

第七、八周作业参考答案

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习题 1

狄拉克方程的平面波解如下：

$$u^s(\vec{p}) = \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ \sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix}, \quad v^s(\vec{p}) = \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ -\sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix} \quad (1)$$

请证明如下关系：

$$\begin{aligned} u^r(\vec{p})^\dagger \cdot u^s(\vec{p}) &= 2\omega_p \delta^{rs}, & v^r(\vec{p})^\dagger \cdot v^s(\vec{p}) &= 2\omega_p \delta^{rs} \\ \bar{u}^r(\vec{p}) \cdot u^s(\vec{p}) &= 2m \delta^{rs}, & \bar{v}^r(\vec{p}) \cdot v^s(\vec{p}) &= -2m \delta^{rs} \\ \bar{u}^s(\vec{p}) \cdot v^r(\vec{p}) &= 0, & u^r(\vec{p})^\dagger \cdot v^s(-\vec{p}) &= 0 \\ \sum_{s=1}^2 u^s(\vec{p}) \bar{u}^s(\vec{p}) &= \not{p} + m, & \sum_{s=1}^2 v^s(\vec{p}) \bar{v}^s(\vec{p}) &= \not{p} - m \end{aligned} \quad (2)$$

解：

$$u^r(\mathbf{p})^\dagger u^s(\mathbf{p}) = \begin{pmatrix} \xi^{r\dagger} \sqrt{p \cdot \sigma}, \xi^{r\dagger} \sqrt{p \cdot \bar{\sigma}} \end{pmatrix} \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ \sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix} = \xi^{r\dagger} p \cdot \sigma \xi^s + \xi^{r\dagger} p \cdot \bar{\sigma} \xi^s = 2\omega_p \delta^{rs} \quad (3)$$

$$v^r(\mathbf{p})^\dagger v^s(\mathbf{p}) = \begin{pmatrix} \xi^{r\dagger} \sqrt{p \cdot \sigma}, -\xi^{r\dagger} \sqrt{p \cdot \bar{\sigma}} \end{pmatrix} \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ -\sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix} = \xi^{r\dagger} p \cdot \sigma \xi^s + \xi^{r\dagger} p \cdot \bar{\sigma} \xi^s = 2\omega_p \delta^{rs} \quad (4)$$

$$\begin{aligned} \bar{u}^r(\mathbf{p}) u^s(\mathbf{p}) &= \begin{pmatrix} \xi^{r\dagger} \sqrt{p \cdot \sigma}, \xi^{r\dagger} \sqrt{p \cdot \bar{\sigma}} \end{pmatrix} \gamma^0 \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ \sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix} \\ &= \begin{pmatrix} \xi^{r\dagger} \sqrt{p \cdot \bar{\sigma}}, \xi^{r\dagger} \sqrt{p \cdot \sigma} \end{pmatrix} \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ \sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix} = \xi^{r\dagger} (\sqrt{(p \cdot \sigma)(p \cdot \bar{\sigma})} + \sqrt{(p \cdot \bar{\sigma})(p \cdot \sigma)}) \xi^s \\ &= 2m \delta^{rs} \end{aligned} \quad (5)$$

$$\begin{aligned}
\bar{v}^r(\mathbf{p})v^s(\mathbf{p}) &= \begin{pmatrix} \xi^{r\dagger}\sqrt{p\cdot\sigma}, -\xi^{r\dagger}\sqrt{p\cdot\bar{\sigma}} \end{pmatrix} \gamma^0 \begin{pmatrix} \sqrt{p\cdot\sigma}\xi^s \\ -\sqrt{p\cdot\bar{\sigma}}\xi^s \end{pmatrix} \\
&= \begin{pmatrix} -\xi^{r\dagger}\sqrt{p\cdot\bar{\sigma}}, \xi^{r\dagger}\sqrt{p\cdot\sigma} \end{pmatrix} \begin{pmatrix} \sqrt{p\cdot\sigma}\xi^s \\ -\sqrt{p\cdot\bar{\sigma}}\xi^s \end{pmatrix} = -\xi^{r\dagger}(\sqrt{(p\cdot\sigma)(p\cdot\bar{\sigma})} + \sqrt{(p\cdot\bar{\sigma})(p\cdot\sigma)})\xi^s \\
&= -2m\delta^{rs}
\end{aligned} \tag{6}$$

