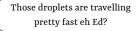
Hats and pancakes in the sky: high-speed droplet dynamics

Presented by Kamran Pentland

28/09/20

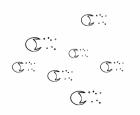
EPSRC & MRC Centre for Doctoral Training in Mathematics for Real-World Systems, University of Warwick

Supervisors: Radu Cimpeanu and Ed Brambley



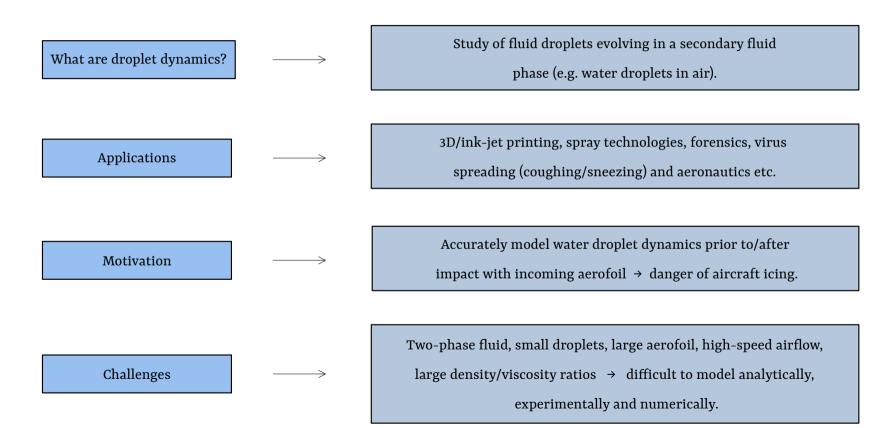








(I) Motivation and aims



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What did we do?

Analysed an existing droplet trajectory and deformation model (simplistic but currently the most advanced analytical model).

Developed a high resolution predictive droplet model

→ via direct numerical simulations (DNS).

Used DNS results to assess validity and accuracy of assumptions in the existing model.

Take home message

Highly challenging to capture/predict pre-impact dynamics in this flow regime -> goal is to develop predictive models (using DNS) that capture trajectory, non-spheroidal deformations and breakup.

Origin

• Developed by Sor et al. $[1] \rightarrow$ experimentally informed force-balance model \rightarrow no fluid mechanics.

Overview

- Two-dimensional flow \rightarrow aerofoil moving at constant speed (flow in front of droplet accelerates).
- Trajectory (x,y) and deformation (a) tracked \rightarrow via force balance equations.
- Taylor analogy → deformation modelled like a mass-spring system (harmonic oscillator).

$$m\frac{d^2x}{dt^2} = -F_{D_x}$$

$$m\frac{d^2y}{dt^2} = F_{D_y} - mg$$

$$m\frac{d^2a}{dt^2} = F_p - F_{st} - F_v$$

Assumptions

- 1. Vertical air flow negligible.
- 2. Droplet deforms as oblate spheroid (see Fig.1).
- 3. No breakup of droplet.

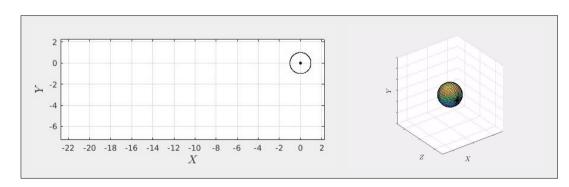


Fig. 1: (Left) Trajectory and deformation of droplet. (Right) 3D rotating view tracking the droplet.

[1] S. Sor, A. García-Magariño, and A. Velazquez, "Model to predict water droplet trajectories in the flow past an airfoil," Aerospace Science and Technology 58, 26–35 (2016).

What did further analysis identify?

- X trajectory varies drastically for small changes in initial velocity.
- Y trajectory can be ignored → gravity has little effect and droplets typically suspended.
- Pressure battles surface tension force → driving the oscillations and deformation.
- Larger droplets oscillate less but deform most uncharacteristic as breakup would occur.

