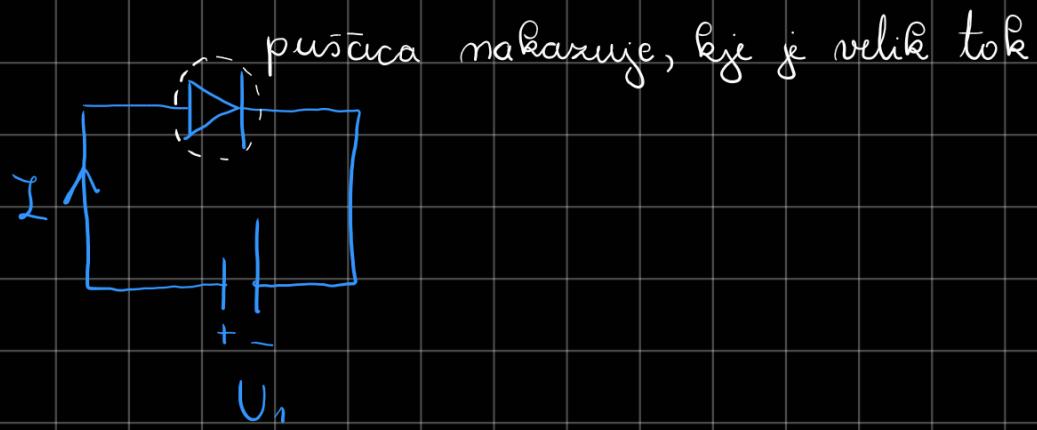


# Nadajevanje od zadnjic

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1) Najprej vzamemo idealno diodo

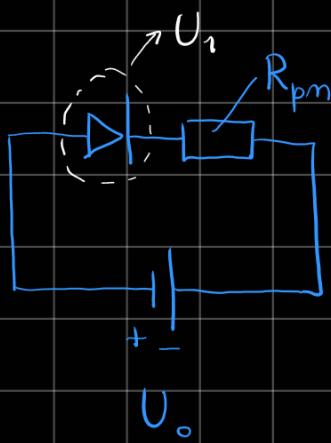


$$I = I_0 (e^{-\beta e_0 U} - 1)$$

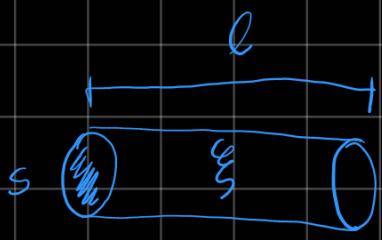
// obmenjujmo za  $U_i$

$$U_i = \frac{1}{\beta e_0} \ln \left[ \frac{I}{I_0} + 1 \right] = 0.24V$$

2)

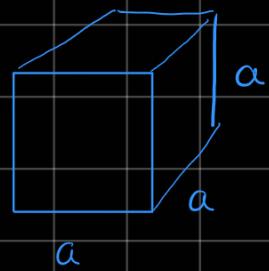


$$U_o = U_i + I R \quad // zavima na  $R$ , ker  $U_o$  in  $U_i$  poravnamo$$



$$R = \frac{\xi l}{s}$$

// za kocko je



$$R = \frac{\xi a}{a^2} = \frac{1}{3a}$$

$$R_{pm} = R_p + R_m = \frac{1}{a} \left[ \frac{1}{R_p} + \frac{1}{R_m} \right] = 7 \Omega$$

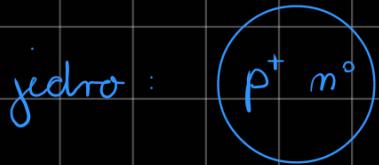
$U_0$  je napetost vrha, ki ga moramo priklopiti

$$U_0 = (\underbrace{U_1}_{\text{"} 0.24V}) + R \underbrace{I}_{\text{"} 10mA} = 0.31V$$

Jedra

Težota v jedru je večinoma konst. Vzamna en. je premo soraznačena s številom gradnikov ( $2 \times \text{št. gradnikov} \rightarrow 2 \times E_v$ )

