

# Quantum Mechanics : : Homework 04

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## Problem 01

**a**

I do the matrix multiplication by hand.

$$\begin{aligned}\hat{\sigma}_1 \hat{\sigma}_2 &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} = i\hat{\sigma}_3 \\ \hat{\sigma}_2 \hat{\sigma}_1 &= \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} -i & 0 \\ 0 & i \end{pmatrix} = -i\hat{\sigma}_3 \\ \hat{\sigma}_2 \hat{\sigma}_3 &= \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} = i\hat{\sigma}_1 \\ \hat{\sigma}_3 \hat{\sigma}_2 &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix} = -i\hat{\sigma}_1 \\ \hat{\sigma}_3 \hat{\sigma}_1 &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = i\hat{\sigma}_2 \\ \hat{\sigma}_1 \hat{\sigma}_3 &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = -i\hat{\sigma}_2\end{aligned}$$

The cross product

$$\begin{aligned}\hat{\vec{\sigma}} \times \hat{\vec{\sigma}} &= \begin{vmatrix} \vec{n}_x & \vec{n}_y & \vec{n}_z \\ \sigma_1 & \sigma_2 & \sigma_3 \\ \sigma_1 & \sigma_2 & \sigma_3 \end{vmatrix} \\ &= \vec{n}_x (\sigma_2 \sigma_3 - \sigma_3 \sigma_2) + \vec{n}_y (\sigma_3 \sigma_1 - \sigma_1 \sigma_3) + \vec{n}_z (\sigma_1 \sigma_2 - \sigma_2 \sigma_1) \\ &= \vec{n}_x (2i\sigma_1) + \vec{n}_y (2i\sigma_2) + \vec{n}_z (2i\sigma_3) \\ &= 2i (\sigma_1 \vec{n}_x + \sigma_2 \vec{n}_y + \sigma_3 \vec{n}_z) \\ &= 2i\hat{\vec{\sigma}}\end{aligned}$$

Note this also is alluring to the Levi-Civita Symbol because

$$\text{cyclic } i \rightarrow j \rightarrow k \rightarrow i \rightarrow j \rightarrow k \implies \sigma_i \sigma_j = i\epsilon_{ij} \sigma_k \quad (\text{where } \{i, j, k\} \in \{1, 2, 3\})$$

**a**

$$(\vec{U} \cdot \hat{\vec{\sigma}}) (\vec{V} \cdot \hat{\vec{\sigma}}) = (U_1 \sigma_1 + U_2 \sigma_2 + U_3 \sigma_3) (V_1 \sigma_1 + V_2 \sigma_2 + V_3 \sigma_3)$$

$$\begin{aligned}
(\vec{U} \cdot \hat{\vec{\sigma}}) (\vec{V} \cdot \hat{\vec{\sigma}}) &= \left( \sum_{i=1}^3 U_i \sigma_i \right) \left( \sum_{j=1}^3 V_j \sigma_j \right) \\
&= \sum_{i,j=1}^3 U_i V_j \sigma_i \sigma_j \\
&= \sum_{i,j=1, i \neq j}^3 U_i V_j \sigma_i \sigma_j + \sum_{n=1}^3 U_n V_n \sigma_n \sigma_n \\
&= \sum_{i,j=1, i \neq j}^3 i U_i V_j \epsilon_{ij} \sigma_k + \sum_{n=1}^3 U_n V_n \sigma_n \sigma_n \\
&= [i(U_1 V_2 - U_2 V_1) \sigma_3 + i(U_2 V_3 - U_3 V_2) \sigma_1 + i(U_3 V_1 - U_1 V_3) \sigma_2] + (U_1 V_1 + U_2 V_2 + U_3 V_3) \hat{I} \\
&= [\vec{U} \times \vec{V}]_3 i \sigma_3 + [\vec{U} \times \vec{V}]_1 i \sigma_1 + [\vec{U} \times \vec{V}]_2 i \sigma_2 + (\vec{U} \cdot \vec{V}) \hat{I} \\
&= i [\vec{U} \times \vec{V}] \cdot \hat{\vec{\sigma}} + (\vec{U} \cdot \vec{V}) \hat{I}
\end{aligned}$$

