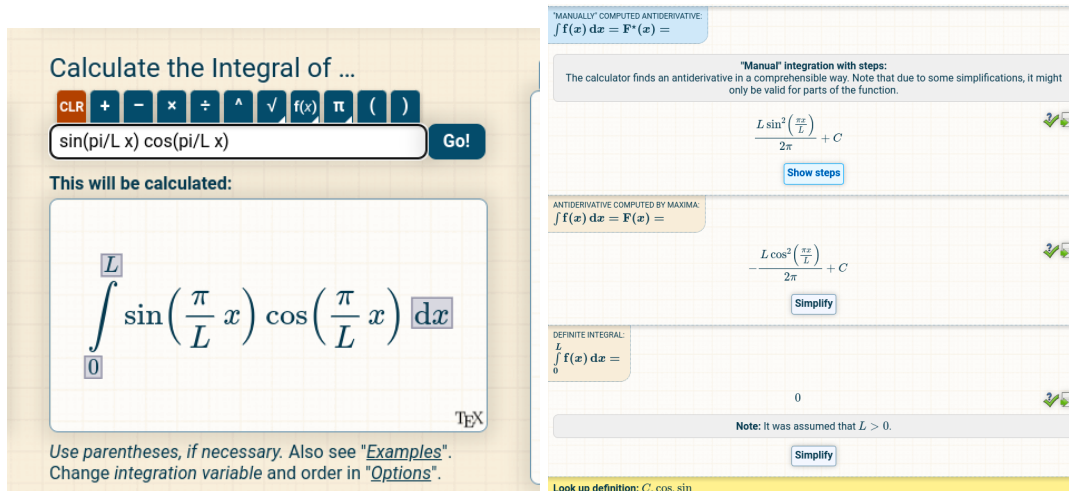


Figure 3: ss/inti04.png

$$\begin{aligned}
 \langle 1 | \hat{K} | 2 \rangle &= \int_0^L dx' \int_0^L dx \langle 1 | x \rangle \langle x | \hat{K} | x' \rangle \langle x' | 2 \rangle \\
 &= \frac{2}{L} \left(-i \frac{2\pi}{L} \right) \int_0^L dx' \int_0^L dx \sin\left(\frac{\pi}{L} x\right) \delta(x - x') \cos\left(\frac{2\pi}{L} x'\right) \\
 &= -i \frac{4\pi}{L^2} \int_0^L dx' \sin\left(\frac{\pi}{L} x'\right) \cos\left(\frac{2\pi}{L} x'\right) \\
 &= -i \frac{4\pi}{L^2} \left(-\frac{2L}{3\pi} \right) \\
 &= i \frac{8}{3L}
 \end{aligned}$$

