

# The Superconducting Heterodyne Approach to Axion Detection

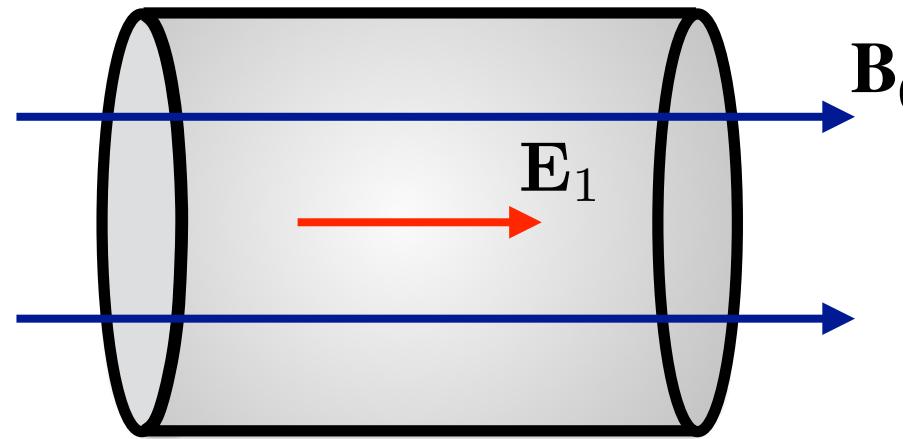
Kevin Zhou



Berkeley Axion Workshop — May 9, 2025

## Driving Cavity Modes

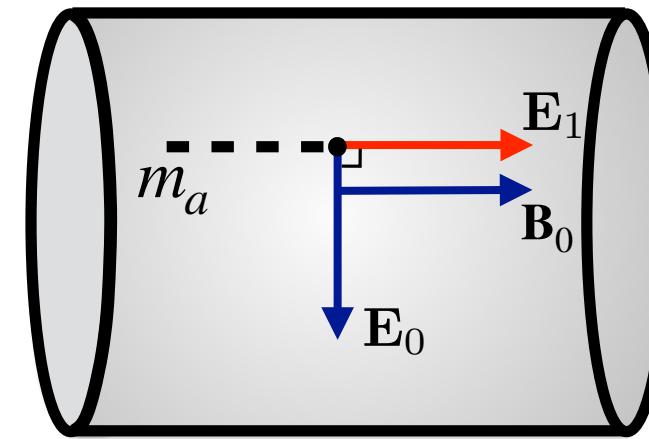
A current  $\mathbf{J}$  drives the amplitude of a cavity mode with profile  $\mathbf{E}_1$  by  $\int_V \mathbf{J} \cdot \mathbf{E}_1$   
Axion dark matter effectively yields  $\mathbf{J} = g_{a\gamma} \dot{a} \mathbf{B}_0$



### traditional cavity haloscope

place cavity into largest  
possible static  $\mathbf{B}_0$

excites mode at  $\omega_1 = m_a$   
probes  $m_a \sim 1/L \sim \text{GHz}$



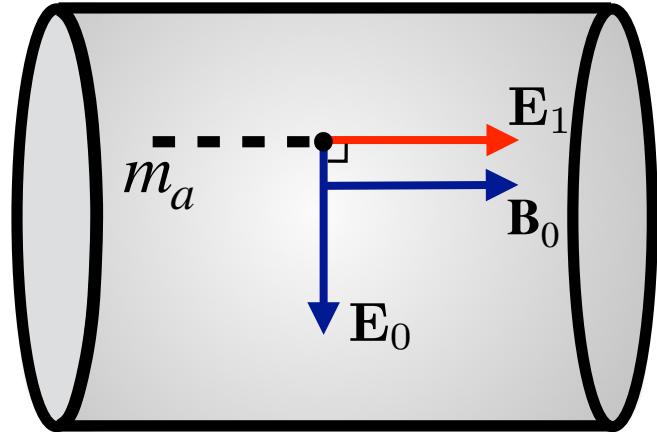
### heterodyne approach

excite a cavity mode at  $\omega_0 \sim \text{GHz}$

axion transfers energy to  
mode at  $\omega_1 = \omega_0 + m_a$

small difference probes  $m_a \ll \text{GHz}$

# The Superconducting Heterodyne Approach



$$\text{driving } g_{a\gamma} \dot{a} \int_V \mathbf{B}_0 \cdot \mathbf{E}_1 \quad P_{\text{sig}} \sim (g_{a\gamma} \dot{a} B_0)^2 V \min(t_r, t_a)$$

cavity decay time  $t_r \sim Q_1/\omega_1 \sim (1 \text{ ns}) Q_1$

axion coherence time  $t_a \sim (1 \text{ sec})(\text{MHz}/m_a)$

Requires high  $B_0$  and  $t_r$ , both enabled by high  $Q \sim 10^{11}$  in superconducting cavities  
Fabrication, calibration, loading, measurement enabled by decades of accelerator R&D

ELECTROMAGNETIC DETECTOR FOR GRAVITATIONAL WAVES

F. PEGORARO, L.A. RADICATI

OBSERVATION OF  $4 \times 10^{-17}$  cm HARMONIC DISPLACEMENT  
USING A 10 GHz SUPERCONDUCTING PARAMETRIC CONVERTER \*

C.E. REED  
Department of Physics  
University of Florida  
Received 15 January 1979



Superconducting Radio Frequency Cavities as Axion Dark Matter Detectors

P. Sikivie  
Department of Physics, University of Florida, Gainesville, FL 32611, USA  
(Dated: September 6, 2011)

suggested for gravitational waves in 1970s  
prototyped in 1980s  
first serious effort (MAGO) in early 2000s  
proposed by Sikivie for axions in 2010  
revived for light axions in 2019