## Problem 02

**a**

Considering the simplest basis in column vector forms

$$|1\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad |2\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad |3\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

We want to compute the following operator  $\hat{H}$

$$\hat{H} = E_0 (|1\rangle\langle 1| + |2\rangle\langle 2| + |3\rangle\langle 3|) - J (|1\rangle\langle 2| + |2\rangle\langle 1| + |2\rangle\langle 3| + |3\rangle\langle 2| + |3\rangle\langle 1| + |1\rangle\langle 3|)$$

Matrices for the first term

$$\begin{aligned} |1\rangle\langle 1| + |2\rangle\langle 2| + |3\rangle\langle 3| &= \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} [1 \ 0 \ 0] + \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix} [0 \ 1 \ 0] + \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} [0 \ 0 \ 1] \\ &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} = \hat{I} \end{aligned}$$

Matrices for the second term

$$\begin{aligned} &|1\rangle\langle 2| + |2\rangle\langle 1| + |2\rangle\langle 3| + |3\rangle\langle 2| + |3\rangle\langle 1| + |1\rangle\langle 3| \\ &= \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} (0 \ 1 \ 0) + \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} (1 \ 0 \ 0) + \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} (0 \ 0 \ 1) + \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} (0 \ 1 \ 0) + \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} (1 \ 0 \ 0) + \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} (0 \ 0 \ 1) \\ &= \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 1 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \\ &= \begin{bmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{bmatrix} \end{aligned}$$

Put them all together

$$\begin{aligned} \hat{H} &= E_0 (|1\rangle\langle 1| + |2\rangle\langle 2| + |3\rangle\langle 3|) - J (|1\rangle\langle 2| + |2\rangle\langle 1| + |2\rangle\langle 3| + |3\rangle\langle 2| + |3\rangle\langle 1| + |1\rangle\langle 3|) \\ &= \begin{bmatrix} E_0 & -J & -J \\ -J & E_0 & -J \\ -J & -J & E_0 \end{bmatrix} \end{aligned}$$

**b (i)**

The  $|E_1\rangle$  in column vector representation

$$|E_1\rangle = \frac{1}{\sqrt{3}} (|1\rangle + |2\rangle + |3\rangle) = \frac{1}{\sqrt{3}} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

Computing  $\hat{H}|E_1\rangle$

$$\begin{aligned}
\hat{H}|E_1\rangle &= \begin{bmatrix} E_0 & -J & -J \\ -J & E_0 & -J \\ -J & -J & E_0 \end{bmatrix} \left( \frac{1}{\sqrt{3}} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} \right) \\
&= \frac{1}{\sqrt{3}} \begin{bmatrix} E_0 - J - J \\ -J + E_0 - J \\ -J - J + E_0 \end{bmatrix} \\
&= \frac{1}{\sqrt{3}} \begin{bmatrix} E_0 - 2J \\ E_0 - 2J \\ E_0 - 2J \end{bmatrix} \\
&= \frac{E_0 - 2J}{\sqrt{3}} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} \\
&= (E_0 - 2J) |E_1\rangle
\end{aligned}$$

We see the eigen-equation for the Hamiltonian

$$\hat{H}|E_1\rangle = E_1|E_1\rangle \implies E_1 = E_0 - 2J$$

**b (ii) and (iii)**

We algebraically initialize the elements of the matrix to minimize the computational effort by hand. Keep in mind that  $\hat{H} - E\hat{I} = \hat{0} = \frac{1}{J}\hat{H} - \frac{1}{J}E\hat{I}$  hence