$$\begin{aligned}
\bar{u}^r(\mathbf{p})v^s(\mathbf{p}) &= \begin{pmatrix} \xi^{r\dagger}\sqrt{p\cdot\sigma}, \xi^{r\dagger}\sqrt{p\cdot\bar{\sigma}} \end{pmatrix} \gamma^0 \begin{pmatrix} \sqrt{p\cdot\sigma}\xi^s \\ \sqrt{p\cdot\bar{\sigma}}\xi^s \end{pmatrix} \\
&= \begin{pmatrix} \xi^{r\dagger}\sqrt{p\cdot\bar{\sigma}}, \xi^{r\dagger}\sqrt{p\cdot\sigma} \end{pmatrix} \begin{pmatrix} \sqrt{p\cdot\sigma}\xi^s \\ \sqrt{p\cdot\bar{\sigma}}\xi^s \end{pmatrix} = \xi^{r\dagger}(\sqrt{(p\cdot\sigma)(p\cdot\bar{\sigma})} + \sqrt{(p\cdot\bar{\sigma})(p\cdot\sigma)})\xi^s \\
&= 2m\delta^{rs}
\end{aligned} \tag{7}$$

$$\begin{aligned}
\bar{u}^r(\mathbf{p})v^s(\mathbf{p}) &= \begin{pmatrix} \xi^{r\dagger}\sqrt{p\cdot\sigma}, \xi^{r\dagger}\sqrt{p\cdot\bar{\sigma}} \end{pmatrix} \gamma^0 \begin{pmatrix} \sqrt{p\cdot\sigma}\xi^s \\ -\sqrt{p\cdot\bar{\sigma}}\xi^s \end{pmatrix} \\
&= \begin{pmatrix} \xi^{r\dagger}\sqrt{p\cdot\bar{\sigma}}, \xi^{r\dagger}\sqrt{p\cdot\sigma} \end{pmatrix} \begin{pmatrix} \sqrt{p\cdot\sigma}\xi^s \\ -\sqrt{p\cdot\bar{\sigma}}\xi^s \end{pmatrix} = \xi^{r\dagger}(\sqrt{(p\cdot\sigma)(p\cdot\bar{\sigma})} - \sqrt{(p\cdot\bar{\sigma})(p\cdot\sigma)})\xi^s \\
&= 0
\end{aligned} \tag{8}$$

$$\begin{aligned}
u^r(\mathbf{p})^\dagger v^s(-\mathbf{p}) &= \begin{pmatrix} \xi^{r\dagger}\sqrt{p\cdot\sigma}, \xi^{r\dagger}\sqrt{p\cdot\bar{\sigma}} \end{pmatrix} \begin{pmatrix} \sqrt{p\cdot\bar{\sigma}}\xi^s \\ -\sqrt{p\cdot\sigma}\xi^s \end{pmatrix} = \xi^{r\dagger}(\sqrt{(p\cdot\sigma)(p\cdot\bar{\sigma})} - \sqrt{(p\cdot\bar{\sigma})(p\cdot\sigma)})\xi^s \\
&= 0
\end{aligned} \tag{9}$$

$$\begin{aligned}
\sum_s u^s(\mathbf{p})\bar{u}^s(\mathbf{p}) &= \sum_s \begin{pmatrix} \sqrt{p\cdot\sigma}\xi^s \\ \sqrt{p\cdot\bar{\sigma}}\xi^s \end{pmatrix} \begin{pmatrix} \xi^{s\dagger}\sqrt{p\cdot\bar{\sigma}}, \xi^{s\dagger}\sqrt{p\cdot\sigma} \end{pmatrix} \\
&= \sum_s \begin{pmatrix} \sqrt{p\cdot\sigma}\xi^s\xi^{s\dagger}\sqrt{p\cdot\bar{\sigma}} & \sqrt{p\cdot\sigma}\xi^s\xi^{s\dagger}\sqrt{p\cdot\sigma} \\ \sqrt{p\cdot\bar{\sigma}}\xi^s\xi^{s\dagger}\sqrt{p\cdot\bar{\sigma}} & \sqrt{p\cdot\bar{\sigma}}\xi^s\xi^{s\dagger}\sqrt{p\cdot\sigma} \end{pmatrix} \\
&= \begin{pmatrix} \sqrt{p\cdot\sigma}\sqrt{p\cdot\bar{\sigma}} & \sqrt{p\cdot\sigma}\sqrt{p\cdot\sigma} \\ \sqrt{p\cdot\bar{\sigma}}\sqrt{p\cdot\bar{\sigma}} & \sqrt{p\cdot\bar{\sigma}}\sqrt{p\cdot\sigma} \end{pmatrix} = \begin{pmatrix} m & p\cdot\sigma \\ p\cdot\bar{\sigma} & m \end{pmatrix} = \not{p} + m
\end{aligned} \tag{10}$$