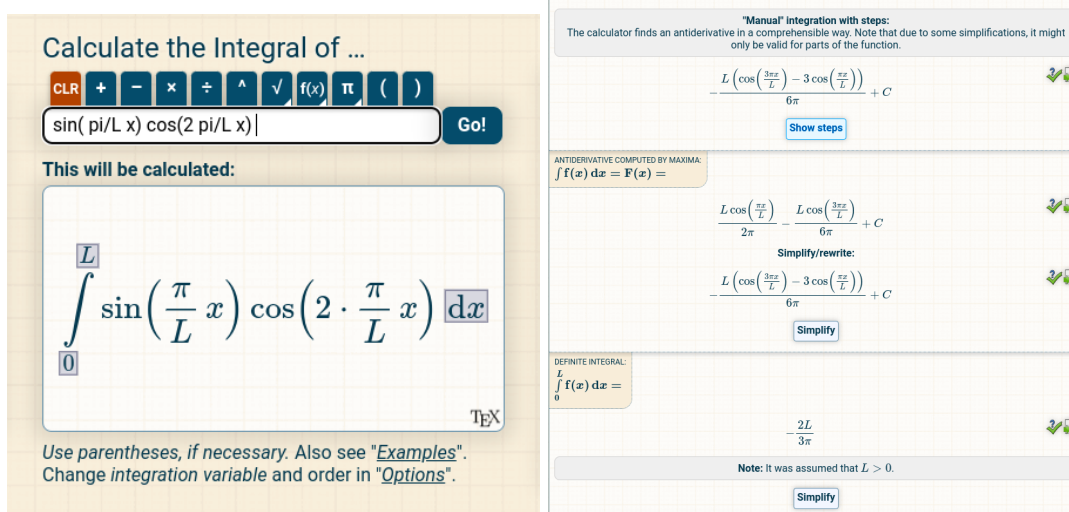


Figure 4: ss/inti07.png

b

$$\begin{aligned}
\langle k_1 | \hat{X} | k_1 \rangle &= \langle k_1 | \hat{I} \hat{X} \hat{I} | k_1 \rangle \\
&= \left(\int_0^L dx' \langle k_1 | x' \rangle \langle x' | \right) \hat{X} \left(\int_0^L dx |x\rangle \langle x | k_1 \rangle \right) \\
&= \int_0^L dx' \int_0^L dx \langle k_1 | x' \rangle \langle x' | \hat{X} | x \rangle \langle x | k_1 \rangle \\
&= \int_0^L dx' \int_0^L dx \langle k_1 | x' \rangle x \delta(x' - x) \langle x | k_1 \rangle \\
&= \int_0^L dx' \int_0^L dx x \left(\frac{1}{\sqrt{L}} e^{-ik_1 x'} \right) \left(\frac{1}{\sqrt{L}} e^{ik_1 x} \right) \delta(x' - x) \\
&= \frac{1}{L} \int_0^L dx x e^{-ik_1 x} e^{ik_1 x} \\
&= \frac{1}{L} \int_0^L dx x \\
&= \frac{1}{L} \left(\frac{L^2}{2} \right) \\
&= \frac{L}{2}
\end{aligned}$$

$$\begin{aligned}
\langle k_1 | \hat{X} | k_2 \rangle &= \left(\int_0^L dx' \langle k_1 | x' \rangle \langle x' | \right) \hat{X} \left(\int_0^L dx |x\rangle \langle x | k_2 \rangle \right) \\
&= \int_0^L dx' \int_0^L dx \langle k_1 | x' \rangle \langle x' | \hat{X} | x \rangle \langle x | k_2 \rangle \\
&= \frac{1}{L} \int_0^L dx' \int_0^L dx x e^{-ik_1 x'} e^{ik_2 x} \delta(x' - x) \\
&= \frac{1}{L} \int_0^L dx x e^{i(k_2 - k_1)x} \\
&= \frac{1}{L} \int_0^L dx x e^{i(2\frac{\pi}{L})x} \\
&= \left(\frac{1}{L} \right) \left(-\frac{iL^2}{2\pi} \right) \\
&= -\frac{iL}{2\pi}
\end{aligned}$$

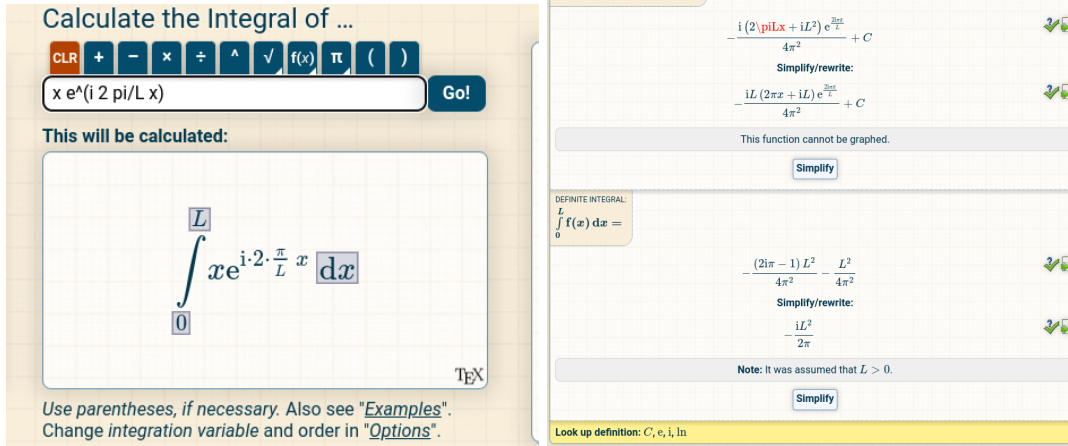


Figure 5: ss/balsal2.png

$$\begin{aligned}
 \langle k_1 | \hat{K} | k_1 \rangle &= \left(\int_0^L dx' \langle k_1 | x' \rangle \langle x' | \right) \hat{K} \left(\int_0^L dx |x\rangle \langle x | k_1 \rangle \right) \\
 &= \int_0^L dx' \int_0^L dx \langle k_1 | x' \rangle \langle x' | \hat{K} | x \rangle \langle x | k_1 \rangle \\
 &= \frac{1}{L} \int_0^L dx' \int_0^L dx e^{-ik_1 x'} \left(-i\delta(x' - x) \frac{d}{dx} \right) e^{ik_1 x} \\
 &= \frac{-i}{L} \int_0^L dx' \int_0^L dx e^{-ik_1 x'} \delta(x' - x) i k_1 e^{ik_1 x} \\
 &= \frac{k_1}{L} \int_0^L dx e^{-ik_1 x} e^{ik_1 x} \\
 &= \frac{k_1}{L} \int_0^L dx \\
 &= k_1 \\
 &= \frac{2\pi}{L}
 \end{aligned}$$

