

# Determination of the beam asymmetry $\Sigma$ in $\eta$ - and $\eta'$ -photoproduction using Bayesian statistics

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I hereby declare that this thesis was formulated by myself and that no sources or tools other than those cited were used.

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## CHAPTER 4

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### Extraction of the beam asymmetries $\Sigma_\eta$ and $\Sigma_{\eta'}$

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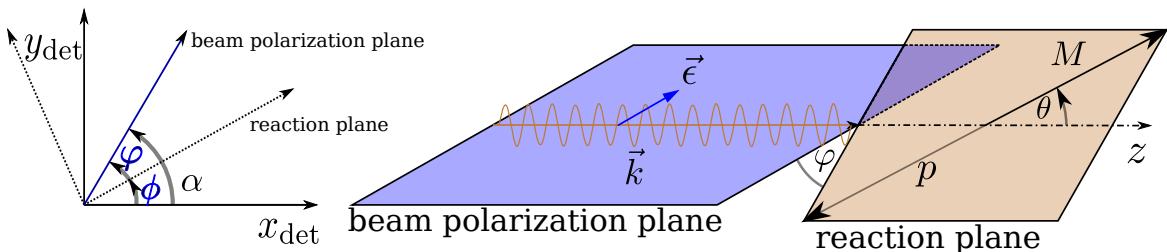
The beam asymmetry  $\Sigma$  is observable when a linearly polarized photon beam and unpolarized liquid hydrogen target are employed. The polarized cross section  $\frac{d\sigma}{d\Omega_{\text{pol}}}$  is not symmetric in the azimuthal angle  $\phi$  anymore as opposed to the unpolarized cross section  $\frac{d\sigma}{d\Omega_0}$ . It is rather modulated by a cosine dependence which scales with the polarization observable  $\Sigma$  and the (linear) beam polarization  $p_\gamma$ , see equation (4.1) [San+11].

$$\frac{d\sigma}{d\Omega_{\text{pol}}} (E_\gamma, \cos \theta, \phi) = \frac{d\sigma}{d\Omega_0} (E_\gamma, \cos \theta) \cdot \left[ 1 - p_\gamma \Sigma (E_\gamma, \cos \theta) \cos (2\phi) \right] \quad (4.1)$$

Since the incident photon beam is polarized, photon momentum  $\vec{k}$  and polarization  $\vec{\epsilon}$  span a plane which is referred to as the beam polarization plane. This plane is tilted by the angle  $\varphi$  with respect to the reaction plane which is defined by the final state momenta. Naturally, this plane builds the angle  $\phi$  in the laboratory system. At the same time the angle of the beam polarization plane in the same reference frame is defined as  $\alpha$ . It holds

$$\varphi = \alpha - \phi. \quad (4.2)$$

Figure 4.1 illustrates definitions of all angles and planes. Theoretically the beam asymmetry can be



**Figure 4.1:** Left: Definition of angles  $\alpha, \phi, \varphi$ . Right: Photon momentum  $\vec{k}$  and polarization  $\vec{\epsilon}$  define the beam polarization plane while the reaction plane is defined by the recoil proton  $p$  and produced meson  $M$ .

68

determined by a measurement of the cross section and a fit using equation (4.1). However, when  
69 calculating polarized cross sections, it is important to have good control over flux normalization and  
70 detector acceptance in three dimensions ( $E_\gamma, \cos \theta, \phi$ ). To avoid this, the measurement of asymmetries  
71

72 can be used to access the polarization observable  $\Sigma$  instead. Particularly, data is taken for two distinct  
73 orthogonal polarization settings corresponding to  $\alpha = \pm 45^\circ$ .

74 This chapter will illustrate the process of determining the beam asymmetry for  $\eta$  and  $\eta'$  photoproduction.  
75 The published results of  $\Sigma_\eta$  [Afz19; Afz+20] are used to check the accuracy and functionality of  
76 employed bayesian methods. Bayesian methods, as well as traditional frequentist approaches are used  
77 afterwards to extract new results for  $\Sigma_{\eta'}$ . First, the used methods will be presented and subsequently  
78 their application for each final state, respectively.

79 **4.1 Methods**

80 The beam asymmetry has to be determined via fits to  $\phi$  distributions obtained from data. These are  
81 performed as either binned or unbinned fits. Both methods allow the application of Bayesian methods  
82 as will be discussed in the following. Additionally the advantages and disadvantages off all methods  
83 are compared.

84 **4.1.1 Event yield asymmetries**

Measurements were made in two distinct polarization settings  $\alpha = \pm 45^\circ = \alpha^{\perp/\parallel}$ . Thus, the polarized cross sections for both settings are given by<sup>1</sup>

$$\frac{d\sigma}{d\Omega_{\text{pol}}}^{\parallel} = \frac{d\sigma}{d\Omega_0} \cdot \left[ 1 - p_\gamma^{\parallel} \Sigma \cos(2(\alpha^{\parallel} - \phi)) \right] \quad (4.3)$$

and

$$\frac{d\sigma}{d\Omega_{\text{pol}}}^{\perp} = \frac{d\sigma}{d\Omega_0} \cdot \left[ 1 - p_\gamma^{\perp} \Sigma \cos(2(\alpha^{\perp} - \phi)) \right] \quad (4.4)$$

$$= \frac{d\sigma}{d\Omega_0} \cdot \left[ 1 + p_\gamma^{\perp} \Sigma \cos(2(\alpha^{\parallel} - \phi)) \right]. \quad (4.5)$$

Note that equation (4.5) holds, because

$$\alpha^{\perp} = \alpha^{\parallel} + \pi/2 \quad \text{and} \quad \cos x = -1 \cdot \cos(x + \pi).$$

Consider now taking the difference of equations (4.3) and (4.5)

$$\frac{d\sigma}{d\Omega_{\text{pol}}}^{\perp} - \frac{d\sigma}{d\Omega_{\text{pol}}}^{\parallel} = \frac{d\sigma}{d\Omega_0} \cdot \left( p_\gamma^{\perp} + p_\gamma^{\parallel} \right) \Sigma \cos(2(\alpha^{\parallel} - \phi)). \quad (4.6)$$

One can further eliminate the unpolarized cross section from this equation by dividing by the polarization weighted sum of equations (4.3) and (4.5)

$$\alpha \cdot \frac{d\sigma}{d\Omega_{\text{pol}}}^{\perp} + \beta \cdot \frac{d\sigma}{d\Omega_{\text{pol}}}^{\parallel} = \frac{d\sigma}{d\Omega_0} \cdot \left[ \alpha + \beta - \left( \alpha p_\gamma^{\perp} - \beta p_\gamma^{\parallel} \right) \Sigma \cos(2(\alpha^{\parallel} - \phi)) \right] \stackrel{!}{=} 2 \frac{d\sigma}{d\Omega_0}. \quad (4.7)$$

<sup>1</sup> The dependencies of polarized and unpolarized cross sections as well as the beam asymmetry like in equation (4.1) are implied

Since

$$\frac{d}{d\phi} \frac{d\sigma}{d\Omega_0} \stackrel{!}{=} 0 \forall \phi,$$

it holds

$$\alpha p_\gamma^{\parallel} - \beta p_\gamma^{\perp} \stackrel{!}{=} 0 \quad \alpha + \beta \stackrel{!}{=} 2, \quad (4.8)$$

such that

$$\alpha = \frac{2p_\gamma^{\parallel}}{p_\gamma^{\perp} + p_\gamma^{\parallel}} \quad \beta = \frac{2p_\gamma^{\perp}}{p_\gamma^{\perp} + p_\gamma^{\parallel}}. \quad (4.9)$$

The beam asymmetry  $\Sigma$  is thus accessible via the asymmetry

$$A(\phi) = \frac{\frac{d\sigma^{\perp}}{d\Omega \text{ pol}} - \frac{d\sigma^{\parallel}}{d\Omega \text{ pol}}}{p_\gamma^{\parallel} \frac{d\sigma^{\perp}}{d\Omega \text{ pol}} + p_\gamma^{\perp} \frac{d\sigma^{\parallel}}{d\Omega \text{ pol}}} = \Sigma \cos(2(\alpha^{\parallel} - \phi)). \quad (4.10)$$

At this point one can now make use of the fact that in any scattering reaction the number of events  $N$  is given by the product of luminosity  $L$  and total cross section  $\sigma$

$$N = L \cdot \sigma = \Phi \cdot N_t \cdot \frac{d\sigma}{d\Omega} \cdot \Delta\Omega,$$

where  $\Phi$  is the beam flux,  $N_t$  the number of target particles and  $\Delta\Omega$  is the solid angle covered by the detector. Substituting this in equation (4.10) one can build the asymmetry  $A(\phi)$  using only the (flux-)normalized event yields  $\tilde{N}^{\parallel/\perp}(E_\gamma, \cos\theta, \phi)^2$

$$A(\phi) = \frac{\tilde{N}^{\perp} - \tilde{N}^{\parallel}}{p_\gamma^{\parallel} \tilde{N}^{\perp} + p_\gamma^{\perp} \tilde{N}^{\parallel}} = \Sigma \cos(2(\alpha^{\parallel} - \phi)). \quad (4.11)$$

Alternatively, the event yields  $N$  can also be normalized by integrating over the total azimuthal angle range in each bin of  $(E_\gamma, \cos\theta)$ . This normalization technique has been used in reference [Afz19] and will also be used in this work. Using appropriate binning in  $\phi$  in addition to beam energy and meson polar angle the asymmetry can be build for all kinematic bins and the beam asymmetry then be extracted via a one-Parameter fit. The statistical errors for  $A(\phi)$  are given by GAUSSIAN error propagation (see appendix A.1).

## 91 Frequentist

92 The beam asymmetry can now be determined via a frequentist fit, where  $\Sigma$  is determined such that the  
93  $\chi^2$  value resulting from the data points and equation 4.11 is minimized. The results are point estimates  
94 with statistical error bars that are also obtained from the fit:  $\hat{\Sigma} \pm \sigma_{\hat{\Sigma}}$ . In addition  $\chi^2/\text{NDF} \approx 1$  may be  
95 verified in order to diagnose the fit itself. Multiple automated minimization and calculation algorithms

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<sup>2</sup> again, arguments  $(E_\gamma, \cos\theta, \phi)$  are implied.

for  $\chi^2$  fitting are available as open source. The *Python* [RP22] module *scipy* [Vir+20] and *ROOT* [BR97] offer e.g. the methods `scipy.optimize.curve_fit` [sci] and `TH1::Fit()` [ROOa] for discrete/binned data, which were used in the analysis.

99 **BAYESIAN**

Following section 1.3, where the basics of BAYESIAN inference were discussed, the goal of a BAYESIAN approach is to sample marginal posterior distributions for each fitted parameter from the joint posterior  $p(\theta|y)$  which depends on the observed data  $y$ . The joint posterior itself is proportional to the product of priors  $\pi(\theta)$  and likelihood  $\mathcal{L}(y|\theta)$  (BAYES’ theorem). This collapses to a one parameter problem in the case of fitting the event yield asymmetries (Eq. (4.11))

$$p(\Sigma|y) \propto \pi(\Sigma) \cdot \mathcal{L}(y|\Sigma). \quad (4.12)$$

However, to be able to sample from a joint posterior, prior and likelihood need to be specified. In order not to bias the fit towards any particular values, the prior is chosen non-informative, realized by a broad GAUSSIAN centered at 0 which is truncated to the physically allowed parameter space of  $\Sigma \in [-1, 1]$ . Furthermore, the likelihood is formulated assuming GAUSSIAN errors  $\epsilon_n$  at each data point  $y_n$ , which should be described by the asymmetry (Eq (4.11)) at bin  $n$   $A(\phi_n; \Sigma)$ , i. e.<sup>3</sup>

$$\Sigma \sim \mathcal{N}(0, 1)_{[-1, 1]} \quad y_n = A(\phi_n; \Sigma) + \epsilon_n \quad \epsilon_n \sim \mathcal{N}(0, \sigma_n), \quad (4.13)$$

which is equivalent to

$$\Sigma \sim \mathcal{N}(0, 1) \quad y_n \sim \mathcal{N}(A(\phi_n; \Sigma), \sigma_n). \quad (4.14)$$

The likelihood of all data points now evaluates to the product of the likelihood at each data point  $y_n$  and the posterior results in

$$p(\Sigma|y) \propto \pi(\Sigma) \cdot \mathcal{L}(y|\Sigma) = \mathcal{N}(\Sigma|0, 1)_{[-1, 1]} \cdot \prod_n \mathcal{N}(y_n|A(\phi_n; \Sigma), \sigma_n) \quad (4.15)$$

$$\Leftrightarrow -\ln p(\Sigma|y) = \frac{1}{2}\Sigma^2 + \frac{1}{2} \sum_n \left( \frac{y_n - A(\phi_n; \Sigma)}{\sigma_n} \right)^2 + \text{constant terms}, \quad (4.16)$$

such that all ingredients are present to form a fully BAYESIAN probabilistic model<sup>4</sup>. This model was implemented in Stan [Sta22], directly giving access to samples from the posterior obtained with the No-U-Turn-Sampler (NUTS) [Sta22; HG14]. Hereby, the sampling is restricted to the allowed parameter region  $\Sigma \in [-1, 1]$ . As a measure of goodness of fit, the  $p$ -values obtained from the posterior predictive distributions, as introduced in section 1.3, are reviewed. To diagnose the convergence of the MCMC fit, sensible values for  $\hat{R}$  and the Monte-Carlo standard error  $\sigma_{\text{MCSE}}$  are verified.

<sup>3</sup> Remember the notation introduced in section 1.3:  $x \sim \mathcal{N}(\mu, \sigma) = \mathcal{N}(x|\mu, \sigma) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{(x-\mu)^2}{2\sigma^2}}$

<sup>4</sup> Note that the sampling aims only to reflect the right proportionality of the (marginal) posterior. Thus, constant terms can be dropped and are of no further interest [Sta22].

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### 107 4.1.2 Event based fit

108 Although intuitive and easily implementable the binned fit –BAYESIAN or not– has one critical  
 109 disadvantage: it is inevitable that information is lost because the asymmetry  $A(\phi)$  is a binned quantity  
 110 and hence, the choice of binning influences the fit results. This is discussed in more detail in  
 111 appendix B. Especially kinematic bins with low statistics show this behavior. To circumvent this  
 112 problem, an *unbinned fit*, based on the likelihood function for each event, can be performed. Also, no  
 113 assumptions on the distribution of statistical errors have to be made since each event is taken into  
 114 account individually. Yet, the event based fit does not provide any measure of goodness of fit, so that  
 115 the study of toy Monte Carlo data is essential when checking the working principle of the method.

