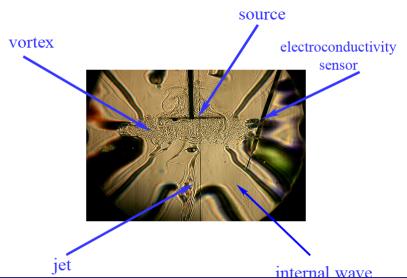
Singularity and limiting cases for oscillations in a viscous continuously stratified fluid problems

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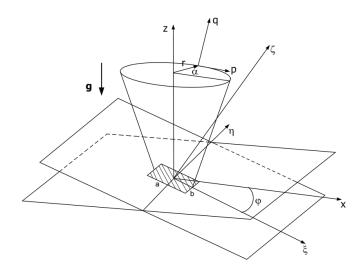


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Intro

Institute for Problems in Mechanics of the RAS

Formulation of the problem.





Governing equations and boundary conditions.

Governing equations

General solutions

$$\frac{\partial \rho}{\partial t} + \mathbf{v} \nabla \rho = 0, \quad \text{div } \mathbf{v} = 0$$

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v}, \nabla) \mathbf{v} = \frac{1}{\rho} \nabla p + \nu \Delta v + \rho \mathbf{g}$$

$$\frac{\partial S}{\partial t} = \kappa_S \Delta S + \frac{v_z}{\Lambda}$$

Boundary conditions

$$|\mathbf{v}|_{\Gamma} = \mathbf{u}_0 e^{-i\omega t}, \quad \kappa_S \left. \frac{\partial S}{\partial n} \right|_{\Gamma} = 0$$

$$v \to 0$$
, $\rho \to \rho_0$, $\partial P/\partial z \to \rho_0(z) g, r \to \infty$

Toroidal-poloidal decomposition

$$\mathbf{v} = \nabla \times \mathbf{e}_z \Phi + \nabla \times (\nabla \times \mathbf{e}_z \Psi)$$

System equation for scalar function Φ and Ψ

General solutions

$$\left(\left(\frac{\partial}{\partial t} - D\Delta \right) \left(\frac{\partial}{\partial t} - \nu \Delta \right) \Delta + N^2 \Delta_{\perp} \right) \Phi = 0$$

$$\left(\frac{\partial}{\partial t} - \nu \Delta \right) \Psi = 0$$

$$\left(\left(\frac{\partial}{\partial t} - D\Delta \right) \left(\frac{\partial}{\partial t} - \nu \Delta \right) \Delta + N^2 \Delta_{\perp} \right) S = 0$$

where

$$\Delta_{\perp} = \partial_{xx}^2 + \partial_{yy}^2$$
 $\Delta = \Delta_{\perp} + \partial_{zz}^2$ $N^2 = \sqrt{\frac{g}{\Lambda}}$



Dispersion equation

$$\left(\nu\kappa_{S}\tilde{k}^{6}-i\omega\left(\nu+\kappa_{S}\right)\tilde{k}^{4}-\omega^{2}\tilde{k}^{2}+N^{2}k_{\perp}^{2}\right)\left(\tilde{k}^{2}+\frac{\omega}{i\nu}\right)=0$$

$$\tilde{k}^2 = 2k_\zeta^2 + k_\perp^2, \quad k_\perp^2 = k_\xi^2 + k_\eta^2$$

Regular solution

$$k_1 = \frac{k_\xi \sin\varphi\cos\varphi \pm \kappa\cos\theta}{\mu_\theta} \pm \delta_N^2 \left(1 + \varepsilon\right) \frac{i \ \tan\theta \mu_\theta^4}{2\kappa\mu^4} + \dots$$

$$\mu = \sin^2 \varphi - \sin^2 \theta, \quad \mu_{\theta} = (k_{\xi} \sin \varphi \cos \varphi \pm \kappa \cos \varphi),$$

$$\varepsilon = Sc^{-1} = \frac{\kappa_S}{\nu}, \quad \delta_N = \sqrt{\frac{\nu}{N}}$$



The viscous stratified fluid taking into account effects of diffusion

Singular solutions

$$k_{2,3} \approx \sqrt{\frac{i\omega\left(\varepsilon + 1 \pm \lambda_{\nu\kappa}\right)}{\varepsilon}}, \quad \lambda_{\nu\kappa} = \frac{2}{\sin\theta}\sqrt{\left(1 + \varepsilon\right)^2 - \frac{4\varepsilon\mu}{\sin^2\theta}}$$

$$k_4 = \sqrt{\frac{2i}{\delta_{\nu}^2} - k^2}, \quad \delta_{\nu} = \delta_N \sqrt{\frac{2}{\sin \theta}}$$

