

Example

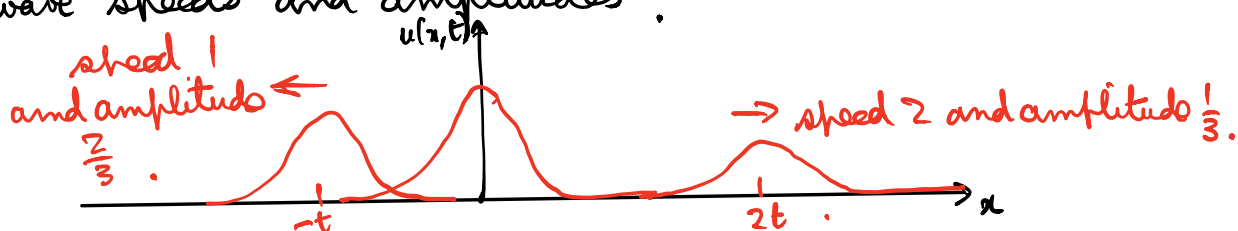
$$\frac{\partial^2 u}{\partial t^2} + \frac{\partial^2 u}{\partial x \partial t} - 2 \frac{\partial^2 u}{\partial x^2} = 0 \quad \text{for } -\infty < x < \infty \\ \text{and } t > 0$$

where $u(x, 0) = e^{-x^2}$ and $\frac{\partial u}{\partial t}(x, 0) = 0$ for all x .

General solution

$$u = \frac{1}{3} e^{-(x-2t)^2} + \frac{2}{3} e^{-(x+t)^2}.$$

Plotting this shows that the parts have different wave speeds and amplitudes.



An alternative way of finding the general solution of:

$$A \cdot \frac{\partial^2 u}{\partial t^2} + B \cdot \frac{\partial^2 u}{\partial x \partial t} + C \cdot \frac{\partial^2 u}{\partial x^2} = 0 \quad (**)$$

is to follow the procedure for the wave equation, and introduce v such that

$$\frac{\partial v}{\partial t} = \frac{\partial u}{\partial x}$$

then it can be shown that satisfy (**), then

$$A \cdot \frac{\partial v}{\partial t} + B \cdot \frac{\partial v}{\partial x} + C \cdot v = 0.$$

The linear PDE can therefore be written as,

$$\frac{\partial v}{\partial t} = \tilde{A} \cdot \frac{\partial v}{\partial x}$$

where $\underline{u} = \begin{bmatrix} u \\ v \end{bmatrix}$ and $\tilde{A} = \begin{bmatrix} -\frac{B}{A} & -\frac{C}{A} \\ 1 & 0 \end{bmatrix}$.

If the eigenvalues λ_1 & λ_2 of \tilde{A} are real and distinct, then they have corresponding eigenvectors $\underline{v}^{(1)}$, $\underline{v}^{(2)}$. Then introducing

$$\underline{u} = P \underline{u}' = \begin{bmatrix} \underline{v}^{(1)} & \underline{v}^{(2)} \end{bmatrix} \underline{u}'$$

the equation for \underline{u} becomes

$$\begin{aligned} \frac{\partial \underline{u}'}{\partial t} &= P^{-1} \tilde{A} P \frac{\partial \underline{u}'}{\partial x} \\ &= \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix} \frac{\partial \underline{u}'}{\partial x} \quad (***). \end{aligned}$$

which is in diagonal form, and each component on \underline{u}' is a linear combination of u, v .

Each component of $(***)$ is linear advection equation with characteristic slope:

$$\frac{dx}{dt} = -\lambda_1 \text{ and } \frac{dx}{dt} = -\lambda_2 \text{ respectively.}$$

For the wave equation we used $\eta = \frac{u+v}{2}$ and $\mathcal{J} = \frac{u-v}{2}$ to derive.

$$\frac{\partial \eta}{\partial t} = a \frac{\partial \eta}{\partial x}, \quad \frac{\partial \mathcal{J}}{\partial t} = -a \frac{\partial \mathcal{J}}{\partial x}$$

so that $\lambda_1 = -a$ and $\lambda_2 = a$ here, and the combinations of u, v for η, \mathcal{J} are given by the eigenvectors of \tilde{A} in this case.