In a polarized experiment the azimuthal angle distribution of events is not isotropic, but modulated by a cosine term coupling to beam asymmetry  $\Sigma$  and beam polarization  $p_\gamma^{\parallel/\perp}$  for each setting  $\alpha^{\parallel/\perp}$ , as is expressed through the respective differential cross sections in Equations 4.3 and 4.4. Since the number of events is proportional to the cross section, the probability  $p(\phi, p_\gamma^{\parallel/\perp} | \Sigma)$  to find an event under the azimuthal angle  $\phi$  for a given bin of  $(E_\gamma, \cos \theta)$  and setting  $\alpha^{\parallel/\perp}$  is<sup>5</sup>

$$p(\phi, p_\gamma^{\parallel/\perp} | \Sigma) = \frac{[1 \mp p_\gamma^{\parallel/\perp} \Sigma \cos(2(\alpha^\parallel - \phi))]}{\frac{1}{2\pi} \int_0^{2\pi} d\phi [1 \mp p_\gamma^{\parallel/\perp} \Sigma \cos(2(\alpha^\parallel - \phi))]} \quad (4.17)$$

This is only true for an idealized experiment with acceptance  $\epsilon = \text{const} \forall \phi$ , so that the acceptance  $\epsilon(\phi)$  has to be included in the probability for each event. As demonstrated in reference [Har17] a FOURIER series truncated after the fourth order is sufficient to model any occurring function

$$\epsilon(\phi) = \sum_{k=0}^4 a_k \sin(k\phi) + b_k \cos(k\phi),$$

where the eight detector coefficients  $a_k$  and  $b_k$  are determined from the fit. With this the measurable probability  $\tilde{p}(\phi, p_\gamma^{\parallel/\perp} | \Sigma, a, b)$  is

$$\tilde{p}(\phi, p_\gamma^{\parallel/\perp} | \Sigma, a, b) \propto [1 \mp p_\gamma^{\parallel/\perp} \Sigma \cos(2(\alpha^\parallel - \phi))] \cdot \epsilon(\phi) \quad (4.18)$$

$$= [1 \mp p_\gamma^{\parallel/\perp} \Sigma \cos(2(\alpha^\parallel - \phi))] \cdot \left( \sum_{k=0}^4 a_k \sin(k\phi) + b_k \cos(k\phi) \right), \quad (4.19)$$

where  $a := \{a_k\}_{k=0}^4$ ,  $b := \{b_k\}_{k=0}^4$ . Finally, normalizing  $\frac{1}{2\pi} \int_0^{2\pi} d\phi \tilde{p}(\phi, p_\gamma^{\parallel/\perp} | \Sigma, a, b) \stackrel{!}{=} 1$ ,

$$\tilde{p}(\phi, p_\gamma^{\parallel/\perp} | \Sigma, a, b) = \frac{[1 \mp p_\gamma^{\parallel/\perp} \Sigma \cos(2(\alpha^\parallel - \phi))] \cdot \left( \sum_{k=0}^4 a_k \sin(k\phi) + b_k \cos(k\phi) \right)}{1 \pm \frac{1}{2} a_2 p_\gamma^{\parallel/\perp} \Sigma}. \quad (4.20)$$

---

<sup>5</sup> Note: Normalizing  $p(\phi, p_\gamma^{\parallel/\perp} | \Sigma)$  to  $2\pi$  (or any other arbitrary constant) is sufficient for the fit as long as the integral does not depend on the fit parameters. The normalization to  $2\pi$  is chosen for better readability. However, to calculate actual probabilities, one must multiply Eq. (4.17) by  $2\pi$ .

The respective polarization setting  $\alpha^{\parallel/\perp}$  determines the sign in the normalizing constant which allows an uncorrelated estimation of the detector coefficient  $a_2$  and the beam asymmetry  $\Sigma$ . To simplify notation and implementation, Equation (4.21) can be written as *one* probability for all events – regardless of polarization setting – if the polarization values  $p_\gamma^\parallel$  are multiplied by  $(-1)$  and summarized as  $p_\gamma$ :

$$\tilde{p}(\phi, p_\gamma | \Sigma, a, b) = \frac{[1 + p_\gamma \Sigma \cos(2(\alpha - \phi))] \cdot (\sum_{k=0}^4 a_k \sin(k\phi) + b_k \cos(k\phi))}{1 - \frac{1}{2}a_2 p_\gamma \Sigma}, \quad (4.21)$$

<sup>116</sup> with  $\alpha = -45^\circ$  fixed.

In section 3.2 the subtraction of uncorrelated time background was discussed via a sideband subtraction. Naturally, without binning the data, this strategy is invalid and prompt peak and sideband events have to be fitted simultaneously. Hereby it is important to consider that the random time background is realized as flat distribution underneath the complete reaction time spectrum (cf. Figure 3.3), *including* the range of the prompt peak. The fraction of true coincident events to random coincidences is given as

$$f = \frac{N_{\text{prompt}} - w \cdot N_{\text{sideband}}}{N_{\text{prompt}}}, \quad (4.22)$$

where  $N_{\text{prompt}}$  and  $N_{\text{sideband}}$  are the number of events where the reaction time lies within the prompt peak or sideband, respectively.  $w$  is the ratio of the widths of the chosen prompt peak and sideband ranges, see section 3.2. It is now assumed that the random coincidences will exhibit an asymmetry  $\Sigma^{\text{bkg}}$  of their own and are not necessarily described by the detector coefficients  $a$  and  $b$  but rather by  $a^{\text{bkg}}$  and  $b^{\text{bkg}}$ . The probability to detect a prompt peak event  $p_{\text{prompt}}$  and the probability to measure a sideband event  $p_{\text{sideband}}$  are then given by

$$p_{\text{prompt}}(\phi, p_\gamma | \Sigma, a, b, \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}) = f \cdot \tilde{p}(\phi, p_\gamma | \Sigma, a, b) + (1 - f) \cdot \tilde{p}(\phi, p_\gamma | \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}) \quad (4.23)$$

$$p_{\text{sideband}}(\phi, p_\gamma | \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}) = \tilde{p}(\phi, p_\gamma | \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}). \quad (4.24)$$

If there are  $n$  prompt peak and  $m$  sideband events, the joint likelihood of all events  $\mathcal{L}$  is thus given by

$$\mathcal{L} = \prod_{i=1}^n p_{\text{prompt}}(\phi_i, p_{\gamma,i} | \Sigma, a, b, \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}) \prod_{j=1}^m p_{\text{sideband}}(\phi_j, p_{\gamma,j} | \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}), \quad (4.25)$$

or, equivalently

$$\begin{aligned} \ln \mathcal{L} &= \sum_{i=1}^n \ln p_{\text{prompt}}(\phi_i, p_{\gamma,i} | \Sigma, a, b, \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}) \\ &\quad + \sum_{j=1}^m \ln p_{\text{sideband}}(\phi_j, p_{\gamma,j} | \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}). \end{aligned} \quad (4.26)$$

<sup>117</sup> Eighteen<sup>6</sup> parameters have to be determined in total, either via a conventional frequentist approach or

<sup>6</sup> The FOURIER series is constructed such that  $a_0 = 0, b_0 = 1$

<sup>118</sup> a BAYESIAN approach to this non-linear fitting problem.

<sup>119</sup> **Frequentist**

<sup>120</sup> Best fit estimates can be derived by maximizing the likelihood  $\mathcal{L}$ , or, for computational convenience,  
<sup>121</sup> by minimizing  $-\ln \mathcal{L}$ . The *ROOT* library [BR97] offers the method *TTree::UnbinnedFit* to perform  
<sup>122</sup> an unbinned maximum likelihood fit on data filled in a *TTree* [ROOb], which was used to perform the  
<sup>123</sup> fit. Minimization and statistical error calculation are performed by *MINUIT* [JR75]. Errors are hereby  
<sup>124</sup> estimated either symmetrical from the paraboloid shape of  $(-\ln \mathcal{L})$  using the *HESSE* algorithm or  
<sup>125</sup> asymmetric from the half-maximum values of  $(-\ln \mathcal{L})$  using the *MINOS* algorithm without making  
<sup>126</sup> assumptions on the shape of the likelihood, if necessary [JR75].

<sup>127</sup> **Bayesian**

The joint posterior of all fit parameters given the data  $\phi, p_\gamma$  is

$$p(\Sigma, a, b, \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}} | \phi, p_\gamma) \propto \mathcal{L}(\phi, p_\gamma | \Sigma, a, b, \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}) \cdot \pi(\Sigma, a, b, \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}), \quad (4.27)$$

where  $\pi(\theta)$  denotes the combined prior of all fit parameters which factors into each individual prior since all parameters are independent:

$$\pi(\Sigma, a, b, \Sigma^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}}) = \pi(\Sigma) \cdot \pi(a) \cdot \pi(b) \cdot \pi(\Sigma^{\text{bkg}}) \cdot \pi(a^{\text{bkg}}) \cdot \pi(b^{\text{bkg}}). \quad (4.28)$$

The priors for  $\Xi \in \{\Sigma, \Sigma^{\text{bkg}}\}$  and  $\xi \in \{a, b, a^{\text{bkg}}, b^{\text{bkg}}\}$  are again chosen non-informative, broadly centered around 0

$$\Xi \sim \mathcal{N}(0, 1)_{[-1, 1]} \quad \xi_k \sim \mathcal{N}(0, 0.1). \quad (4.29)$$

<sup>128</sup> From references [Afz19; Har17] it is expected that  $\xi_k \ll 1$ , therefore the chosen widths of the priors  
<sup>129</sup> resemble a, relatively speaking, broad distribution.

<sup>130</sup> This model, consisting of the likelihood in Eq. (4.26) and the priors in Eq. (4.29), is implemented  
<sup>131</sup> in Stan [Sta22], so that samples from the posterior in the physically allowed region ( $\Xi \in [-1, 1]$ ) are  
<sup>132</sup> again obtained via NUTS [HG14].

<sup>133</sup> **4.2 Determination of  $\Sigma_\eta$  using Bayesian statistics**

<sup>134</sup> This section will now demonstrate the application of the discussed methods to obtain the beam  
<sup>135</sup> asymmetry  $\Sigma$  for  $\eta$  photoproduction with selected data provided from reference [Afz19]. For each  
<sup>136</sup> method only the respective BAYESIAN approach will be used and compared to the results from [Afz19]  
<sup>137</sup> to confirm that it is a valid method. As an additional sanity check toy Monte Carlo samples are  
<sup>138</sup> generated and analyzed.

<sup>139</sup> **4.2.1 Application of methods to toy Monte Carlo data**

<sup>140</sup> Although results can be compared to the already accomplished ones in reference [Afz19], verifying  
<sup>141</sup> the correct functioning of the fitting methods is still useful. This is done by generating events that  
<sup>142</sup> follow the expected distributions of  $N^{\parallel/\perp}$  with fixed and known parameters. With the simulated event  
<sup>143</sup> yields the binned and unbinned fit are performed as described previously. Repeating this for a large  
<sup>144</sup> number of times should reproduce the input parameters if the methods work as intended.

<sup>145</sup> **Event yield asymmetries**

The asymmetry  $A(\phi)$  is built from the event yields  $N^{\parallel/\perp}$ , which are distributed according to

$$N^{\parallel} = N_0 \left[ 1 - p_{\gamma}^{\parallel} \Sigma \cos(2(\alpha^{\parallel} - \phi)) \right], \quad (4.30)$$

$$N^{\perp} = N_0 \left[ 1 - p_{\gamma}^{\perp} \Sigma \cos(2(\alpha^{\perp} - \phi)) \right], \quad (4.31)$$

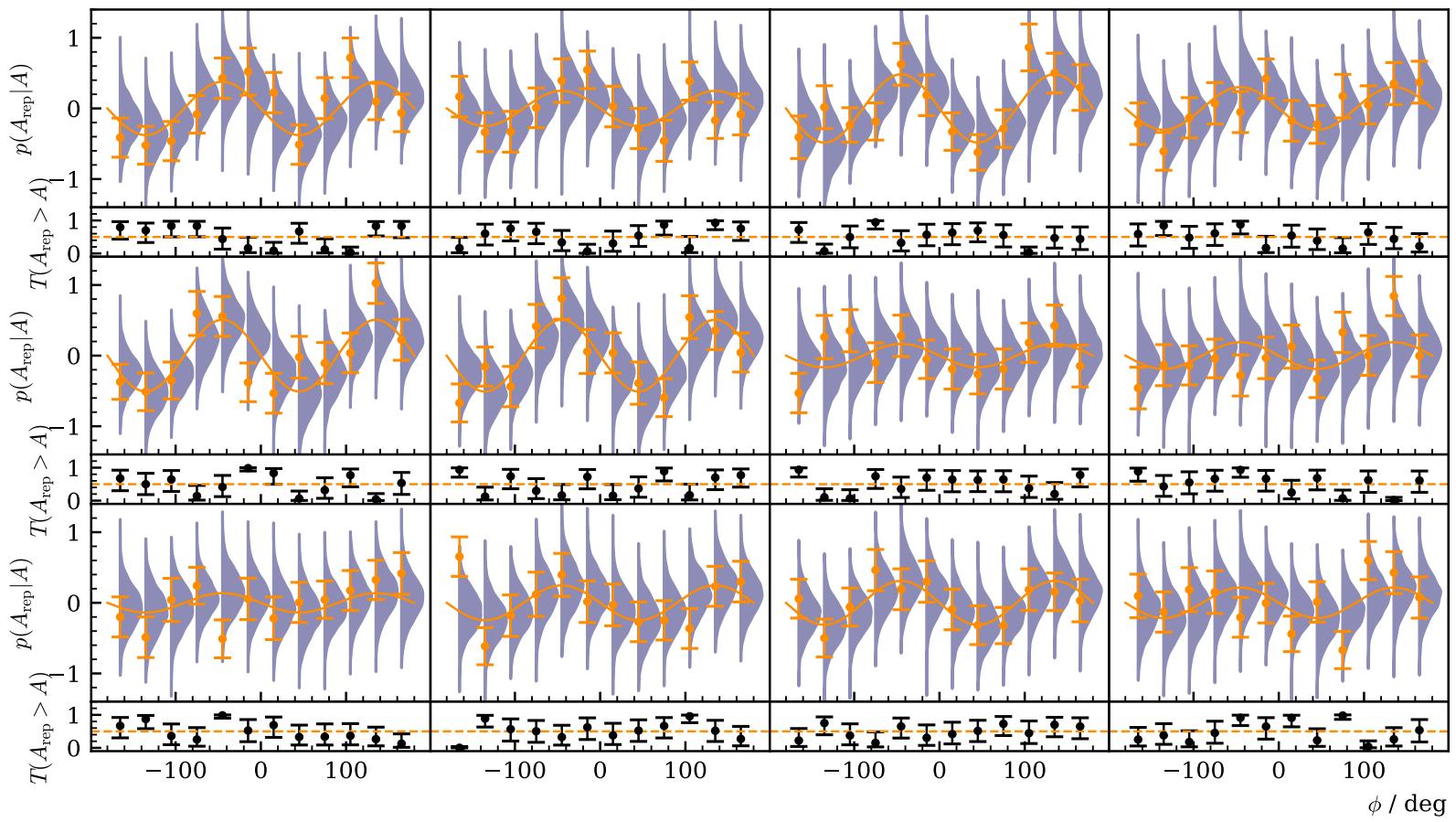
where the parameters are chosen as  $\Sigma = 0.3$ ,  $p_{\gamma}^{\parallel} = 0.25$ ,  $p_{\gamma}^{\perp} = 0.3$ . In each toy Monte Carlo bin the number of generated samples per setting  $N_{\text{total}}^{\parallel/\perp}$  is given by a Poisson distribution

$$N_{\text{total}}^{\parallel} \sim \mathcal{P}(800) \quad N_{\text{total}}^{\perp} \sim \mathcal{P}(1000), \quad (4.32)$$