$$\begin{aligned}
\langle k_1 | \hat{K} | k_2 \rangle &= \left(\int_0^L dx' \langle k_1 | x' \rangle \langle x' | \right) \hat{K} \left(\int_0^L dx | x \rangle \langle x | k_2 \rangle \right) \\
&= \int_0^L dx' \int_0^L dx \langle k_1 | x' \rangle \langle x' | \hat{K} | x \rangle \langle x | k_2 \rangle \\
&= \frac{1}{L} \int_0^L dx' \int_0^L dx e^{-ik_1 x'} \left(-i\delta(x' - x) \frac{d}{dx} \right) e^{ik_2 x} \\
&= \frac{-i}{L} \int_0^L dx' \int_0^L dx e^{-ik_1 x'} \delta(x' - x) i k_2 e^{ik_2 x} \\
&= \frac{k_2}{L} \int_0^L dx e^{-ik_1 x} e^{ik_2 x} \\
&= \frac{k_2}{L} \int_0^L dx e^{i(2\pi/L)x} \\
&= 0
\end{aligned}$$

Problem 05

0.1 a

$$\begin{aligned}
\langle k | \hat{X} | k' \rangle &= \left(\int_{-\infty}^{\infty} dx \langle k | x \rangle \langle x | \right) \hat{X} \left(\int_{-\infty}^{\infty} dx' | x' \rangle \langle x' | k' \rangle \right) \\
&= \int dx \int dx' \langle k | x \rangle \langle x | \hat{X} | x' \rangle \langle x' | k' \rangle \\
&= \int dx \int dx' \frac{1}{\sqrt{2\pi}} e^{-ikx} \langle x | \hat{X} | x' \rangle \frac{1}{\sqrt{2\pi}} e^{ik'x} \\
&= \int dx \int dx' \frac{1}{\sqrt{2\pi}} e^{-ikx} x' \delta(x - x') \frac{1}{\sqrt{2\pi}} e^{ik'x'} \\
&= \int dx \int dx' \frac{1}{\sqrt{2\pi}} e^{-ikx} \delta(x - x') \left(x' \frac{1}{\sqrt{2\pi}} e^{ik'x'} \right) \\
&= \int dx' \frac{1}{\sqrt{2\pi}} e^{-ikx'} \left(x' \frac{1}{\sqrt{2\pi}} e^{ik'x'} \right) \\
&= \int dx' \frac{1}{\sqrt{2\pi}} e^{-ikx'} \left(x' \frac{1}{\sqrt{2\pi}} e^{ik'x'} \right) \\
&= \left(\frac{1}{\sqrt{2\pi}} \right)^2 \int dx' x' e^{i(k' - k)x'} \\
&= (-1)(-i) \left(\frac{1}{\sqrt{2\pi}} \right)^2 \int dx' \frac{d}{dk} e^{i(k' - k)x'} \\
&= i \frac{d}{dk} \left(\frac{1}{2\pi} \int_{-\infty}^{\infty} dx e^{i(k' - k)x} \right) \\
&= i \frac{d}{dk} \delta(k' - k) = i \frac{d}{dk} \delta(k - k')
\end{aligned}$$

b

$$\begin{aligned}\langle k|\hat{X}|f\rangle &= \int_{-\infty}^{\infty} dk' \langle k|\hat{X}|k'\rangle \langle k'|f\rangle \\&= \int_{-\infty}^{\infty} dk' \left(i \frac{d}{dk} \delta(k-k') \right) \langle k'|f\rangle \\&= \int_{-\infty}^{\infty} dk' \left(i \delta(k-k') \frac{d}{dk'} \right) \langle k'|f\rangle \\&= \int_{-\infty}^{\infty} dk' \left(i \delta(k-k') \frac{d}{dk'} \right) \left(\int_{-\infty}^{\infty} dx \frac{e^{-ik'x}}{\sqrt{2\pi}} f(x) \right) \\&= \int_{-\infty}^{\infty} dk' \left(i \delta(k'-k) \frac{d}{dk'} \right) \left(\int_{-\infty}^{\infty} dx \frac{e^{-ik'x}}{\sqrt{2\pi}} f(x) \right) \\&= \left(i \frac{d}{dk} \right) \left(\int_{-\infty}^{\infty} dx \frac{e^{-ikx}}{\sqrt{2\pi}} f(x) \right) \\&= i \frac{d}{dk} \tilde{f}(k)\end{aligned}$$