# Unique Features of the Heterodyne Approach

The heterodyne approach upends all conventional wisdom about axion experiments

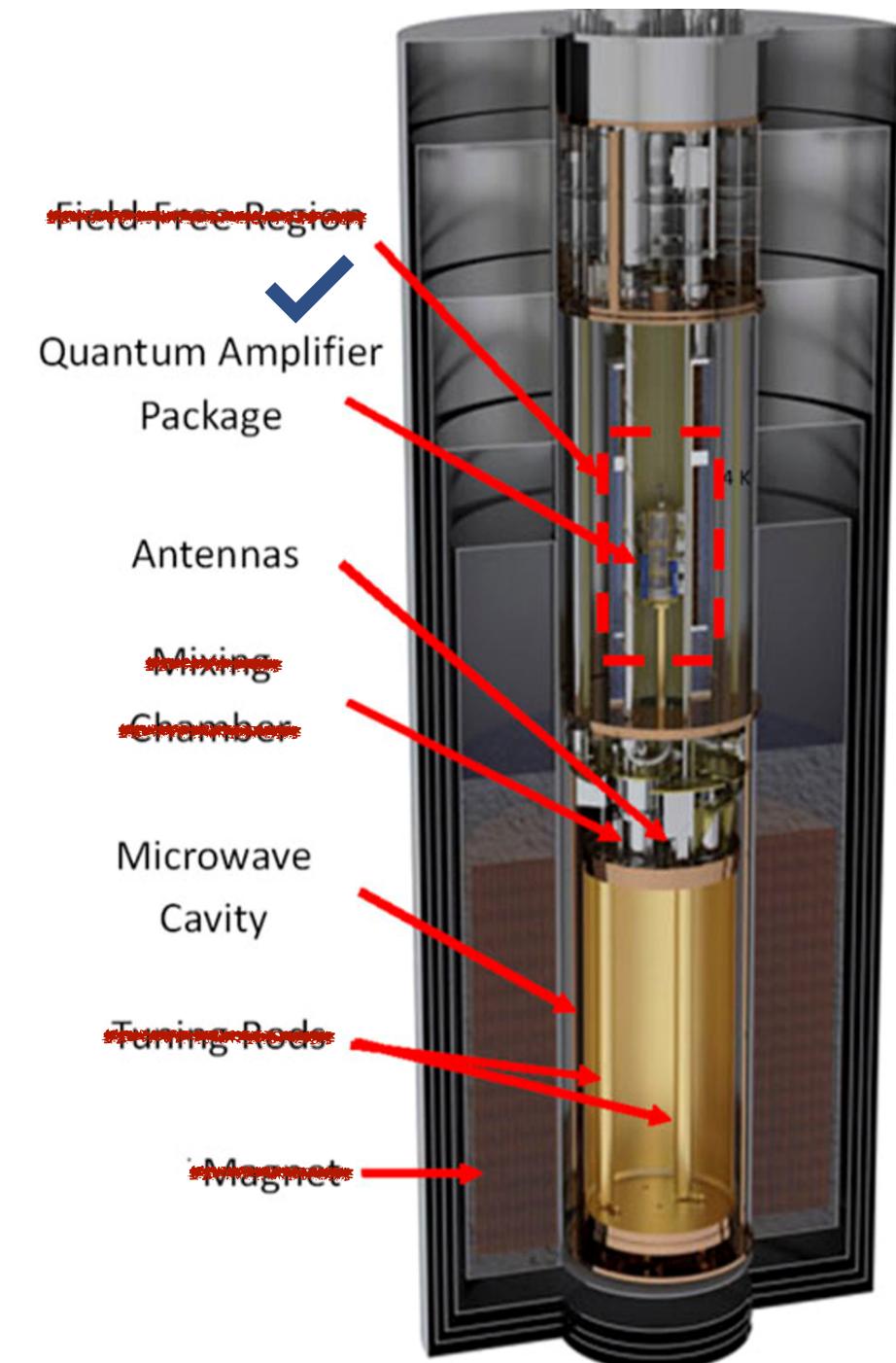
**No large magnet:** field excited in cavity ( $B_0 \sim 0.2$  T)

**No dilution fridge:** use helium cryostat ( $T \sim 2$  K)

(typical operating points for these cavities)

Same output frequency as ADMX, so use same amplifiers; **no new “quantum tricks”**

Linear tuning: a small shift  $\Delta\omega_i/\omega_i \sim 10^{-3}$  in one cavity covers **all**  $m_a$  from zero to MHz



# Parametric Advantage at Low Axion Mass

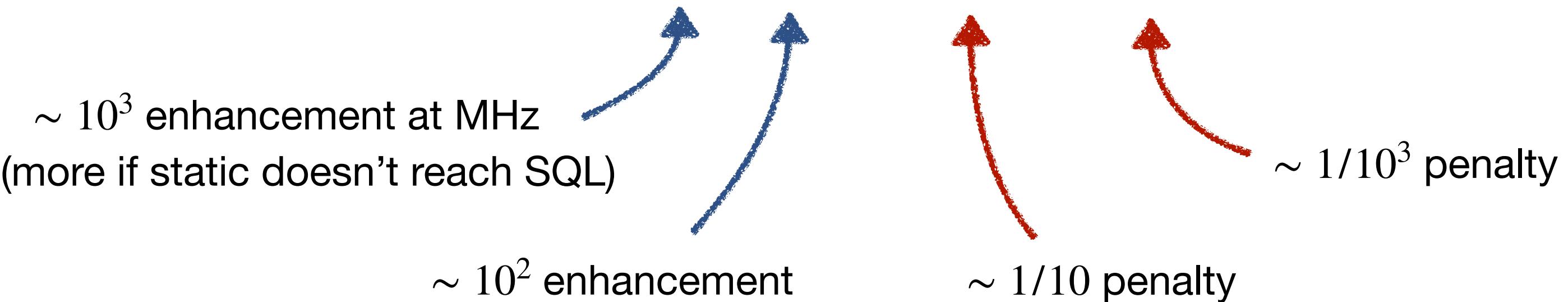
At low  $m_a$ , using a static  $\mathbf{B}_0$  would excite a “magnetoquasistatic” mode

cavity mode:  $E_1 \sim B_1$   
magnetoquasistatic:  $E_1 \sim (m_a L) B_1$

$$\frac{P_{\text{sig}}^{(\text{het.})}}{P_{\text{sig}}^{(\text{static})}} \propto \frac{1}{(m_a L)^2} \sim 10^6 \left( \frac{\text{MHz}}{m_a} \right)^2$$

Assuming thermal + standard quantum limited amplifier noise:

$$\frac{\text{SNR}^{(\text{het})}}{\text{SNR}^{(\text{static})}} \sim \frac{\omega_1}{m_a} \left( \frac{Q}{Q_{\text{static}}} \right)^{1/2} \left( \frac{T_{\text{static}}}{T} \right)^{1/2} \left( \frac{B}{B_{\text{static}}} \right)^2$$