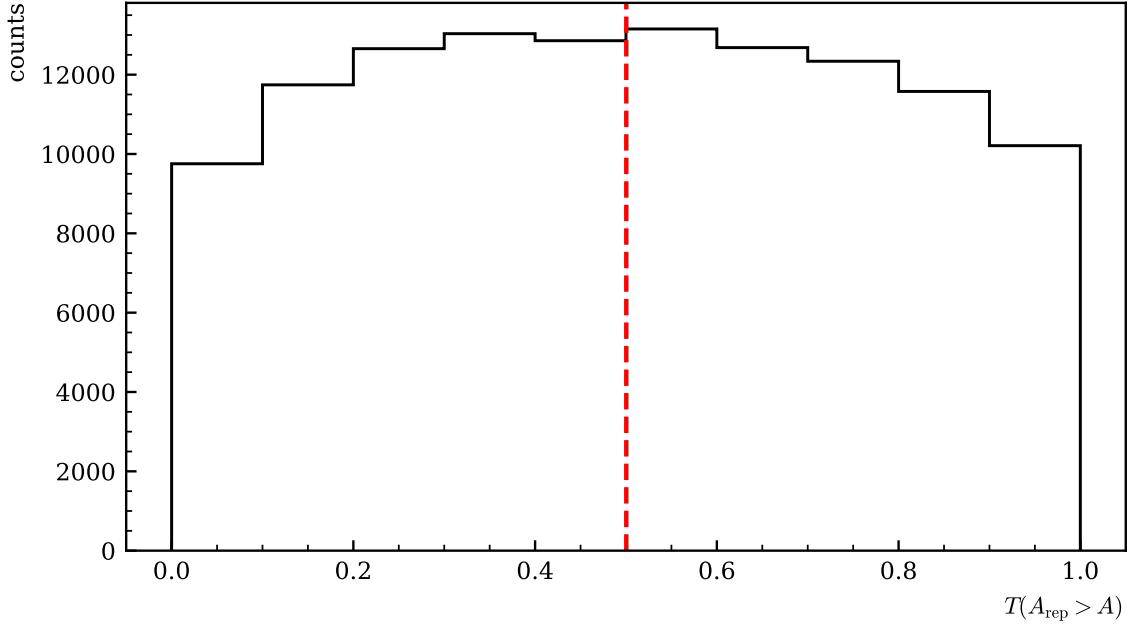
to simulate the statistics of the  $\gamma p \rightarrow p\eta$  final state as accurately as possible [Afz19]. The samples from the distributions (4.30),(4.31) are drawn using the *TH1::GetRandom* [ROOc] function provided by *ROOT* [BR97]. The function from which samples should be drawn is integrated point wise and then normalized. The normalized integral is approximated by a parabola for each bin. A random number between 0 and 1 is generated and assigned to the according bin, where the respective parabola is evaluated to give the desired random value. [ROOc]. In total, 10000 toy Monte Carlo bins were simulated and the asymmetry built for 12 bins in  $\phi$ , to conform with the binning chosen in reference [Afz19]. The resulting asymmetry  $A(\phi)$  is shown in Figure 4.2 for several bins as the orange data points with statistical errors according to GAUSSIAN error propagation. Additionally shown is a  $\chi^2$  fit (orange line) to the asymmetry together with posterior predictive checks as obtained from a fully BAYESIAN fit according to the introduced model (blue distributions). This BAYESIAN fit was performed employing  $n_{\text{chain}} = 4$  MARKOV chains with  $n_{\text{samples}} = 1000$  samples each. The warm-up period for each chain has the same length of  $n_{\text{warm up}} = 1000$

$$n_{\text{chain}} = 4 \quad n_{\text{samples}} = 1000 \quad n_{\text{warm up}} = 1000. \quad (4.33)$$

<sup>146</sup> The goodness of fit is checked via the introduced  $p$ -values  $p = T(A_{\text{rep}} > A)$ , and are shown as black  
<sup>147</sup> points with propagated error bars on the bottom. The optimal value of  $p = 0.5$  is marked by the dashed  
<sup>148</sup> line and realizes the mean of the distribution of all  $p$ -values, so that one can assume good description  
<sup>149</sup> of the data by the fits, see Figure 4.3. This replaces the investigation of the  $\chi^2/\text{NDF}$  distribution in  
<sup>150</sup> the case of a frequentist fit, which should have a mean of 1.



**Figure 4.2:** Posterior predictive checks  $p(A_{\text{rep}} | A)$  from a BAYESIAN fit to the event yield asymmetries for six toy Monte Carlo bins are shown as distributions. The data points in the upper plot are the asymmetry  $A(\phi)$ , which was additionally fitted using a  $\chi^2$  fit (solid line). The goodness of fit is shown using  $p$ -values, which give the fraction  $T(A_{\text{rep}} > A)$  of replicated samples greater than the original measured value, with propagated statistical error bars on the bottom of each plot. The expected mean value of  $T(A_{\text{rep}} > A) = 0.5$  is indicated by the dashed line.



**Figure 4.3:**  $p$  values of all toy Monte Carlo bins. They are centered around their mean at 0.5, which is indicated by the dashed line, and show no bias towards higher or lower values, thus confirming an adequate fit.

To check whether the fit is unbiased and provides correct error estimation one can investigate the normalized residuals

$$\xi = \frac{\Sigma^{\text{fit}} - \Sigma^{\text{true}}}{\Delta\Sigma^{\text{fit}}} \quad (4.34)$$

in the case of a least-squares fit. Here  $\Sigma^{\text{fit}}$  and  $\Delta\Sigma^{\text{fit}}$  are the value and corresponding statistical error for the beam asymmetry as obtained from the fit and  $\Sigma^{\text{true}}$  is the true value that was used to throw the toy MC experiments. An unbiased fit with right estimation of errors yields [Bar89]

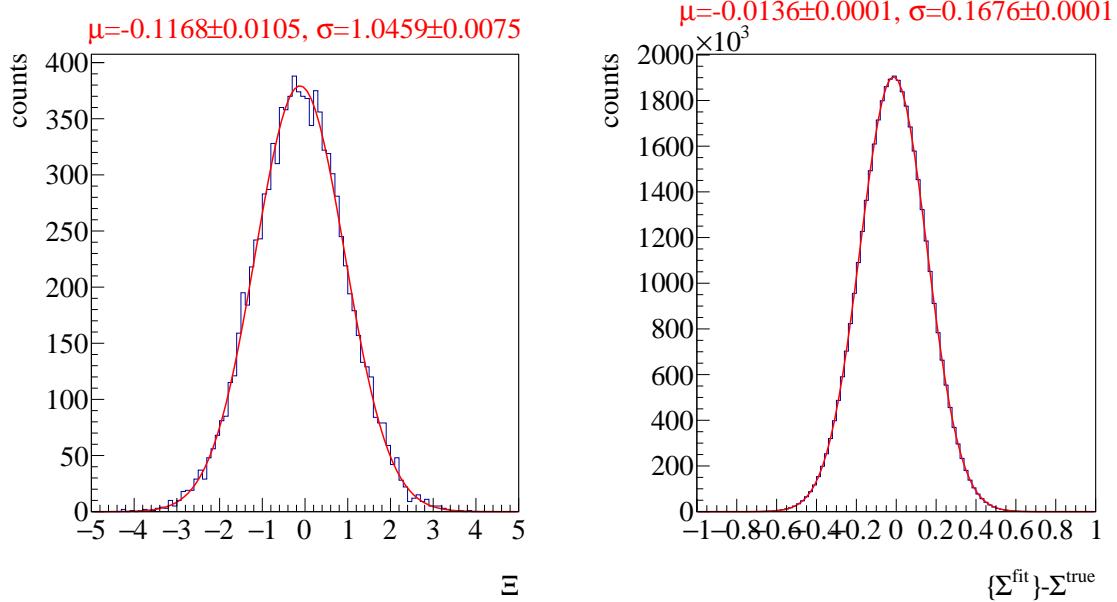
$$\xi \sim \mathcal{N}(0, 1). \quad (4.35)$$

This criterion obviously cannot be applied in the same way to a BAYESIAN fit. The fit results are distributions and therefore lack point estimates  $\Sigma^{\text{fit}}$  and errors  $\Delta\Sigma^{\text{fit}}$ . However, one can modify Equation (4.34) to the needs of a BAYESIAN fit to assess its performance:

$$\Xi = \frac{\text{median}(\{\Sigma^{\text{fit}}\}) - \Sigma^{\text{true}}}{\text{std}(\{\Sigma^{\text{fit}}\})} \sim \mathcal{N}(0, 1). \quad (4.36)$$

<sup>151</sup> Instead of the point estimates  $\Sigma^{\text{fit}}$  the median of the set of all draws from the marginal posteriors  
<sup>152</sup>  $\{\Sigma^{\text{fit}}\}$  is shifted by the true value  $\Sigma^{\text{true}}$  and normalized by the standard deviation as an error estimate  
<sup>153</sup> for each fit. Although this is not a rigorously derived quantity it allows to identify possible bias  
<sup>154</sup> if mean ( $\Xi$ )  $\neq 0$ . Furthermore, checking that  $\sigma \approx 1$  will affirm whether the width of the marginal

posterior distributions is sensible. Figure 4.4 shows the combined  $\Xi$ -distributions of all fits and the unaltered combined posteriors shifted by the true value. As expected, both distributions are normal as a GAUSSIAN fit proves. The width of the marginal posterior distributions can be regarded as sensible



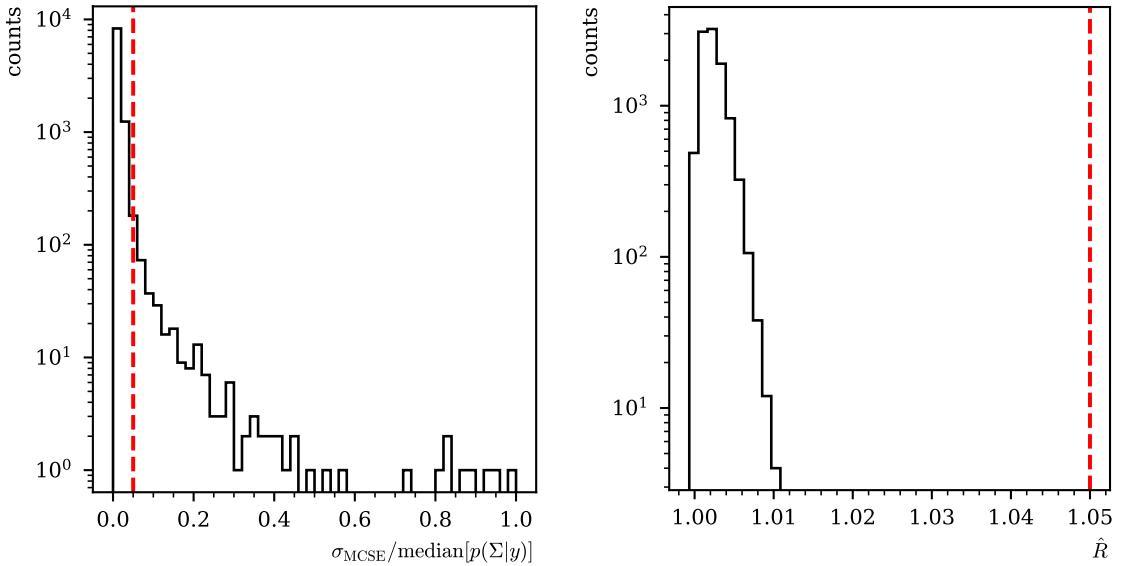
**Figure 4.4:** Left: Combined posterior distributions of all 10000 fits normalized by their respective standard deviation. Right: Unaltered combined posterior distributions of all 10000 fits. A GAUSSIAN fit was performed to determine mean  $\mu$  and standard deviation  $\sigma$  of the distributions with results given on top.

since the standard deviation of the  $\Xi$  distribution  $\sigma_\Xi \approx 1^7$ . This is furthermore confirmed by the fact that  $\{\Sigma^{\text{fit}}\} - \Sigma^{\text{true}} = 0$  within one standard deviation. Yet, the fit results tend to underestimate the beam asymmetry, because both distributions are centered out of their statistical error at  $\mu \pm \Delta\mu < 0$ . This indicates bias towards smaller values of the beam asymmetry. Remarkably, this bias is *not* introduced by the BAYESIAN fit, but by the choice of binning and available statistics as an extensive investigation of binned fits (least squares and BAYESIAN) to toy Monte Carlo data showed. A detailed discussion thereof is given in appendix B. For now it suffices to note the origin of this bias lies in the choice of binning and is thus inevitable. And, most importantly, the introduced bias is negligible compared to the width of the posterior distributions, so that a binned fit will produce valid fit results in any case. Nevertheless this emphasizes the advantage of unbinned fitting, which is discussed in the next paragraph.