Velikost je redar  $\sim fm$



A ... oznaka št. gradnikov

$\xrightarrow{\text{mimoš.}} A \rightarrow Z + N$ ,  $N$  je št. protonov)

Jedro elementa



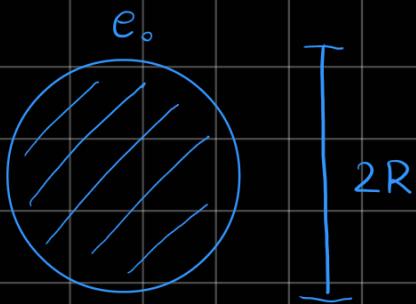
$\Rightarrow$  vrsta št. (# protonov)

// predpostavimo, da je jedro iz trdih movi (ni gradniki so tudi // empirično lahko  $\propto g = \text{konst}$ . zapisimo radij jedra kroglice)

$$r_0 = r_1 A^{\frac{1}{2}}, \quad r_1 = 1.1 \text{ fm}$$

## Elektrostatika

### Čopac VI / 2



// nabolj je enakomerno razmejan

$$E_{el} = ? \quad g(r=0) = ?$$

// z Gaussovim zakonom dočimmo  $\vec{E}$  zunaj in unutrašnj  
kroglice

$$1) \text{Gauss} \rightarrow \vec{E}(\vec{r})$$

$$2) \vec{E}(\vec{r}) \rightarrow g(\vec{r}) \quad // \text{potencial}$$

$$3) g(\vec{r}) \rightarrow E_{el} \quad // \text{elektrostatska en.}$$

$$1) e = \epsilon_0 \oint \vec{E} d\vec{S} \quad d\vec{S} \text{ so sfere (amo v 3D)}$$

$$E(r) = \frac{e(r)}{4\pi r^2 \epsilon_0}$$

$$a) r \leq R$$

$$e = \rho e V \quad // \bar{u} \& e = e(r) \text{ integral } \int_V g(r) d^3 \vec{r}$$

$$C = \frac{\epsilon_0}{\frac{4}{3}\pi R_0^3} \cdot \frac{4}{3}\pi r^3$$

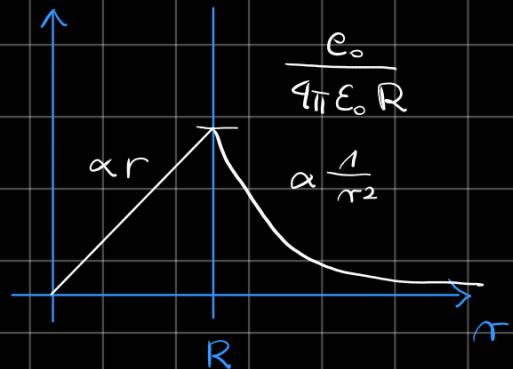
$$E_{<R}(e) = \frac{\epsilon_0}{4\pi\epsilon_0 R^3} r$$

b)  $r > R$

točkast náboj,  $C = \text{konst.}$

$$E(r) = \frac{e(r)}{4\pi\epsilon_0 r^2} \Rightarrow E_{>R}(e) = \frac{\epsilon_0}{4\pi\epsilon_0 r^2}$$

$$E(r) = \begin{cases} E_{<R}, & r \leq R \\ E_{>R}, & r > R \end{cases}$$



2) potencial  $\phi(\vec{r}) = ?$

$$\vec{E}(\vec{r}) = -\nabla \phi(\vec{r})$$

$$\phi(\vec{r}) = - \int E(\vec{r}) d\vec{s}$$

$\nabla$  v kartezičních  $\rightarrow$   $\nabla$  v sfemických

$$\nabla = \left( \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \rightarrow \nabla = \frac{\partial}{\partial r} \hat{e}_r + \frac{1}{r} \frac{\partial}{\partial \theta} \hat{e}_\theta + \frac{1}{r \sin \theta} \frac{\partial}{\partial \varphi} \hat{e}_\varphi$$