3.5. The wave equation with reflection at boundaries

Consider the wave equation

$$\frac{\partial^2 u}{\partial t^2} = a^2 \frac{\partial^2 u}{\partial x^2}$$

in a domain $-\infty < x < 0$ with a boundary at $x=0$.
We seek a solution for $t > 0$ with given initial conditions

$$\left. \begin{aligned} u(x, 0) &= f(x) \\ \frac{\partial u}{\partial t}(x, 0) &= g(x) \end{aligned} \right\} \text{ for } x \leq 0,$$

and a stated boundary condition at $x=0$ for $t > 0$.
We know that the general solution of the wave equation is

$$u(x, t) = F(x+at) + G(x-at)$$

for any functions F, G .

Applying the initial conditions at $t=0$ gives.

$$\left. \begin{aligned} u(x, 0) &= F(x) + G(x) = f(x) \\ \frac{\partial u}{\partial t}(x, 0) &= aF'(x) - aG'(x) = g(x) \end{aligned} \right\} \text{ for } x \leq 0.$$

and hence we find $F(x)$ and $G(x)$ for given f, g when $x \leq 0$.

Once $t > 0$, notice that $x+at > 0$ when $t > -\frac{x}{a}$ for some negative x . Hence $u(x, t)$

is undetermined when $x+at > 0$ and F has some positive argument.

If we specify a boundary condition at $x=0$,
e.g.

$$u(0, t) = 0 \quad \text{for all } t > 0$$

then from the general solution .

$$F(0+at) + G(0-at) = 0 \quad \text{for all } t > 0.$$

Let $\eta = at$ then this tells us that .

$F(\eta) = -G(-\eta)$ for all $\eta > 0$,
and hence the argument of F is known for all positive values .

Example

Consider when $u(x,0) = f(x)$ for any given f when $x \leq 0$ and $\frac{\partial u}{\partial t}(x,0) = 0$ for all $x \leq 0$, with $u(0,t) = 0$ for $t > 0$.

From earlier, we showed that .

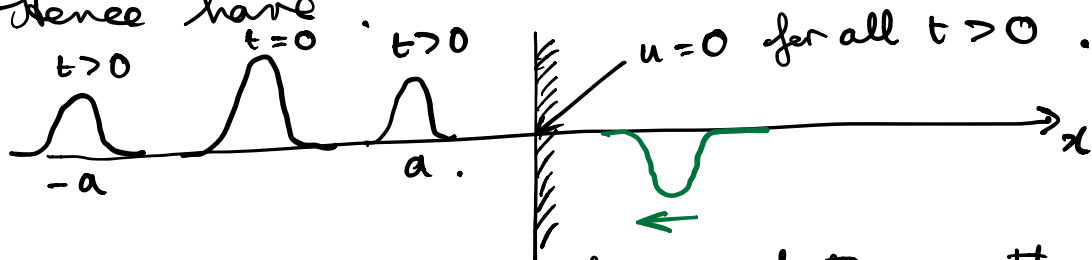
$$F(x) = \frac{1}{2}f(x) \text{ and } G(x) = \frac{1}{2}f(x) \text{ when } x \leq 0.$$

and hence

$$u(x,t) = \frac{1}{2}f(x+at) + \frac{1}{2}f(x-at)$$

provided that $x+at \leq 0$.

Hence have



Once $x+at > 0$ we also need to use that $u(0,t) = 0$ so that

$$u(0,t) = F(at) + G(-at) = 0 \quad \text{for } t > 0$$

and hence

$$F(\eta) = -G(-\eta) = -\frac{1}{2}f(-\eta) \text{ for } \eta > 0.$$

Hence

$$F(\eta) = \begin{cases} \frac{1}{2} f(\eta) & \text{for } \eta \leq 0 \\ -\frac{1}{2} f(-\eta) & \text{for } \eta > 0. \end{cases}$$

with $G(\eta) = \frac{1}{2} f(\eta)$ for all $\eta < 0$. This leads to a second component of the solution which travels to left (i.e. F part) which becomes important once its argument is positive. This corresponds to the reflection of the G part of the solution in the boundary $x = 0$.