

`mfsw_test` is a test directory for running a program, `test.cpp`, built by applying the mean-field theory and spin-wave theory to the general localized interaction model,
$$\sum_{\{i,j\}} \sum_{\{x|x_i \neq x_j\}} \mathcal{H}_{ij} \mathcal{O}_i^{x_i} \mathcal{O}_j^{x_j}$$

- $\sum_{\{i\}} \sum_{\{x_i\}} H_i^{x_i} \mathcal{O}_{\{i\}}^{x_i} = \mathcal{H}^{\text{MF}} + \mathcal{H}'$. $\$$

- Local state
- Spin-wave dispersion
- Dynamical structure factor $S(\mathbf{q}, \omega)$
- Moment contraction (or SW moment)
- Quasiparticle damping rate

As for the dynamical structure factor, in the linear spin-wave theory, $S_{\xi\xi'}(\mathbf{q}, \omega)$ is expressed by $\sum_{\eta}^N \tilde{W}_{\{\eta, \mathbf{q}\}}^{\{\xi\}} \tilde{W}_{\{\eta, \mathbf{q}\}}^{\{\xi'\}} \delta(\omega - \varepsilon_{\{\eta, \mathbf{q}\}})$, where, N is a number of spin-wave bands and $\varepsilon_{\{\eta, \mathbf{q}\}}$ is the excitation energy for the η -th band.

On the program, all $\tilde{W}_{n,\mathbf{g}}^\xi$ are calculated and written to a file opened by `std::ofstream`.

Below is the code to calculate the mean-field solution, spin-wave dispersion, SW moments, and dynamical structure factor.

```
#include <iostream>
#include <fstream>
#include <complex>
#include <string>
#include "cpplapack/cpplapack.h"

using namespace std;

#include "mfsw.hpp"
#include "makedata.hpp"

using exch_type = CPPL::dgematrix;
struct param
{
    vector<tuple<vector<double>, exch_type>> J;
    vector<vector<double>> H;
```

```

    double T;
};

void bond(int n, int Nc, int i, int j, const vector<double> &xs,
mfsw<exch_type>& ms,
        param &p)
{
    int m;
    exch_type Js;
    for (int i = 0; i < p.J.size(); i++)
    {
        vector<double> x = std::get<0>(p.J[i]);
        double sum = 0.;
        for (int k = 0; k < xs.size(); k++)
            sum += abs(x[k] - xs[k]);
        double eps = 1e-4;
        if (sum < eps)
        {
            m = i;
            Js = std::get<1>(p.J[m]);
            break;
        }
        else
        {
            Js.resize(Nc, Nc); Js.zero();
        }
    }
    ms.set_J(n) = std::tuple<int, int, exch_type>{i, j, Js};
}

int main(int argc, char *argv[])
{
    if(argc!=1)
        exit(1);

    std::string fn_lattice =
"data/input/spin_wave_230203/cubic2subbcc_lattice.dat";
    std::string fn_iij =
"data/input/spin_wave_230203/cubic2subbcc_iij_1_1.dat";
    std::string fn_base =
"data/input/spin_wave_230203/cubic2subbcc_base_1_1.dat";

    makedata<exch_type> md(fn_lattice, fn_iij, fn_base);

    complex<double> im(0, 1);

    int Ns = md.make_Ns();           // 局所状態の数 (generator の次元)
    int Nc = md.make_Nc();           // サイトあたりの平均場の数
    int M = 2;                       // 副格子の数
    int N = M * Nc;                  // 平均場の数
    int Z = 8;                       // 1つの注目サイトと相互作用するサイトの合計 (最近
// 接間でのみ相互作用する bcc なら 8)
    int Nb = M * Z/2;                // ユニットセルに含まれるボンドの総数

```

```

int Nr = 5; // 初期値を（ランダムに）入力する回数
int Ni = 100000; // ループ回数

/////*generator*////
vector<CPPL::zhematrix> Oop(Nc); // generator
for (int i = 0; i < Oop.size(); i++)
    Oop[i] = md.make_generator()[i];

param p;
md.make_J(p.J);
md.make_H(p.H, M);

mfsw<exch_type> ms(N, Nc, Nb, Ni, Nr, 1e-8);

for (size_t i = 0; i < N; i += Nc)
{
    for (size_t j = 0; j < Nc; j++)
        ms.set_mat(i + j) = Oop[j];
}

for(size_t i = 0; i < N; i += Nc)
{
    for (size_t j = 0; j < Nc; j++)
    {
        ms.set_H(i + j) = p.H[i / Nc][j];
    }
}

vector<vector<double>> xs(Nb); // ユニットセル内の独立なボンド間の相対座標

/* bond input */
{
    ms.set_bond(xs, Nb, md); // 元データと同じ副格子構造を仮定していれば
md.set_bond を呼び出すだけで良い
}

ms.exec_mf(); // Nr 回のランダムな初期値で解を探す
cout << "# ";
cout << ms.mf_out() << endl; // 平均場解の出力（Nr回実行した内の最安定解のみ
出力される）

p.T = 1e-4; // System の温度
ms.set_T() = p.T;
std::function<complex<double>(vector<double>, vector<double>)> g_;
g_ = [&im](vector<double> x, vector<double> k)
{
    complex<double> pd(1);