The viscous stratified fluid taking into account effects of diffusion

$$v_{\xi} \approx -\frac{1-i}{2}\delta_{\varphi}G_1 + \frac{i\delta_N^2}{\sqrt{|\mu|}}\tan^2\varphi \ G_2 - G_3,$$

$$Q_n = u_0 \int_{-\infty}^{+\infty} \frac{g_n}{k_\eta^2 \cos \varphi + k_\xi \beta} dk_\xi dk_\eta, \quad \beta = k_\xi \cos \varphi - k_1 \sin \varphi,$$

$$g_0 = \gamma k_\xi k_\eta e_1, \quad g_1 = k_\xi e_1 \left(k_\eta^2 \sin \varphi - k_1 \beta \right),$$

$$g_2 = k_{\eta} Q \exp\left(\frac{i-1}{\delta_{\nu}\xi}\right) \left[k_{\eta}^2 \cos\varphi + k_{\xi} \left(k_1 - \gamma\right)\right], \quad Q = e^{i\left(k_{\xi}\xi + k_{\eta}\eta\right)},$$

$$g_3 = k_{\eta} Q \exp\left(-\frac{i+1}{\delta_{\nu}\xi}\right) \left(k_{\eta}^2 \sin\varphi + k_{\xi}\beta \sin\varphi\right), \quad \gamma = k_{\xi}\sin\varphi + k_1\cos\varphi$$

The viscous stratified fluid taking into account effects of diffusion

$$v_{\eta} \approx (1-i) \, \delta_{\varphi} G_0 + \frac{1-i}{2} \delta_{\nu} \tan \varphi \, G_2 - (1-i) \, \delta_{\varphi} \frac{\cos \varphi}{\sin^2 \varphi} G_3$$

$$v_{\zeta} \approx \frac{i-1}{2} \delta_{\varphi} \int_{-\infty}^{+\infty} u_0 k_{\xi} e_1 dk_{\xi} dk_{\eta} + \frac{i \delta_N^2}{2\sqrt{|\mu|}} \tan \varphi G_2 - \frac{1+i}{\delta_{\varphi} \sin \varphi} G_3$$

$$v_{\xi} = \int_{-\infty}^{+\infty} k_{\eta}^{2} L_{3} dk_{\xi} dk_{\eta} + \int_{-\infty}^{+\infty} k_{\xi}^{2} L_{1} dk_{\xi} dk_{\eta},$$

$$v_{\eta} = \int_{-\infty}^{+\infty} L_3 dk_{\xi} dk_{\eta} + \int_{-\infty}^{+\infty} k_{\xi} k_{\eta} L_1 dk_{\xi} dk_{\eta}, \quad v_{\zeta} = \int_{-\infty}^{+\infty} k_{\eta} k_{\perp}^2 L_2 dk_{\xi} dk_{\eta}$$

$$L_m = u_0 Q \frac{k_1^{2-m} \exp(ik_1\zeta) + k_2^{2-m} \exp(ik_2\zeta)}{\left(k_\eta^2 - k_\xi^2\right) (k_2 - k_1)}, \quad m = 1, 2,$$

$$L_3 = \frac{u_0 Q}{k_\eta^2 - k_\xi^2} \exp\left(-\frac{1 - i}{\delta_\nu}\zeta\right)$$



$$v_{\xi} = \frac{u_0}{\pi^2} \int_{-\infty}^{+\infty} \frac{1}{k_{\xi} (k_1 - k_2)} \sin \frac{k_{\xi} a}{2} \left(k_1 e^{ik_1 \zeta} - k_2 e^{ik_2 \zeta} \right) e^{ik_{\xi} \xi} dk_{\xi},$$

$$v_{\eta} = 0,$$

$$v_{\zeta} = \frac{u_0}{\pi^2} \int_{-\infty}^{+\infty} \frac{1}{k_2 - k_1} \sin \frac{k_{\xi} a}{2} \left(e^{ik_1 \zeta} + e^{ik_2 \zeta} \right) e^{ik_{\xi} \xi} dk_{\xi},$$

$$k_1 = k_{\xi} \cot(\varphi + \theta) \pm \frac{ik_{\xi}^3 \delta_N^2}{2\cos\theta \sin^4(\varphi - \theta)}, \quad k_2 = \frac{1+i}{\delta_{\varphi}}$$



