

Quantum Mechanics Homework 01

September 7, 2024

Ahmed Saad Sabit, Rice University

Problem 1 (a)

Studying the problem

The initial configuration of the string

$$q(x, 0) = f(x) = \begin{cases} \frac{2h}{L}x & 0 \leq x \leq \frac{L}{2} \\ 2h - \frac{2h}{L}x & \frac{L}{2} \leq x \leq L \end{cases}$$

$$\frac{\partial q(x, t)}{\partial t} \Big|_{t=0} = \sum_{n=1}^{\infty} d_n \Omega_n \phi_n(x) = g(x) = 0 \implies \boxed{d_n = 0}$$

The general solution to the string equation (assumed solution is separable between time and position)

$$q(x, t) = \sum_{n=1}^{\infty} [c_n \cos(\Omega_n t) + d_n \sin(\Omega_n t)] \phi_n(x)$$

For $t = 0$ we get,

$$q(x, 0) = f(x) = \sum_{n=1}^{\infty} c_n \phi_n(x)$$

We are interested on finding the general solution of $q(x, t)$ that will hold for the future given this initial condition. The variables of our equation are obviously x, t and what we need to find out is c_n, d_n . The next sub-section will find out a solution for c_n (d_n is trivially zero given zero initial velocity).

Solving for c_n

Let us do the following computation now. Let us multiply both sides of the above equation with $\phi_p(x)$ where p represents the p -th term while we take a summation over the index of n .

$$f(x) \phi_p(x) = \sum_{n=1}^{\infty} c_n \phi_n(x) \phi_p(x)$$

Just so that we can invoke the inner product between orthonormal bases, we can take an integral with the following way

$$\begin{aligned} \int_0^L dx f(x) \phi_p(x) &= \int_0^L dx \left(\sum_{n=1}^{\infty} c_n \phi_n(x) \phi_p(x) \right) \\ &= \sum_{n=1}^{\infty} c_n \int_0^L dx \phi_n(x) \phi_p(x) \\ &= \sum_{n=1}^{\infty} c_n \delta_{np} \frac{L}{2} \\ &= c_p \frac{L}{2} \end{aligned}$$

This above gives us the p -th term

$$c_p = \frac{2}{L} \int_0^L dx f(x) \phi_p(x)$$

Using the explicit equation for the bases and also looking at the piecewise function, we can write,

$$\begin{aligned} c_p &= \frac{2}{L} \int_0^L dx f(x) \sin\left(\frac{p\pi x}{L}\right) \\ &= \frac{2}{L} \left(\int_0^{\frac{L}{2}} f(x) \sin\left(\frac{p\pi x}{L}\right) + \int_{\frac{L}{2}}^L f(x) \sin\left(\frac{p\pi x}{L}\right) \right) \\ &= \frac{2}{L} \left(\int_0^{\frac{L}{2}} \frac{2h}{L} x \sin\left(\frac{p\pi x}{L}\right) + \int_{\frac{L}{2}}^L \left(2h - \frac{2h}{L} x\right) \sin\left(\frac{p\pi x}{L}\right) \right) \\ &= \frac{2}{L} \left(\frac{hL}{\pi^2 p^2} \left[2 \sin\left(p \frac{\pi}{2}\right) - \pi p \cos\left(p \frac{\pi}{2}\right) \right] - \frac{hL}{\pi^2 p^2} \left[2 \sin(\pi p) - 2 \sin\left(p \frac{\pi}{2}\right) - \pi p \cos\left(p \frac{\pi}{2}\right) \right] \right) \\ &= \frac{8h}{\pi^2 p^2} \sin\left(p \frac{\pi}{2}\right) \left[1 - \cos\left(p \frac{\pi}{2}\right) \right] \end{aligned}$$

Hence if I write this huge mess properly

$$c_p = \frac{8h}{\pi^2 p^2} \sin\left(\frac{p\pi}{2}\right) \left[1 - \cos\left(\frac{p\pi}{2}\right) \right]$$