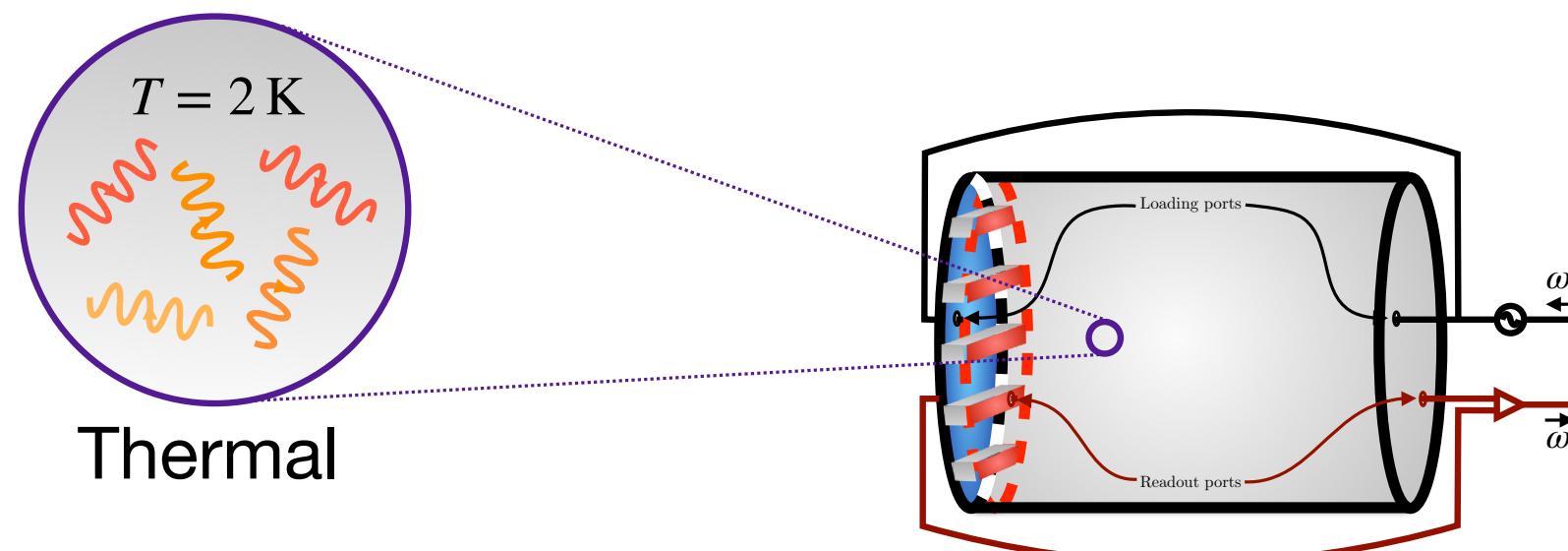


# Estimating Noise: Thermal

$$S(m_a) \sim T \sim 10^{-4} \text{ eV}$$

Higher temperature than other axion experiments,  
signal mode has  $\sim 100$  thermal photons

But compensated by much higher signal power



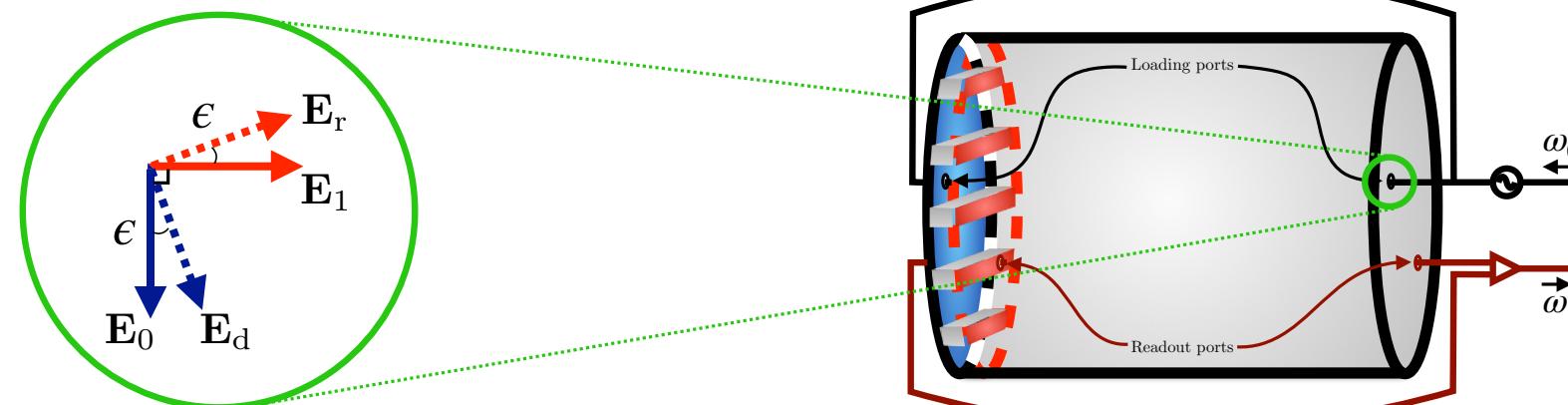
# Estimating Noise: Crosstalk

$P_{\text{in}} \sim (1 \text{ kW})(10^{11}/Q) \sim 10^{20} P_{\text{sig}}$ , isolated by geometric and frequency separation

$$S(m_a) \sim P_{\text{in}} \epsilon^2 \times \begin{cases} S_\varphi(m_a) & \text{oscillator phase noise} \\ S_\delta(m_a) & \text{fractional wall deformation} \end{cases}$$

cross-coupling factor,  
 $\epsilon = 10^{-7}$  achieved by MAGO

for commercial oscillators:  $S_\varphi(\text{kHz}) \sim 10^{-12}/\text{Hz}$   
DarkSRF measurement:  $S_\delta(100 \text{ Hz}) \sim (10^{-9})^2/(100 \text{ Hz})$   
both rapidly fall off for higher frequency



Crosstalk

1912.11048, 1912.11056, 2007.15656

Thermal noise larger above:  
kHz:  $\epsilon = 10^{-7}$   
MHz:  $\epsilon = 10^{-4}$

