Quantum Mechanics Homework 01

December 5, 2024

Ahmed Saad Sabit, Rice University

Problem 1 (a)

Studying the problem

The initial configuration of the string

$$q(x,0) = f(x) = \begin{cases} \frac{2h}{L}x & 0 \le x \le \frac{L}{2} \\ 2h - \frac{2h}{L}x & \frac{L}{2} \le x \le L \end{cases}$$
$$\frac{\partial q(x,t)}{\partial t}_{t=0} = \sum_{n=1}^{\infty} d_n \Omega_n \phi_n(x) = g(x) = 0 \implies \boxed{d_n = 0}$$

The general solution to the string equation (assumed solution is separable between time and position)

$$q(x,t) = \sum_{n=1}^{\infty} [c_n \cos(\Omega_n t) + d_n \sin(\Omega_n t)] \phi_n(x)$$

For t = 0 we get,

$$q(x,0) = f(x) = \sum_{n=1}^{\infty} c_n \phi_n(x)$$

We are interested on finding the general solution of q(x,t) that will hold for the future given this initial condition. The variables of our equation are obviously x,t and what we need to find out is c_n,d_n . The next sub-section will find out a solution for c_n (d_n is trivially zero given zero initial velocity).

Solving for c_n

Let us do the following computation now. Let us multiply both sides of the above equation with $\phi_p(x)$ where p represents the p-th term while we take a summation over the index of n.

$$f(x)\phi_p(x) = \sum_{n=1}^{\infty} c_n \phi_n(x)\phi_p(x)$$

Just so that we can invoke the inner product between orthonormal bases, we can take an integral with the following way

$$\int_0^L dx f(x)\phi_p(x) = \int_0^L dx \left(\sum_{n=1}^\infty c_n \phi_n(x)\phi_p(x) \right)$$
$$= \sum_{n=1}^\infty c_n \int_0^L dx \, \phi_n(x)\phi_p(x)$$
$$= \sum_{n=1}^\infty c_n \delta_{np} \frac{L}{2}$$
$$= c_p \frac{L}{2}$$

This above gives us the p-th term

$$c_p = \frac{2}{L} \int_0^L \mathrm{d}x \, f(x) \phi_p(x)$$

Using the explicit equation for the bases and also looking at the piecewise function, we can write,

$$c_{p} = \frac{2}{L} \int_{0}^{L} dx f(x) \sin\left(\frac{p\pi x}{L}\right)$$

$$= \frac{2}{L} \left(\int_{0}^{\frac{L}{2}} f(x) \sin\left(\frac{p\pi x}{L}\right) + \int_{\frac{L}{2}}^{L} f(x) \sin\left(\frac{p\pi x}{L}\right)\right)$$

$$= \frac{2}{L} \left(\int_{0}^{\frac{L}{2}} \frac{2h}{L} x \sin\left(\frac{p\pi x}{L}\right) + \int_{\frac{L}{2}}^{L} \left(2h - \frac{2h}{L} x\right) \sin\left(\frac{p\pi x}{L}\right)\right)$$

$$= \frac{2}{L} \left(\frac{hL}{\pi^{2}p^{2}} \left[2 \sin\left(p\frac{\pi}{2}\right) - \pi p \cos\left(p\frac{\pi}{2}\right)\right] - \frac{hL}{\pi^{2}p^{2}} \left[2 \sin\left(\pi p\right) - 2 \sin\left(p\frac{\pi}{2}\right) - \pi p \cos\left(p\frac{\pi}{2}\right)\right]\right)$$

$$= \frac{8h}{\pi^{2}p^{2}} \sin\left(p\frac{\pi}{2}\right) \left[1 - \cos\left(p\frac{\pi}{2}\right)\right]$$

Hence if I write this huge mess properly

$$c_p = \frac{8h}{\pi^2 p^2} \sin\left(\frac{p\pi}{2}\right) \left[1 - \cos\left(\frac{p\pi}{2}\right)\right]$$

