# Quantum Mechanics: : Homework 03

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# Problem 01

(a)

I did the first matrix multiplication computation by hand and it aligns with what I got from Matlab. There's way too many multiplication to do so I am instead opting for matlab solution.

$$\sigma_1 \kappa_1 = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}$$

$$\kappa_1 \sigma_1 = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}$$

I have thought enough how to convince the grader that I know how to do Matrix Multiplication, the most tedious approach being I show each and every single individual dot product on paper. That's too much work.

We just wanna prove that the commutator is zero

$$[\sigma_1 \kappa_1] = \sigma_1 \kappa_1 - \kappa_1 \sigma_1 = 0$$

Very similarly we can keep doing it, one by one in MATLAB

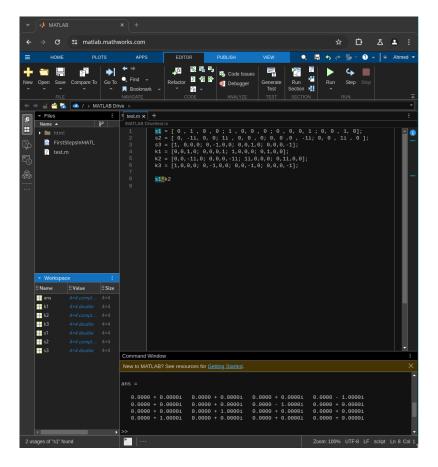


Figure 1: Computing  $\sigma_1 \kappa_2$ 

We find

$$\sigma_1 \kappa_2 = \kappa_2 \sigma_1 = \begin{pmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & -i & 0 \\ 0 & i & 0 & 0 \\ i & 0 & 0 & 0 \end{pmatrix} \quad \to \quad [\sigma_1, \kappa_2] = 0$$

$$\sigma_1 \kappa_3 = \kappa_3 \sigma_1 = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 \end{pmatrix} \quad \rightarrow \quad [\sigma_1, \kappa_3] = 0$$

$$\sigma_2 \kappa_2 = \kappa_2 \sigma_2 = \begin{pmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} \quad \rightarrow \quad [\sigma_2, \kappa_2] = 0$$

$$\sigma_2 \kappa_3 = \kappa_3 \sigma_2 = \begin{pmatrix} 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \\ 0 & 0 & 0 & i \\ 0 & 0 & -i & 0 \end{pmatrix} \quad \rightarrow \quad [\sigma_2, \kappa_3] = 0$$

#### b

Take  $\sigma_1$  and do pen and paper calculation to find out the eigenvalues. I used mathematica to find out the determinant safely and we get the eigenvalues from the characteristic polynomial

$$\lambda^4 - 2\lambda^2 + 1 = 0$$

Which gives,

$$\lambda = -1, -1, 1, 1$$

Set of eigenvectors for  $\sigma_1$  being

$$\begin{pmatrix} 0 \\ 0 \\ -1 \\ 1 \end{pmatrix}, \begin{pmatrix} -1 \\ 1 \\ 0 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 0 \\ 1 \\ 1 \end{pmatrix}, \begin{pmatrix} 1 \\ 1 \\ 0 \\ 0 \end{pmatrix}$$

Now let's do this for the  $\kappa_2$  and we get the characteristic equation

$$\lambda^4 - 2\lambda^2 + 1$$

And we get

$$\lambda = -1, -1, 1, 1$$

With Eigenvectors

$$\begin{pmatrix} 0 \\ i \\ 0 \\ 1 \end{pmatrix}, \begin{pmatrix} i \\ 0 \\ 1 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ -i \\ 0 \\ 1 \end{pmatrix}, \begin{pmatrix} -i \\ 0 \\ 1 \\ 0 \end{pmatrix}$$

By theory if there are common eigenvectors,

$$A|\phi\rangle = \phi_A|\phi\rangle$$

$$B|\phi\rangle = \phi_B|\phi\rangle$$

Then

$$AB|\phi\rangle = \phi_A \phi_B |\phi\rangle$$

Taking  $\sigma_1 \kappa_2$ ,

$$\sigma_1 \kappa_2 = \begin{pmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & -i & 0 \\ 0 & i & 0 & 0 \\ i & 0 & 0 & 0 \end{pmatrix}$$

The eigenvectors of these are

$$\begin{pmatrix} i \\ 0 \\ 0 \\ 1 \end{pmatrix}, \begin{pmatrix} 0 \\ i \\ 1 \\ 0 \end{pmatrix}, \begin{pmatrix} -i \\ 0 \\ 0 \\ 1 \end{pmatrix}, \begin{pmatrix} 0 \\ -i \\ 1 \\ 0 \end{pmatrix}$$

Looking at the eigenvectors we know there are NO common eigenvectors.