# Estimating Noise: Vibrations

As usual for axion experiments, vibrational noise greatly dominates at kHz

$$S(m_a) \sim P_{\text{in}} \eta^2 Q_1^2 S_\delta(m_a)$$

mode mixing form factor

$$\eta = \int_S (\mathbf{u} \cdot d\mathbf{S}) (\mathbf{B}_0 \cdot \mathbf{B}_1 - \mathbf{E}_0 \cdot \mathbf{E}_1)$$

ideally zero for certain cavity geometries

often set by  $\sim 10^{-4}$  precision of manufacture

falls off for higher mech. modes,  $\eta \propto \omega_{\min}/m_a$

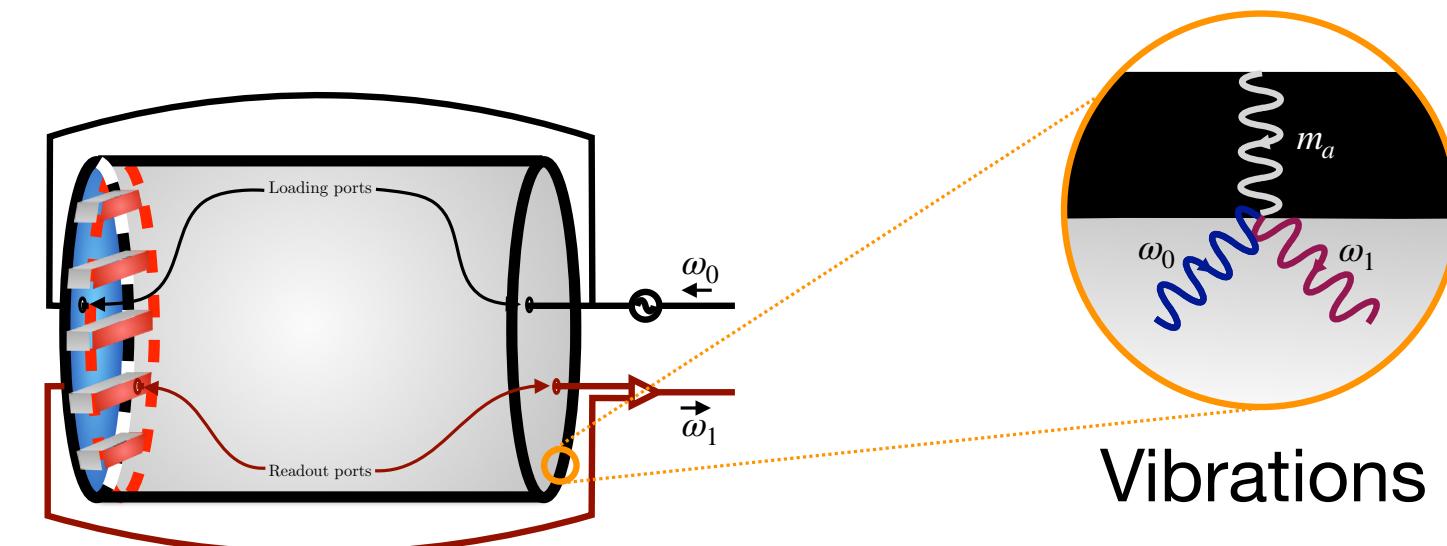
(reducible by stiffening walls)

largely due to cooling system

pessimistic assumption:  $S_F(m_a) \propto 1/m_a$

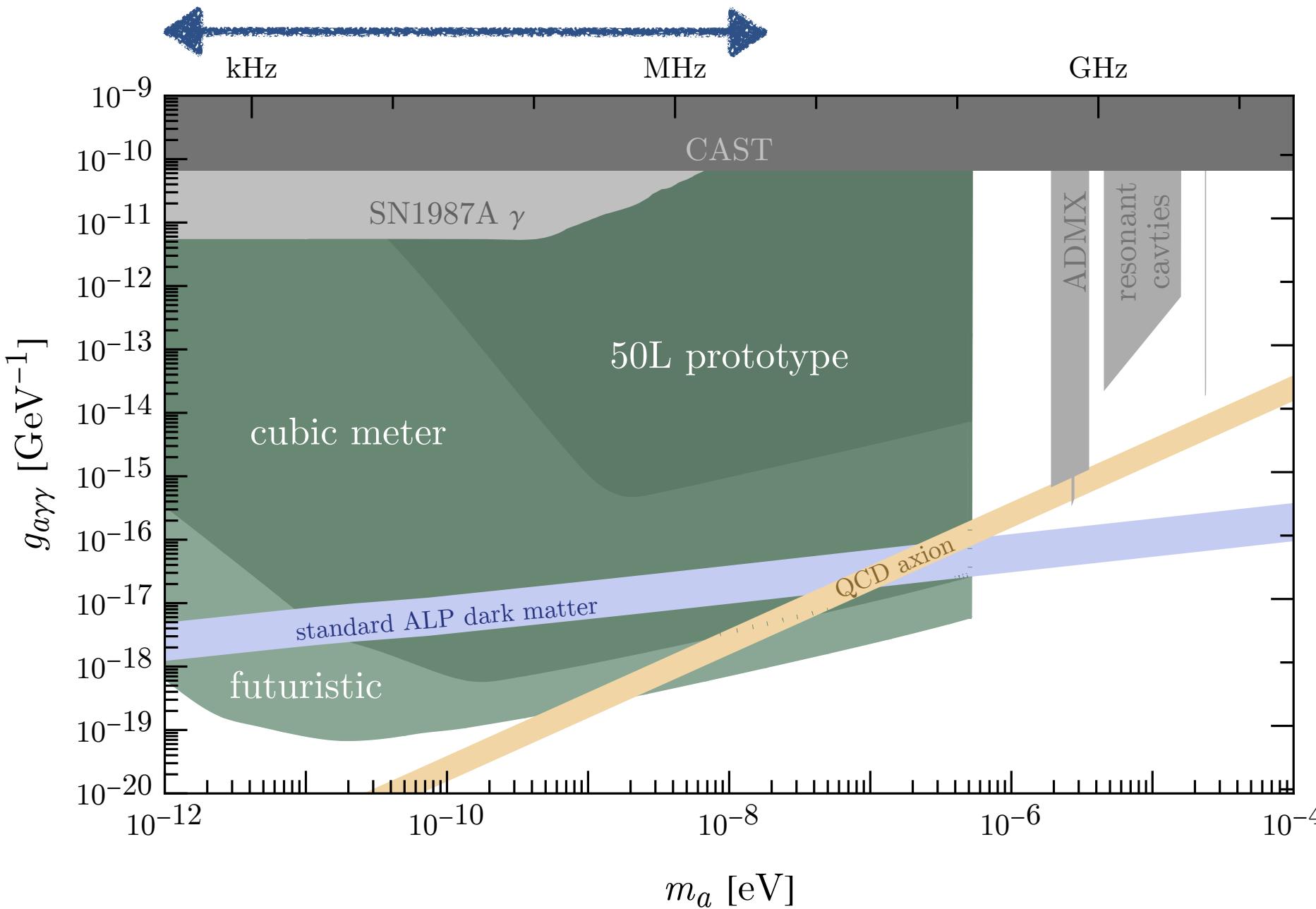
at high frequency,  $S_\delta(m_a) \propto S_F(m_a)/m_a^4$

subdominant at MHz if  $\eta$  reasonably small



# Potential Sensitivity

tuning range of a single cavity



High mass: thermal limited

$$B_0 = 0.2 \text{ T}, T = 2 \text{ K}$$

$$Q = \begin{cases} 10^{10} & \text{prototype} \\ 10^{12} & \text{cubic meter} \end{cases}$$