It has been established that the fit produces valid results with sensible statistical errors, i.e. widths of posterior distributions, and the description of the data is in agreement with the expectations. It remains to diagnose the convergence of the MARKOV chains to finally deem the binned fit method appropriate. For this the  $\hat{R}$  value and the relative Monte Carlo standard error (MCSE)  $\frac{\sigma_{\text{MCSE}}}{\text{median}[\rho(\Sigma|y)]}$  that were introduced in section 1.3 are investigated. If  $\hat{R} < 1.05$  one can assume that the different chains have converged and explored the same parameter space [Veh+19]. Furthermore the accuracy

<sup>7</sup> One cannot expect to fulfill Eq. (4.35) exactly, since the errors are after all only estimates that aim to create comparability to the least squares fit.

supplied by the number of draws from the posteriors can be evaluated with the relative MCSE. It  
 is demanded that  $\frac{\sigma_{\text{MCSE}}}{\text{median}[P(\Sigma|y)]} \lesssim 5\%$ . This is the case for 97% of all fits, see Figure 4.5 on the  
 left. When the fit result for the beam asymmetry is close to 0 larger values for the relative MCSE  
 are observed. The  $\hat{R}$  values (Figure 4.5 on the right) clearly indicate good within and between chain  
 convergence. This means that the hyper parameters  $n_{\text{chain}}$ ,  $n_{\text{samples}}$  and  $n_{\text{warm up}}$  as chosen are well  
 tuned, although one may increase the number of sample draws from the posterior for the fit from  
 measured data to further suppress the relative MCSE. Then, only  $11 \cdot 12$  as opposed to 10000 fits have  
 to be performed keeping the computational cost reasonable.



**Figure 4.5:** Left: relative error  $\frac{\sigma_{\text{MCSE}}}{\text{median}[P(\Sigma|y)]}$  Right:  $\hat{R}$  associated with the fit parameter  $\Sigma$ . Both are shown for all 10000 fits. The critical values that should not be exceeded are marked by dashed lines.

182

### 183 Event based fit

Not only signal events but also contributions from random time background as well as the imperfect  
 detector efficiency  $\epsilon(\phi)$  have to be simulated in order to test the method of event based fitting. For each  
 polarization setting prompt peak and sideband events are drawn from the theoretical  $\phi$ -distributions  
 $p_{\text{prompt}}$  and  $p_{\text{sideband}}$  which are given in Equations (4.23) and (4.24). The total number of draws per  
 setting and bin is again given by Poisson distributions and the ratio of prompt peak to sideband events  
 is given by the time cut weights  $\{w_i\}_{i=1}^7$  employed for the actual analysis of the  $p\eta$  final state data  
 [Afz19]. This means that seven times prompt peak and sideband events need to be simulated. The  
 fraction of background events within the prompt peak  $f$  was set to  $f = 0.95$ . Further, the values  
 $p_\gamma^\parallel = 0.25$ ,  $p_\gamma^\perp = 0.3$ ,  $\Sigma = 0.5$ ,  $\Sigma^{\text{bkg}} = -0.5$  and lastly, a random efficiency function as already chosen  
 in reference [Afz19] were appointed. Finally the hyperparameters  $n_{\text{chain}}$ ,  $n_{\text{samples}}$ ,  $n_{\text{warm up}}$  are chosen  
 the same as previously. Table 4.1 shows a summary of all toy Monte Carlo properties. In total, 1000  
 toy Monte Carlo experiments are thrown. To evaluate the fit results only the posterior distributions  
 are available. As before, the residuals  $\Xi$  are built from all fits, now for  $\Sigma$  as well as  $\Sigma^{\text{bkg}}$ . They are

<b>chosen parameters</b>	$p_\gamma^{\parallel} = 0.25, p_\gamma^{\perp} = 0.3, \Sigma = 0.5, \Sigma^{\text{bkg}} = -0.5, f = 0.95,$ $w_1 = \frac{15}{210}, w_2 = \frac{8}{210}, w_3 = \frac{4}{210}, w_4 = \frac{10}{210}, w_5 = \frac{14}{210}, w_6 = \frac{6}{210}, w_7 = \frac{11}{210}$
<b>simulation draws</b>	$7 \cdot N_{\text{total},i}^{\parallel} \sim \mathcal{P}(800), 7 \cdot N_{\text{total},i}^{\perp} \sim \mathcal{P}(1000) \quad  _{i=1}^7$
	signal in prompt      background in prompt      sideband $N_{\text{total}}^{\parallel/\perp} \cdot f \quad N_{\text{total}}^{\parallel/\perp} \cdot (1-f) \quad N_{\text{total}}^{\parallel/\perp} \cdot (1-f) \cdot 1/w_i$
<b>efficiency function</b>	$\epsilon(\phi) = 1/10.5 \cdot (9.3 + 0.28 \cdot \cos \phi + 0.24 \cdot \sin 3\phi)$
<b>hyperparameters</b>	$n_{\text{samples}} = 1000, n_{\text{chain}} = 4, n_{\text{warm up}} = 1000$

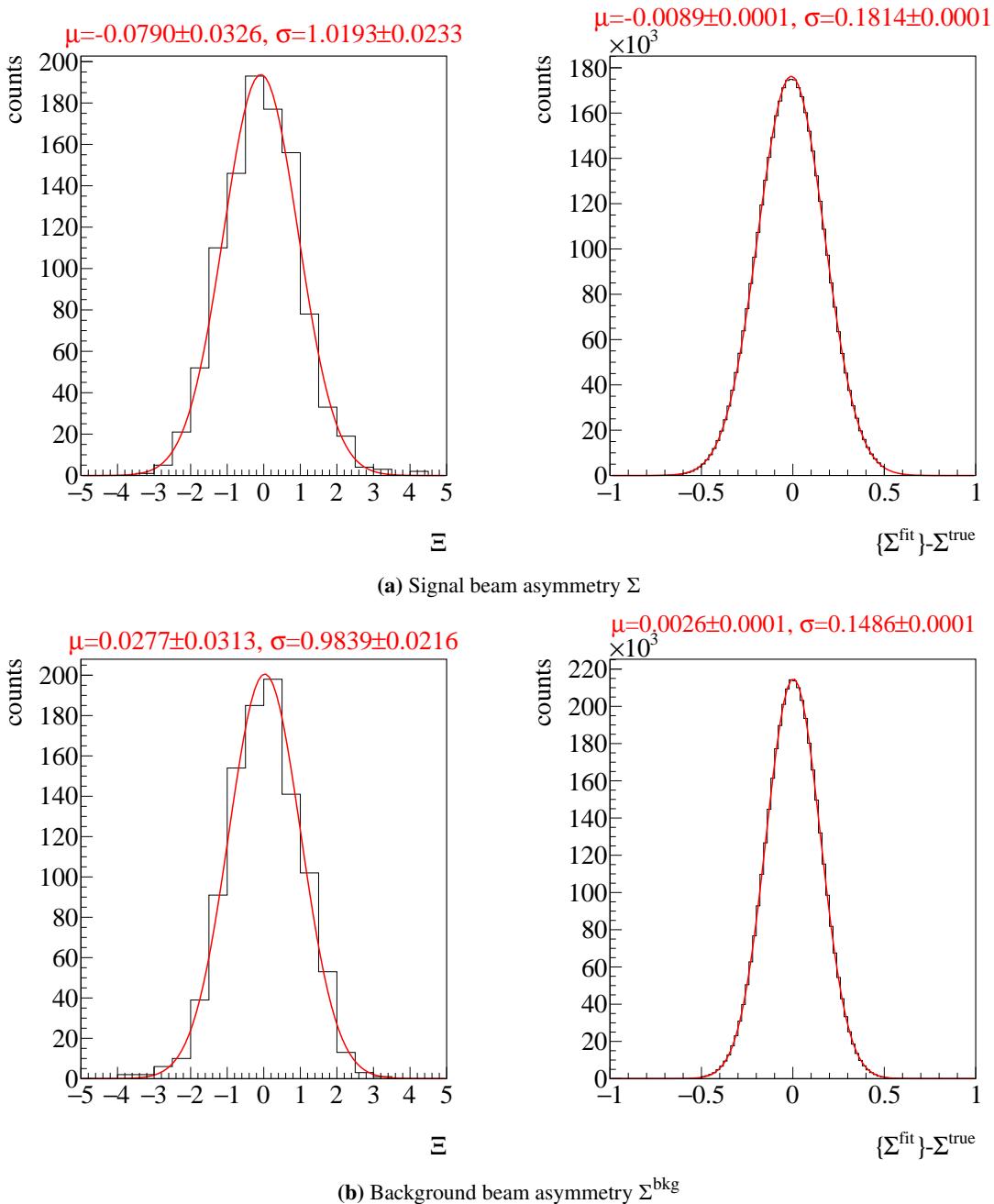
**Table 4.1:** Summary of the complete setting of all toy Monte Carlo experiments for the event based fit. Values and table layout adapted from [Afz19].

197 shown in Figure 4.6 together with the unnormalized posterior distributions. Note that the distributions  
 198 are truncated towards the tails at  $\Sigma \rightarrow 1$  and  $\Sigma^{\text{bkg}} \rightarrow -1$  as implemented in the BAYESIAN fit in Stan.  
 199 Also the event based fit is able to reproduce the true values within one standard deviation of the  
 200 posterior distribution  $\sigma_{\text{posterior}}$ . Only slight bias towards smaller values of the signal beam asymmetry  
 201 is seen, but again with negligible effect on the results in comparison to the widths of the distributions.  
 202 Additionally, the mean of the distributions are 0 within  $3\sigma$  of their statistical error and the errors are  
 203 estimated correctly, since  $\sigma_{\text{posterior}} = 1$  within the statistical error. Note that the background beam  
 204 asymmetry is estimated closer to its true value than the signal beam asymmetry because there are  
 205 more random background events than signal events.

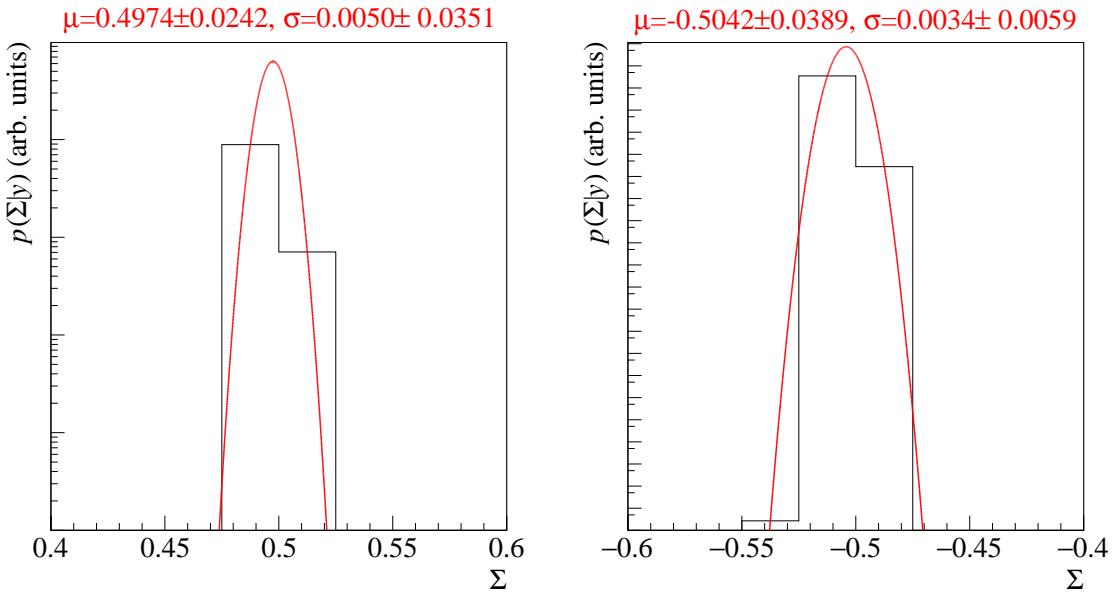
206 If the total amount of posterior distributions is not too large one might also combine them in a  
 207 *pooled likelihood* approach (see subsection 1.3.3). This is shown in Figure 4.7 for both determined  
 208 asymmetries. Because the true asymmetries are within  $1\sigma$  of the resulting distribution one can finally  
 209 conclude that the event based fit estimates positions and widths of the posterior distributions correctly.

210 A check of the relative MCSE and  $\widehat{R}$  values as shown in Figure 4.8 reveals also that the hyperpara-  
 211 meters  $n_{\text{chain}}, n_{\text{samples}}, n_{\text{warm up}}$  are correctly tuned. It is noteworthy that the event based fit creates  
 212 more precise results although less Monte Carlo experiments in total were performed than for the  
 213 binned fit.

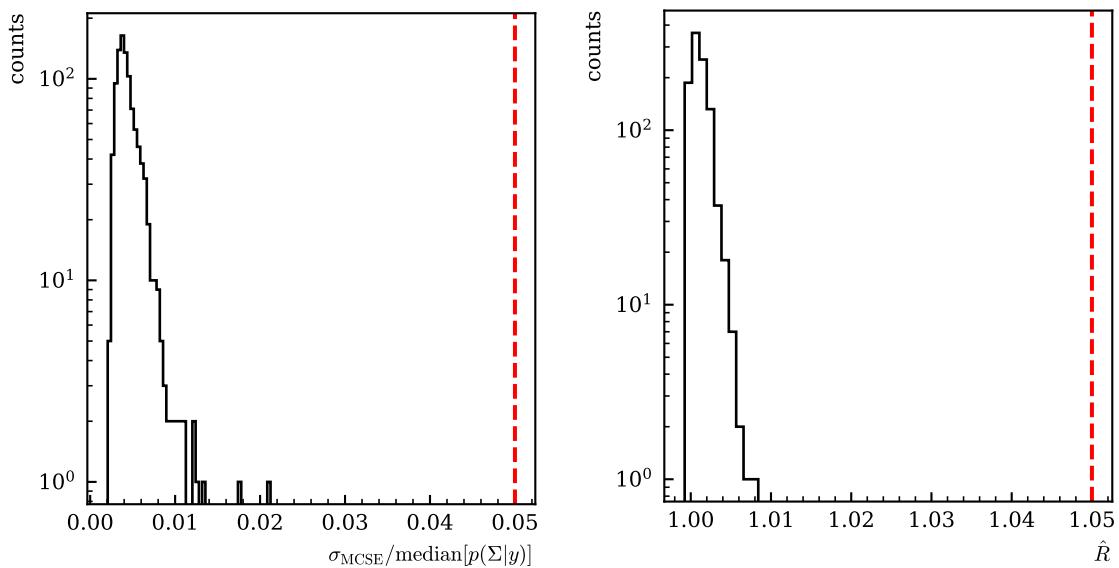
214 Next to signal and background beam asymmetry the fit also estimates an efficiency function  $\epsilon(\phi)$   
 215 that is described by a FOURIER series. The polarization weighted sum of event yields allows to confirm  
 216 whether a good description of the efficiency is achieved because it is given by the efficiency function  
 217 modulated only by a constant, see previous section 4.1  $\alpha \tilde{N}^{\parallel} + \beta \tilde{N}^{\perp} = c \cdot \epsilon(\phi)$ . Figure 4.9 shows the  
 218 weighted sum of event yields with the same binning as used in the binned fit together with a posterior  
 219 predictive check using the draws of all parameters distributions  $a, b$ . As opposed to the binned fit it is  
 220 not only performed at each data point but for significantly more points. This mimics a continuous  
 221 function as is appropriate for an unbinned fit. However, for better visibility the posterior predictive  
 222 distributions are not built at every point but the efficiency function is plotted for each set of draws  
 223  $a^{(s)}, b^{(s)}, s = 1 \dots 4000$ . No  $p$  values are estimated since they would not be representative of the  
 224 whole fit. The mean of the posterior predictive distributions is given by the dashed line and agrees  
 225 within statistical error bars with the simulated data.