$$\begin{aligned}
\frac{1}{J}\hat{H} &= \begin{bmatrix} \frac{E_0}{J} & -1 & -1 \\ -1 & \frac{E_0}{J} & -1 \\ -1 & -1 & \frac{E_0}{J} \end{bmatrix} \\
\left( \frac{1}{J}\hat{H} \right) - \frac{E}{J}\hat{I} &= \begin{bmatrix} \frac{E_0-E}{J} & -1 & -1 \\ -1 & \frac{E_0-E}{J} & -1 \\ -1 & -1 & \frac{E_0-E}{J} \end{bmatrix} \\
&= (-1) \begin{bmatrix} \lambda & 1 & 1 \\ 1 & \lambda & 1 \\ 1 & 1 & \lambda \end{bmatrix} = 0 \quad (\lambda = (E - E_0)/J) \\
\implies \det \begin{bmatrix} \lambda & 1 & 1 \\ 1 & \lambda & 1 \\ 1 & 1 & \lambda \end{bmatrix} &= 0 \\
\implies \lambda^3 + 3\lambda + 2 &= 0 \\
\implies (\lambda - 1)^2(\lambda + 2) &= 0 \\
\implies (\lambda^2 - 2\lambda + 1)(\lambda + 2) &= 0 \\
\implies \lambda = 1, 1, -2 &
\end{aligned}$$

As stated in the problem statement,  $0 = (\lambda - \lambda_1)(\lambda^2 + b\lambda + c) \implies 0 = (\lambda - (-2))(\lambda^2 + (-2)\lambda + 1)$  giving us  $(b, c) = (-2, 1)$

$$\text{For } \lambda = -2 \text{ the eigen-energy is } \frac{E - E_0}{J} = -2 \implies E = E_0 - 2J$$

$$\text{For } \lambda = 1 \text{ the eigen-energy is } \frac{E - E_0}{J} = 1 \implies E = E_0 + J$$

The eigenvalues are

$$\lambda_1 = -2, \quad \lambda_2 = 1, \quad \lambda_3 = 1$$

And corresponding eigen-energy

$$E_1 = E_0 - 2J, \quad E_2 = E_0 + J, \quad E_3 = E_0 + J$$

## Problem 03

**a**

$$\begin{aligned}
 \hat{T}|n\rangle = |n+1\rangle &\implies \langle n|T^\dagger = \langle n+1| \\
 &\text{now, } \langle n|T^\dagger|n+1\rangle = \langle n+1|n+1\rangle = 1 \\
 &\text{or, } \langle n|T^\dagger T|n\rangle = 1 \\
 &\text{and, } \langle n|T^\dagger|k+1\rangle = \langle n+1|k+1\rangle = 0 \quad (n \neq k) \\
 &\implies T^\dagger T = \hat{I}
 \end{aligned}$$

**b**

To compute  $[\hat{H}, \hat{T}]$  we compute the two operator terms individually

$$\begin{aligned}
 \hat{H}\hat{T} &= \left( \sum_{n=-\infty}^{\infty} [E_0|n\rangle\langle n| + J(|n+1\rangle\langle n| + |n\rangle\langle n+1|)] \right) \hat{T} \\
 \hat{H}\hat{T}|n\rangle &= \left( \sum_{n=-\infty}^{\infty} [E_0|n\rangle\langle n| + J(|n+1\rangle\langle n| + |n\rangle\langle n+1|)] \right) |n+1\rangle \\
 &= [E_0|n\rangle\langle n| + J(|n+1\rangle\langle n| + |n\rangle\langle n+1|)] |n+1\rangle + [E_0|n+1\rangle\langle n+1| + J(|n+2\rangle\langle n+1| + |n+1\rangle\langle n+2|)] |n+1\rangle \\
 &= J|n\rangle + E_0|n+1\rangle + J(|n+2\rangle) \\
 \hat{T}\hat{H}|n\rangle &= \hat{T} \left( [E_0|n\rangle\langle n| + J(|n+1\rangle\langle n| + |n\rangle\langle n+1|)] |n\rangle + [E_0|n-1\rangle\langle n-1| + J(|n\rangle\langle n-1| + |n-1\rangle\langle n|)] |n\rangle \right) \\
 &= \hat{T}(E_0|n\rangle + J|n+1\rangle + J|n-1\rangle) \\
 &= E_0|n+1\rangle + J|n+2\rangle + J|n\rangle = \hat{H}\hat{T}|n\rangle
 \end{aligned}$$

$$\therefore (\hat{H}\hat{T} - \hat{T}\hat{H})|n\rangle = |0\rangle \implies [\hat{H}, \hat{T}] = 0$$