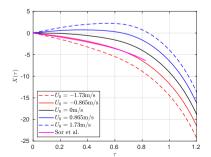


Fig. 2: Effect of initial horizontal velocity on X trajectory vs. time.

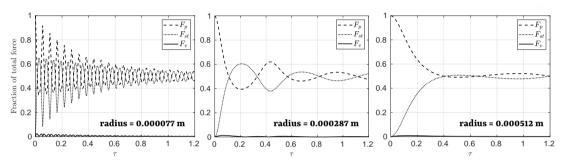


Fig. 4: Relative force contributions to deformation vs. time for various droplet radii.

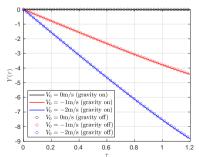


Fig 3: Effect of initial vertical velocity (and gravity) on Y trajectory vs. time.

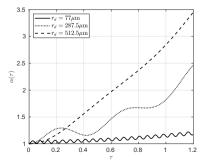


Fig. 5: Vertical deformation vs. time for varying droplet radii.

Overview

• Axisymmetric half-droplet → governed by two-phase Navier-Stokes equations + interface conditions → open-source C library Basilisk.

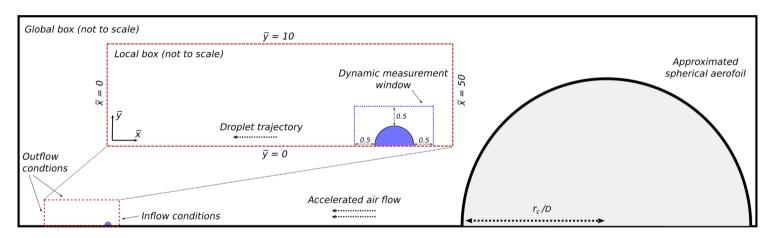


Fig. 6: Global and local computational boxes.

- Global model \rightarrow too computationally heavy \rightarrow requires $> O(10^7)$ grid points.
- Local model → still multi-scale, requiring tens of CPU hours solving in parallel but tractable → requires outflow and inflow conditions.
- Inflow conditions found solving global model *without* droplet → computationally inexpensive (alternatively use analytical potential flow around sphere).

Trajectory and deformation results

- Droplet accelerates \rightarrow initially oblate spheroidal shape \rightarrow assumption breaks down at later times.
- No oscillations observed \rightarrow calls Taylor (mass-spring) analogy into question in this particular case.

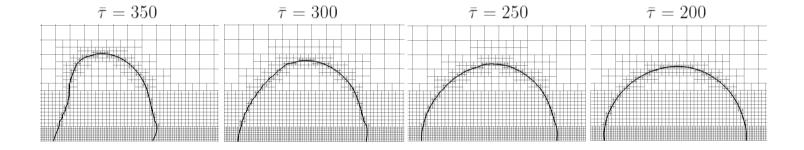


Fig. 7: Trajectory and deformation observed in the DNS model with horizontal velocity field (red = slow flow, blue = fast flow).

What did the flow analysis reveal?

- Flow measured in front, behind and above droplet.
- Non-spheroidal deformation driven by increasing pressure gradient across droplet.
- Negligible vertical velocity assumption verified above droplet center of mass.

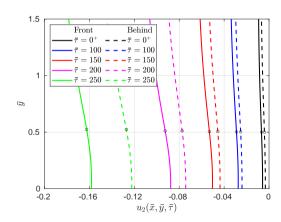


Fig. 8: Horizontal air velocity profile in front/behind droplet.

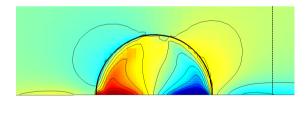
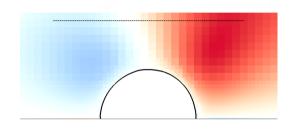


Fig. 9: Pressure field around/inside the droplet prior to deformation.



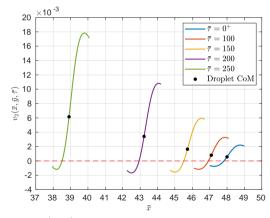


Fig. 10: (Top) Vertical velocity field around droplet prior to deformation. (Bottom) Velocity profiles above droplet.