$$\begin{aligned}
\sum_s v^s(\mathbf{p}) \bar{v}^s(\mathbf{p}) &= \sum_s \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ -\sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix} \begin{pmatrix} -\xi^{s\dagger} \sqrt{p \cdot \bar{\sigma}}, \xi^{s\dagger} \sqrt{p \cdot \sigma} \end{pmatrix} \\
&= \sum_s \begin{pmatrix} -\sqrt{p \cdot \sigma} \xi^s \xi^{s\dagger} \sqrt{p \cdot \bar{\sigma}} & \sqrt{p \cdot \sigma} \xi^s \xi^{s\dagger} \sqrt{p \cdot \sigma} \\ \sqrt{p \cdot \bar{\sigma}} \xi^s \xi^{s\dagger} \sqrt{p \cdot \bar{\sigma}} & -\sqrt{p \cdot \bar{\sigma}} \xi^s \xi^{s\dagger} \sqrt{p \cdot \sigma} \end{pmatrix} \\
&= \begin{pmatrix} -\sqrt{p \cdot \sigma} \sqrt{p \cdot \bar{\sigma}} & \sqrt{p \cdot \sigma} \sqrt{p \cdot \sigma} \\ \sqrt{p \cdot \bar{\sigma}} \sqrt{p \cdot \bar{\sigma}} & -\sqrt{p \cdot \bar{\sigma}} \sqrt{p \cdot \sigma} \end{pmatrix} = \begin{pmatrix} -m & p \cdot \sigma \\ p \cdot \bar{\sigma} & -m \end{pmatrix} = \not{p} - m
\end{aligned} \tag{11}$$

习题 2

计算如下场算符分别在 P,T,C 作用下的变换：

$$\begin{aligned}
&i\bar{\psi}\gamma^5\psi \\
&\bar{\psi}\gamma^\mu\psi \\
&\bar{\psi}\gamma^\mu\gamma^5\psi
\end{aligned} \tag{12}$$

解：

$$\begin{aligned}
P i\bar{\psi}\gamma^5\psi P^{-1} &= i\bar{\psi}(t, -\vec{x}) \gamma^0 \gamma^5 \gamma^0 \psi(t, -\vec{x}) = -i\bar{\psi}\gamma^5\psi(t, -\vec{x}) \\
T i\bar{\psi}\gamma^5\psi T^{-1} &= -i (\gamma^1 \gamma^3 \psi(-t, \vec{x}))^\dagger \gamma^{0*} \gamma^{5*} \gamma^1 \gamma^3 \psi(-t, \vec{x}) = -i\psi^\dagger \gamma^0 \gamma^5 \psi(-t, \vec{x}) = -i\bar{\psi}\gamma^5\psi(-t, \vec{x}) \\
C i\bar{\psi}\gamma^5\psi C^{-1} &= i (-i\gamma^0 \gamma^2 \psi)^\text{T} \gamma^5 (-\bar{\psi} i\gamma^0 \gamma^2)^\text{T} = -i (\bar{\psi} i\gamma^0 \gamma^2 \gamma^{5\text{T}} i\gamma^0 \gamma^2 \psi)^\text{T} = \bar{\psi} i\gamma^5 \psi
\end{aligned} \tag{13}$$