**c**

Using the form given for the energy, and the eigen-equation for  $\hat{T}$ , we determine the general formula for  $\psi_{E,n}$  in terms of  $\psi_{E,0}$

$$\begin{aligned}
|E\rangle &= \sum_{n=-\infty}^{\infty} |n\rangle \psi_{E,n} \\
T|E\rangle &= \sum_{n=-\infty}^{\infty} T|n\rangle \psi_{E,n} \\
e^{-i\phi}|E\rangle &= \sum_{n=-\infty}^{\infty} |n+1\rangle \psi_{E,n} \\
\sum_{n=-\infty}^{\infty} e^{-i\phi}|n\rangle \psi_{E,n} &= \sum_{n=-\infty}^{\infty} |n+1\rangle \psi_{E,n} \\
\implies e^{-i\phi}|n+1\rangle \psi_{E,n+1} &= |n+1\rangle \psi_{E,n} \\
\therefore e^{-i\phi} &= \frac{\psi_{E,n}}{\psi_{E,n+1}}
\end{aligned}$$

The inductive relation between two coefficient is

$$\begin{aligned}
\psi_{E,n+1} = e^{i\phi} \psi_{E,n} &\implies \psi_{E,1} = e^{i\phi} \psi_{E,0} \implies \psi_{E,2} = e^{i\phi} \psi_{E,1} = e^{2i\phi} \psi_{E,0} \\
&\implies \psi_{E,n} = e^{in\phi} \psi_{E,0}
\end{aligned}$$

The general form of energy is then

$$|E\rangle = \psi_{E,0} \sum_{n=-\infty}^{\infty} e^{in\phi} |n\rangle$$

$$\begin{aligned}
\hat{H}|E_n\rangle &= \psi_{E,n} \hat{H}|n\rangle = \psi_{E,n} (E_0|n\rangle + J|n+1\rangle + J|n-1\rangle) \\
\hat{H}|E\rangle &= \hat{H} \left( \sum_{n=-\infty}^{\infty} \psi_{E,n} |n\rangle \right) = \sum_{n=-\infty}^{\infty} \psi_{E,n} (E_0|n\rangle + J|n+1\rangle + J|n-1\rangle) \\
\langle n'|\hat{H}|E\rangle &= \sum_{n=-\infty}^{\infty} \psi_{E,n} (E_0\langle n'|n\rangle + J\langle n'|n+1\rangle + J\langle n'|n-1\rangle) \\
&= \psi_{E,n'} + \psi_{E,n'-1} + \psi_{E,n'+1} \\
&= \psi_{E,n'} + \frac{\psi_{E,n'}}{e^{i\phi}} + e^{i\phi} \psi_{E,n'} \\
&= \psi_{E,n'} (1 + e^{-i\phi} + e^{i\phi}) \\
&= \psi_{E,n'} (1 + 2\cos(\phi))
\end{aligned}$$

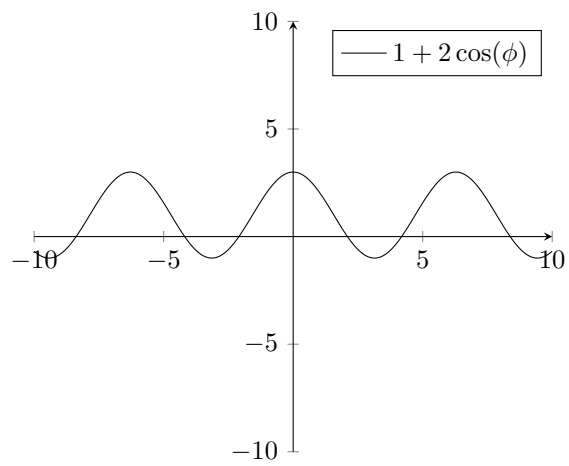


Figure 1: Simple pgfplot for aesthetic purposes

$$[1 + 2 \cos(\phi)]_{\max} = 3$$

$$[1 + 2 \cos(\phi)]_{\min} = -1$$