Discussion on odd and even modes

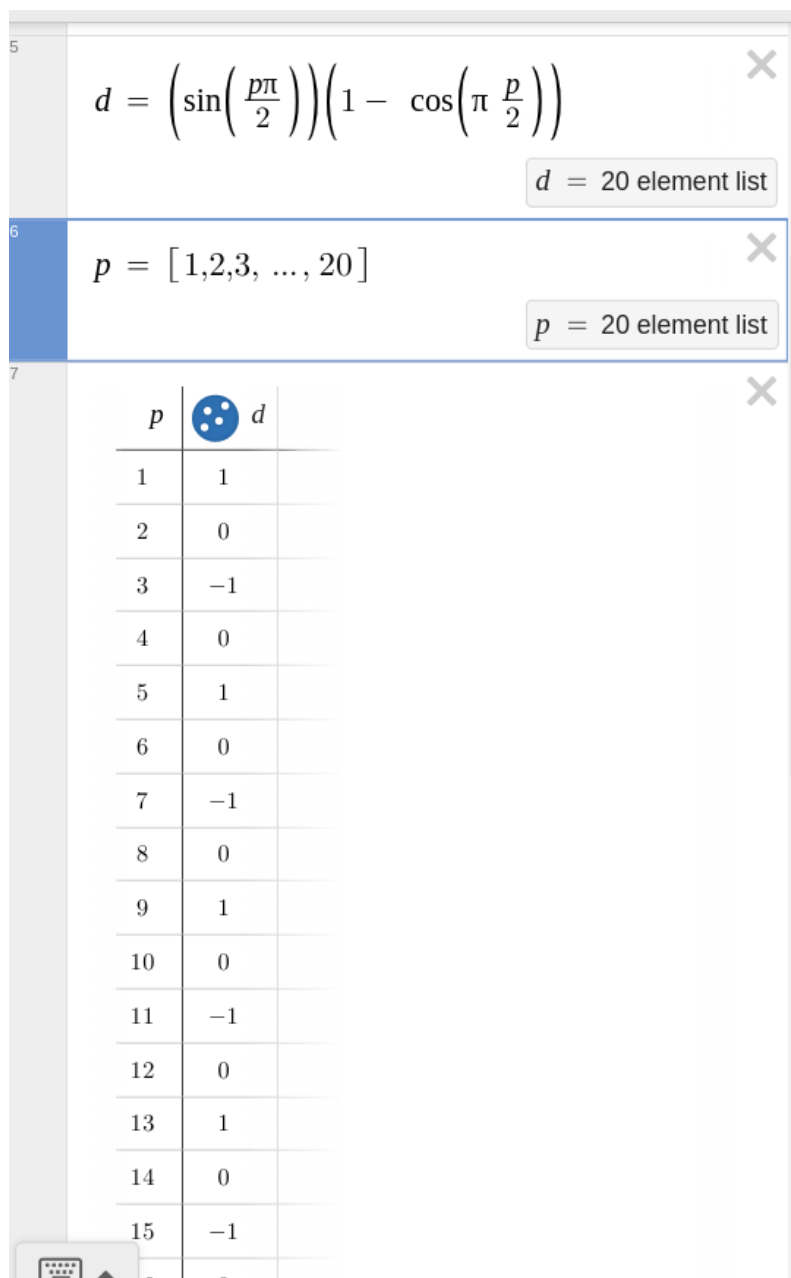


Figure 1: ss/graph-cn.png

We can see that every *even* numbers is giving us a 0 for c_p (d term in graph). The *odd* terms alter between -1 and 1 . The contribution of even and odd modes is obvious from here.

Plotting the displacement with time

Let's set $L = 1$ and we will look at $0 \leq x \leq 1$.

As given $\Omega_n = k_n = \frac{n\pi}{L} = n\pi$.

For h let's pick $h = 0.4$.

Using the c_p we derived above, the displacement function $q(x, t)$ is then,

$$q(x, t)_{(10)} = \sum_{n=1}^{10} c_n \cos(n\pi t) \sin(n\pi x)$$

Expanding c_n ,

$$q(x, t)_{(10)} = \frac{8(0.4)}{\pi^2 n^2} \sum_{n=1}^{10} \sin\left(\frac{n\pi}{2}\right) \left[1 - \cos\left(\frac{n\pi}{2}\right)\right] \cos(n\pi t) \sin(n\pi x)$$

And $\tau = \frac{2\pi}{\Omega_1} = \frac{2\pi}{\pi} = 2$

Plots of $q(x, t)$ where $\tau = 2$

$$q(x, t)_{(10)} = \frac{8(0.4)}{\pi^2} \sum_{n=1}^{10} \frac{1}{n^2} \sin\left(\frac{n\pi}{2}\right) \left[1 - \cos\left(\frac{n\pi}{2}\right)\right] \cos(n\pi t) \sin(n\pi x) \quad \text{and } x \in [0, 1]$$

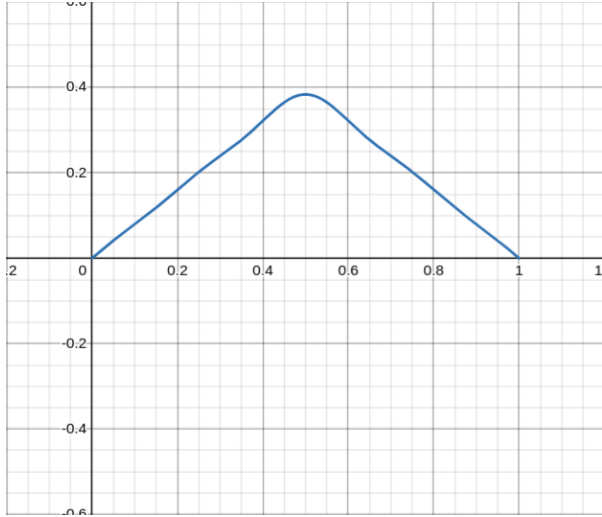


Figure 2: $t = 0$ and also $t = \tau$

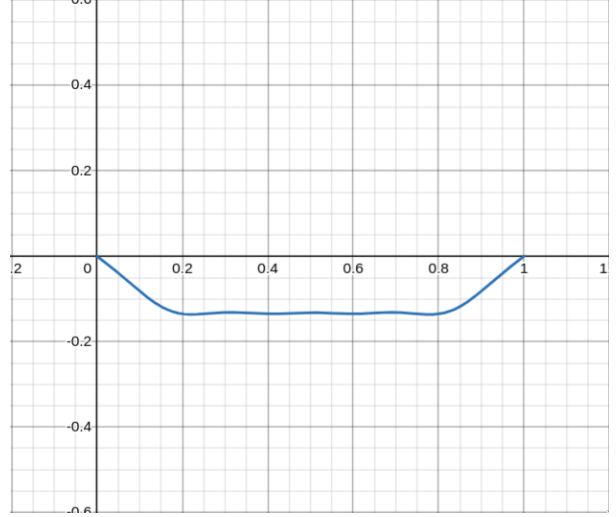


Figure 4: $t = \frac{\tau}{3} = \frac{2}{3}$

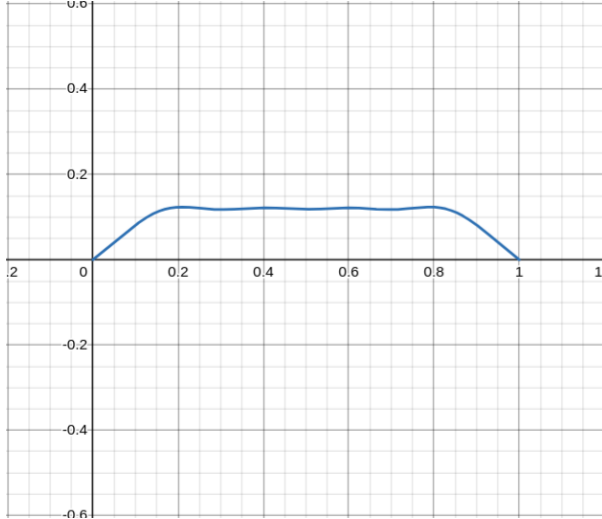


Figure 3: $t = 0.35$

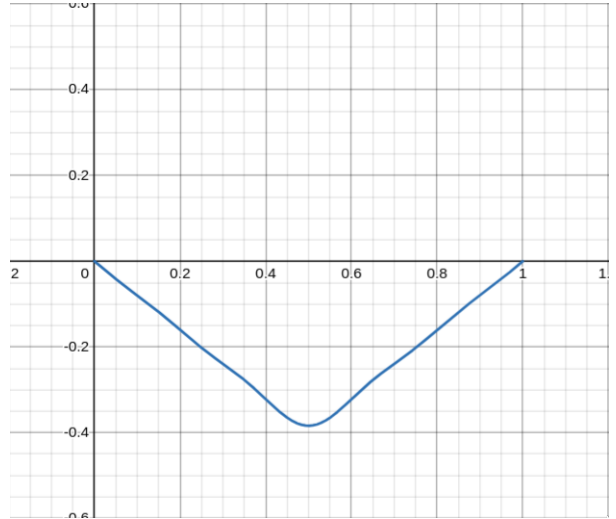


Figure 5: $t = \frac{\tau}{2} = 1$

Problem 1 (b)