```

```

    for (int i = 0; i < x.size(); i++)
        pd = pd * exp(im * x[i] * k[i]);

    return pd;
};

vector<std::function<complex<double>(vector<double>, vector<double>)>>
g(Nb);
for (size_t i = 0; i < Nb; i++)
    g[i] = g_;

vector<vector<double>> a(3);          // 基本並進ベクトル
md.make_unitvec(a);

std::string fns = "data/output/spec.txt";
std::ofstream sw(fns);

{
    double x = 0.;
    double t = 0.;

    /* (-2π, -2π, -2π) --- (2π, 2π, 2π) line */
    for (x = -2.; x < 2.001; x += 0.005)
    {
        vector<double> k(3);
        k[0] = x;
        k[1] = x;
        k[2] = x;
        k[0] *= M_PI;
        k[1] *= M_PI;
        k[2] *= M_PI;
        vector<complex<double>> gamma(g.size());
        for (int i = 0; i < Nb; i++)
            gamma[i] = g[i](xs[i], k);
        sw << x << " " << ms.exec_sw_out(gamma);    // スピン波エネルギーおよび動
        的構造因子の書き込み
        sw << endl;
    }
    t = x;
}

std::string fnmg = "data/output/swmag.txt";
std::ofstream mg(fnmg);
{
    double dk = 0.05;
    mg << " " << ms.exec_sw_mag_out(xs, g, a, dk) << endl;    // SW モーメント
    の書き込み
}

}

```

Note the following:

1. `cout << ms.mf_out() << endl;` outputs in the following order:

$$\begin{aligned} & \mathcal{O}^{\{x_i=0\}}\{A\}, | \mathcal{O}^{\{x_i=1\}}\{A\}, \cdots, \mathcal{O}^{\{x_i=x_{\text{max}}\}}\{A\}, | \mathcal{O}^{\{x_i=0\}}\{B\}, \cdots, \mathcal{O}^{\{x_i=x_{\text{max}}\}}\{B\}, \\ & \mathcal{O}^{\{x_i=0\}}\{C\}, \cdots \end{aligned}$$
 where, capital alphabets represent sublattices. As an example, for the data of "cubic2subbcc", the output is as follows.

```
# -1.00000e+00
0.7071067973 0.5845336644 -0.0056875930 -0.3978543406 0.7071066436
-0.5845336470 0.0056875929 0.3978543288
```

2. `sw << x << " " << ms.exec_sw_out(gamma);` write out in `fsw` in the following order:

$$\begin{aligned} & \epsilon_0, \epsilon_1, \cdots, \epsilon_N, | \mathcal{O}^{\{x_i=0\}}\{A\}, | \mathcal{O}^{\{x_i=1\}}\{A\}, \cdots, \mathcal{O}^{\{x_i=x_{\text{max}}\}}\{A\}, | \\ & \mathcal{O}^{\{x_i=0\}}\{B\}, \cdots, \mathcal{O}^{\{x_i=x_{\text{max}}\}}\{B\}, | \mathcal{O}^{\{x_i=0\}}\{C\}, \cdots \end{aligned}$$
 where, ϵ_i is the expectation value in Bogoliubov vacuum.

3. `mg << " " << ms.exec_sw_mag_out(xs, g, a, dk) << endl;` write out in `fnmg` in the following order:

$$\begin{aligned} & \langle \mathcal{O}^{\{x_i=0\}}\{A\} \rangle, | \langle \mathcal{O}^{\{x_i=1\}}\{A\} \rangle, \cdots, \langle \mathcal{O}^{\{x_i=x_{\text{max}}\}}\{A\} \rangle, | \\ & \langle \mathcal{O}^{\{x_i=0\}}\{B\} \rangle, \cdots, \langle \mathcal{O}^{\{x_i=x_{\text{max}}\}}\{B\} \rangle, | \langle \mathcal{O}^{\{x_i=0\}}\{C\} \rangle, \cdots \end{aligned}$$
 where, $\langle \mathcal{O}^{\{x_i=0\}}\{A\} \rangle$ is the expectation value in Bogoliubov vacuum.

Interfaces

`makedata.hpp` and `mfsfsw.hpp` provide one interface each.

- `makedata`

The primary role of `makedata` class is to produce $\mathcal{O}_i, \mathcal{H}_i$, and the primitive lattice vector from the original data.

- `mfsfsw`

'mfsfsw' class computes the mean-field solution and above quantities. In the mean-field solution, we can change the initial values to file read.

For example, we prepare the initial value input file

```
0.7071067812 0 0 0.7071067812 0.7071067812 0 0 -0.7071067812
```

and change `ms.exec_mf()` to the following.

```
ms.exec_mf("data/input/init.txt");
```

As a result, the output is as follows.

```
# -1.000000e+00  
0.7071067812 0.00000000000 0.00000000000 0.7071067812 0.7071067812  
0.00000000000 0.00000000000 -0.7071067812
```

Note that if you put some initial values, it will solve a self-consistent equation for the number of initial values.

Graphics

The spin-wave dispersion and dynamical structure factor are plotted by `sw.py` and `spec.py`. These python codes use command-line arguments. Run these codes as follows:

```
$ python sw.py M Ne  
$ python spec.py M Ne
```

Here, M is total number of sublattices and Ne is total number of local excited states. For "`cubic2subbcc`", we type as

```
$ python sw.py 2 1  
$ python spec.py 2 1
```

Requirement

- C++11 compatible environment
- `cpplapack-2015.05.11` (header-only library)