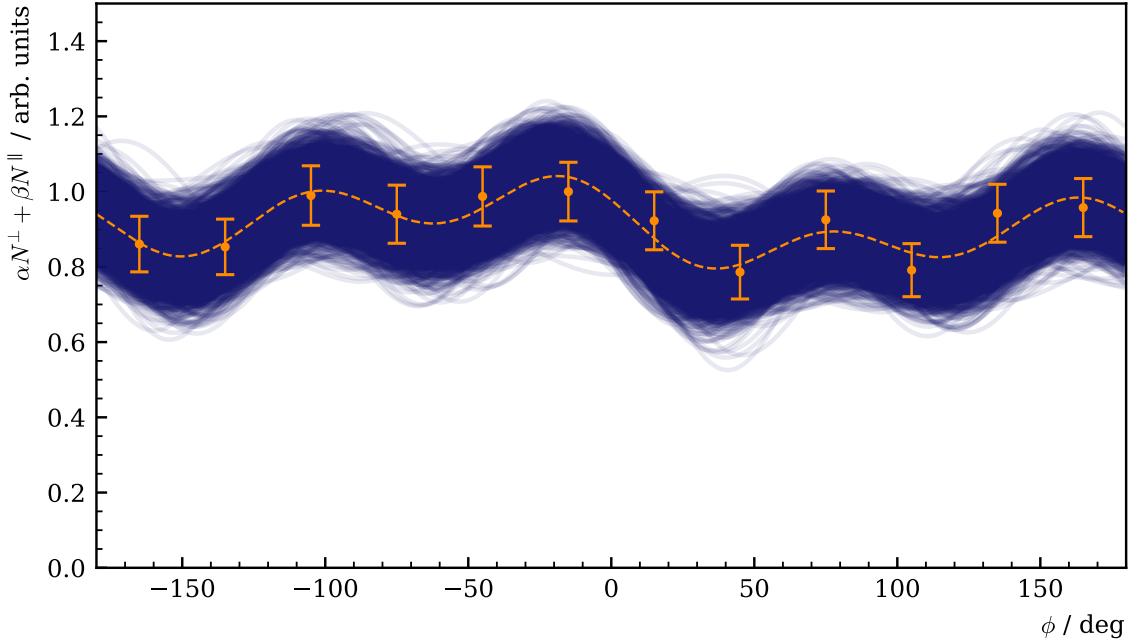
**Figure 4.6:** Combined posteriors for the beam asymmetries  $\Sigma$  and  $\Sigma^{\text{bkg}}$  from all 1000 event based fits. Left: Residuals  $\Xi$  Right: Unnormalized posterior distributions. A GAUSSIAN fit is performed on the distributions with results for mean  $\mu$  and standard deviation  $\sigma$  on top.



**Figure 4.7:** Combined posterior probabilities using the *pooled likelihood* approach. Left: Signal beam asymmetry, Right: background beam asymmetry. Mean and standard deviation as obtained from a Gaussian fit are shown on top



**Figure 4.8:** Left: relative error  $\frac{\sigma_{\text{MCSE}}}{\text{median}[p(\Sigma|y)]}$ . Right:  $\hat{R}$  associated with the fit parameter  $\Sigma$ . Both are shown for all 1000 fits. The critical values that should not be exceeded are marked by dashed lines.



**Figure 4.9:** Posterior predictive check using the draws of the detector coefficients  $a$  and  $b$ . Points with error bars are the polarization weighted sum of event yields. The dashed line is the mean of the predictive values while the solid opaque lines are representative of one simulation draw  $a^{(s)}, b^{(s)}$ .

### 226 4.2.2 Application of methods to data

Since both methods were tested successfully using toy Monte Carlo data, both are subsequently applied to real data to extract the beam asymmetry. Since the number of fits is small compared to the amount of fits performed during testing, the amount of posterior draws was increased to 5000 per chain. All other hyperparameters were left the same

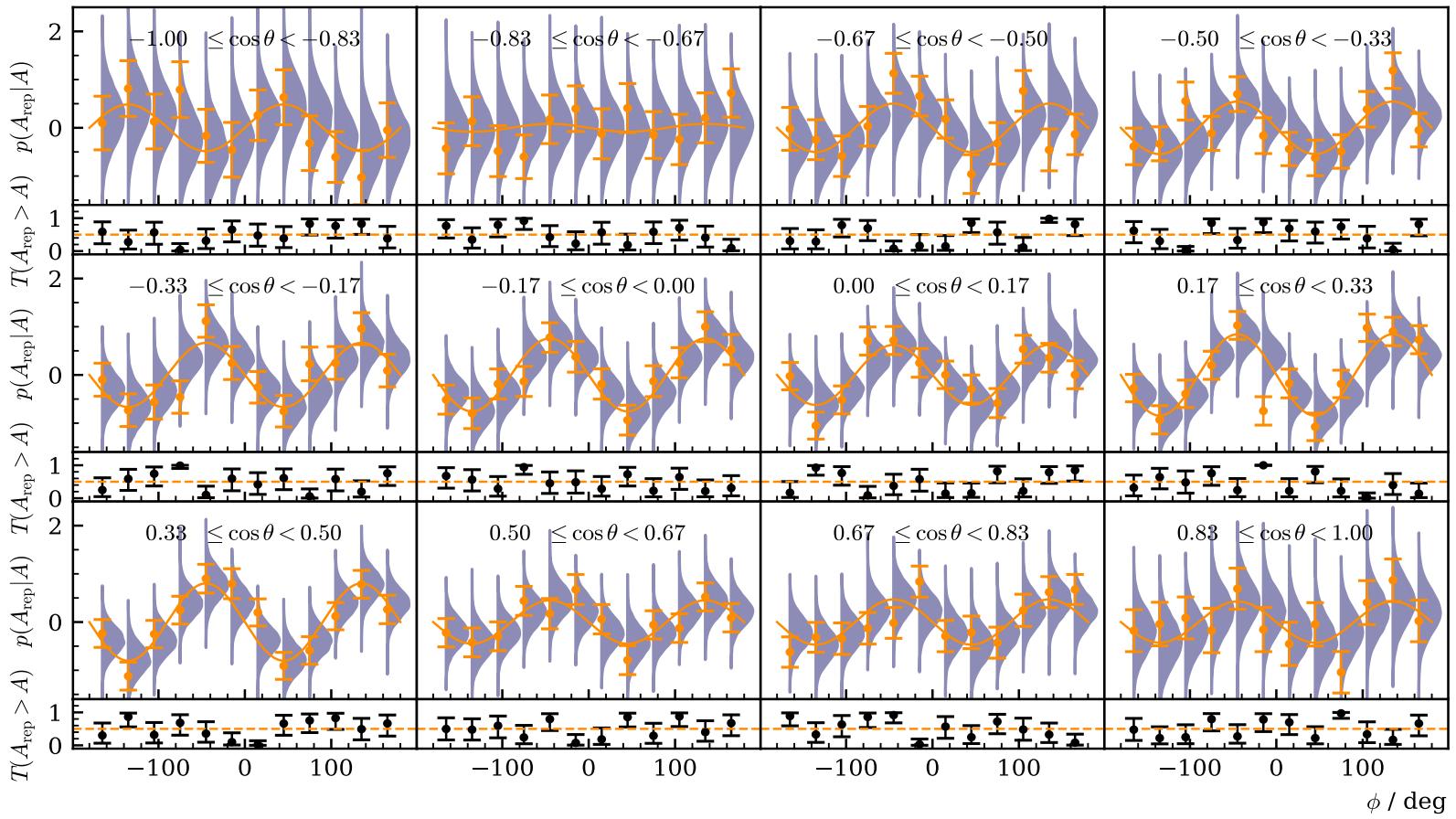
$$n_{\text{samples}} = 5000 \quad n_{\text{chain}} = 4 \quad n_{\text{warm up}} = 1000. \quad (4.37)$$

### 227 Event yield asymmetries

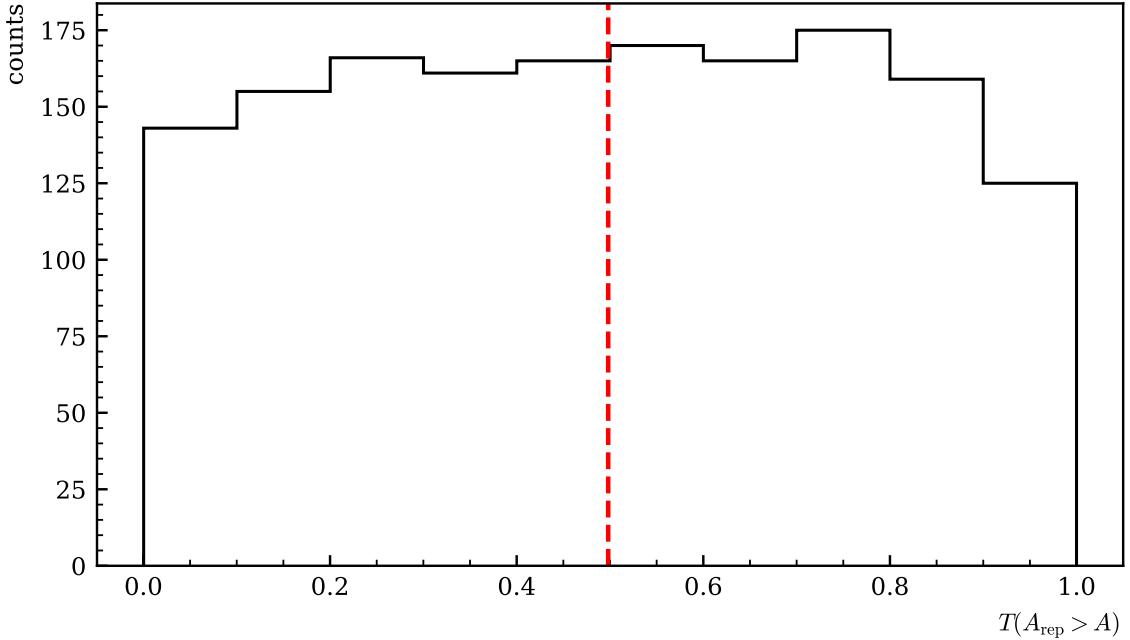
228 In all bins of beam energy and meson polar angle the binned fit was performed. Similar to the Monte  
229 Carlo samples, 12 bins in phi are built to comply with the binning in reference [Afz19], and guarantee  
230 comparability.

231 Figure 4.10 shows the resulting asymmetry  $A(\phi)$  for all angular bins for the energy bin  $1250 \text{ MeV} \leq$   
232  $E_\gamma < 1310 \text{ MeV}$  as orange data points with statistical error bars. Additionally shown is a  $\chi^2$  fit  
233 (orange line) to the asymmetry together with posterior predictive checks as obtained from a fully  
234 BAYESIAN fit according to the introduced model (blue distributions). The goodness of fit is checked  
235 via the introduced  $p$ -values  $p = T(A_{\text{rep}} > A)$ , and are shown as black points with propagated error  
236 bars on the bottom. The optimal value of  $p = 0.5$  is marked by the dashed line and realizes the mean  
237 of the distribution of all  $p$ -values, so that one can assume good description of the data by the fits, see  
238 Figure 4.11. This replaces the investigation of the  $\chi^2/\text{NDF}$  distribution in the case of a frequentist fit.

## 4.2 Determination of $\Sigma_\eta$ using Bayesian statistics



**Figure 4.10:** Posterior predictive checks  $p(A_{\text{rep}}|A)$  from a BAYESIAN fit to the event yield asymmetries for all angular bins of the energy bin  $1250 \text{ MeV} \leq E_\gamma < 1310 \text{ MeV}$ . The data points in the upper plot are the asymmetry  $A(\phi)$ , which was additionally fitted using a  $\chi^2$  fit (solid line). The goodness of fit is shown using  $p$ -values, which give the fraction  $T(A_{\text{rep}} > A)$  of replicated samples greater than the original measured value, with propagated statistical error bars on the bottom of each plot. The expected mean value of  $T(A_{\text{rep}} > A) = 0.5$  is indicated by the dashed line.



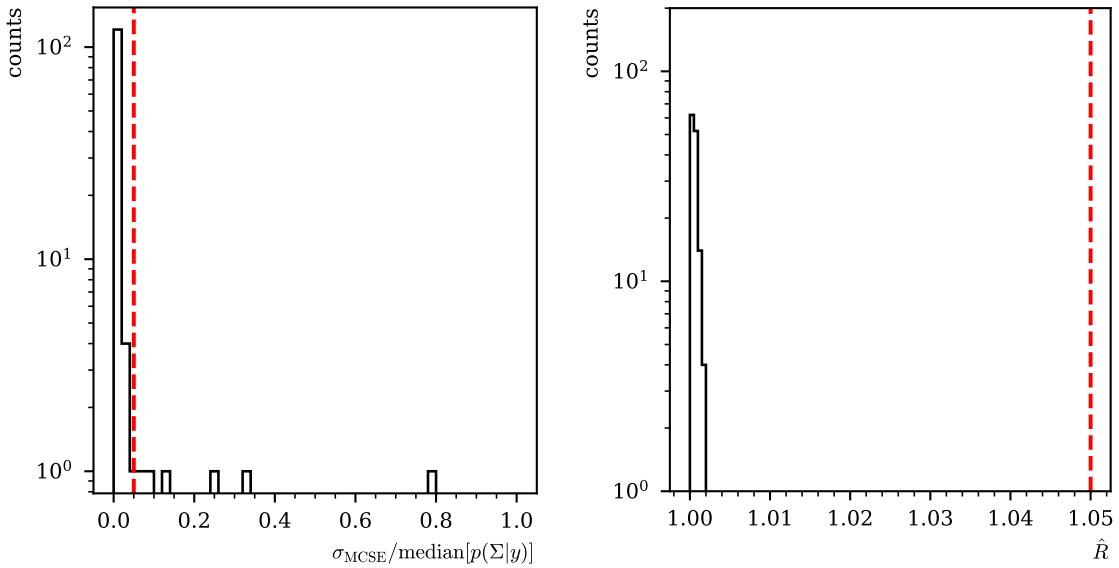
**Figure 4.11:**  $p$  values generated using all fits. They are centered around their mean at 0.5, which is indicated by the dashed line, and show no bias towards higher or lower values, thus confirming an adequate fit.