Studying the Problem

Initially the string is tight and straight hence

$$q(x, 0) = \sum_{n=1}^{\infty} c_n \phi_n(x) = f(x) = 0 \implies \boxed{c_n = 0}$$

$$\frac{\partial q(x, t)}{\partial t} \Big|_{t=0} = g(x) = v_0 \theta \left(a - \left| x - \frac{L}{2} \right| \right)$$

Where $\theta(x)$ is a Heaviside step function (outputs 1 whenever input is 0 or positive). I am not going to waste my and graders time by re-writing everything I wrote above, the procedure we are going to follow is same as above.

Computation of d_n

$$\begin{aligned} g(x) &= \sum_{n=1}^{\infty} d_n \Omega_n \phi_n(x) \\ \int_0^L dx g(x) \phi_p(x) &= \sum_{n=1}^{\infty} \int_0^L dx d_n \Omega_n \phi_n(x) \phi_p(x) \\ \int_{\frac{L}{2}-a}^{\frac{L}{2}+a} v_0 \phi_p(x) &= d_p \Omega_p \frac{L}{2} \\ \int_{\frac{L}{2}-a}^{\frac{L}{2}+a} v_0 \sin(k_p x) &= d_p k_p \frac{L}{2} \\ \frac{v_0}{k_p} [-\cos(k_p x)]_{x=\frac{L}{2}-a}^{x=\frac{L}{2}+a} &= d_p k_p \frac{L}{2} \\ \cos\left(k_p \frac{L}{2} - k_p a\right) - \cos\left(k_p \frac{L}{2} + k_p a\right) &= \frac{d_p k_p^2 L}{2v_0} \\ \cos\left(\frac{n\pi}{2} - \frac{n\pi a}{L}\right) - \cos\left(\frac{n\pi}{2} + \frac{n\pi a}{L}\right) &= \frac{d_p n^2 \pi^2}{2v_0 L} \end{aligned}$$

This gives us

$$d_p = \frac{2v_0 L}{n^2 \pi^2} \left[\cos\left(\frac{n\pi}{2} - \frac{n\pi a}{L}\right) - \cos\left(\frac{n\pi}{2} + \frac{n\pi a}{L}\right) \right] = \frac{4v_0 L}{n^2 \pi^2} \sin\left(\frac{n\pi}{2}\right) \sin\left(\frac{n\pi a}{L}\right)$$

Discussion on Even and Odd modes

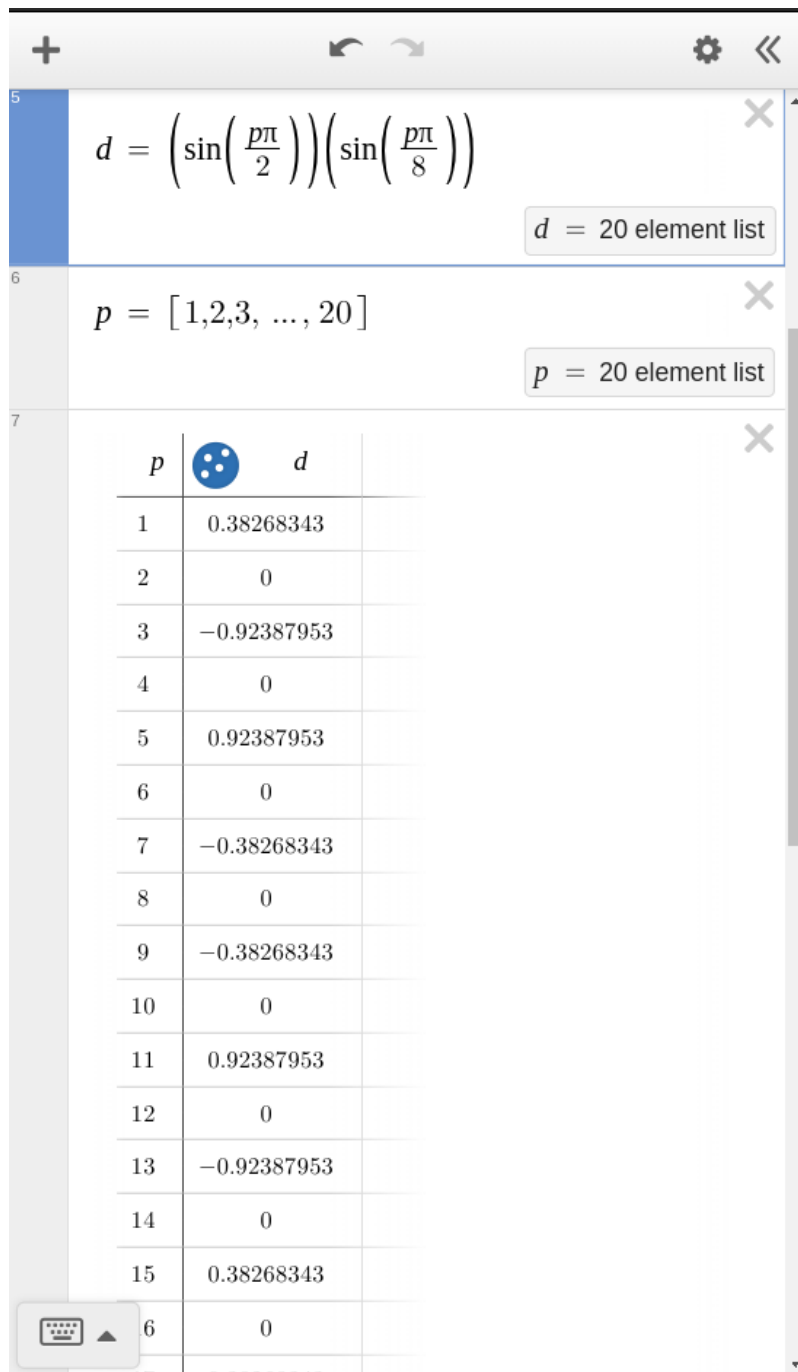


Figure 6: ss/graph-dn

Like before we can see every *even* modes are 0.