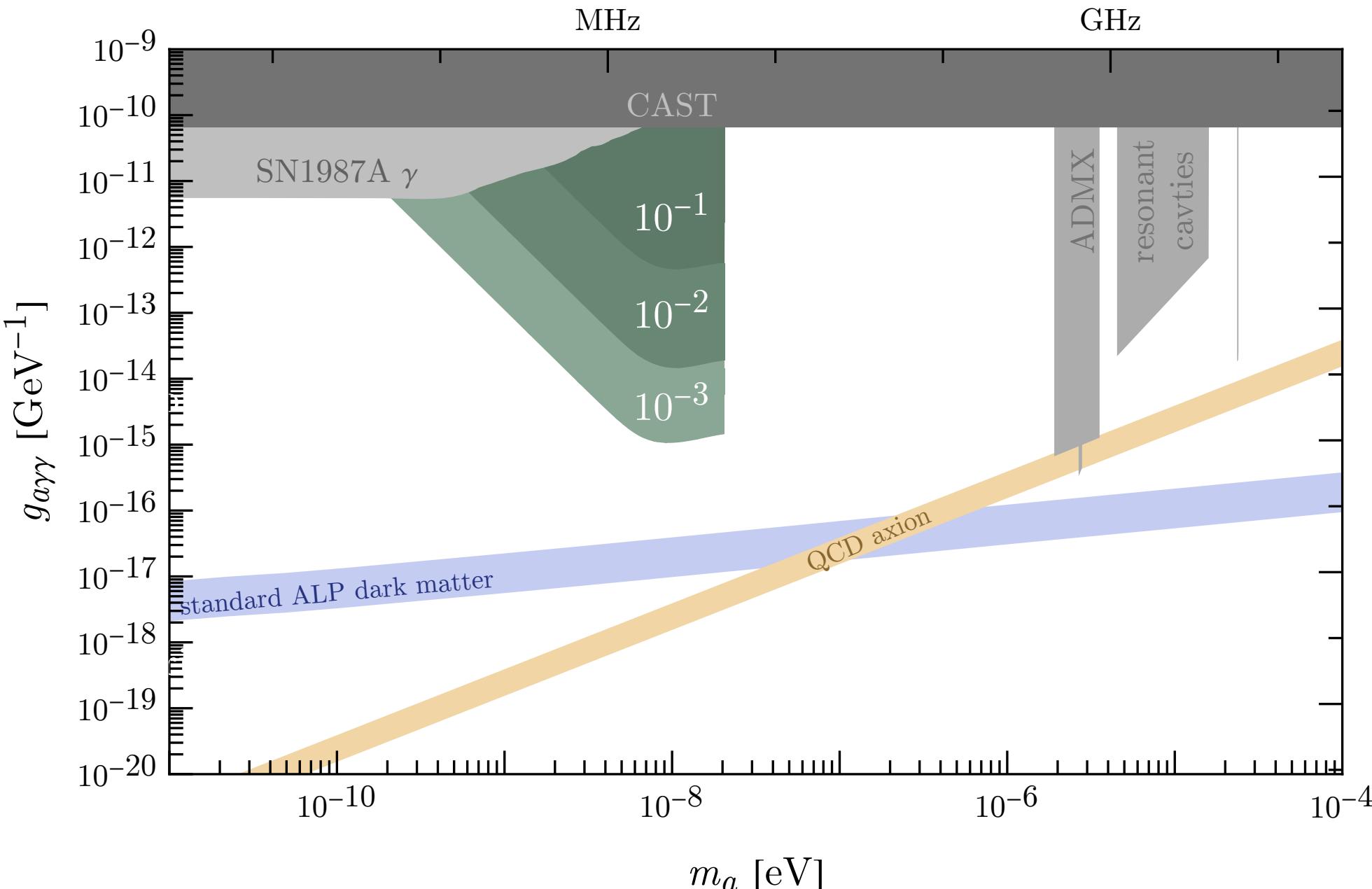
Lower mass: crosstalk/vibration

$$\epsilon = \begin{cases} 10^{-3} & \text{prototype} \\ 10^{-7} & \text{cubic meter} \end{cases}$$

$$\eta = \begin{cases} 10^{-4} & \text{prototype} \\ 10^{-5} & \text{cubic meter} \end{cases}$$

Probes QCD axion at  $\sim 10 \text{ MHz}$

# “Off the Shelf” Prototype Sensitivity



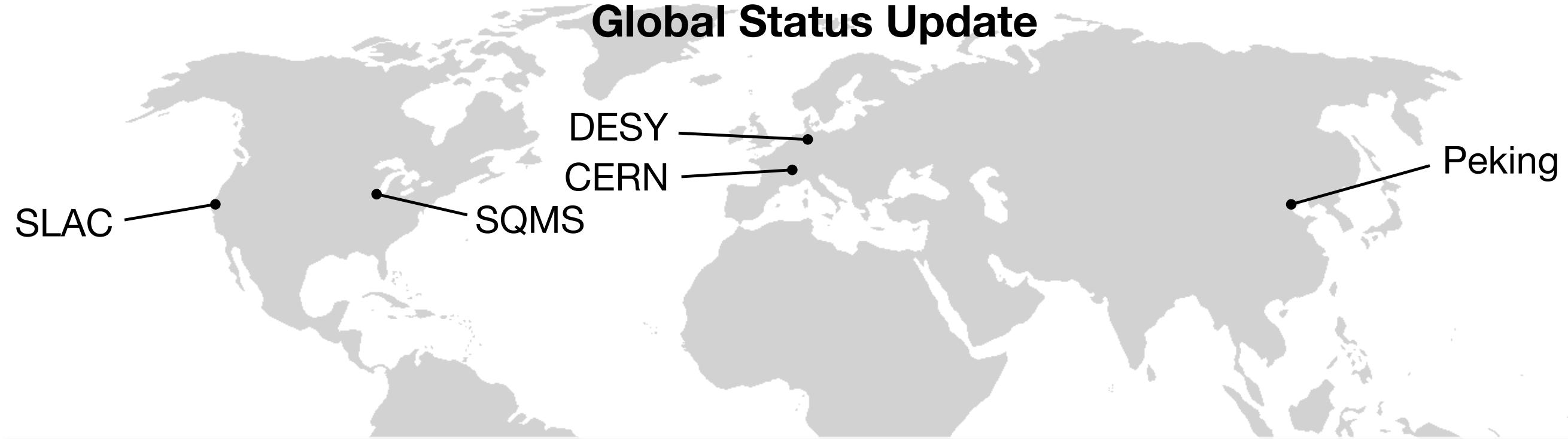
Projection for existing  
**unoptimized** cavities