239 The MCMC diagnostics reveal that all MARKOV chains have converged and explored the available  
 240 parameter space completely which is characterized by  $\widehat{R} = 1.00$  for all fits. The increment of posterior  
 241 draws removed any fluctuations in the  $\widehat{R}$  value that was present before. The relative MCSE is below  
 242 the (self-)defined threshold of  $\frac{\sigma_{\text{MCSE}}}{\text{median}[P(\Sigma|y)]} < 5\%$  for 95% of all fits. The remaining 5% are fits that  
 243 estimated the beam asymmetry close to 0 such that further incrementing the number of posterior  
 244 draws will not amend this. Both, relative MCSE and  $\widehat{R}$  associated with the fit parameter  $\Sigma$  of the  
 245 binned fit, are shown in Figure 4.12. Thus, one can conclude the successful tuning of hyperparameters  
 246 concerning the BAYESIAN fit.

#### 247 Event based fit

248 The discussed BAYESIAN event based fit is also performed for all kinematic bins. The MCMC  
 249 diagnostics reveal good choice of hyperparameters, as Figure 4.13 shows. As already observed during  
 250 the toy Monte Carlo experiments, the unbinned fit produces more precise results; 98% of all fits exhibit  
 251 an MCSE below threshold, those bins with error above threshold are again explained by values of the  
 252 beam asymmetry  $\Sigma \approx 0$ . Good within and between chain convergence is reached, as  $\widehat{R} = 1.00$  for all  
 253 fits shows.

254 The only illustrative measure regarding the performance of the unbinned fit is the posterior predictive  
 255 distribution of the detector coefficients  $a, b$ . Figure 4.14 shows the weighted sum of event yields with  
 256 the same binning as used in the binned fit together with a posterior predictive check using the draws of all  
 257 parameters distributions  $a, b$  for the kinematic bin  $1250 \text{ MeV} \leq E_\gamma < 1310 \text{ MeV}, 0 \leq \cos \theta < 0.17$ .  
 258 As opposed to the binned fit it is not only performed at each data point but for significantly more  
 259 points. This mimics a continuous function as is appropriate for an unbinned fit. However, for better



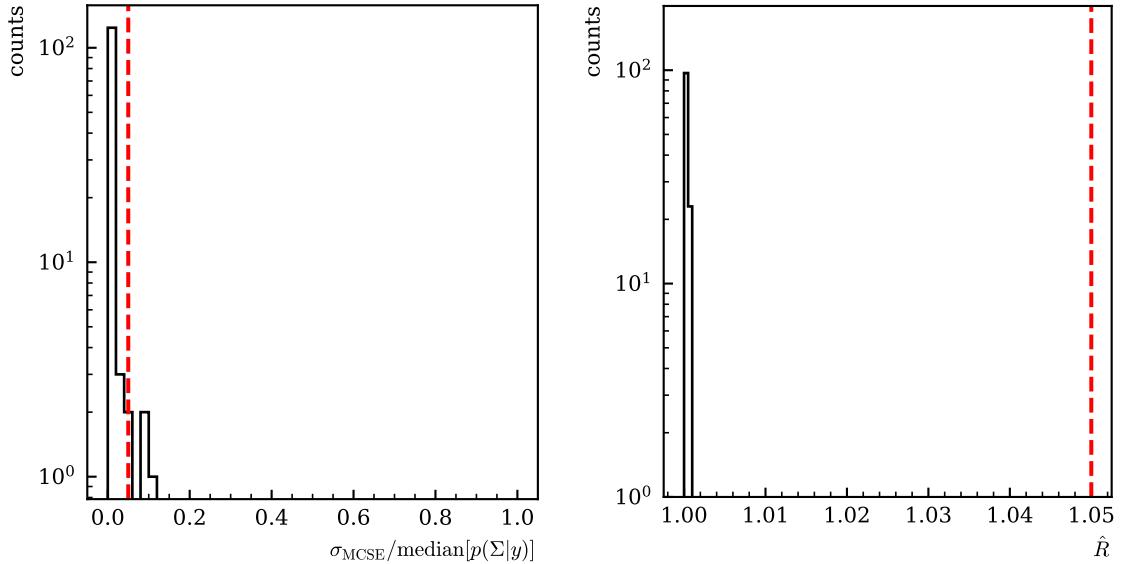
**Figure 4.12:** Left: relative error  $\frac{\sigma_{\text{MCSE}}}{\text{median}[p(\Sigma|y)]}$  Right:  $\hat{R}$  associated with the fit parameter  $\Sigma$ . Both are shown for all  $11 \cdot 12$  binned fits to the asymmetry  $A(\phi)$ . The critical values that should not be exceeded are marked by dashed lines.

visibility the posterior predictive distributions are not built at every point but the efficiency function is plotted for each set of draws  $a^{(s)}, b^{(s)}, s = 1 \dots 20000$ . No  $p$  values are estimated since they would not be representative of the whole fit. Note that the detector efficiency  $\epsilon(\phi)$  is almost flat as opposed to the simulated efficiency in subsection 4.2.1. This is an overall observation and not specific to the shown kinematic bin and further confirmed by the fact that the sums of all fit coefficients  $a, b$  are distributed around zero. Regardless of the significance of detection inefficiencies the unbinned fit is nevertheless able to obtain good agreement between fit and data; the mean of all posterior predictive distributions (dashed line) agrees within statistical error bars with the data points.

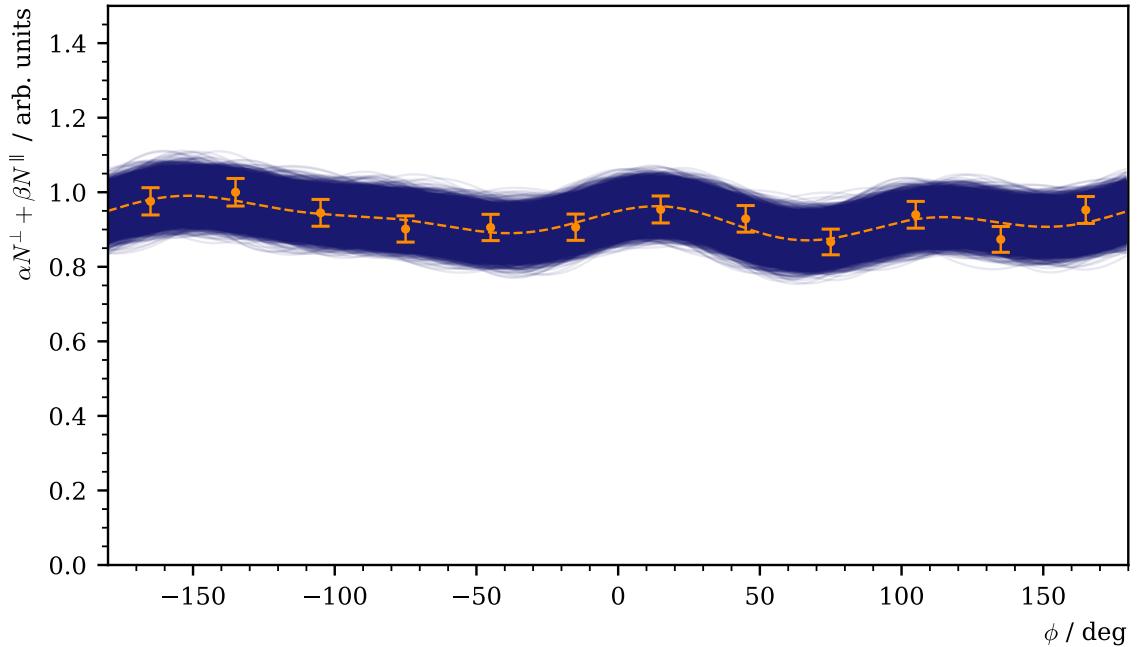
### 4.2.3 Discussion

The final results for the beam asymmetry obtained with both methods are shown in Figure 4.15 for all bins in beam energy and meson polar angle. Along with the marginal posterior distributions obtained from both introduced methods point estimates from a  $\chi^2$  fit and an unbinned fit [Afz19] are shown. An important intermediate result is that all results agree with each other within statistical error bars or within the width of the marginal posterior distributions. Any effects that lead to slightly distinguishable point estimates from a binned or unbinned fit are propagated directly throughout the BAYESIAN fit; then each distribution is centered around the point estimate of the respective fitting method. The results from the binned and unbinned fit with error bars on average cover 68% of the respective marginal posterior distributions, corresponding to a  $1\sigma$  interval. Thus, in general, the application of BAYESIAN methods has proven to give the same results and no systematic error is introduced by the BAYESIAN approach.

The effort of implementing the binned or unbinned fit in a BAYESIAN approach is comparable to



**Figure 4.13:** Left: relative error  $\frac{\sigma_{\text{MCSE}}}{\text{median}[p(\Sigma|y)]}$ . Right:  $\hat{R}$  associated with the fit parameter  $\Sigma$ . Both are shown for all  $11 \cdot 12$  unbinned fits. The critical values that should not be exceeded are marked by dashed lines.

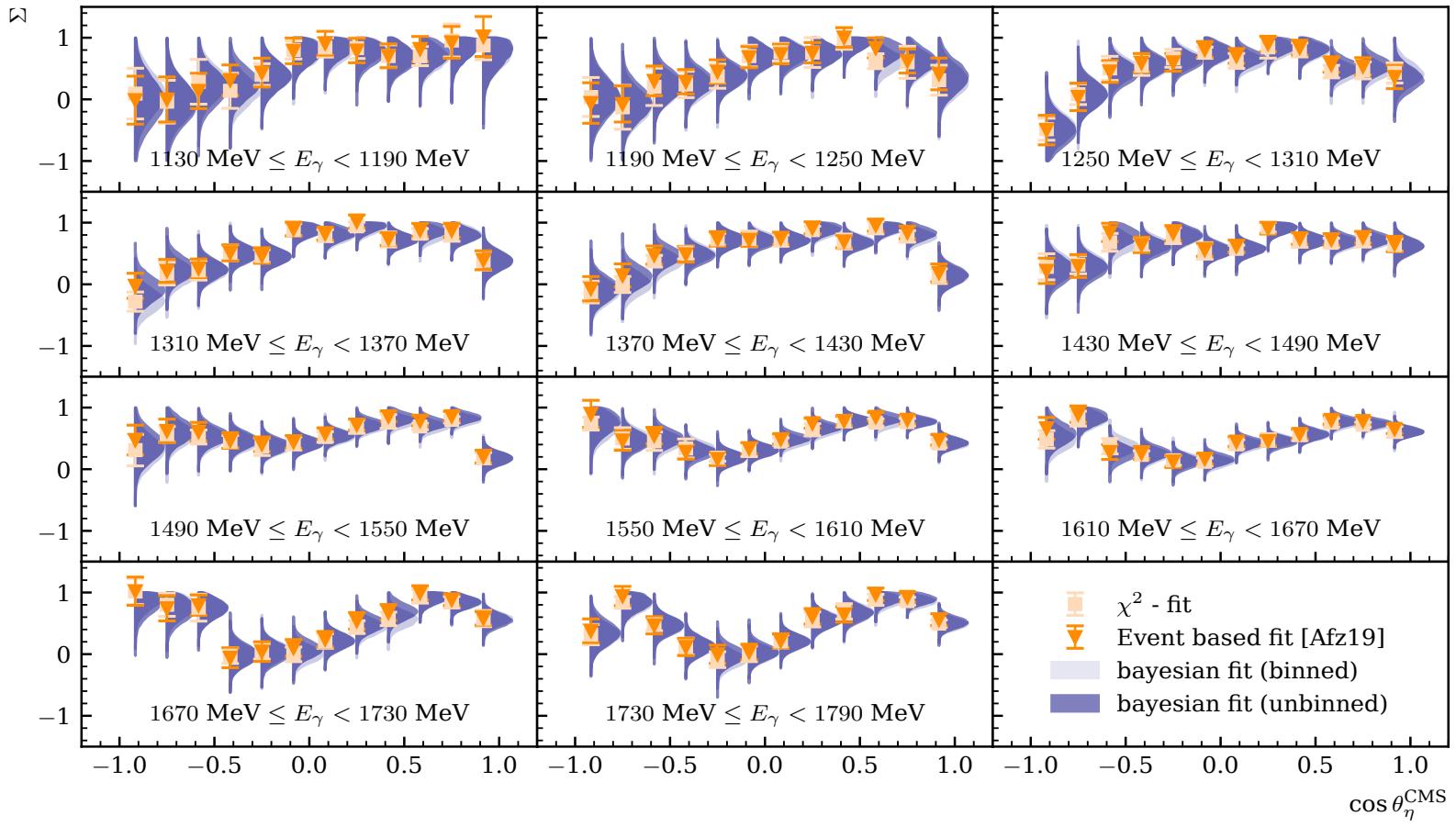


**Figure 4.14:** Posterior predictive check using the draws of the detector coefficients  $a$  and  $b$  for the kinematic bin  $1250 \text{ MeV} \leq E_\gamma < 1310 \text{ MeV}, 0 \leq \cos \theta < 0.17$ . Points with error bars are the polarization weighted sum of event yields. The dashed line is the mean of the predictive values while the solid opaque lines are representative of one simulation draw  $a^{(s)}, b^{(s)}$ .

281 using standard libraries provided by e.g. *python* [RP22] or *ROOT* [BR97]. Due to the probabilistic  
282 structure of the Stan language [Sta22] implementation of likelihood and prior models is straightforward  
283 and the sampling algorithms can be accessed or modified to ones needs intuitively. However, the  
284 BAYESIAN fit requires more careful preparation and also diagnostics. On one hand, the choice of  
285 priors has to be made. On the other hand, the fitting procedure is inherently different; not only the  
286 goodness of fit compared with data has to be checked but also the convergence of the MARKOV chains  
287 themselves.

288 The computational cost also increases greatly in comparison to the traditional frequentist approaches.  
289 Especially the unbinned BAYESIAN fit requires a lot of computation time, where the number of data  
290 points is of order  $O(10^4)$  as opposed to 12 data points for the binned fit. This can be understood  
291 considering the underlying algorithms; the least squares fit and also the unbinned maximum likelihood  
292 fit both aim to minimize a given test statistic which is achievable by (numeric) differentiation.  
293 Determining marginal posterior distributions from BAYESIAN inference on the other hand requires  
294 (numeric) integration, making it the more complex task which is observable in the consumed  
295 computation time.