Discussion on odd and even modes

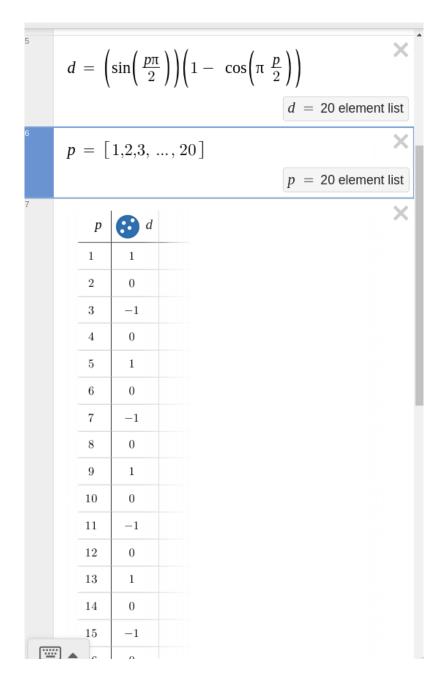


Figure 1: ss/graph-cn.png

We can see that every *even* numbers is giving us a 0 for c_p (d term in graph). The *odd* terms alter between -1 and 1. The contribution of even and odd modes is obvious from here.

Plotting the displacement with time

Let's set L=1 and we will look at $0 \le x \le 1$.

As given $\Omega_n = k_n = \frac{n\pi}{L} = n\pi$.

For h let's pick h = 0.4.

Using hte c_p we derived above, the displacement function q(x,t) is then,

$$q(x,t)_{(10)} = \sum_{n=1}^{10} c_n \cos(n\pi t) \sin(n\pi x)$$

Expanding c_n ,

$$q(x,t)_{(10)} = \frac{8(0.4)}{\pi^2 n^2} \sum_{n=1}^{10} \sin\left(\frac{n\pi}{2}\right) \left[1 - \cos\left(\frac{n\pi}{2}\right)\right] \cos(n\pi t) \sin(n\pi x)$$

And
$$\tau = \frac{2\pi}{\Omega_1} = \frac{2\pi}{\pi} = 2$$

Plots of q(x,t) where $\tau=2$

$$q(x,t)_{(10)} = \frac{8(0.4)}{\pi^2} \sum_{n=1}^{10} \frac{1}{n^2} \sin\left(\frac{n\pi}{2}\right) \left[1 - \cos\left(\frac{n\pi}{2}\right)\right] \cos(n\pi t) \sin(n\pi x) \quad \text{and } x \in [0,1]$$

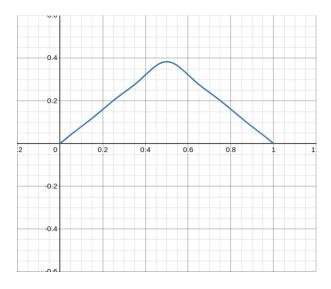


Figure 2: t=0 and also $t=\tau$

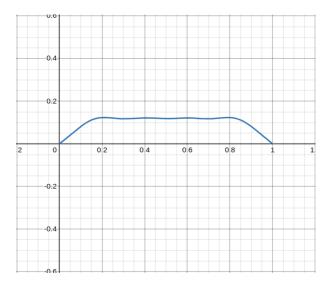


Figure 3: t = 0.35

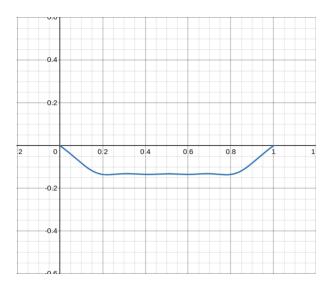


Figure 4: $t = \frac{\tau}{3} = \frac{2}{3}$

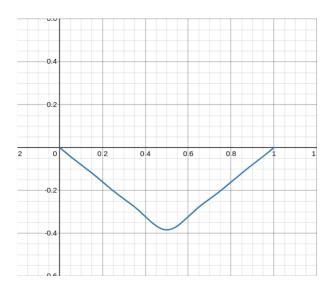


Figure 5: $t = \frac{\tau}{2} = 1$

Problem 1 (b)

Studying the Problem

Initially the string is tight and straight hence

$$q(x,0) = \sum_{n=1}^{\infty} c_n \phi_n(x) = f(x) = 0 \implies \left[c_n = 0 \right]$$
$$\frac{\partial q(x,t)}{\partial t}_{t=0} = g(x) = v_0 \theta \left(a - \left| x - \frac{L}{2} \right| \right)$$

Where $\theta(x)$ is a Heaviside step function (outputs 1 whenever input is 0 or positive). I am not going to waste my and graders time by re-writing everything I wrote above, the procedure we are going to follow is same as above.