### Problem 02

 $\mathbf{a}$ 

Total force on  $m_1$  and  $m_2$ 

$$m_1 \ddot{x}_1 = -k_A x_1 + k_B (x_2 - x_1)$$
  
$$m_2 \ddot{x}_2 = -k_B (x_2 - x_1) + k_C (L - x_2 - x_1)$$

The equilibriums are for which net force is zeroed

$$0 = -k_A x_1 + k_B (x_2 - x_1)$$
$$0 = -k_B (x_2 - x_1) + k_C (L - x_2 - x_1)$$

Trying to solve this by hand I get,

$$x_1 = \frac{k_B}{k_A + k_B} x_2$$

Individual solution

$$x_{2}^{0} = \frac{k_{A}k_{C} + k_{B}k_{C}}{(k_{B} + k_{C})(k_{A} + k_{B}) - (k_{B} - k_{C})k_{B}}L$$

$$x_{1}^{0} = \frac{k_{B}k_{C}}{(k_{B} + k_{C})(k_{A} + k_{B}) - (k_{B} - k_{C})k_{B}}L$$

$$-ax_1 + b(x_2 - x_1) = 0, -b(x_2 - x_1) + c(L(x) - x_2 - x_1) = 0$$
Solution 
$$x_1 = \frac{bcL}{ab + ac + 2bc}, x_2 = -\frac{-acL - bcL}{ab + ac + 2bc}, \quad a \neq -\frac{2bc}{b + c}, a \neq -b, a + b \neq 0$$

Figure 2: Verifying that I haven't really messed up.

b

Apparently  $\ddot{x}_i = \delta \ddot{x}_i$ . And

$$x_2 - x_1 = x_2^0 - x_1^0 + \delta x_2 - \delta x_1$$

As we know the denominator of  $x_1^0, x_2^0$  are the same, calling them as D,

$$x_2 - x_1 = \frac{k_A k_C L}{D} + \delta x_2 - \delta x_1$$

$$k_A k_C + 2k_B k_C L + \delta x_1 + \delta x_2 + \delta x_3$$

$$x_2 + x_1 = \frac{k_A k_C + 2k_B k_C}{D} L + \delta x_1 + \delta x_2$$

Thus through substituting our new formulas for  $x_2 - x_1$  and  $x_1 + x_2$  with equilibrium position in the newton's equations we can get,

$$m_1 \delta \ddot{x}_1 = -\frac{k_A k_B k_C L}{D} - k_A \delta x_1 + \frac{k_A k_B k_C L}{D} + k_B (\delta x_2 - \delta x_1) = \boxed{-k_A \delta x_1 + k_B (\delta x_2 - \delta x_1)}$$

For the second equation,

$$m_2 \delta \ddot{x}_2 = -\frac{k_A k_B k_C L}{D} - k_B (\delta x_2 - \delta x_1) + k_C L - \frac{k_A k_C^2 L + 2k_B k_C^2 L}{D} - k_C (\delta dx_2 + \delta dx_1)$$

$$m_2 \delta \ddot{x}_2 = -\frac{k_A k_B k_C L}{D} - k_B (\delta x_2 - \delta x_1) + \frac{k_C D L - k_A k_C^2 L - 2k_B k_C^2 L}{D} + k_C L - k_C (\delta dx_2 + \delta dx_1)$$

Note that  $D = (k_B + k_C)(k_A + k_B) - (k_B - k_C)k_B = k_A k_B + k_A k_C + 2k_B k_C$ , hence

$$\frac{k_C D L - k_A k_C^2 L - 2k_B k_C^2 L}{D} = \frac{k_A k_B k_C + k_A k_C^2 + 2k_B k_C^2 - k_A k_C^2 - 2k_B k_C^2}{D} L = \frac{k_A k_B k_C}{D} L$$

$$m_2\delta\ddot{x}_2 = -\frac{k_Ak_Bk_CL}{D} - k_B(\delta x_2 - \delta x_1) + \frac{k_Ak_Bk_C}{D}L + k_CL - k_C(\delta dx_2 + \delta dx_1)$$

With  $m_1 = m_2 = m$ , we finalize

$$m\delta\ddot{x}_2 = -k_B(\delta x_2 - \delta x_1) - k_C\delta x_2$$

$$m\delta\ddot{x}_1 = -k_A\delta x_1 + k_B(\delta x_2 - \delta x_1)$$

 $\mathbf{c}$ 

We can divide both sides of equations with m and name  $k_I/m = \omega_I^2$ 

$$\delta \ddot{x}_1 = -\omega_A^2 \delta x_1 + \omega_B^2 (\delta x_2 - \delta x_1)$$

$$\delta \ddot{x}_2 = -\omega_B^2 (\delta x_2 - \delta x_1) - \omega_C^2 \delta x_2$$