296 Yet, the additional effort is rewarded by the fact that the BAYESIAN fit will yield *distributions* for all  
297 parameters as opposed to point estimates with error bars. This is especially useful for polarization  
298 observables which may be used as input for PWA calculations. Here error estimates can be derived  
299 from the distributions e.g. as (multiple) standard deviations, the full width at half maximum or similar.  
300 If furthermore a BAYESIAN approach is pursued, also the whole distributions may be used. Another  
301 advantage of the BAYESIAN fit is the ability to truncate the posterior samples to the relevant or allowed  
302 parameter space. In the very first energy bin in Figure 4.15 this is visible clearly where several point  
303 estimates exceed the physical limit of  $\Sigma = 1$  as opposed to the obtained marginal posteriors. Lastly,  
304 the BAYESIAN fit allows much more flexible changes to the fitted model, as will be discussed in the  
305 next section 4.3.



**Figure 4.15:** Final results for the beam asymmetry  $\Sigma$  in  $\eta$  photoproduction off the proton for all kinematic bins obtained with BAYESIAN methods. They are compared with the results of a least squares fit and an unbinned fit as given in reference [Afz19]. All results agree within statistical error bars or within the widths of marginal posterior distributions.

### 306 4.3 Determination of $\Sigma_{\eta'}$

307 After successfully testing the application of BAYESIAN methods this section will demonstrate the  
 308 extraction of new results for the beam asymmetry in  $\eta'$  photoproduction off the proton obtained at the  
 309 CBELSA/TAPS experiment. Only the event based fitting method is applied and the binned fitting of  
 310 event yield asymmetries is discarded because on one hand the available statistics will on average give  
 311 only few counts per bin  $(E_\gamma, \cos \theta, \phi)$  of order  $O(10^1)$ , depending on the specific number of  $\phi$ -bins.  
 312 This will result in very large statistical errors per bin. On the other hand this could be compensated by  
 313 choosing fewer bins in  $\phi$  which would however amplify the systematic underestimation of the beam  
 314 asymmetry which is discussed in appendix B.

315 The event based fits are performed as unbinned maximum likelihood fit and as BAYESIAN fit. New  
 316 toy Monte Carlo experiments are thrown adapting to the significantly decreased statistics. Also, tests  
 317 regarding the background contamination of Monte Carlo samples are made.

#### 318 4.3.1 Application of event based fit to toy Monte Carlo data

After selecting all suitable event candidates for  $\gamma p \rightarrow p\eta' \rightarrow p\gamma\gamma$  reactions there still remain significant background contributions which make up roughly 20% of all events averaged over all kinematic bins. Investigation of Monte Carlo simulations of different mesonic final states revealed they are mainly realized by the false reconstruction of  $2\pi^0$  photoproduction events (see chapter 3). It was thus investigated how the presence of two beam asymmetries  $\Sigma_1$  and  $\Sigma_2$  affects the fit. The combined differential cross section is then given by

$$\frac{d\sigma}{d\Omega_{\text{pol}}}(\phi) = \frac{d\sigma^{(1)}}{d\Omega_{\text{pol}}}(\phi) + \frac{d\sigma^{(2)}}{d\Omega_{\text{pol}}}(\phi) \quad (4.38)$$

$$= \frac{d\sigma^{(1)}}{d\Omega_0} \cdot (1 - p_\gamma \Sigma_1 \cos(2\phi)) + \frac{d\sigma^{(2)}}{d\Omega_0} \cdot (1 - p_\gamma \Sigma_2 \cos(2\phi)). \quad (4.39)$$

If now the fraction  $\delta$  of events in the final state with beam asymmetry  $\Sigma_2$  is known one can express the unpolarized cross sections via a common combined unpolarized cross section as

$$\frac{d\sigma^{(1)}}{d\Omega_0} = (1 - \delta) \cdot \frac{d\sigma}{d\Omega_0} \quad \frac{d\sigma^{(2)}}{d\Omega_0} = \delta \cdot \frac{d\sigma}{d\Omega_0}. \quad (4.40)$$

Plugging Equation (4.40) into Equation (4.39) then yields

$$\frac{d\sigma}{d\Omega_{\text{pol}}}(\phi) = \frac{d\sigma}{d\Omega_0} \cdot (1 - p_\gamma [(1 - \delta) \cdot \Sigma_1 + \delta \cdot \Sigma_2] \cos(2\phi)). \quad (4.41)$$

In the presence of background  $\Sigma^{\text{bkg}}$ <sup>8</sup> all measured values for the beam asymmetry  $\Sigma^{\text{meas}}$  are systematically shifted from their true value  $\Sigma^{\text{true}}$ , depending on the amount of background, if the methods described in section 4.1 are used without modification, i.e.

$$\Sigma^{\text{meas}} = (1 - \delta) \cdot \Sigma^{\text{true}} + \delta \cdot \Sigma^{\text{bkg}}. \quad (4.42)$$

---

<sup>8</sup> For clarity, the beam asymmetry associated with random time background will in this section be denoted as  $\Sigma_t^{\text{bkg}}$  and the beam asymmetry associated with background reactions as  $\Sigma^{\text{bkg}}$

319 This is an important intermediate result: If the background asymmetry  $\Sigma^{\text{bkg}}$  and the corresponding  
 320 fraction of events  $\delta$  that exhibit this asymmetry are known, the true values  $\Sigma^{\text{true}}$  can be obtained. If  
 321  $\Sigma^{\text{bkg}}$  is not known this is a source of systematic error.

322 Table 4.2 shows all characteristics of the investigated toy Monte Carlo experiments, the values for  
 323 polarization and statistics per bin were chosen to match the respective average values in the selected  
 324 data for the  $\gamma p \rightarrow p\eta' \rightarrow p\gamma\gamma$  final state. Furthermore  $\delta = 20\%$  of all events were simulated as  
 325 background with beam asymmetry  $\Sigma_2$  while the remainder of events is simulated with beam asymmetry  
 326  $\Sigma_1$ . The amount of random time background in the prompt peak is set to  $(1 - f) = 0.05$  while  
 327 the sideband events are weighted by  $w = \frac{13}{200}$  (see section 3.2). The same efficiency function and  
 328 hyperparameters as previously are employed with the exception of the number of posterior draws  
 329 which is raised to  $n_{\text{samples}} = 5000$  from the beginning. Because of the smaller amount of data this  
 does not increase the computational cost irrationally.

<b>chosen parameters</b>	$p_\gamma^\parallel = 0.3, p_\gamma^\perp = 0.3, \Sigma_1 = 0.5, \Sigma_2 = -0.3, \Sigma_t^{\text{bkg}} = -0.5, f = 0.95, \delta = 0.2$ $w = \frac{13}{200}$
<b>simulation draws</b>	$N_{\text{total}}^\parallel \sim \mathcal{P}(200), N_{\text{total}}^\perp \sim \mathcal{P}(200)$
	signal in prompt      background in prompt      sideband
	$N_{\text{total}}^{\parallel/\perp} \cdot f$ $N_{\text{total}}^{\parallel/\perp} \cdot (1 - f)$ $N_{\text{total}}^{\parallel/\perp} \cdot (1 - f) \cdot 1/w$
<b>efficiency function</b>	$\epsilon(\phi) = 1/10.5 \cdot (9.3 + 0.28 \cdot \cos \phi + 0.24 \cdot \sin 3\phi)$
<b>hyperparameters</b>	$n_{\text{samples}} = 5000, n_{\text{chain}} = 4, n_{\text{warm up}} = 1000$

**Table 4.2:** Summary of the complete setting of all toy Monte Carlo experiments for the event based fit. Table layout adapted from [Afz19].

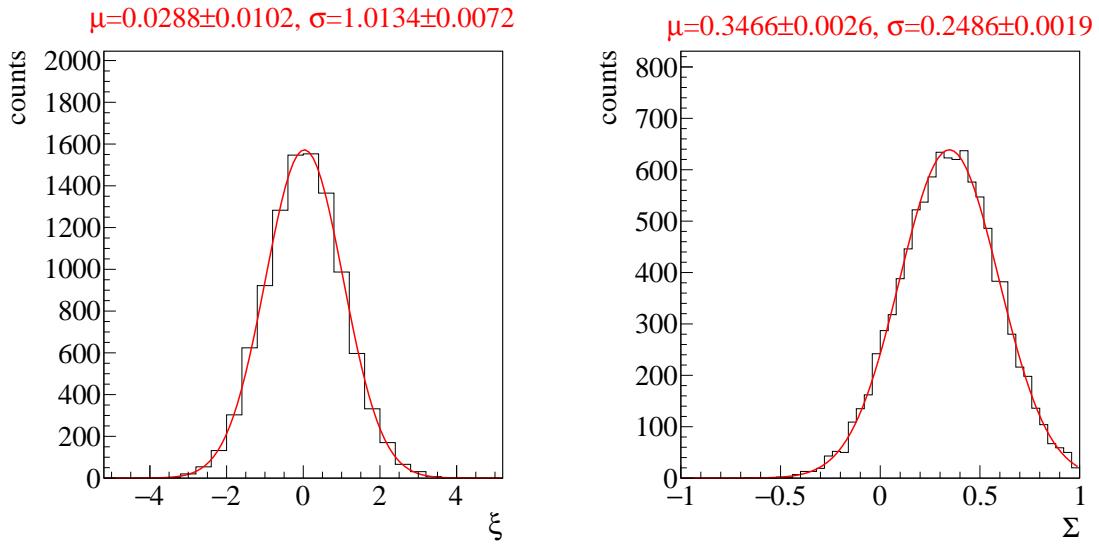
330

### 331 Unbinned maximum likelihood fit

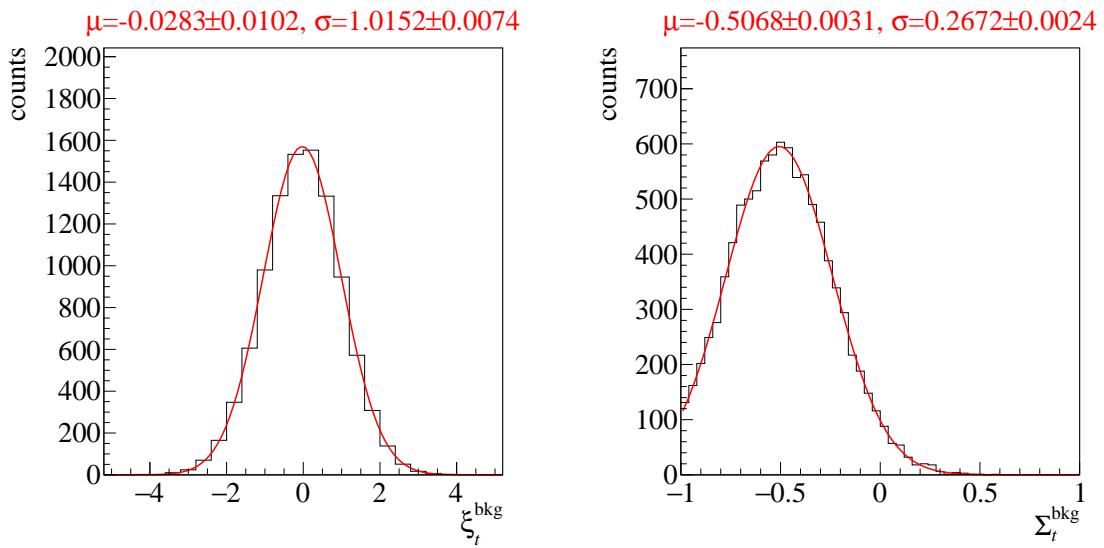
To test the unbinned maximum likelihood fit and its response to the presence of background, 10000 toy Monte Carlo bins are created according to table 4.2. The expected beam asymmetry is

$$\Sigma^{\text{meas}} = 0.8 \cdot 0.5 - 0.2 \cdot 0.3 = 0.34. \quad (4.43)$$

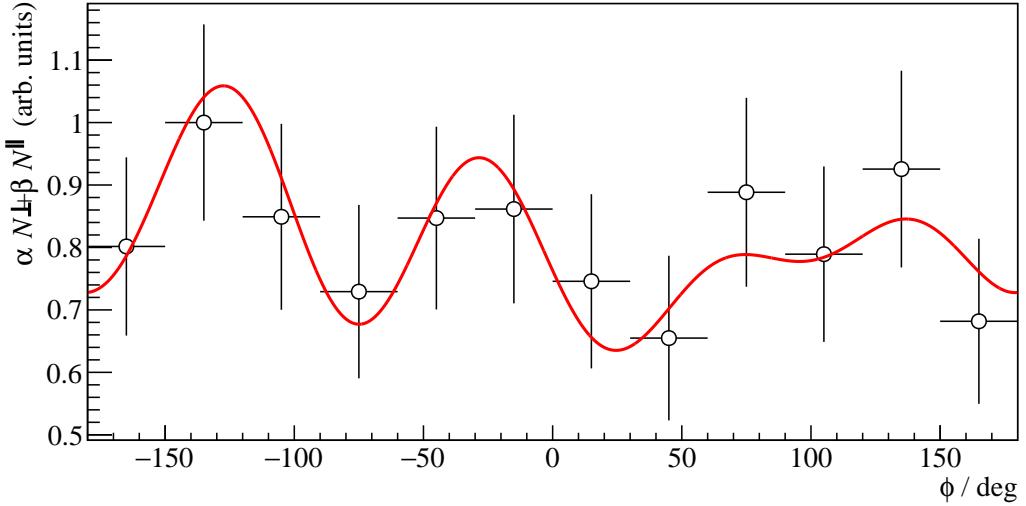
332 Building the normalized residuals with respect to  $\Sigma^{\text{meas}}$  and the random time background asymmetry  
 333  $\Sigma_t^{\text{bkg}}$  yields standard normal distributions, see Figures 4.16 and 4.17, thus fulfilling the expectations  
 334 entirely. This allows Equation (4.42) to be applied on results obtained with this method provided the  
 335 background beam asymmetry is known. Any errors on  $\Sigma^{\text{bkg}}$  may be propagated in a GAUSSIAN manner.  
 336 The beam asymmetry is estimated without bias and the errors are correctly estimated as confirmed by  
 337  $\mu = 0$  and  $\sigma = 1$  for the normalized residuals within few widths of the statistical errors. Note that the  
 338 unnormalized distributions are significantly broader than the ones examined in the previous section  
 339 which is due to the decreased statistics in each bin which already indicates larger statistical errors.  
 340 The decreased statistics also cause greater variation in the fitted efficiency function that is seen in  
 341 Figure 4.18, which still describes the data very well.



**Figure 4.16:** Normalized residuals (left) and unaltered distribution (right) of all 10000 fits for the beam asymmetry  $\Sigma = (1 - \delta) \cdot \Sigma_1 + \delta \cdot \Sigma_2$ . GAUSSIAN fits are performed with results given on top of each plot.