Computation of d_n

$$g(x) = \sum_{n=1}^{\infty} d_n \Omega_n \phi_n(x)$$

$$\int_0^L dx \, g(x) \phi_p(x) = \sum_{n=1}^{\infty} \int_0^L dx \, d_n \Omega_n \phi_n(x) \phi_p(x)$$

$$\int_{\frac{L}{2} - a}^{\frac{L}{2} + a} v_0 \phi_p(x) = d_p \Omega_p \frac{L}{2}$$

$$\int_{\frac{L}{2} - a}^{\frac{L}{2} + a} v_0 \sin(k_p x) = d_p k_p \frac{L}{2}$$

$$\frac{v_0}{k_p} \left[-\cos(k_p x) \right]_{x = \frac{L}{2} - a}^{x = \frac{L}{2} + a} = d_p k_p \frac{L}{2}$$

$$\cos\left(k_p \frac{L}{2} - k_p a\right) - \cos\left(k_p \frac{L}{2} + k_p a\right) = \frac{d_p k_p^2 L}{2v_0}$$

$$\cos\left(\frac{n\pi}{2} - \frac{n\pi a}{L}\right) - \cos\left(\frac{n\pi}{2} + \frac{n\pi a}{L}\right) = \frac{d_p n^2 \pi^2}{2v_0 L}$$

This gives us

$$d_p = \frac{2v_0L}{n^2\pi^2} \left[\cos\left(\frac{n\pi}{2} - \frac{n\pi a}{L}\right) - \cos\left(\frac{n\pi}{2} + \frac{n\pi a}{L}\right) \right] = \frac{4v_0L}{n^2\pi^2} \sin\left(\frac{n\pi}{2}\right) \sin\left(\frac{n\pi a}{L}\right)$$

Discussion on Even and Odd modes

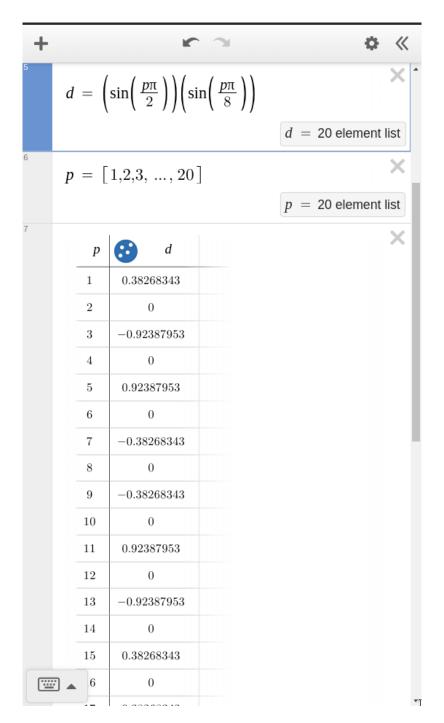


Figure 6: ss/graph-dn

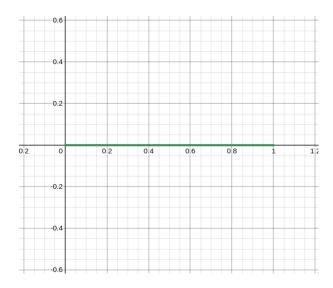
Like before we can see every *even* modes are 0.

Plotting the displacement with time

Let's set L=1, then a=1/8. Let's set $v_0=1.5$. As before $\tau=2$.