$$\begin{aligned}
P \bar{\psi}\gamma^\mu\psi P^{-1} &= (\gamma^0 \psi)^\dagger \gamma^0 \gamma^\mu \gamma^0 \psi(t, -\vec{x}) = \bar{\psi} \gamma^0 \gamma^\mu \gamma^0 \psi(t, -\vec{x}) = \begin{cases} \bar{\psi} \gamma^0 \psi(t, -\vec{x}) \\ -\bar{\psi} \gamma^i \psi(t, -\vec{x}) \end{cases} \\
T \bar{\psi}\gamma^\mu\psi T^{-1} &= \bar{\psi} (-\gamma^1 \gamma^3) \gamma^{\mu*} \gamma^1 \gamma^3 \psi(-t, \vec{x}) = \begin{cases} \bar{\psi} \gamma^0 \psi(-t, \vec{x}) \\ -\bar{\psi} \gamma^i \psi(-t, \vec{x}) \end{cases} \\
C \bar{\psi}\gamma^\mu\psi C^{-1} &= (-i\gamma^0 \gamma^2 \psi)^\text{T} \gamma^\mu (-\bar{\psi} i\gamma^0 \gamma^2)^\text{T} = -(\bar{\psi} i\gamma^0 \gamma^2 \gamma^{\mu\text{T}} i\gamma^0 \gamma^2 \psi)^\text{T} = -\bar{\psi} \gamma^\mu \psi
\end{aligned} \tag{14}$$

$$\begin{aligned}
P \bar{\psi}\gamma^\mu \gamma^5 \psi P^{-1} &= \bar{\psi} \gamma^0 \gamma^\mu \gamma^0 \gamma^0 \gamma^5 \gamma^0 \psi = \begin{cases} -\bar{\psi} \gamma^0 \psi(t, -\vec{x}) \\ \bar{\psi} \gamma^i \psi(t, -\vec{x}) \end{cases} \\
T \bar{\psi}\gamma^\mu \gamma^5 \psi T^{-1} &= (\gamma^1 \gamma^3 \psi)^\dagger \gamma^0 \gamma^{\mu*} \gamma^{5*} \gamma^1 \gamma^3 \psi = \begin{cases} \bar{\psi} \gamma^0 \psi(-t, \vec{x}) \\ -\bar{\psi} \gamma^i \psi(-t, \vec{x}) \end{cases} \\
C \bar{\psi}\gamma^\mu \gamma^5 \psi C^{-1} &= -(\bar{\psi} i\gamma^0 \gamma^2 \gamma^{5\text{T}} \gamma^{\mu\text{T}} i\gamma^0 \gamma^2 \psi)^\text{T} = \bar{\psi} \gamma^\mu \gamma^5 \psi
\end{aligned} \tag{15}$$

习题 3

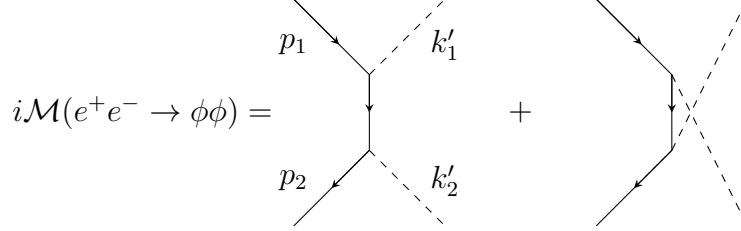
在标量 Yukawa 理论中，拉氏量为：

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m^2 \phi^2 + \bar{\psi} (i\gamma^\mu \partial_\mu - m) \psi - g \phi \bar{\psi} \psi \tag{16}$$

请分别计算 $e^+e^- \rightarrow \phi\phi$ 和 $e^-\phi \rightarrow e^-\phi$ 的最低阶非平庸 S 矩阵元模方 $\sum_s |M|^2$ ，需要对矩阵元自旋求和，及对初态自旋求平均，结果以 Mandelstam 变量表达。如果可能的话，请用交叉对称性（crossing symmetry）解释这两个结果的关系。

解：

对于 $e^+e^- \rightarrow \phi\phi$



即，

$$i\mathcal{M} = \bar{v}_{s_2}(p_2)(-ig)\frac{i(\not{p}_1 - \not{k}'_1 + m)}{t - m^2}(-ig)u_{s_1}(p_1) + \bar{v}_{s_2}(p_2)(-ig)\frac{i(\not{p}_1 - \not{k}'_2 + m)}{u - m^2}(-ig)u_{s_1}(p_1) \quad (17)$$