Plotting the displacement with time

Let's set $L = 1$, then $a = 1/8$. Let's set $v_0 = 1.5$. As before $\tau = 2$.

Plots of $q(x, t)$ where $\tau = 2$

$$q(x, t)_{(10)} = \frac{4(1.5)}{\pi^2} \sum_{n=1}^{10} \frac{1}{n^2} \sin\left(\frac{n\pi}{2}\right) \sin(n\pi(1/8)) \sin(n\pi t) \sin(n\pi x) \quad \text{and } x \in [0, 1]$$

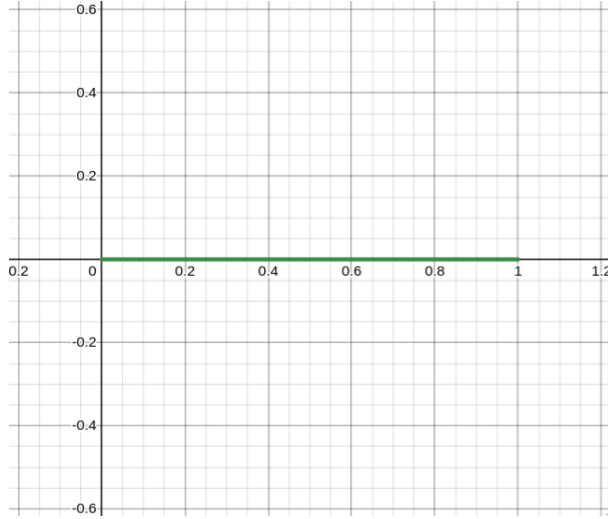


Figure 7: $t = 0$ and also $t = \tau$



Figure 9: $t = \frac{\tau}{3} = \frac{2}{3}$

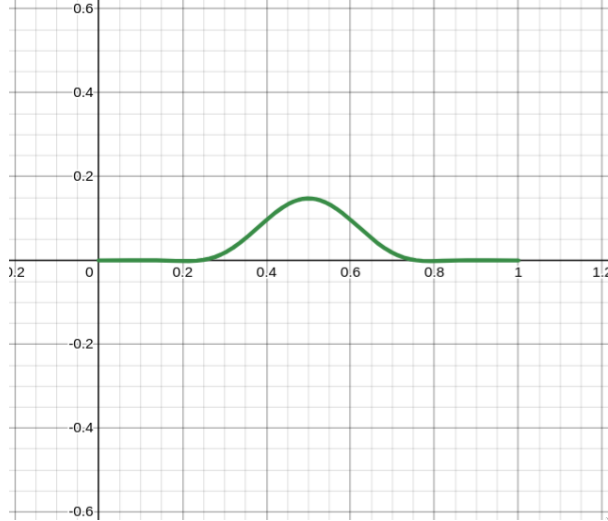


Figure 8: $t = \frac{\tau}{20} = 0.1$

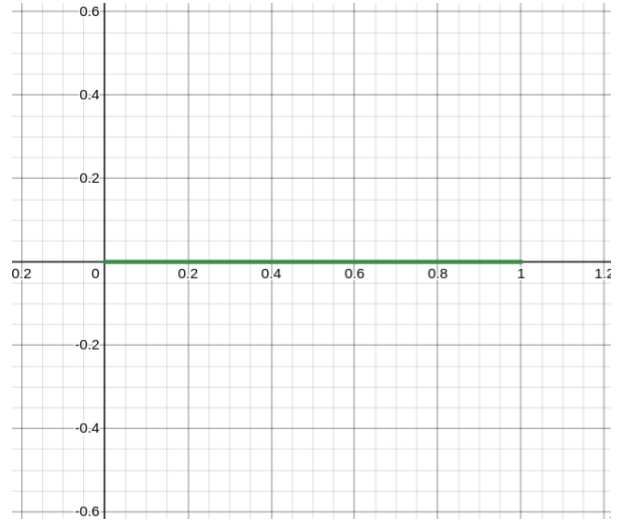


Figure 10: $t = \frac{\tau}{2} = 1$

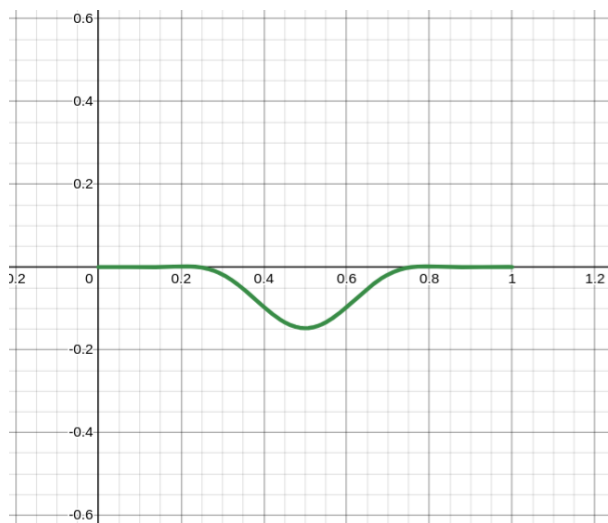


Figure 11: $t = 1.1$

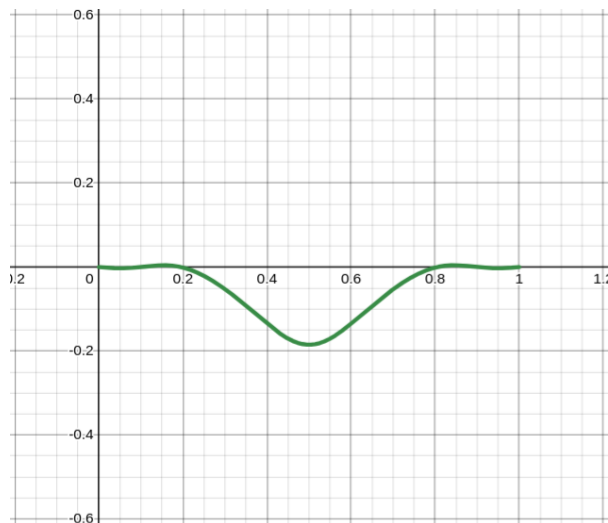


Figure 13: $t = 1.85$

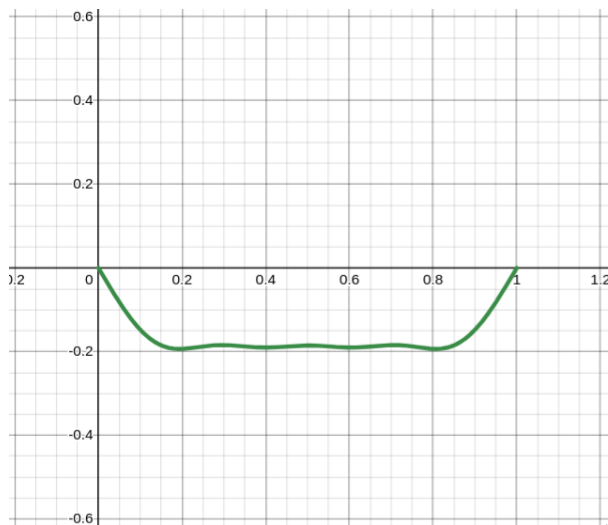


Figure 12: $t = 1.5$

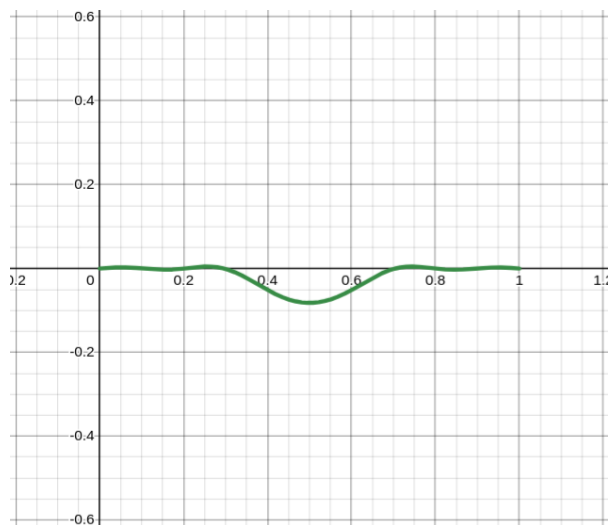


Figure 14: $t = 1.95$

Plotting the Initial Velocity

We can simply take the time derivative and make a plot. The more iterations we make, we should get close to a heaviside function as stated in the initial velocity profile.