Plots of q(x,t) where $\tau=2$

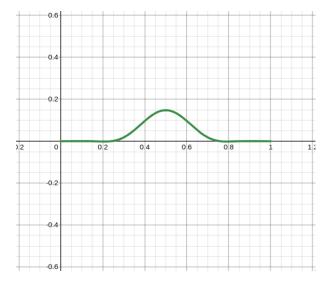
$$q(x,t)_{(10)} = \frac{4(1.5)}{\pi^2} \sum_{n=1}^{10} \frac{1}{n^2} \sin\left(\frac{n\pi}{2}\right) \sin\left(n\pi(1/8)\right) \sin\left(n\pi t\right) \sin\left(n\pi x\right) \quad \text{and } x \in [0,1]$$



0.6

Figure 7: t=0 and also $t=\tau$

Figure 9: $t = \frac{\tau}{3} = \frac{2}{3}$



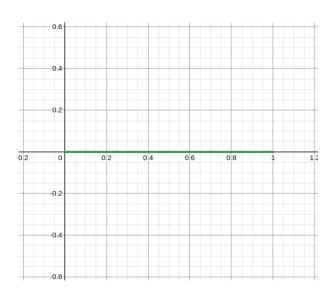
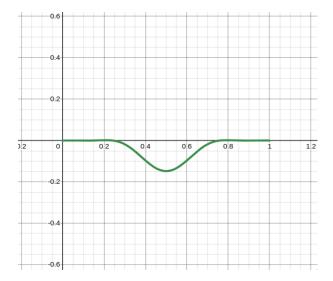


Figure 8: $t = \frac{\tau}{20} = 0.1$

Figure 10: $t = \frac{\tau}{2} = 1$



0.4
0.2
0.2
0.2
0.2
0.4
0.6
0.8
1
1

Figure 11: t = 1.1

Figure 13: t = 1.85

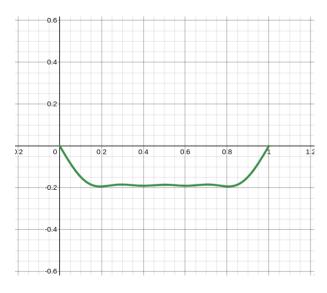




Figure 12: t = 1.5

Figure 14: t = 1.95

Plotting the Initial Velocity

We can simply take the time derivative and make a plot. The more iterations we make, we should get close to a heaviside function as stated in the initial velocity profile.