// dla  $\hat{e}_\theta$  i  $\hat{e}_\varphi$  odpadita, kur  $\vec{E} = E(\hat{e}_r)$

a)  $r > R$

// nádolocíni int

$$\varphi_{>R}(r) = - \int \frac{e_0}{4\pi\epsilon_0 r^2} dr = \frac{e_0}{4\pi\epsilon_0 r} + A$$

b)  $r \leq R$

$$\varphi_{\leq R}(r) = - \int \frac{e_0 r}{4\pi\epsilon_0 R^3} dr = - \frac{e_0 r^2}{8\pi\epsilon_0 R^3} + B$$

a. 1)  $\varphi(r \rightarrow \infty) \rightarrow 0 \Rightarrow A = \emptyset$

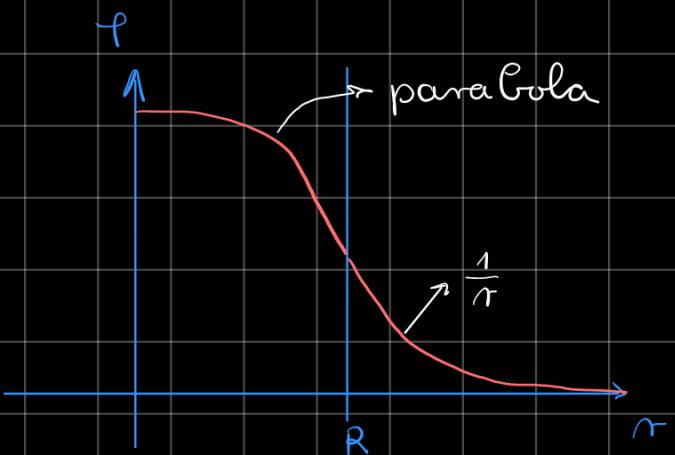
b. 1)  $\varphi_{\leq R}(r=R) = \varphi_{>R}(r=R)$

$$-\frac{e_0 r^2}{8\pi\epsilon_0 R} + B = \frac{e_0}{4\pi\epsilon_0 R}$$

$$\Rightarrow B = \frac{3 e_0}{8\pi\epsilon_0 R}$$

$$\varphi(r=0) = \frac{3 e_0}{8\pi\epsilon_0 R}$$

$$\varphi_{>R}(r) = \frac{e_0}{8\pi\epsilon_0 R} \left[ 3 - \left( \frac{r}{R} \right)^2 \right]$$



### 3) Elektrostatska energija

$$\begin{array}{c} \text{---} & \varphi(\vec{r}) \\ \text{---} & \text{točkasto maboj} \\ \text{---} & \bullet \vec{r} = \vec{r}_0 \end{array}$$

U nivau točkastega maboga, ampak nutočkastega, razbijemo na enj niteximalni deli

$$E_{el} = e \varphi(\vec{r} = \vec{r}_0)$$

$$\text{// ka jedro} \quad E_{el} = \frac{1}{2} \int_V g(\vec{r}) de$$

parska interakcija, deluje in izračunava  
je ma par, mas pa xauima rame  
ka kroglico

$$e(r) = \frac{e_0 r^3}{R^3} \Rightarrow \frac{de}{dr} = \frac{3e_0 r^2}{R^3}$$

$$\begin{aligned} E_{el} &= \frac{1}{2} \frac{e_0}{8\pi\epsilon_0 R} \frac{3e_0}{R^3} \int_0^R dr r^2 \left[ 3 - \left( \frac{r}{R} \right)^2 \right] \\ &= \frac{3e_0^2}{16\pi\epsilon_0 R^4} \left[ 3 \frac{x^3}{3} \Big|_0^1 - \frac{x^5}{5} \Big|_0^1 \right] \end{aligned}$$

$$\begin{aligned} x &= \frac{r}{R} \\ \frac{dx}{dr} &= \frac{1}{R} \end{aligned}$$

$$E_{el} = \frac{3e_0^2}{20\pi\epsilon_0 R}$$

// za obutek  $R \sim fm \Rightarrow E_{el} \sim MeV$

$$E_{el} = \frac{3}{5} \frac{e_0^2}{4\pi\epsilon_0\hbar c} \frac{\hbar c}{R}$$

"  $\alpha$  "