(IV) Conclusions and future work

Existing model – what did we find?

- Vertical trajectory can be ignored → droplets typically suspended in clouds.
- Oblate spheroidal shape → holds up to certain time → cannot capture non-uniform deformation thereafter.
- Taylor analogy no shape oscillations found in this regime → needs further investigation.
- Verified negligible vertical background flow in stagnation region.
- Difficult to re-create results → heavy reliance on experimental parameters → hinders predictive power.

Numerical model - what did we achieve?

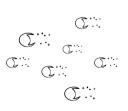
- Good first attempt at predicting pre-impact dynamics in violent flow regime → self contained predictive model.
- High-resolution flow → detailed deformation and flow quantities close to/within droplet.
- Efficient coupling of global and local domains \rightarrow solvable on realistic timescales ($O(10^2)$ CPU hours).

(IV) Conclusions and future work

Future work

- Further numerical validation of DNS over range of droplet sizes/flow conditions.
- Relax negligible vertical airflow assumption → considers droplets away from stagnation region.
- Investigate droplet breakup in accelerating flow vs. constant background flow.





Thank you for listening, questions?



Extra Slides

Experimental setup



Fig: Experimental rotating arm facility used to verify the analytical model (taken from [1]).

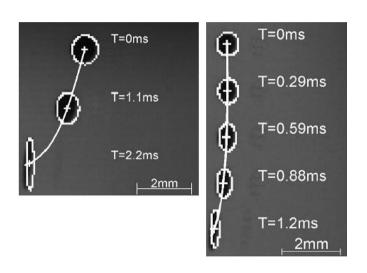
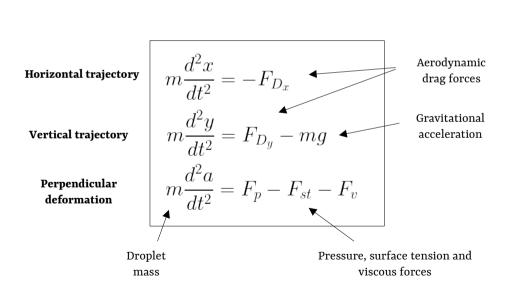


Fig: Analytical model results overlaid on images from the experiments [1].

[1] S. Sor, A. García-Magariño, and A. Velazquez, "Model to predict water droplet trajectories in the flow past an airfoil," Aerospace Science and Technology 58, 26-35 (2016).

Governing equations

- Newton's second law → approximated forces incorporate: accelerating flow, drag laws, surface area change etc.
- Solve numerically using RK4 method.



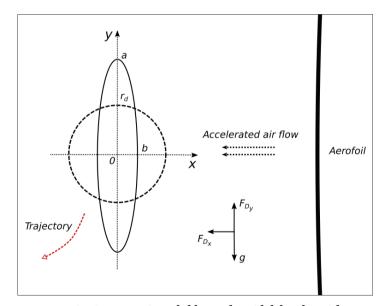
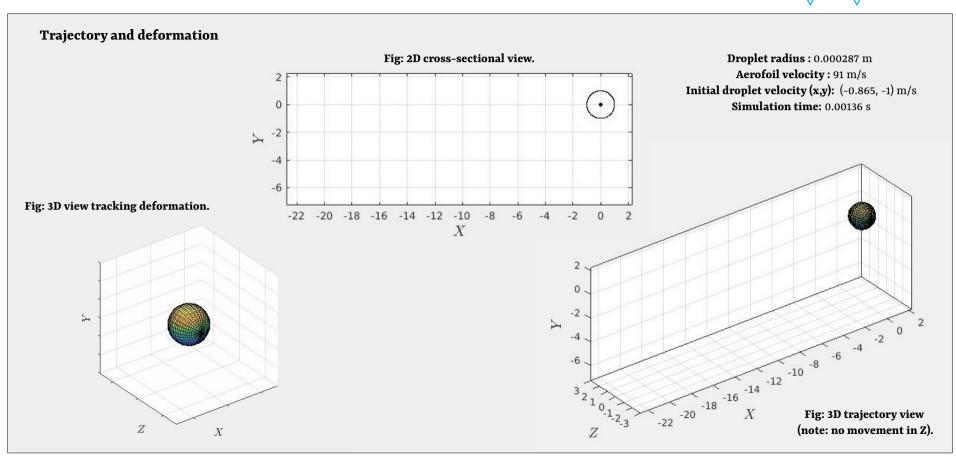


Fig: Cross-section of oblate spheroidal droplet with incoming aerofoil.