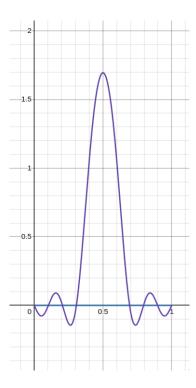


Figure 15: Plotted for n = 10

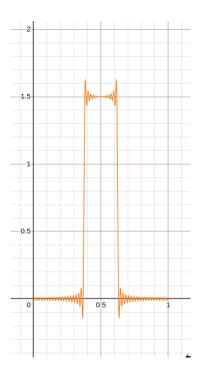


Figure 18: Plotted for n = 100

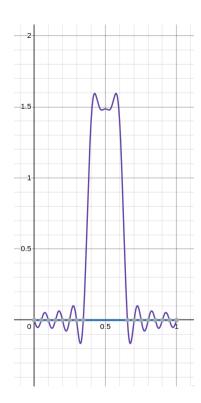


Figure 16: Plotted for n = 20

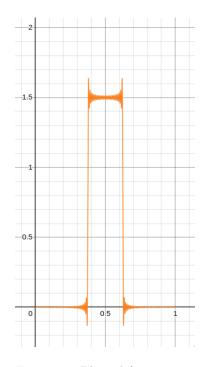


Figure 19: Plotted for n = 200

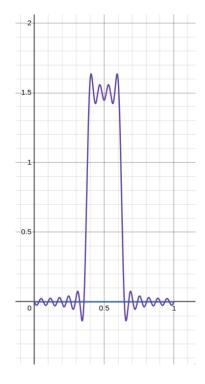


Figure 17: Plotted for n = 30

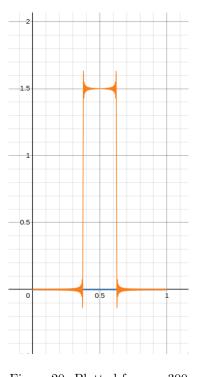


Figure 20: Plotted for n=300

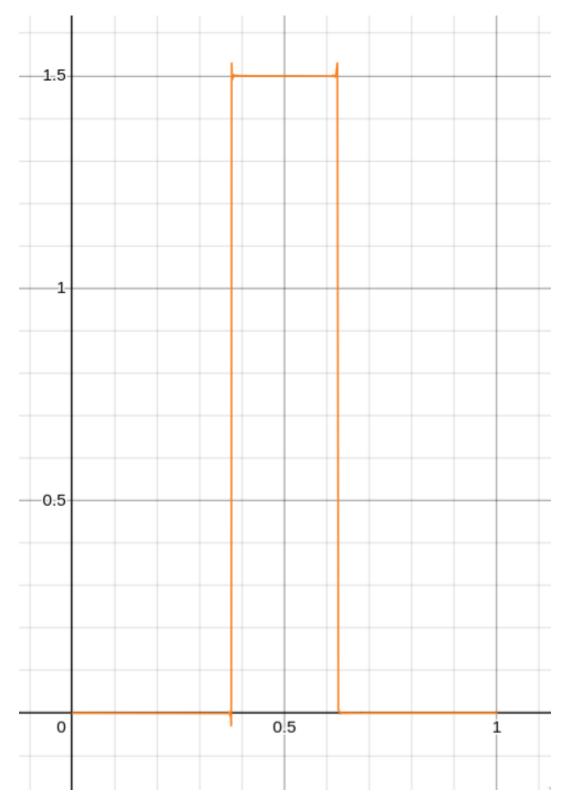


Figure 21: Waited 2 minutes to try n = 10,000. We have the exact initial conditions we wanted. This is absolutely beautiful to do this myself.

Problem 2

Linear Dependence equation

$$a\begin{pmatrix}0&1\\0&0\end{pmatrix}+b\begin{pmatrix}1&1\\0&1\end{pmatrix}+c\begin{pmatrix}-2&-1\\0&-2\end{pmatrix}=0$$

If the three above equations are linearly independent we are going to have the only possible solution of the equality being zero as $\{a, b, c\} = \{0, 0, 0\}$. We get three different equations

$$b-2c = 0$$

$$a+b-c = 0$$

$$b-2c = 0$$

$$0+0 = 0$$

Where the I didn't count the last one as it's not really anything. If we solve the system of equation above we are left with (did the computation by hand on paper)

$$a = -c$$

$$b = 2c$$

$$c = \text{can be anything}$$

a, b, c doesn't necessarily need to be 0 to yield above equation to be zero. This system is NOT linearly independent. For example, setting c = 1

$$-\begin{pmatrix}0&1\\0&0\end{pmatrix}+2\begin{pmatrix}1&1\\0&1\end{pmatrix}+\begin{pmatrix}-2&-1\\0&-2\end{pmatrix}=\begin{pmatrix}0&0\\0&0\end{pmatrix}$$

Problem 3

Deduction

$$|v+w|^2 = \langle v+w|v+w \rangle$$

$$= \langle v|v+w \rangle + \langle w|v+w \rangle$$

$$= \langle v|v \rangle + \langle v|w \rangle + \langle w|v \rangle + \langle w|w \rangle$$

$$= |v|^2 + |w|^2 + \langle v|w \rangle + \langle v|w \rangle^*$$

$$= |v|^2 + |w|^2 + 2\operatorname{Re}(\langle v|w \rangle)$$

$$\leq |v|^2 + |w|^2 + 2|\langle v|w \rangle|$$

$$\leq |v|^2 + |w|^2 + 2|v||w|$$

$$\leq (|v| + |w|)^2$$

This shows that

$$|v + w| \le |v| + |w|$$

Equality condition

Above we can see

$$|v + w|^2 = \langle v + w|v + w\rangle$$
$$= |v|^2 + |w|^2 + 2\operatorname{Re}(\langle v|w\rangle)$$

For the equality to hold we strictly need

$$2\operatorname{Re}(\langle v|w\rangle) = 2|v||w|$$

Let's try $|v\rangle = a|w\rangle$.

$$2\operatorname{Re}\left(\langle aw|w\rangle\right) = 2a\operatorname{Re}\left(\langle w|w\rangle\right) = 2|aw||w| = 2a|w||w|$$

 $|v\rangle=a|w\rangle$ satisfies this. If $|v\rangle$ has component perpendicular to $|w\rangle$ (in a way that $|v\rangle$ can be broken down into constituent factors of basis vectors such that inner product with $|w\rangle$ gives zero), then some of the value of $2\text{Re}(\langle v|w\rangle)$ is getting lost. Having the two vectors parallel gives equal to their norm. This is more of a physical vector-like intuition.