The homogenous set of equation,

$$\delta \ddot{x}_1 = (-\omega_A^2 - \omega_B^2) \delta x_1 + \omega_B^2 \delta x_2$$
  
$$\delta \ddot{x}_2 = \omega_B^2 \delta x_1 + (-\omega_B^2 - \omega_C^2) \delta x_2$$

This is very equivalently

$$\begin{pmatrix} \delta \ddot{x}_1 \\ \delta \ddot{x}_2 \end{pmatrix} = \begin{pmatrix} (-\omega_A^2 - \omega_B^2) & \omega_B^2 \\ \omega_B^2 & (-\omega_B^2 - \omega_C^2) \end{pmatrix} \begin{pmatrix} \delta x_1 \\ \delta x_2 \end{pmatrix}$$

Simplify notation to avoid myself to getting hospitalized.

$$\begin{pmatrix} \delta \ddot{x}_1 \\ \delta \ddot{x}_2 \end{pmatrix} = \begin{pmatrix} -A - B & B \\ B & -B - C \end{pmatrix} \begin{pmatrix} \delta x_1 \\ \delta x_2 \end{pmatrix}$$

The characteristic equation is

$$\frac{1}{m^2}(k_A k_B + k_A k_C + k_B k_C) + \frac{1}{m}(k_A + 2k_B + k_C)\lambda + \lambda^2 = 0$$

$$\lambda = \frac{1}{2m} \left( \pm \sqrt{4k_B^2 - (k_A - k_C)^2} - k_A - 2k_B - k_C \right)$$

$$\lambda_1 = -\omega_1^2 = \frac{1}{2m} \left( k_D - k_A - 2k_B - k_C \right)$$

$$\lambda_2 = -\omega_2^2 = \frac{1}{2m} \left( -k_D - k_A - 2k_B - k_C \right)$$

$$\omega_{1} = \sqrt{\frac{(k_{A} + k_{C} + 2k_{B}) - k_{D}}{2m}}$$

$$\omega_{2} = \sqrt{\frac{(k_{A} + k_{C} + 2k_{B}) + k_{D}}{2m}}$$

The reason why we can equate the eigenfrequency with eigenvalue in such a way is given in the following box

#### Showing eigenfrequency of the matrix here itself is also the eigenvalue

$$\begin{pmatrix} \delta \ddot{x}_1 \\ \delta \ddot{x}_2 \end{pmatrix} = \begin{pmatrix} -A - B & B \\ B & -B - C \end{pmatrix} \begin{pmatrix} \delta x_1 \\ \delta x_2 \end{pmatrix}$$

$$\delta \ddot{x}_1 = (-\omega_A^2 - \omega_B^2) \delta x_1 + \omega_B^2 \delta x_2$$
  
$$\delta \ddot{x}_2 = \omega_B^2 \delta x_1 + (-\omega_B^2 - \omega_C^2) \delta x_2$$

Let's say  $\delta x_i = A_i e^{-i\lambda t}$  (using  $\lambda$  instead of  $\omega$  to reduce eyesore).

$$(A_1 e^{-i\lambda t})(-1)\lambda^2 = (-\omega_A^2 - \omega_B^2)\delta x_1 + \omega_B^2 \delta x_2 (A_2 e^{-i\lambda t})(-1)\lambda^2 = \omega_B^2 \delta x_1 + (-\omega_B^2 - \omega_C^2)\delta x_2$$

$$-(\delta x_1)\lambda^2 = (-\omega_A^2 - \omega_B^2)\delta x_1 + \omega_B^2 \delta x_2$$
$$-(\delta x_2)\lambda^2 = \omega_B^2 \delta x_1 + (-\omega_B^2 - \omega_C^2)\delta x_2$$

$$0 = (-\omega_A^2 - \omega_B^2 + \lambda^2)\delta x_1 + \omega_B^2 \delta x_2$$
$$0 = \omega_B^2 \delta x_1 + (-\omega_B^2 - \omega_C^2 + \lambda^2)\delta x_2$$

Re-writing this whole mess is basically

$$0 = \begin{pmatrix} (-\omega_A^2 - \omega_B^2) - H & \omega_B^2 \\ \omega_B^2 & (-\omega_B^2 - \omega_C^2) - H \end{pmatrix} \begin{pmatrix} \delta x_1 \\ \delta x_2 \end{pmatrix}$$