High values of  $\epsilon$  and  $\eta$

Low signal form factor,  
multiples power by 0.04

Vibrational force PSD  
scaled up by  $\sim 10^6$

## Global Status Update



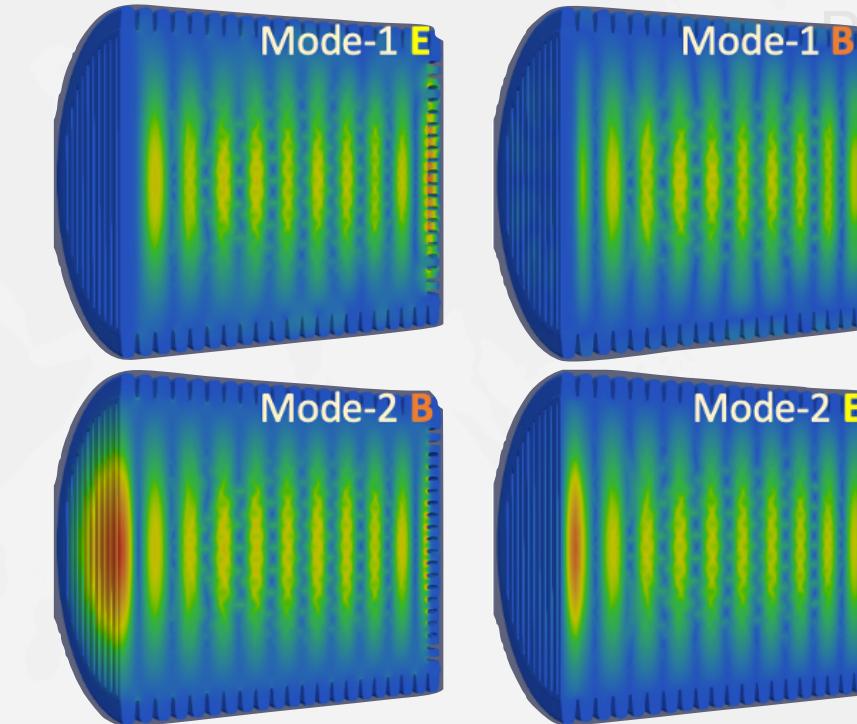
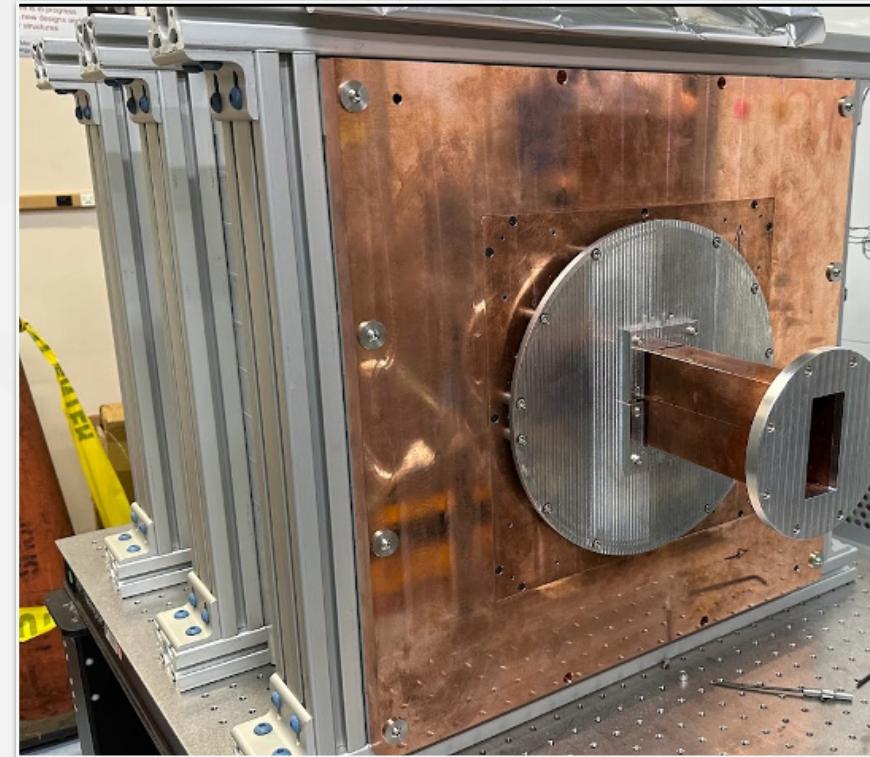
Several groups, at labs with expertise in superconducting cavities,  
are currently designing prototypes and taking measurements

Current level of funding is sufficient to demonstrate proof of principle  
and probe far beyond astrophysical bounds

Scaling to QCD axion sensitivity requires qualitative increase in funding



## LDRD effort at SLAC (2022-2024)

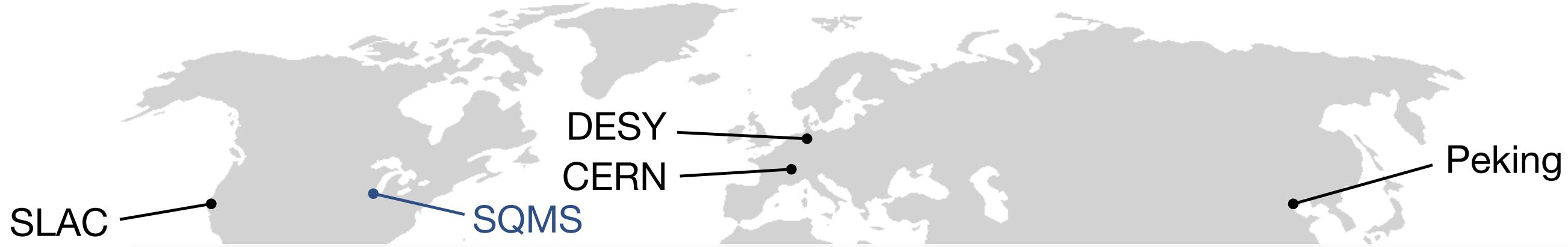


made corrugated copper cavity with two orthogonal modes

moving tuning fins gives range  $\Delta f = 4 \text{ MHz}$

mode profiles enable:  $\eta < 10^{-3}$  for **any** vibration

$\epsilon \leq 2 \times 10^{-4}$  by just hand tuning



SHADE collaboration at SQMS (started 2021)