Problem 4(a)

I will write $\hat{\sigma}^n$ as simply σ^n for this problem. I did the multiplication by hand.

$$(\sigma^1)^2 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$(\sigma^2)^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$(\sigma^3)^2 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

We had been already defined

$$\sigma^0 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

All of these result matrix above is the identity matrix that validates

$$(\sigma^1)^2 = (\sigma^2)^2 = (\sigma^3)^2 = \sigma^0$$

Problem 4(b)

We are required to solve for $A^{\mu,\nu}$ where

$$A^{\mu,\nu} = \sigma^{\mu}\sigma^{\nu} + \sigma^{\nu}\sigma^{\mu}$$

Note that it is obvious

$$A^{\mu,\nu} = A^{\nu,\mu}$$

Computing each of the matrix multiplications, and also referring to previous computations

$$A^{k,k} = A^{1,1} = A^{2,2} = A^{3,3} = 2(\sigma^k)^2 = 2\sigma^0 = 2\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$A^{1,2} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$
$$= \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} + \begin{pmatrix} -i & 0 \\ 0 & i \end{pmatrix}$$
$$= \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

$$A^{2,3} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$
$$= \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} + \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix}$$
$$= \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

$$A^{1,3} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$
$$= \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$
$$= \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

Different indexes cause a zero-matrix, and similar causes a double of identity matrix. From here we can easily figure out that

$$A^{\mu,\nu} = 2\sigma^0 \delta_{\mu,\nu}$$

Problem 4(c)

I am going to borrow the computations I did last problem

$$\begin{split} \operatorname{Tr}[\sigma^1\sigma^1] &= \operatorname{Tr}\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = 2 \\ \operatorname{Tr}[\sigma^2\sigma^2] &= \operatorname{Tr}\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = 2 \\ \operatorname{Tr}[\sigma^3\sigma^3] &= \operatorname{Tr}\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = 2 \\ \operatorname{Tr}[\sigma^1\sigma^2] &= \operatorname{Tr}\begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} = 0 = \operatorname{Tr}[\sigma^2\sigma^1] \\ \operatorname{Tr}[\sigma^2\sigma^3] &= \operatorname{Tr}\begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} = 0 = \operatorname{Tr}[\sigma^3\sigma^2] \\ \operatorname{Tr}[\sigma^1\sigma^3] &= \operatorname{Tr}\begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = 0 = \operatorname{Tr}[\sigma^3\sigma^1] \end{split}$$

From this we can see that same-index gives 2 and different gives 0. From this it's obvious

$$Tr[\sigma^{\mu}\sigma^{\nu}] = 2\delta_{\mu\nu}$$

Problem 4(d)

Expanding the equation of the operator

$$\hat{V} = \sum_{i=1}^{3} V_i \sigma^i = V_1 \sigma^1 + V_2 \sigma^2 + V_3 \sigma^3$$

Multiply σ^p where $p \in \{1, 2, 3\}$

$$\hat{V}\sigma^p = V_1\sigma^1\sigma^p + V_2\sigma^2\sigma^p + V_3\sigma^3\sigma^p$$

Taking the trace and using the property Tr(A + B) = Tr(A) + Tr(B)

$$\operatorname{Tr}(\hat{V}\sigma^{p}) = V_{1}(2\delta_{1p}) + V_{2}(2\delta_{2p}) + V_{3}(2\delta_{3p})$$

From this using the definition of the $\delta_{\mu,p}$ we can simply write,

$$V_p = \frac{1}{2} \text{Tr}(\hat{V} \sigma^p)$$

Using the form

$$\hat{V} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

we can find the coefficients V_p

$$\hat{V}\sigma^{0} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \implies V_{0} = \frac{a+d}{2}$$

$$\hat{V}\sigma^{1} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} b & a \\ d & c \end{pmatrix} \implies V_{1} = \frac{b+c}{2}$$

$$\hat{V}\sigma^{2} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} ib & -ia \\ id & -ic \end{pmatrix} \implies V_{2} = i\frac{b-c}{2}$$

$$\hat{V}\sigma^{3} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} a & -b \\ c & -d \end{pmatrix} \implies V_{3} = \frac{a-d}{2}$$

So our representation is

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \frac{a+d}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \frac{b+c}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + \frac{ib-ic}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + \frac{a-d}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

I've checked this in Wolfram Alpha and it seems to work.