$\rightarrow 197 MeV fm$

$\sim fm$

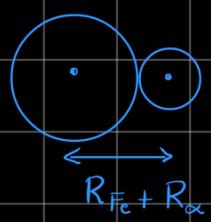
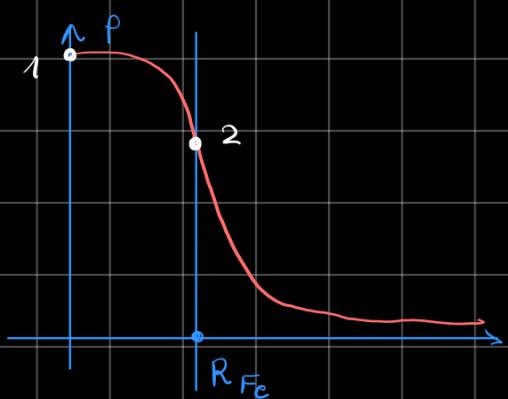
$\alpha \doteq \frac{1}{137}$  // konstanta fine strukture

VI / 3

$\alpha \rightarrow {}^{56}_{26} Fe$  // losuvu kinetiku cu.  $\alpha$  delca

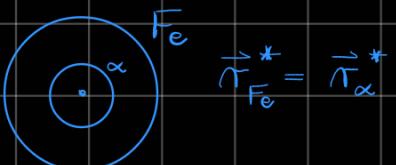
$\alpha = {}^4_2 He^{2+}$  // bez  $e^-$

ykica // pri 2 x jedri dotikata



// tukaj bomo  
vzeli potencial, ki  
je xunaj jedra

// pri 1 x medisici prenikata



$F_e$ :

$$A = 56, Z = 26$$

$\alpha$ :

$$A = 4, Z = 2$$

// empirična formula od prej  $r_\delta = r_1 A^{\frac{1}{2}}$

$$R_{Fe} = 1'75 \text{ fm}$$

$$R_\alpha = 4'21 \text{ fm}$$

1)  $E_1 = \varphi (r = 0) \cdot e_\alpha //$  predpostavimo, da je  $\alpha$  točkasto

$$= \frac{3 Z_{Fe} Z_\alpha e^2}{8\pi \epsilon_0 R_{Fe}} = \frac{3}{2} Z_{Fe} Z_\alpha \alpha \cdot \frac{\frac{e^2}{4\pi \epsilon_0} c}{R_{Fe}} = 26'6 \text{ MeV}$$

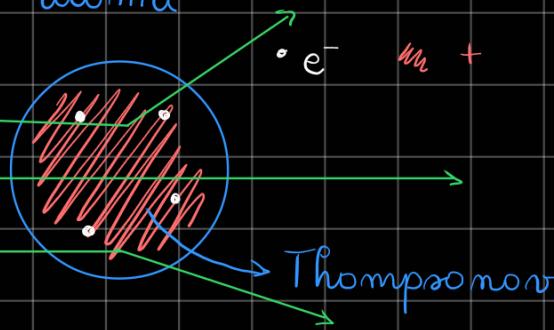
$$2) E_2 = \frac{Z_{Fe} Z_\alpha e^2}{4\pi \epsilon_0 (R_{Fe} + R_\alpha)} = 12'5 \text{ MeV}$$

## Rutherfordovo razprava

Pred Rutherfordom je bil Thompsonov puding (homogena mona točkasti načrt)

Po poskusu: jedro je majhen del atoma

Thompsonov model atoma



// maksimalni kot razpinaja  
 $\theta = 20^\circ$

Thompsonov puding