Homogeneous fluid $N \to 0$

Solutions of the dispersion equation

$$k_1 = k_{\perp} \quad k_2 = \frac{1+i}{\delta_{\varphi}} \quad k_3 = \frac{1+i}{\delta_{\kappa}} \quad k_4 = k_2$$

$$M_j = a_j \exp\left(-\sigma_{\kappa}\xi + i\frac{\zeta}{\sqrt{2}\delta_{\kappa}}\right) \sin^2\varphi \int_{-\infty}^{+\infty} A_3 Q b_j dk_{\xi} dk_{\eta}$$

$$\sigma_{\kappa} = \frac{\delta_N}{\delta_{\nu}\delta_{\kappa}}$$

Velocity components

$$v_{\xi} = \int_{-\infty}^{+\infty} A_1 e_1 \left(k_{\eta}^2 \sin \varphi + k_1 \beta_1 \right) dk_{\xi} dk_{\eta}$$

$$+ i \exp \left(-\frac{1-i}{\delta_{\nu}} \zeta \right) \int_{-\infty}^{+\infty} \left(\frac{2}{\delta_{\nu}^2} + B \cos \varphi \right) Q dk_{\xi} dk_{\eta} + M_1$$

$$v_{\eta} = \int_{-\infty}^{+\infty} A_1 \gamma_1 e_1 \kappa_1 dk_{\xi} dk_{\eta}$$

$$- \frac{1-i}{\delta_{\nu}} \exp \left(\frac{1-i}{\delta_{\nu}} \zeta \right) \int_{-\infty}^{+\infty} \left(A_2 \kappa_{\eta} - B \sin \varphi \right) Q dk_{\xi} dk_{\eta} - M_2$$

$$v_{\zeta} = \int_{-\infty}^{+\infty} A_1 e_1 \left(k_{\eta}^2 \sin \varphi - k_{\eta} \beta \right) dk_{\xi} dk_{\eta}$$

$$+ i \exp \left(-\frac{1-i}{\delta_{\nu}} \zeta \right) \int_{-\infty}^{+\infty} \left(\frac{1-i}{\delta_{\varphi}} A_2 k_{\xi} - B \sin \varphi \right) Q dk_{\xi} dk_{\eta} - M_3$$

$$v_{\xi} = \int_{-\infty}^{+\infty} \left(k_{\eta}^2 L_3 dk_{\xi} dk_{\eta} + k_{\eta}^2 L_1 \right) dk_{\xi} dk_{\eta}$$
$$v_{\eta} = \int_{-\infty}^{+\infty} \left(L_3 + k_{\xi} k_{\eta} L_1 \right) dk_{\xi} dk_{\eta}$$
$$v_{\zeta} = \int_{-\infty}^{+\infty} k_{\eta} k_{\perp}^2 L_2 dk_{\xi} dk_{\eta}$$

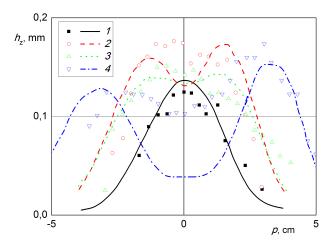
$$L_m = u_0 Q \frac{k_1^{2-m} e^{ik_1 \zeta} + k_1^{2-m} e^{ik_2 \zeta}}{(k_\eta^2 - k_\xi^2)(k_2 - k_1)}, \quad L_3 = \frac{u_0 Q}{k_\eta^2 - k_\xi^2} \exp\left(-\frac{1-i}{\delta_\nu} \zeta\right)$$



Type of sources	Plane(2D)	Plane(3D)	Disk(3D)
Friction	$\frac{u_0}{N} \frac{a}{\delta_N^{1/3}} \frac{1}{q^{2/3}}$	$\frac{u_0}{N} \frac{S}{\delta_N^{2/3} q^{4/3}}$	$\frac{u_0}{\pi N} \frac{S}{\delta_N^{2/3} q^{4/3}}$
Piston	$\frac{u_0}{N} \frac{a}{\delta_N^{2/3}} \frac{1}{q^{1/3}}$	$\frac{u_0}{N} \frac{S}{\delta_N q}$	$\frac{u_0}{\pi N} \frac{S}{\delta_N q}$
Composite	$\frac{u_0}{N} \frac{a}{\delta_N^{4/3}} \frac{a}{q^{2/3}}$	$\frac{u_0}{N} \frac{S}{\delta_N^{5/3}} \frac{b}{q^{4/3}}$	



Comparison of the laboratory experiment (point) and calculations (lines) of the vertical displacement.





Conclusions:

- 1 In general case there are the two types of solutions: regular (internal waves) and three types of singular components of flow (boundary layers). Two of them have no analogue in a homogeneous fluid, their thickness is defined by dissipative factors:
- 2 Near the source viscosity and diffusion is basic factors;
- 3 The obtained results show that it is necessary to consider influence of dissipative factors (viscosity, stratification, diffusion).

Thank you for your attention!
Any quesitions?