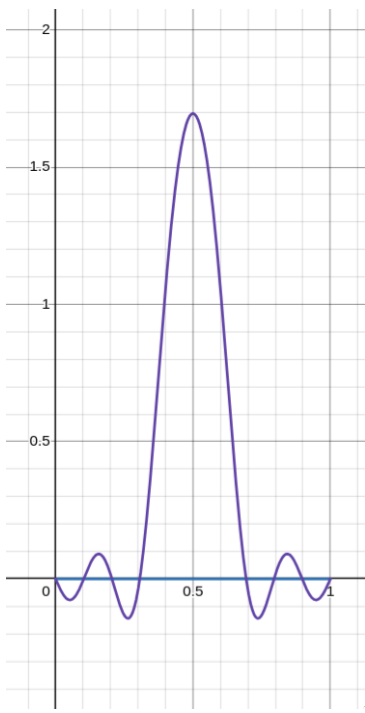


Figure 15: Plotted for $n = 10$

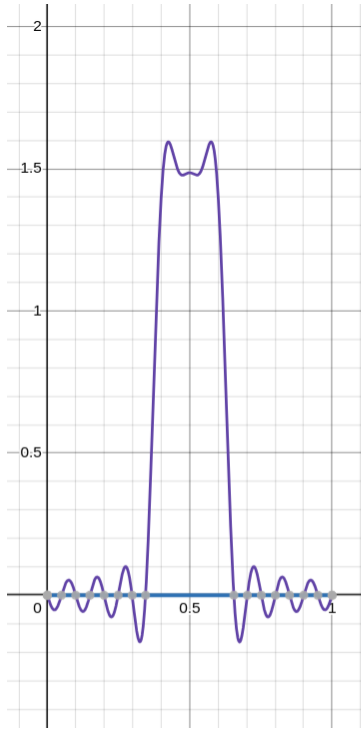


Figure 16: Plotted for $n = 20$

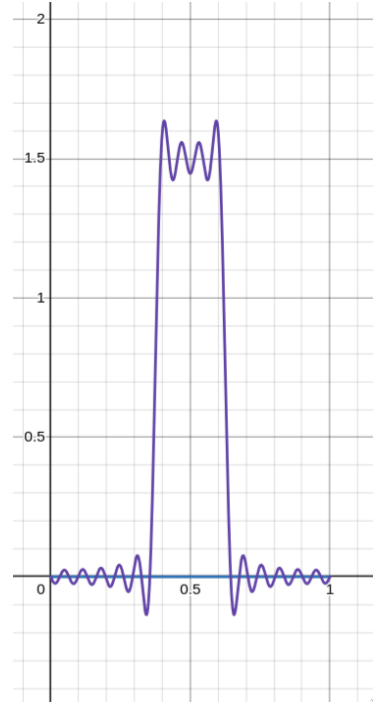


Figure 17: Plotted for $n = 30$

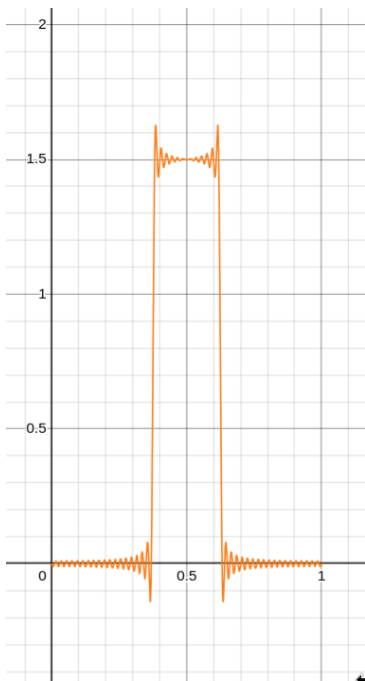


Figure 18: Plotted for $n = 100$