Hence solving the eigenvalue  $H=-\lambda^2$  for the matrix  $\begin{pmatrix} -A-B & B \\ B & -B-C \end{pmatrix}$  basically yields us with the required  $\lambda^2$  in  $A_i e^{-i\lambda t}$ 

Corresponding eigenvectors are

$$|\omega_1\rangle = \begin{pmatrix} \frac{k_C - k_A - k_D}{2k_B} \\ 1 \end{pmatrix}$$

$$|\omega_2\rangle = \begin{pmatrix} \frac{k_C - k_A + k_D}{2k_B} \\ 1 \end{pmatrix}$$

 $\mathbf{d}$ 

For ridiculous  $k_C$ ,

$$k_D = \sqrt{4k_B^2 + (k_C - k_A)^2} = \sqrt{4k_B^2 + k_C^2 \left(1 - \frac{k_A}{k_C}\right)^2} \approx k_C \left(1 - \frac{k_A}{k_C}\right) = k_C - k_A$$

Then

$$\omega_1 = \sqrt{\frac{k_A + k_C + 2k_B - k_C + k_A}{2m}} = \sqrt{\frac{k_A + k_B}{m}}$$

$$\omega_2 = \sqrt{\frac{k_A + k_C + 2k_B + k_C - k_A}{2m}} = \sqrt{\frac{k_C + k_B}{m}}$$

$$|\omega_1\rangle = \begin{pmatrix} 0\\1 \end{pmatrix}$$

$$|\omega_2\rangle = \begin{pmatrix} \frac{k_C - k_A}{k_B}\\1 \end{pmatrix}$$

 $\mathbf{e}$ 

$$k_A = k_C = 0$$
 hence,

$$\omega_1 = \omega_2 = \sqrt{\frac{k_B}{m}}$$

## Problem 03

 $\mathbf{a}$ 

For unitary matrix we know

$$\det(I) = \det(U^t U) = \det(U^t) \det(U) = 1$$

Hence,

$$\det(U^t \Omega U) = \det(U^t) \det(\Omega U) = \det(U^t) \det(\Omega) \det(U) = \det(\Omega)$$

Proven.

b

Note that what we have here is a diagonal matrix. Determinant of a diagonal matrix is given by product of all diagonal elements of the matrix. Hence, for the given matrix,

$$\det U = e^{i\omega_1} e^{i\omega_2} \cdots e^{i\omega_n} = e^{i(\sum_{i=1}^n \omega_i)}$$

Now take a look at  $\log U$ , it is,

$$\begin{pmatrix}
i\omega_1 & 0 & 0 & 0 & \cdots & 0 \\
0 & i\omega_2 & 0 & 0 & \cdots & 0 \\
0 & 0 & i\omega_3 & 0 & \cdots & 0 \\
0 & 0 & 0 & i\omega_4 & \cdots & 0 \\
\vdots & \vdots & \vdots & \vdots & \ddots & 0 \\
0 & 0 & 0 & 0 & 0 & i\omega_n
\end{pmatrix}$$

Trace of this matrix is simply

$$\operatorname{Tr}\log U = i\left(\sum_{i=1}^n \omega_i\right)$$

For this,

$$\det U = e^{\operatorname{Tr} \log U}$$

Proven.

## Problem 04

 $\mathbf{a}$ 

$$(\vec{n} \cdot \vec{\sigma})^2 = (\vec{n} \cdot \vec{\sigma}) (\vec{n} \cdot \vec{\sigma}) = \sum_{i=1}^{3} \sum_{i=1}^{3} n_i n_j \sigma_i \sigma_j$$

We will break the sum between i = j and  $i \neq j$  cases, we get

$$= \sum_{\mu \neq \nu} n_{\mu} n_{\nu} \sigma_{\mu} \sigma_{\nu} + \sum_{k=1}^{3} n_{k}^{2} \sigma_{k} \sigma_{k}$$

Because of being real numbers, obviously  $n_{\nu}n_{\mu} = n_{\mu}n_{\nu}$ . The first sum goes over both  $(\mu, \nu)$  indices and also  $(\nu, \mu)$  indices. Hence because  $\sigma_{\mu}\sigma_{\nu} + \sigma_{\nu}\sigma_{\mu} = 2\delta_{\mu\nu}I$ , and as  $\mu \neq \nu$  we have the first term equal 0.