简记， $u_1 \equiv u_{s_1}(p_1)$ 和 $v_2 \equiv v_{s_2}(p_2)$ 。使用 $\not{p}_1 u_1 = m u_1$ 简化，

$$\mathcal{M} = g^2 \bar{v}_2 \left[\frac{\not{k}'_1 - 2m}{t - m^2} + \frac{\not{k}'_2 - 2m}{u - m^2} \right] u_1 \quad (18)$$

$$\mathcal{M}^* = g^2 \bar{u}_1 \left[\frac{\not{k}'_1 - 2m}{t - m^2} + \frac{\not{k}'_2 - 2m}{u - m^2} \right] v_2 \quad (19)$$

因此，

$$\begin{aligned} |\mathcal{M}|^2 = & + \frac{g^4}{(m^2 - t)^2} \text{Tr}[(v_2 \bar{v}_2)(\not{k}'_1 - 2m)(u_1 \bar{u}_1)(\not{k}'_1 - 2m)] \\ & + \frac{g^4}{(m^2 - u)^2} \text{Tr}[(v_2 \bar{v}_2)(\not{k}'_2 - 2m)(u_1 \bar{u}_1)(\not{k}'_2 - 2m)] \\ & + \frac{g^4}{(m^2 - t)(m^2 - u)} \text{Tr}[(v_2 \bar{v}_2)(\not{k}'_1 - 2m)(u_1 \bar{u}_1)(\not{k}'_2 - 2m)] \\ & + \frac{g^4}{(m^2 - t)(m^2 - u)} \text{Tr}[(v_2 \bar{v}_2)(\not{k}'_2 - 2m)(u_1 \bar{u}_1)(\not{k}'_1 - 2m)]. \end{aligned} \quad (20)$$

对初态自旋求平均，末态求和，

$$\sum_s |\mathcal{M}|^2 = g^4 \left[\frac{\langle \Phi_{tt} \rangle}{(m^2 - t)^2} + \frac{\langle \Phi_{uu} \rangle}{(m^2 - u)^2} + \frac{\langle \Phi_{tu} \rangle + \langle \Phi_{ut} \rangle}{(m^2 - t)(m^2 - u)} \right], \quad (21)$$

其中，

$$\begin{aligned} \langle \Phi_{tt} \rangle &= \frac{1}{4} \text{Tr}[(\not{p}_2 - m)(\not{k}'_1 - 2m)(\not{p}_1 + m)(\not{k}'_1 - 2m)], \\ \langle \Phi_{uu} \rangle &= \frac{1}{4} \text{Tr}[(\not{p}_2 - m)(\not{k}'_2 - 2m)(\not{p}_1 + m)(\not{k}'_2 - 2m)], \\ \langle \Phi_{tu} \rangle &= \frac{1}{4} \text{Tr}[(\not{p}_2 - m)(\not{k}'_1 - 2m)(\not{p}_1 + m)(\not{k}'_2 - 2m)], \\ \langle \Phi_{ut} \rangle &= \frac{1}{4} \text{Tr}[(\not{p}_2 - m)(\not{k}'_2 - 2m)(\not{p}_1 + m)(\not{k}'_1 - 2m)]. \end{aligned} \quad (22)$$

而,

$$\begin{aligned}
\langle \Phi_{tt} \rangle &= \frac{1}{4} \text{Tr} [\not{p}_2 \not{k}'_1 \not{p}_1 \not{k}'_1] + \frac{1}{4} m^2 \text{Tr} [4 \not{p}_1 \not{p}_2 + 2 \not{p}_1 \not{k}'_1 + 2 \not{p}_1 \not{k}'_1 - 2 \not{p}_2 \not{k}'_1 - 2 \not{p}_2 \not{k}'_1 - \not{k}'_1 \not{k}'_1] - m^4 \text{Tr} 1 \\
&= 2 (p_1 k'_1) (p_2 k'_1) - (p_1 p_2) k_1'^2 - m^2 (4 p_1 p_2 + 4 p_1 k'_1 - 4 p_2 k'_1 - k_1'^2) - 4 m^4 \\
&= \frac{1}{2} (t - m^2 - m^2) (u - m^2 - M^2) - \frac{1}{2} (s - 2 m^2) m^2 \\
&\quad - m^2 \left[4 \left(m^2 - \frac{1}{2} s \right) + 2 (t - m^2 - m^2) - 2 (u - m^2 - M^2) + m^2 \right] - 4 m^4 \\
&= -\frac{1}{2} [-tu + m^2(9t + u)]
\end{aligned} \tag{23}$$