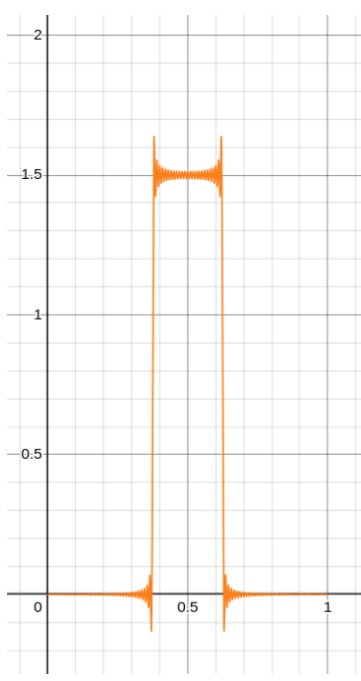


Figure 19: Plotted for $n = 200$

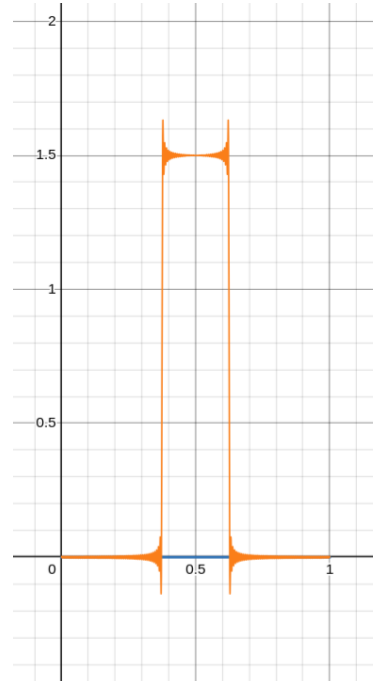


Figure 20: Plotted for $n = 300$

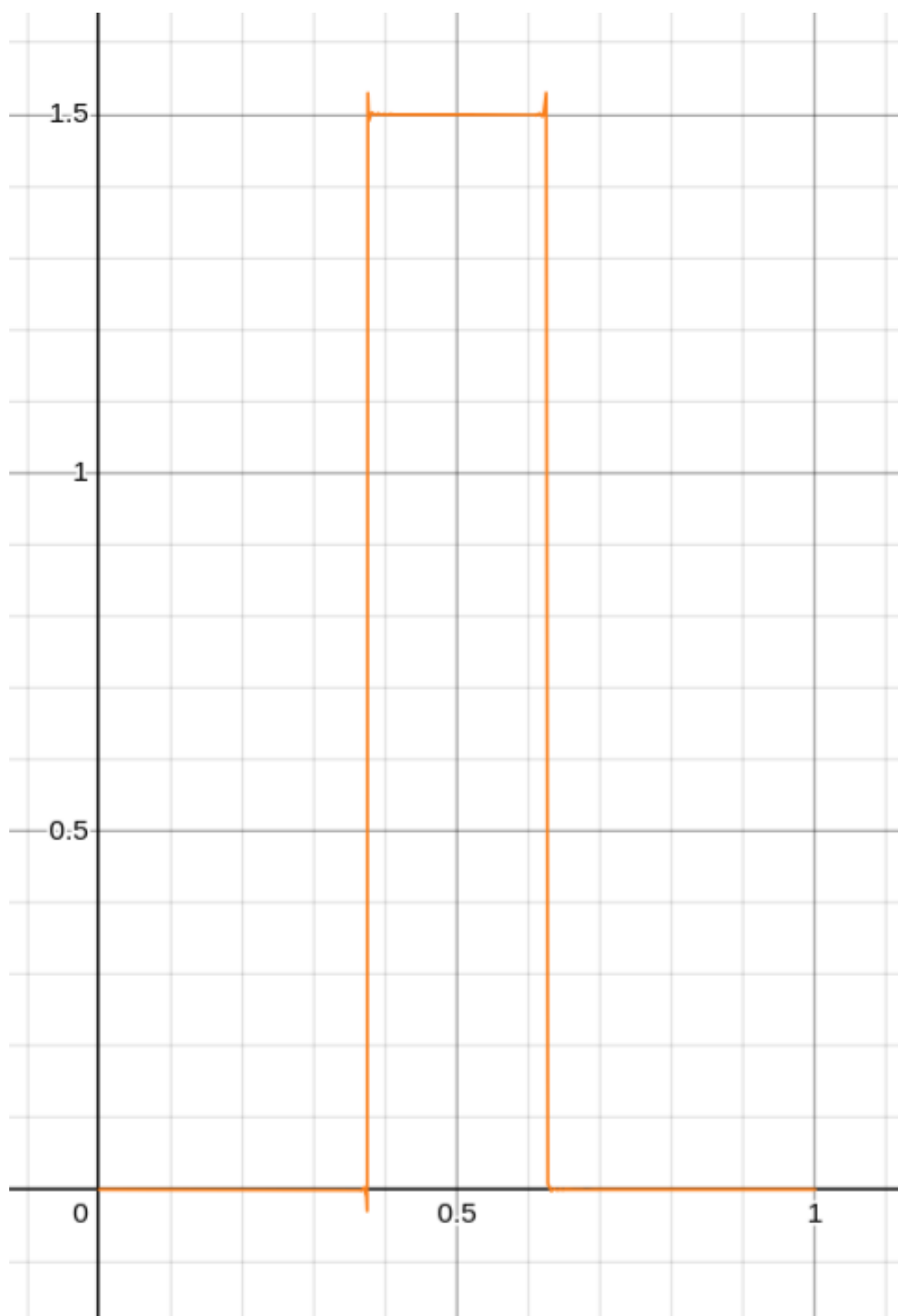


Figure 21: Waited 2 minutes to try $n = 10,000$. We have the exact initial conditions we wanted. This is absolutely beautiful to do this myself.