Problem 5(a)

From the question, my understanding of the basis is

$$|1\rangle = \begin{pmatrix} 1\\0 \end{pmatrix} \quad |2\rangle = \begin{pmatrix} 0\\1 \end{pmatrix}$$

Projection operator

$$P_1 = |1\rangle\langle 1| = \begin{pmatrix} 1\\0 \end{pmatrix} \begin{pmatrix} 1&0 \end{pmatrix} = \begin{pmatrix} 1&0\\0&0 \end{pmatrix}$$

$$P_2 = |2\rangle\langle 2| = \begin{pmatrix} 0\\1 \end{pmatrix}\begin{pmatrix} 0&1 \end{pmatrix} = \begin{pmatrix} 0&0\\0&1 \end{pmatrix}$$

We've determined the elements of the matrices.

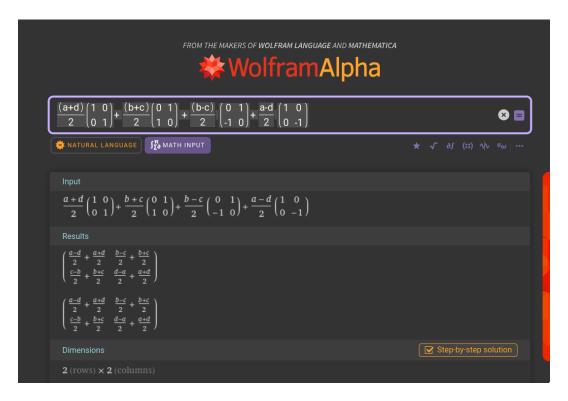


Figure 22: ss/pauli-basis.png

Problem 5(b)

Doing the matrix multiplication by hand

$$P_{1}P_{1} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$P_{2}P_{2} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$P_{1}P_{2} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

$$P_{2}P_{1} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

My intuition for the projection operator is basically

$$P_i P_j = (|i\rangle\langle i|)(|j\rangle\langle j|) = |i\rangle\langle i|j\rangle\langle j| = |i\rangle\delta_{ij}\langle j| = \delta_{ij}|i\rangle\langle j| = |k\rangle\langle k|$$

where $k = \{i, j\}$

Problem 5(c)

Computation

$$\begin{split} |V\rangle &= \frac{1}{\sqrt{2}}(|1\rangle + |2\rangle) \\ P_V &= |V\rangle\langle V| = \frac{1}{2}(|1\rangle + |2\rangle)(\langle 1| + \langle 2|) = \frac{1}{2}(|1\rangle\langle 1| + |1\rangle\langle 2| + |2\rangle\langle 1| + |2\rangle\langle 2|) \end{split}$$

$$P_V = \frac{1}{2} \begin{bmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \end{bmatrix} = \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$$

Now computing

$$P_1 P_V P_2 = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{bmatrix} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \end{bmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$
$$= \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix}$$
$$= \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

We could have also done this directly like

$$\begin{split} P_1 P_V P_2 &= \frac{1}{2} \left(|1\rangle\langle 1| \right) \left(|1\rangle\langle 1| + |1\rangle\langle 2| + |2\rangle\langle 1| + |2\rangle\langle 1| \right) \left(|2\rangle\langle 2| \right) \\ &= \frac{1}{2} \left(|1\rangle\langle 1| \right) \left[|1\rangle\langle 1| 2\rangle\langle 2| + |1\rangle\langle 2| 2\rangle\langle 2| + |2\rangle\langle 1| 2\rangle\langle 1| + |2\rangle\langle 2| 2\rangle\langle 2| \right] \\ &= \frac{1}{2} \left(|1\rangle\langle 1| \right) \left(|1\rangle\langle 2| + |2\rangle\langle 2| \right) \\ &= \frac{1}{2} \left(|1\rangle\langle 1| 1\rangle\langle 2| + |1\rangle\langle 1| 2\rangle\langle 2| \right) \\ &= \frac{1}{2} |1\rangle\langle 2| \\ &= \frac{1}{2} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \end{split}$$

Analysis

We can keep three polarizers, each at a increasing angle, then components of lights get filtered as the pass through, remaining light falls on the screen behind.