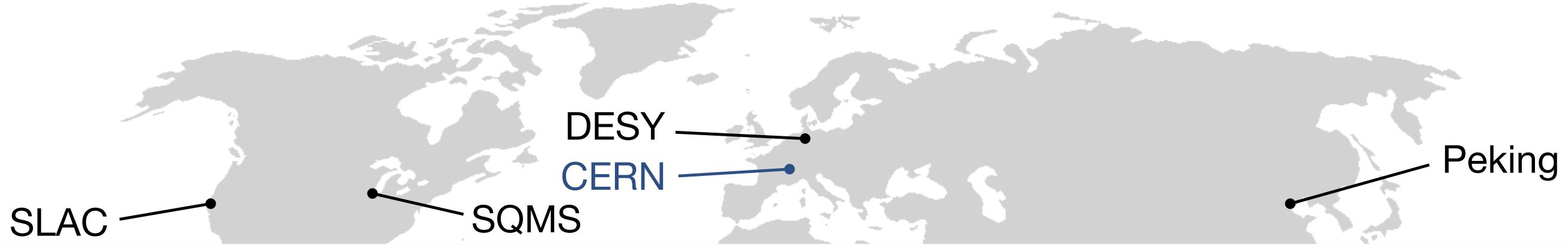
longest running and most developed effort

2022: measurement of existing 9-cell cavity at  
 $T = 2 \text{ K}$ , no exotic noise observed

2024: new cavities with small  $\eta$  instrumented

2025: planned cold measurement, tuning test

further details in Bianca Giaccone's talk



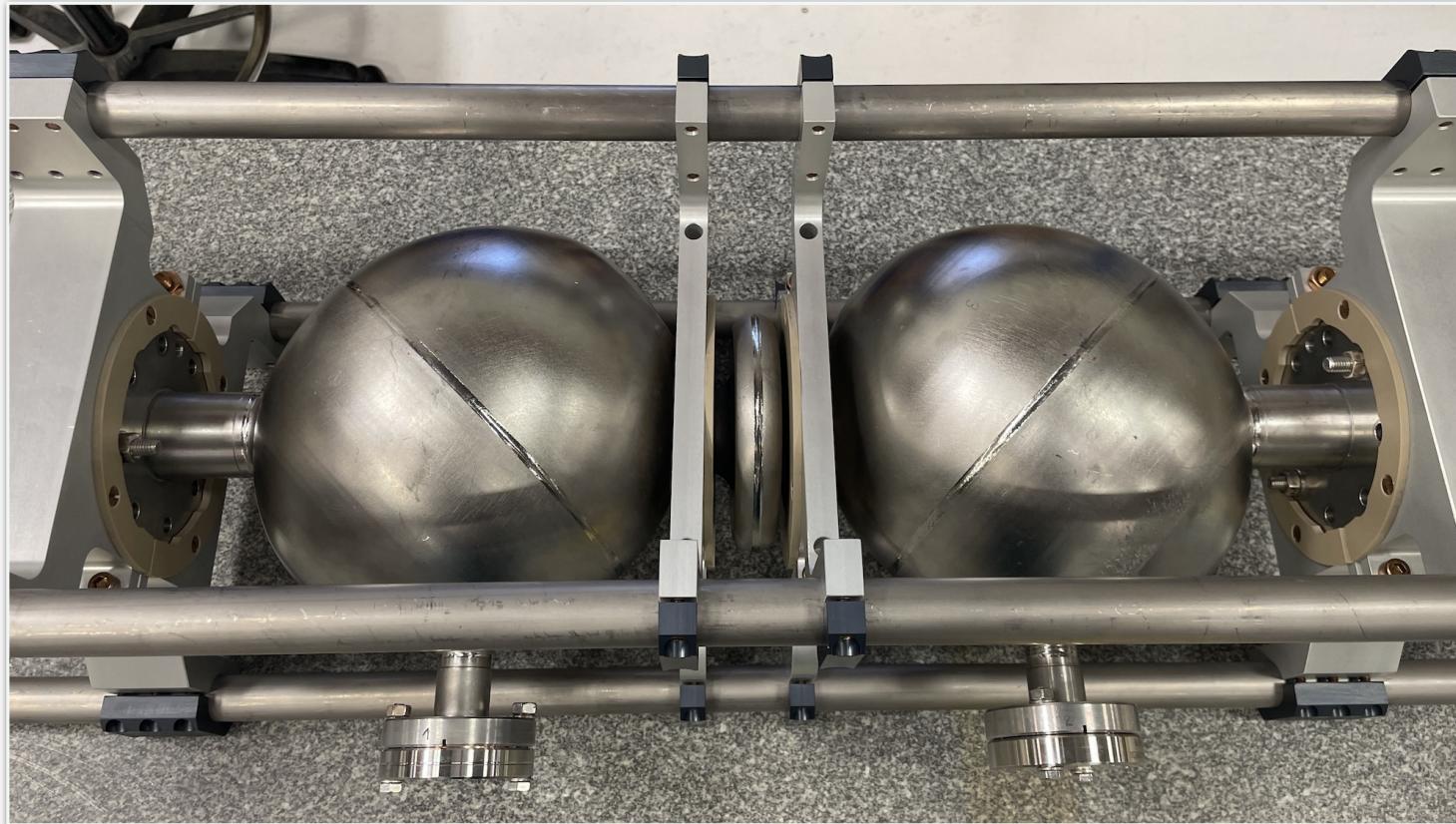
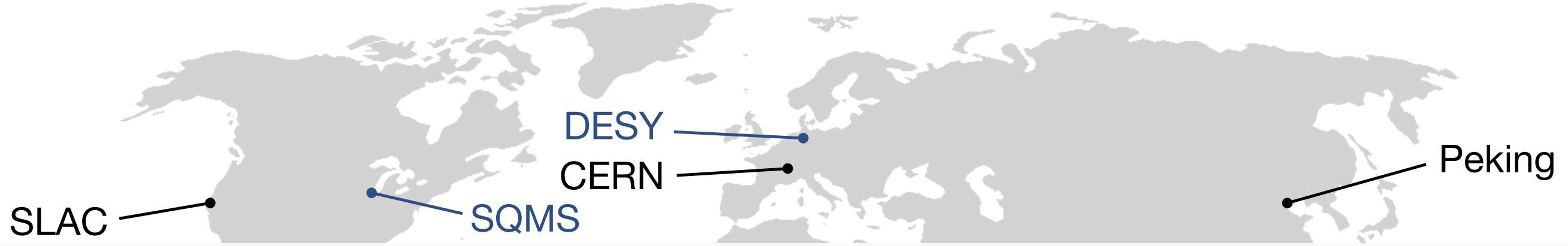
New effort at CERN (started 2025)  
Funded by QTI in affiliation with PBC