Problem 2

Linear Dependence equation

$$a \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} + b \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} + c \begin{pmatrix} -2 & -1 \\ 0 & -2 \end{pmatrix} = 0$$

If the three above equations are linearly independent we are going to have the only possible solution of the equality being zero as $\{a, b, c\} = \{0, 0, 0\}$. We get three different equations

$$\begin{aligned} b - 2c &= 0 \\ a + b - c &= 0 \\ b - 2c &= 0 \\ 0 + 0 &= 0 \end{aligned}$$

Where the I didn't count the last one as it's not really anything. If we solve the system of equation above we are left with (did the computation by hand on paper)

$$\begin{aligned} a &= -c \\ b &= 2c \\ c &= \text{can be anything} \end{aligned}$$

a, b, c doesn't necessarily need to be 0 to yield above equation to be zero. This system is *NOT linearly independent*. For example, setting $c = 1$

$$- \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} + 2 \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} + \begin{pmatrix} -2 & -1 \\ 0 & -2 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

Problem 3

Deduction

$$\begin{aligned} |v + w|^2 &= \langle v + w | v + w \rangle \\ &= \langle v | v + w \rangle + \langle w | v + w \rangle \\ &= \langle v | v \rangle + \langle v | w \rangle + \langle w | v \rangle + \langle w | w \rangle \\ &= |v|^2 + |w|^2 + \langle v | w \rangle + \langle v | w \rangle^* \\ &= |v|^2 + |w|^2 + 2\text{Re}(\langle v | w \rangle) \\ &\leq |v|^2 + |w|^2 + 2|\langle v | w \rangle| \\ &\leq |v|^2 + |w|^2 + 2|v||w| \\ &\leq (|v| + |w|)^2 \end{aligned}$$

This shows that

$$|v + w| \leq |v| + |w|$$

Equality condition

Above we can see

$$\begin{aligned} |v + w|^2 &= \langle v + w | v + w \rangle \\ &= |v|^2 + |w|^2 + 2\text{Re}(\langle v | w \rangle) \end{aligned}$$

For the equality to hold we strictly need

$$2\text{Re}(\langle v|w\rangle) = 2|v||w|$$

Let's try $|v\rangle = a|w\rangle$.

$$2\text{Re}(\langle aw|w\rangle) = 2a\text{Re}(\langle w|w\rangle) = 2|aw||w| = 2a|w||w|$$

$|v\rangle = a|w\rangle$ satisfies this. If $|v\rangle$ has component perpendicular to $|w\rangle$ (in a way that $|v\rangle$ can be broken down into constituent factors of basis vectors such that inner product with $|w\rangle$ gives zero), then some of the value of $2\text{Re}(\langle v|w\rangle)$ is getting lost. Having the two vectors parallel gives equal to their norm. This is more of a physical vector-like intuition.

Problem 4(a)

I will write $\hat{\sigma}^n$ as simply σ^n for this problem. I did the multiplication by hand.

$$(\sigma^1)^2 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$(\sigma^2)^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$(\sigma^3)^2 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

We had been already defined

$$\sigma^0 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

All of these result matrix above is the identity matrix that validates

$$(\sigma^1)^2 = (\sigma^2)^2 = (\sigma^3)^2 = \sigma^0$$

Problem 4(b)

We are required to solve for $A^{\mu,\nu}$ where

$$A^{\mu,\nu} = \sigma^\mu \sigma^\nu + \sigma^\nu \sigma^\mu$$

Note that it is obvious

$$A^{\mu,\nu} = A^{\nu,\mu}$$

Computing each of the matrix multiplications, and also referring to previous computations

$$A^{k,k} = A^{1,1} = A^{2,2} = A^{3,3} = 2(\sigma^k)^2 = 2\sigma^0 = 2 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\begin{aligned} A^{1,2} &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ &= \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} + \begin{pmatrix} -i & 0 \\ 0 & i \end{pmatrix} \\ &= \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} \end{aligned}$$

$$\begin{aligned}
A^{2,3} &= \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \\
&= \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} + \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix} \\
&= \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}
\end{aligned}$$

$$\begin{aligned}
A^{1,3} &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\
&= \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \\
&= \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}
\end{aligned}$$

Different indexes cause a zero-matrix, and similar causes a double of identity matrix. From here we can easily figure out that

$$A^{\mu,\nu} = 2\sigma^0\delta_{\mu,\nu}$$

Problem 4(c)

I am going to borrow the computations I did last problem

$$\text{Tr}[\sigma^1\sigma^1] = \text{Tr}\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = 2$$

$$\text{Tr}[\sigma^2\sigma^2] = \text{Tr}\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = 2$$

$$\text{Tr}[\sigma^3\sigma^3] = \text{Tr}\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = 2$$

$$\text{Tr}[\sigma^1\sigma^2] = \text{Tr}\begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} = 0 = \text{Tr}[\sigma^2\sigma^1]$$

$$\text{Tr}[\sigma^2\sigma^3] = \text{Tr}\begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} = 0 = \text{Tr}[\sigma^3\sigma^2]$$

$$\text{Tr}[\sigma^1\sigma^3] = \text{Tr}\begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = 0 = \text{Tr}[\sigma^3\sigma^1]$$

From this we can see that same-index gives 2 and different gives 0. From this it's obvious

$$\text{Tr}[\sigma^\mu\sigma^\nu] = 2\delta_{\mu\nu}$$

Problem 4(d)

Expanding the equation of the operator

$$\hat{V} = \sum_{i=1}^3 V_i \sigma^i = V_1 \sigma^1 + V_2 \sigma^2 + V_3 \sigma^3$$

Multiply σ^p where $p \in \{1, 2, 3\}$

$$\hat{V} \sigma^p = V_1 \sigma^1 \sigma^p + V_2 \sigma^2 \sigma^p + V_3 \sigma^3 \sigma^p$$

Taking the trace and using the property $\text{Tr}(A + B) = \text{Tr}(A) + \text{Tr}(B)$

$$\text{Tr}(\hat{V} \sigma^p) = V_1 (2\delta_{1p}) + V_2 (2\delta_{2p}) + V_3 (2\delta_{3p})$$

From this using the definition of the $\delta_{\mu,p}$ we can simply write,

$$V_p = \frac{1}{2} \text{Tr}(\hat{V} \sigma^p)$$

Using the form

$$\hat{V} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

we can find the coefficients V_p

$$\begin{aligned} \hat{V} \sigma^0 &= \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \implies V_0 = \frac{a+d}{2} \\ \hat{V} \sigma^1 &= \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} b & a \\ d & c \end{pmatrix} \implies V_1 = \frac{b+c}{2} \\ \hat{V} \sigma^2 &= \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} ib & -ia \\ id & -ic \end{pmatrix} \implies V_2 = i \frac{b-c}{2} \\ \hat{V} \sigma^3 &= \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} a & -b \\ c & -d \end{pmatrix} \implies V_3 = \frac{a-d}{2} \end{aligned}$$

So our representation is

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \frac{a+d}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \frac{b+c}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + \frac{ib-ic}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + \frac{a-d}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

I've checked this in Wolfram Alpha and it seems to work.