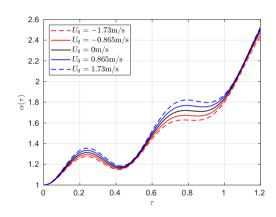


Fig: Deformation vs. time for droplets with varying initial horizontal velocity.

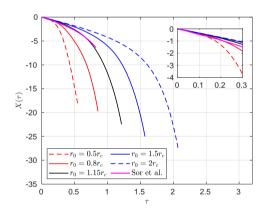


Fig: X displacement vs. time for droplets with varying initial distance from aerofoil.

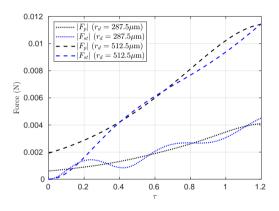


Fig: Force vs. time for droplets of varying radii.

Governing equations

- Two-phase fluid (water/air) + droplet interface between them.
- Requires robust numerical integrator \rightarrow open-source library $Basilisk^{[2]} \rightarrow$ second-order accurate solutions in space/time on adaptive meshes.

Dimensionless Navier-Stokes equations (in each fluid)

$$u_{1\bar{\tau}} + (u_1 \cdot \nabla)u_1 = -\nabla p_1 + \frac{1}{\text{Re}_1} \nabla^2 u_1 - \frac{1}{\text{Fr}^2} F_g$$

$$\nabla \cdot u_1 = 0$$

$$\rho \Big(u_{2\bar{\tau}} + (u_2 \cdot \nabla)u_2 \Big) = -\nabla p_2 + \frac{\mu}{\text{Re}_1} \nabla^2 u_2 - \frac{\rho}{\text{Fr}^2} F_g$$

$$\nabla \cdot u_2 = 0$$

Interface equations on $\bar{y} = h(\bar{x}, \bar{\tau})$

$$u_{1} = u_{2}$$

$$v_{1} = h_{\bar{\tau}} + u_{1}h_{\bar{x}}, \quad v_{2} = h_{\bar{\tau}} + u_{2}h_{\bar{x}}$$

$$\left[4\frac{\mu_{i}}{\mu}h_{\bar{x}}u_{i\bar{x}} + \frac{\mu_{i}}{\mu}(h_{\bar{x}}^{2} - 1)(u_{i\bar{y}} + v_{i\bar{x}})\right]_{2}^{1} = 0$$

$$\left[-p_{i}(1+h_{\bar{x}}^{2}) + \frac{2}{\operatorname{Re}_{1}}\frac{\mu_{i}}{\mu}(h_{\bar{x}}^{2}u_{i\bar{x}} + v_{i\bar{y}} - h_{\bar{x}}(u_{i\bar{y}} + v_{i\bar{x}}))\right]_{2}^{1} = \frac{1}{\operatorname{We}}\frac{h_{\bar{x}\bar{x}}^{2}}{\sqrt{1 + h_{\bar{x}}^{2}}}$$

[2] S. Popinet, "An accurate adaptive solver for surface-tension-driven interfacial flows," Journal of Computational Physics 228, 5838-5866 (2009).

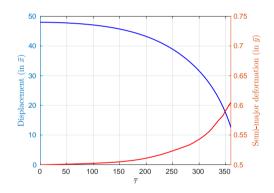


Fig: Displacement and deformation vs. time for droplet from DNS model.

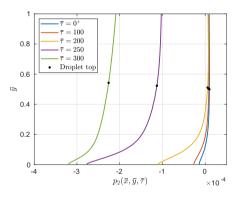


Fig: Air pressure profiles in front of droplet at increasing times.