Aims to develop and run superconducting prototype in next 2-3 years

Developing optimizations for heterodyne detection:

non-mechanical tuning,  
to operate cryogenically

new cavity designs to  
reduce  $\epsilon$  and  $\eta$



Revival of MAGO (started 2024)

Joint effort of DESY and SQMS

Two-sphere cavity recovered, tuned

Vibrational modes simulated

Optimized for gravitational waves,  
but shares noise sources

# SHANHE collaboration at Peking (started 2023)



2023: dark photon search, no driving

early 2025: calibration run at  
 $T = 4\text{ K}$ , no exotic noise observed

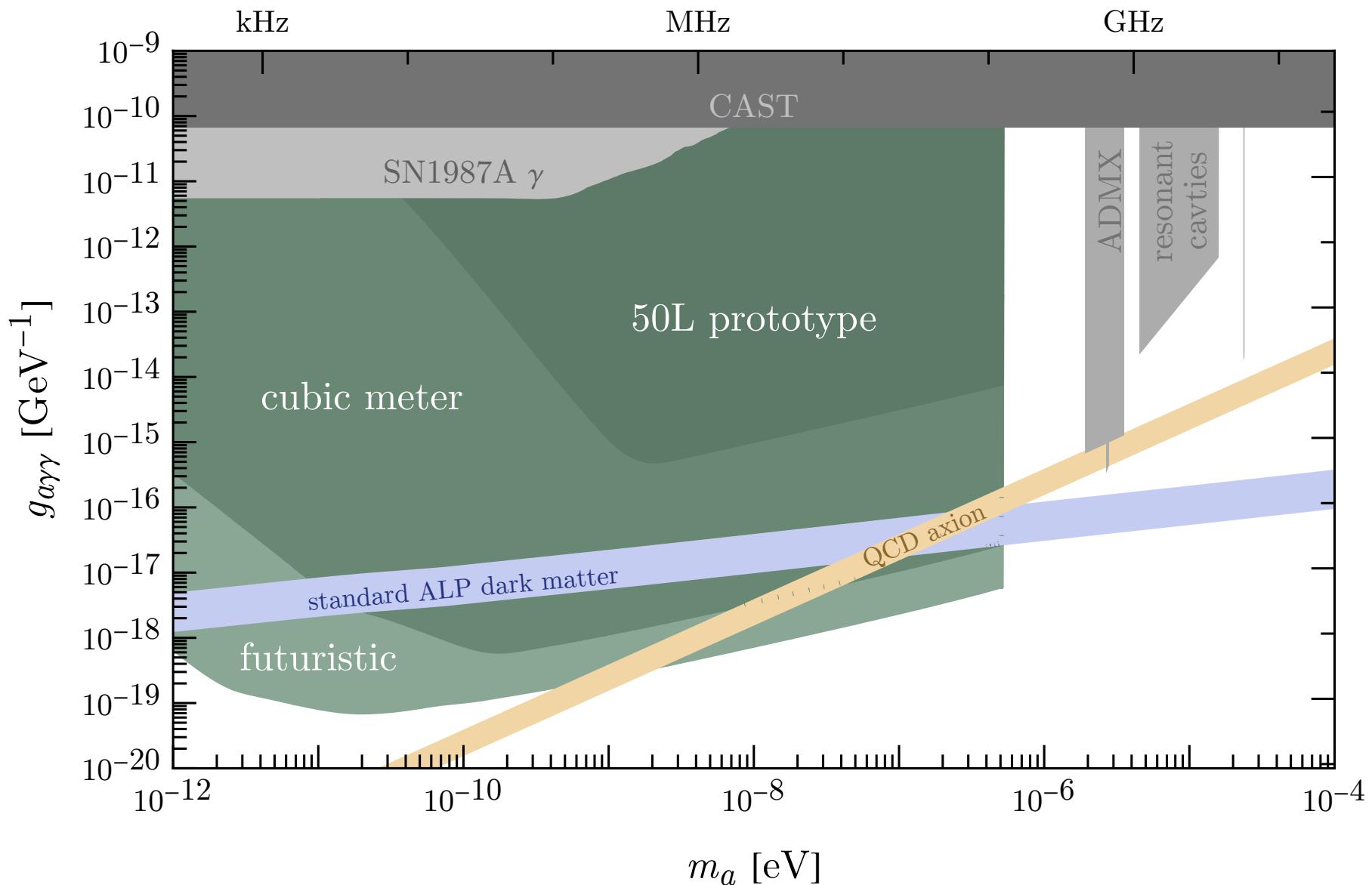
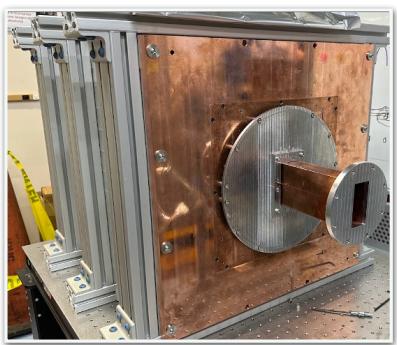
mid 2025: data taking run ongoing

2026: planned run with new cavity  
designed to reduce  $\epsilon$  and  $\eta$



# Outlook

Superconducting cavities are **not** science experiments  
They are mass-produced, practical technology



But they have the potential  
to transform the search for  
light axion dark matter

Multiple ongoing efforts will  
demonstrate feasibility in  
next ~2 years