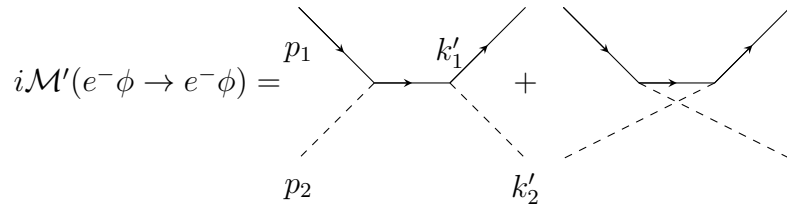
$$\begin{aligned}
\langle \Phi_{tu} \rangle &= \frac{1}{4} \text{Tr} [\not{p}_2 \not{k}'_1 \not{p}_1 \not{k}'_2] + \frac{1}{4} m^2 \text{Tr} [4 \not{p}_1 \not{p}_2 + \not{p}_1 \not{k}'_1 + 2 \not{p}_1 \not{k}'_2 - 2 \not{p}_2 \not{k}'_1 - 2 \not{p}_2 \not{k}'_2 - \not{k}'_1 \not{k}'_2] - m^4 \text{Tr} 1 \\
&= (p_1 k'_1) (p_2 k'_2) + (p_1 k'_2) (p_2 k'_1) - (p_1 p_2) (k'_1 k'_2) \\
&\quad - m^2 (4 p_1 p_2 + 2 p_1 k'_1 + 2 p_1 k'_2 - 2 p_2 k'_1 - 2 p_2 k'_2 - k'_1 k'_2) - 4 m^4 \\
&= \frac{1}{4} (t - m^2 - m^2)^2 + \frac{1}{4} (u - m^2 - m^2)^2 - \frac{1}{4} (s - 2 m^2) (s - 2 m^2) \\
&\quad - m^2 \left[4 \left(m^2 - \frac{1}{2} s \right) - \left(m^2 - \frac{1}{2} s \right) \right] - 4 m^4 \\
&= -\frac{1}{2} [tu + 3m^2(t + u)].
\end{aligned} \tag{24}$$

通过 $t \leftrightarrow u$,

$$\langle \Phi_{tt} \rangle = -\frac{1}{2} [-tu + m^2(9u + t)] \tag{25}$$

$$\langle \Phi_{tu} \rangle = -\frac{1}{2} [tu + 3m^2(t + u)]. \tag{26}$$

对于 $e^- \phi \rightarrow e^- \phi$,



即,

$$i\mathcal{M}' = \bar{u}_{s_2}(k'_1) (-ig) \frac{i(\not{p}_1 + \not{p}_2 + m)}{s - m^2} (-ig) u_{s_1}(p_1) + \bar{u}_{s_2}(k'_1) (-ig) \frac{i(\not{p}_1 - \not{k}'_2 + m)}{u - m^2} (-ig) u_{s_1}(p_1) \tag{27}$$

同理,

$$\sum_s |\mathcal{M}'|^2 = g^4 \left[\frac{\langle \Phi_{ss} \rangle}{(m^2 - s)^2} + \frac{\langle \Phi_{uu} \rangle}{(m^2 - u)^2} + \frac{\langle \Phi_{su} \rangle + \langle \Phi_{ut} \rangle}{(m^2 - s)(m^2 - u)} \right], \tag{28}$$

其中,

$$\begin{aligned}
\langle \Phi_{ss} \rangle &= \frac{1}{2} \text{Tr} [(k'_1 + m) (\not{p}_2 + 2m) (\not{p}_1 + m) (\not{p}_2 + 2m)] = -su + m^2(9s + u), \\
\langle \Phi_{uu} \rangle &= \frac{1}{2} \text{Tr} [(k'_1 + m) (\not{k}'_2 - 2m) (\not{p}_1 + m) (\not{k}'_2 - 2m)] = -su + m^2(9u + s), \\
\langle \Phi_{su} \rangle &= \frac{1}{2} \text{Tr} [(k'_1 + m) (\not{p}_2 + 2m) (\not{p}_1 + m) (\not{k}'_2 - 2m)] = su + 3m^2(s + u), \\
\langle \Phi_{us} \rangle &= \frac{1}{2} \text{Tr} [(k'_1 + m) (\not{k}'_2 - 2m) (\not{p}_1 + m) (\not{p}_2 + 2m)] = su + 3m^2(s + u).
\end{aligned} \tag{29}$$