Problem 5(a)

From the question, my understanding of the basis is

$$|1\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad |2\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Projection operator

$$P_1 = |1\rangle\langle 1| = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$P_2 = |2\rangle\langle 2| = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \begin{pmatrix} 0 & 1 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

We've determined the elements of the matrices.

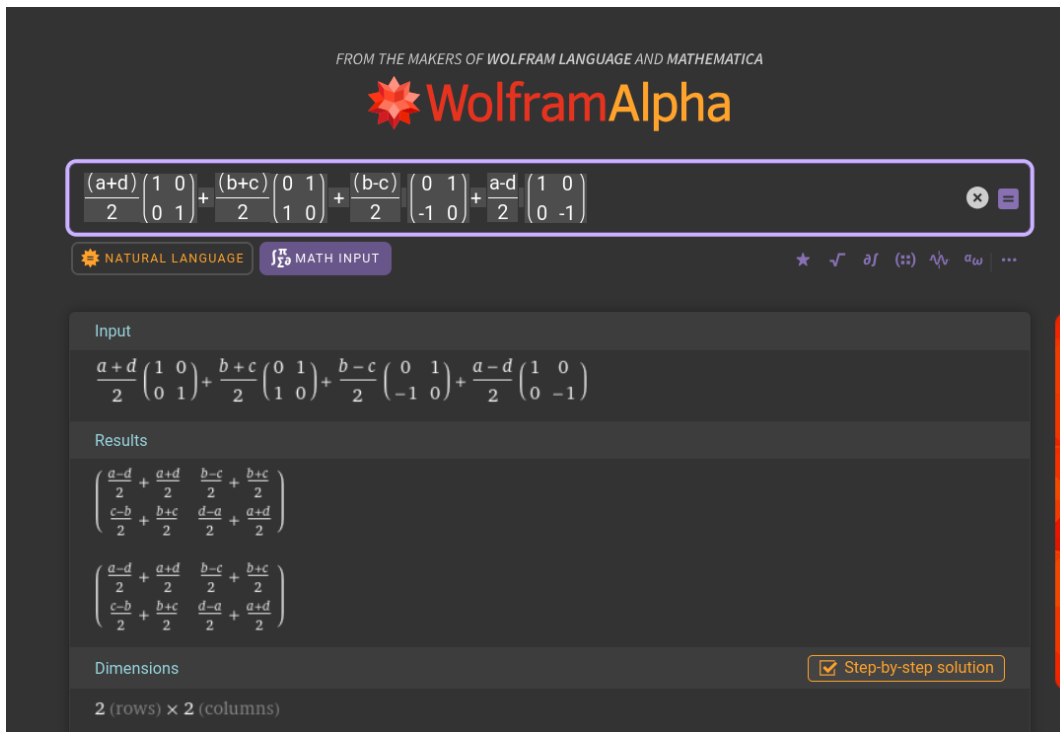


Figure 22: ss/pauli-basis.png

Problem 5(b)

Doing the matrix multiplication by hand

$$P_1 P_1 = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$P_2 P_2 = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$P_1 P_2 = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

$$P_2 P_1 = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

My intuition for the projection operator is basically

$$P_i P_j = (|i\rangle\langle i|)(|j\rangle\langle j|) = |i\rangle\langle i|j\rangle\langle j| = |i\rangle\delta_{ij}\langle j| = \delta_{ij}|i\rangle\langle j| = |k\rangle\langle k|$$

where $k = \{i, j\}$

Problem 5(c)

Computation

$$|V\rangle = \frac{1}{\sqrt{2}}(|1\rangle + |2\rangle)$$

$$P_V = |V\rangle\langle V| = \frac{1}{2}(|1\rangle + |2\rangle)(\langle 1| + \langle 2|) = \frac{1}{2}(|1\rangle\langle 1| + |1\rangle\langle 2| + |2\rangle\langle 1| + |2\rangle\langle 2|)$$

$$P_V = \frac{1}{2} \left[\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \right] = \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$$

Now computing

$$\begin{aligned} P_1 P_V P_2 &= \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \left[\frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \right] \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \\ &= \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix} \\ &= \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \end{aligned}$$

We could have also done this directly like

$$\begin{aligned} P_1 P_V P_2 &= \frac{1}{2} (|1\rangle\langle 1|) (|1\rangle\langle 1| + |1\rangle\langle 2| + |2\rangle\langle 1| + |2\rangle\langle 2|) (|2\rangle\langle 2|) \\ &= \frac{1}{2} (|1\rangle\langle 1|) [|1\rangle\langle 1|2\rangle\langle 2| + |1\rangle\langle 2|2\rangle\langle 2| + |2\rangle\langle 1|2\rangle\langle 1| + |2\rangle\langle 2|2\rangle\langle 2|] \\ &= \frac{1}{2} (|1\rangle\langle 1|) (|1\rangle\langle 2| + |2\rangle\langle 2|) \\ &= \frac{1}{2} (|1\rangle\langle 1|1\rangle\langle 2| + |1\rangle\langle 1|2\rangle\langle 2|) \\ &= \frac{1}{2} |1\rangle\langle 2| \\ &= \frac{1}{2} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \end{pmatrix} \\ &= \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \end{aligned}$$

Analysis

We can keep three polarizers, each at a increasing angle, then components of lights get filtered as the pass through, remaining light falls on the screen behind.