For the second term with same indices we see  $2\sigma_k\sigma_k=2I$  hence  $\sigma_k\sigma_k=I$ . Hence the sum overall ends up being,

$$(\vec{n} \cdot \vec{\sigma})^2 = (n_1^2 + n_2^2 + n_3^2)I = I$$

As  $(n_1, n_2, n_3)$  is a unit vector. Proven.

b

I will do this first,

$$i = i \quad i^{2} = -1 \quad i^{3} = i^{2}i = -i \quad i^{4} = 1$$

$$U_{1} = \exp(-i\vec{\phi} \cdot \vec{\sigma}) = 1 + (-i\phi\vec{n}_{\phi} \cdot \vec{\sigma}) + \frac{1}{2!}(-i\phi\vec{n}_{\phi} \cdot \vec{\sigma})^{2} + \frac{1}{3!}(-i\phi\vec{n}_{\phi} \cdot \vec{\sigma})^{3} + \frac{1}{4!}(-i\phi\vec{n}_{\phi} \cdot \vec{\sigma})^{4} + \frac{1}{5!}(-i\phi\vec{n}_{\phi}\vec{\sigma})^{5} + \cdots$$

Isolating the terms individually, firstly we note that even power on  $\vec{n}_{\phi} \cdot \vec{\sigma}$  is going to be I (last problem).

Every even terms hence become the series,

$$1 - \frac{1}{2!}(\phi^2)I + \frac{1}{4!}(\phi^4)I + \cdots$$

For every odd terms of power n, the previous term n-1 is even so,  $(\vec{n}_{\phi} \cdot \vec{\sigma})^{(n-1)+1} = (\vec{n}_{\phi} \cdot \vec{\sigma})$ . Like this every odd terms become,

$$i\left(-\phi + \frac{1}{3!}\phi^3 - \frac{1}{5!}\phi^5\right)(\vec{n}_\phi \cdot \vec{\sigma})$$

So the two even and odd terms together combine to form,

$$U_1 = \cos(\phi)\hat{I} - i\sin(\phi)(\vec{n}_{\phi} \cdot \hat{\vec{\sigma}})$$

Proven.

 $\mathbf{c}$ 

 $\vec{\phi} = \phi \vec{n}_{\phi}$  so using that

$$\begin{split} U_1 &= \exp(-i\phi \vec{n}_\phi \cdot \vec{\sigma}) \\ &\log(U_1) = -i\phi \left( \vec{n}_\phi \cdot \vec{\sigma} \right) \\ &\log(U_1) = -i\phi (n_\phi^1 \sigma^1 + n_\phi^2 \sigma^2 + n_\phi^3 \sigma^3) \\ &= -i\phi \left( n_\phi^1 \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + n_\phi^2 \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + n_\phi^3 \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right) \end{split}$$

$$\operatorname{Tr}(\log U_1) = -i\phi \left( n_\phi^1 \operatorname{Tr} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + n_\phi^2 \operatorname{Tr} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + n_\phi^3 \operatorname{Tr} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right) = 0$$

As we know

$$\det U_1 = e^{\text{Tr}\log(U_1)} = e^0 = 1$$

Proven.

# Problem 05

The core idea (after a lot of thinking) is that the introduction of a  $\lambda$  factor and taking derivative with respect to that brings down the operator  $\hat{A}$  from  $e^{\lambda \hat{A}}$ . It's like applying a trick to get the stubborn kid off the tree house.

Let's start with  $g(\lambda) = e^{\lambda A} B e^{-\lambda A}$  as instructed. Taking derivatives at  $\lambda = 0$ 

$$\begin{split} g(0) &= B \\ g'(0) &= (e^{\lambda A}ABe^{-\lambda A} + e^{\lambda A}B(-A)e^{-\lambda A})_{\lambda = 0} = AB - BA = [A, B] = 0 \\ g''(0) &= \left(e^{\lambda A}A[A, B]e^{-\lambda A} + e^{\lambda A}[A, B](-A)e^{-\lambda A}\right)_{\lambda = 0} = e^{\lambda A}[A, [A, B]]e^{-\lambda A}\big|_{\lambda = 0} = [A, [A, B]] \\ &\vdots \quad \text{(induction)} \\ g^{(k)}(0) &= [A, [A, \cdots (k - \text{times}) \cdots, [A, B] \cdots]] = [A, \cdot]^k B \end{split}$$

As given in the series,

$$g(\lambda) = \sum_{k=0}^{\infty} \frac{\lambda^k}{k!} g^k(0)$$

But we just solved for the particular form of  $g^k(0)$ , which gives us,

$$g(1) = e^A B e^{-A} = B + \frac{1}{1!} [A, B] + \frac{1}{2!} [A, [A, B]] + \frac{1}{3!} [A, [A, A, B]] + \cdots$$

Proven.