除了自旋平均多出的 $\frac{1}{2}$ 因子外, cross symmetry 满足的 $p_2 \leftrightarrow -k'_1$ 在 (29) 和 (22) 公式中显然体现出来。

习题 4

如下相互作用拉氏量描述了 $\pi^- \rightarrow \mu^- \bar{\nu}_\mu$ 的衰变:

$$\mathcal{L}_I = 2c_1 G_F f_\pi \partial_\mu \phi \bar{\psi}_{mu} \gamma^\mu P_L \psi_{\nu_\mu} + h.c.$$

其中 ϕ 是复标量场, ψ_{mu} 是 muon 场, ψ_{ν_μ} 是 muon 中微子场。 $G_F = 1.166 \times 10^{-5} \text{GeV}^{-2}$, $G_F = 1.166 \times 10^{-5} \text{GeV}^{-2}$ 是 Cabibbo 角, f_π 称作 pion 衰变常数。 muon 质量为 105.7MeV , pion 质量为 139.6MeV 。 请计算 $\pi^- \rightarrow \mu^- \bar{\nu}_\mu$ 的总衰变宽度 Γ , 并与实验测量得到的 pion 寿命 $T = \frac{\hbar}{\Gamma} \sim 2.603 \times 10^{-8}$ 比较抽取 f_π 的值。 对 $\pi^- \rightarrow e^- \bar{\nu}_e$ 也做同样计算, 所有参数不变, 仅将电子质量设为 0.511MeV 。 并尝试解释两个过程总衰变宽度的差异。

解:

记 $g \equiv c_1 G_F f_\pi$; 顶点为 $(ig) (ik_\mu) \gamma^\mu (1 - \gamma_5) = -g \not{k} (1 - \gamma_5)$, k 为介子动量。 则

$$i\mathcal{M}(\pi^-(k) \rightarrow \mu^-(p_1) \bar{\nu}_\mu(p_2)) = -ig \bar{u}_1 \not{k} (1 - \gamma_5) v_2 \tag{30}$$

对初态自旋求平均, 末态求和,

$$\begin{aligned}
\sum |\mathcal{M}|^2 &= g^2 m_\mu^2 \text{Tr} [(\not{p}_1 + m_\mu) (1 - \gamma_5) (\not{p}_2) (1 + \gamma_5)] \\
&= g^2 m_\mu^2 \text{Tr} [(\not{p}_1 + m_\mu) (\not{p}_2) (1 + \gamma_5) (1 + \gamma_5)] \\
&= 2g^2 m_\mu^2 \text{Tr} [(\not{p}_1 + m_\mu) (\not{p}_2) (1 + \gamma_5)] \\
&= 2g^2 m_\mu^2 \text{Tr} [\not{p}_1 \not{p}_2] \\
&= 2g^2 m_\mu^2 (-4p_1 p_2) \\
&= 4g^2 m_\mu^2 [-(p_1 + p_2)^2 + p_1^2 + p_2^2] \\
&= 4g^2 m_\mu^2 (-k^2 + p_1^2 + p_2^2) \\
&= 4g^2 m_\mu^2 (m_\pi^2 - m_\mu^2).
\end{aligned} \tag{31}$$

那么得到 $\Gamma = \sum |\mathcal{M}|^2 |\mathbf{p}_1| / 8\pi m_\pi^2$, 和 $|\mathbf{p}_1| = (m_\pi^2 - m_\mu^2) / 2m_\pi$, 则

$$\Gamma = \frac{g^2 m_\mu^2}{4\pi m_\pi^3} (m_\pi^2 - m_\mu^2)^2 \approx 2.528 \times 10^{-17} \text{GeV} \tag{32}$$

由此可以得到 $f_\pi = 93.14\text{MeV}$ 。同理对衰变道 $\pi^- \rightarrow e^- \bar{\nu}_e$ ，可以得到 $\Gamma = 5.200 \times 10^{-31}\text{GeV}$ 。此差异是由于弱衰变只有和左手费米子的耦合项，而手征对称性由 $m\bar{\psi}\psi$ 破坏，同时衰变宽度正比于破坏程度导致。