**Figure 4.17:** Normalized residuals (left) and unaltered distribution (right) of all 10000 fits for the background beam asymmetry  $\Sigma_t^{\text{bkg}}$ . GAUSSIAN fits are performed with results given on top of each plot.



**Figure 4.18:** Fitted efficiency function (red line) applied to the polarization weighted sum of event yields (data points) for one toy Monte Carlo bin. 12 bins in  $\phi$  are built for demonstration.

### 342 Unbinned BAYESIAN fit

When performing a BAYESIAN fit a systematic shift like Equation (4.42) as result of background contributions is not easily applicable. The posterior distributions may be shifted by a constant value but including statistical errors of the background beam asymmetry is not possible. To circumvent this, the background beam asymmetry can be included inherently in the likelihood function that is part of the model being fitted. In practice,  $\delta$  and  $\Sigma_2$  have to be part of the data that is fed into the probabilistic model. In general though  $\Sigma_2$  will have a statistical error, which can be included in the fashion of modeling *missing data* [Sta22]. Doing so, the value for  $\Sigma_2$  is regarded as an estimate for the true, unknown (or missing) value  $\Sigma_2^{\text{true}}$ . The statistical error  $\tau$  of  $\Sigma_2$  specifies the measurement model assuming GAUSSIAN errors

$$\Sigma_2^{\text{true}} \sim \mathcal{N}(\Sigma_2, \tau). \quad (4.44)$$

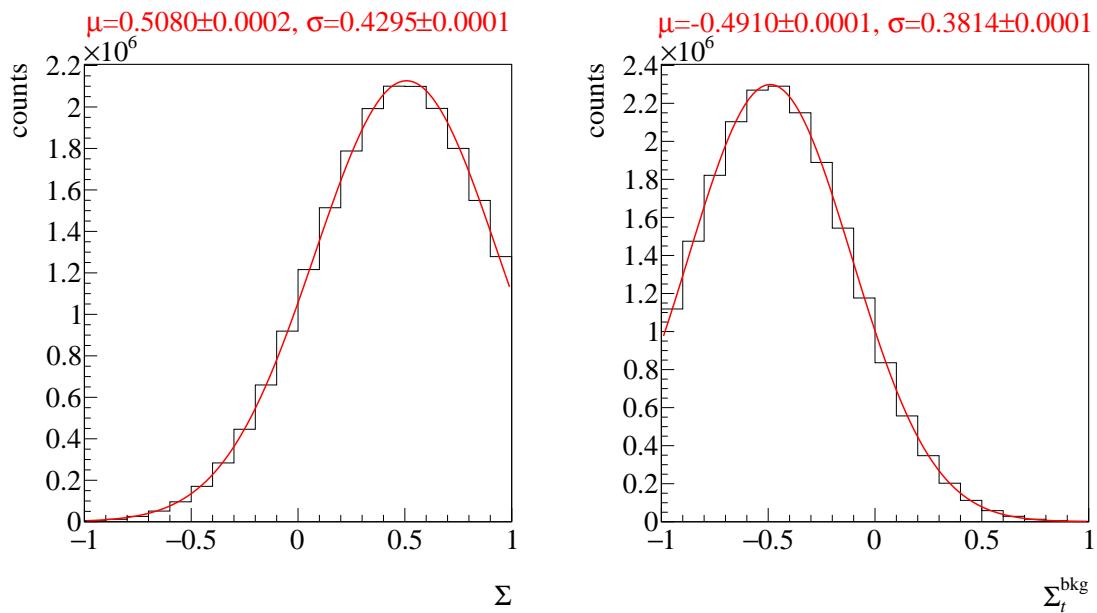
With this  $\Sigma_2^{\text{true}}$  is added as an additional fit parameter to the probabilistic model, so that the likelihood is modified as

$$\begin{aligned} \ln \mathcal{L} = & \sum_{i=1}^n \ln p_{\text{prompt}} \left( \phi_i, p_{\gamma,i} \mid (1 - \delta) \cdot \Sigma_1 + \delta \cdot \Sigma_2^{\text{true}}, a, b, \Sigma_t^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}} \right) \\ & + \sum_{j=1}^m \ln p_{\text{sideband}} \left( \phi_j, p_{\gamma,j} \mid \Sigma_t^{\text{bkg}}, a^{\text{bkg}}, b^{\text{bkg}} \right). \end{aligned} \quad (4.45)$$

<sup>343</sup> For a complete inference  $\Sigma_2^{\text{bkg}}$  also needs to be assigned a prior, which is chosen to be uniform to keep the number of fit parameters minimal.

In total, 1000 bins in accordance with Table 4.2 are fitted with the adapted model that includes a background asymmetry  $\Sigma_2$  with statistical error  $\tau$  which are used to estimate the additional fit parameter  $\Sigma_2^{\text{true}}$ . For the sake of the toy Monte Carlo experiments the measurement error is arbitrarily chosen to be  $\tau = 0.05$ .

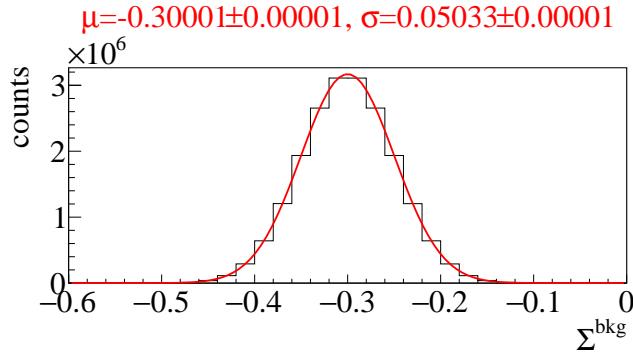
Figure 4.19 shows the combined posteriors for the fit parameters  $\Sigma_1$  and  $\Sigma_t^{\text{bkg}}$  of the event based BAYESIAN fit for all generated toy Monte Carlo bins. The posteriors are added up since the truncation of the posteriors biases the results when combining them in a pooled likelihood model or normalized residuals  $\Xi$ . Although previous experiments already proved sensible error estimation, i.e. distribution widths, when using the unbinned BAYESIAN fit, this is further confirmed by removing the truncation and repeating the investigation of posteriors. This is discussed in appendix C. The input parameters are very well reproduced within the statistical error, which is given by the standard deviation of the GAUSSIAN. The posterior distribution of  $\Sigma_1$  is wider than the posterior of the background beam asymmetry  $\Sigma_t^{\text{bkg}}$  because an additional measurement of  $\Sigma_2$  with measurement error  $\tau$  is used to estimate  $\Sigma_1$ . This is in complete analogy to shifting the point estimates  $\Sigma^{\text{meas}}$  to their true values  $\Sigma^{\text{true}}$ , as described previously, which will also result in larger statistical error bars according to GAUSSIAN error propagation. Since the signal beam asymmetry  $\Sigma_1$  was recognized correctly by the fit it is evident



**Figure 4.19:** Combined (added) posteriors of all 1000 fits. Left: Signal beam asymmetry  $\Sigma_1$  Right: Background beam asymmetry  $\Sigma_t^{\text{bkg}}$ . A GAUSSIAN fit is performed with results given on top.

that the combined posterior distributions for  $\Sigma_2$  are in full agreement with the expectations, see figure 4.20. Mean  $\mu$  and standard deviation  $\sigma$  are reproduced exactly as they were modeled as a GAUSSIAN fit to the combined posteriors reveals.

It remains to investigate the properties of the employed MARKOV chains. First of all one observes that the accuracy of results, i.e. the relative MCSE, suffers from the available statistics which are significantly smaller in comparison with the previously investigated case for the  $p\eta \rightarrow \gamma\gamma$  final state. Only 97% of all unbinned fits fulfill the self set criterion of  $\sigma_{\text{MCSE}}/\text{median}(p(\Sigma|y)) \leq 0.05$  whereas

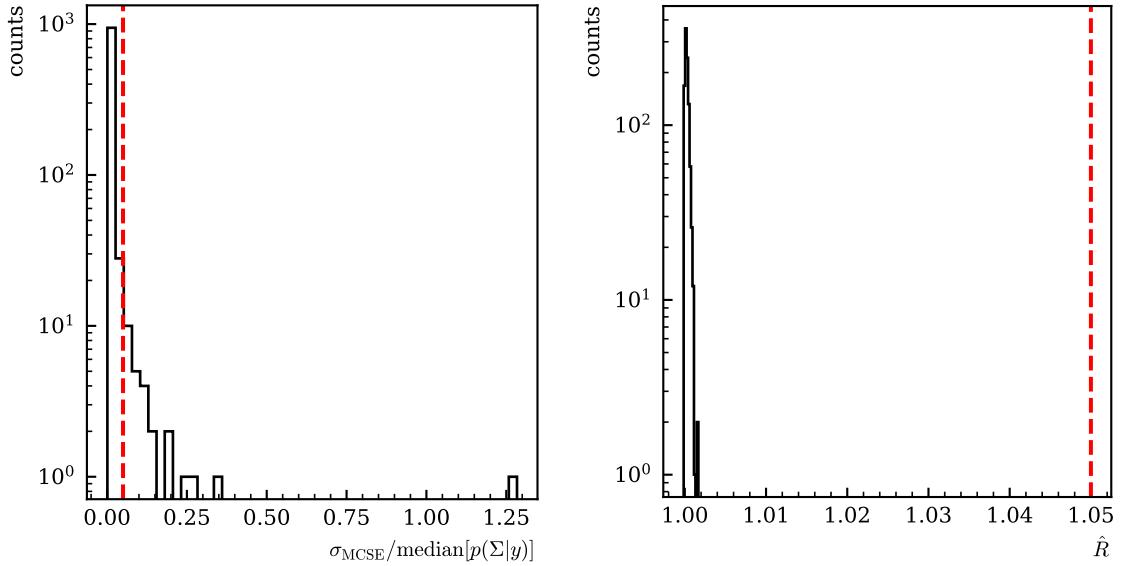


**Figure 4.20:** Combined (added) posteriors of all fits for the fit parameter  $\Sigma_2^{\text{true}}$ . A GAUSSIAN fit is performed which reproduces exactly the values that were used for the simulations.

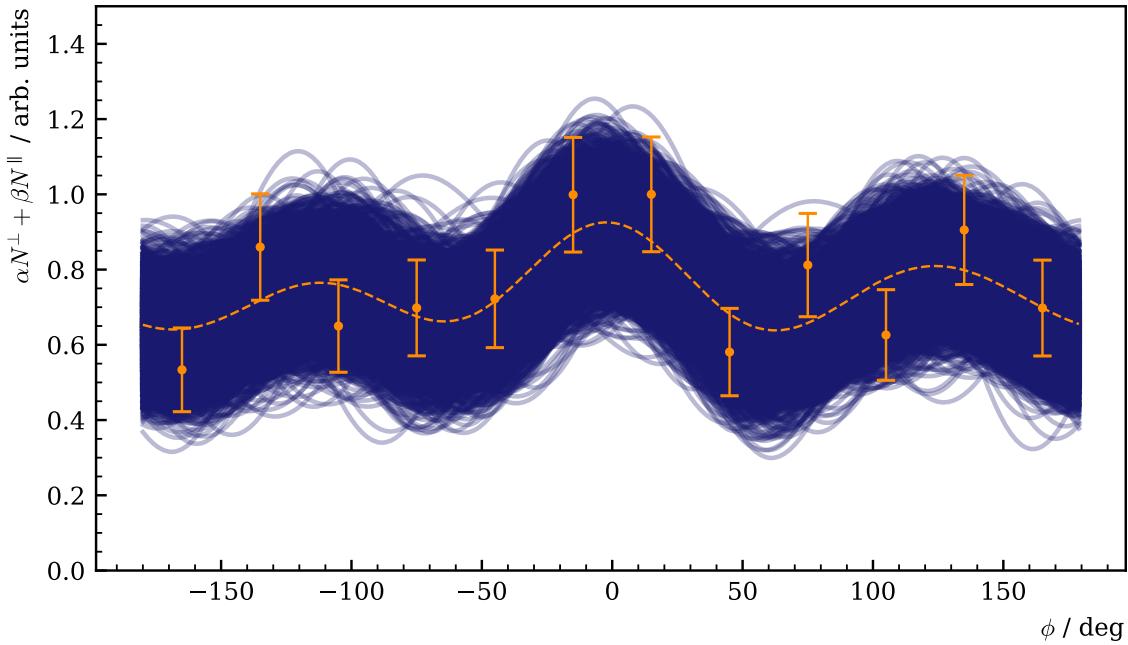
for the toy Monte Carlo experiments in the previous section 4.2 this was the case for *all* unbinned fits. This observation however is not resembling a failing fit but rather the possibility for the data to mimic beam asymmetries close to zero because of the fewer available data points. In these cases the relative MCSE exceeds the threshold as has also been observed before. Convergence and full exploration of the available parameter space on the other hand is guaranteed for all fits because all  $\widehat{R}$  values for all parameters and fits are within the demanded thresholds.

Lastly, consistency of the fit results with the data that is fitted can be checked as before by a posterior predictive check of the detector coefficients  $a, b$  plugged into the efficiency function  $\epsilon(\phi) = \frac{1}{c} \cdot (\alpha \tilde{N}^{\parallel} + \beta \tilde{N}^{\perp})$  as can be seen in Figure 4.22. Each draw  $a^{(s)}, b^{(s)}, s = 1, \dots, 20000$  is plotted as an opaque blue line. The mean over the posterior draws is indicated by the dashed line and follows the polarization weighted sum of event yields which is shown as data points. The mean agrees within statistical error bars at each point indicating finally a successful fit in every respect.

It has been demonstrated that the chosen BAYESIAN model can describe data generated with two beam asymmetries very well. Note that expanding the existing model from before requires minimal changes since *Stan* [Sta22] allows modular implementations; only the new parameter  $\Sigma_2^{\text{true}}$  and the new data points  $(\Sigma_2, \tau, \delta)$  have to be added to give correctly centered posteriors that are not anymore systematically shifted. More effort would have to be made to get similar results with the traditional unbinned fitting method described before.



**Figure 4.21:** MCMC diagnostics for the event based BAYESIAN fit. Left: MCSE, Right:  $\hat{R}$ -value. The critical values not to be exceeded are marked by the dashed lines.



**Figure 4.22:** Posterior predictive checks of one toy Monte Carlo bin using the draws from the marginal posteriors of the detector coefficients  $a, b$  (opaque blue lines). The mean values are marked by the dashed line and follow the distribution of the data points which are the polarization weighted sum of event yields, using 12  $\phi$  bins.

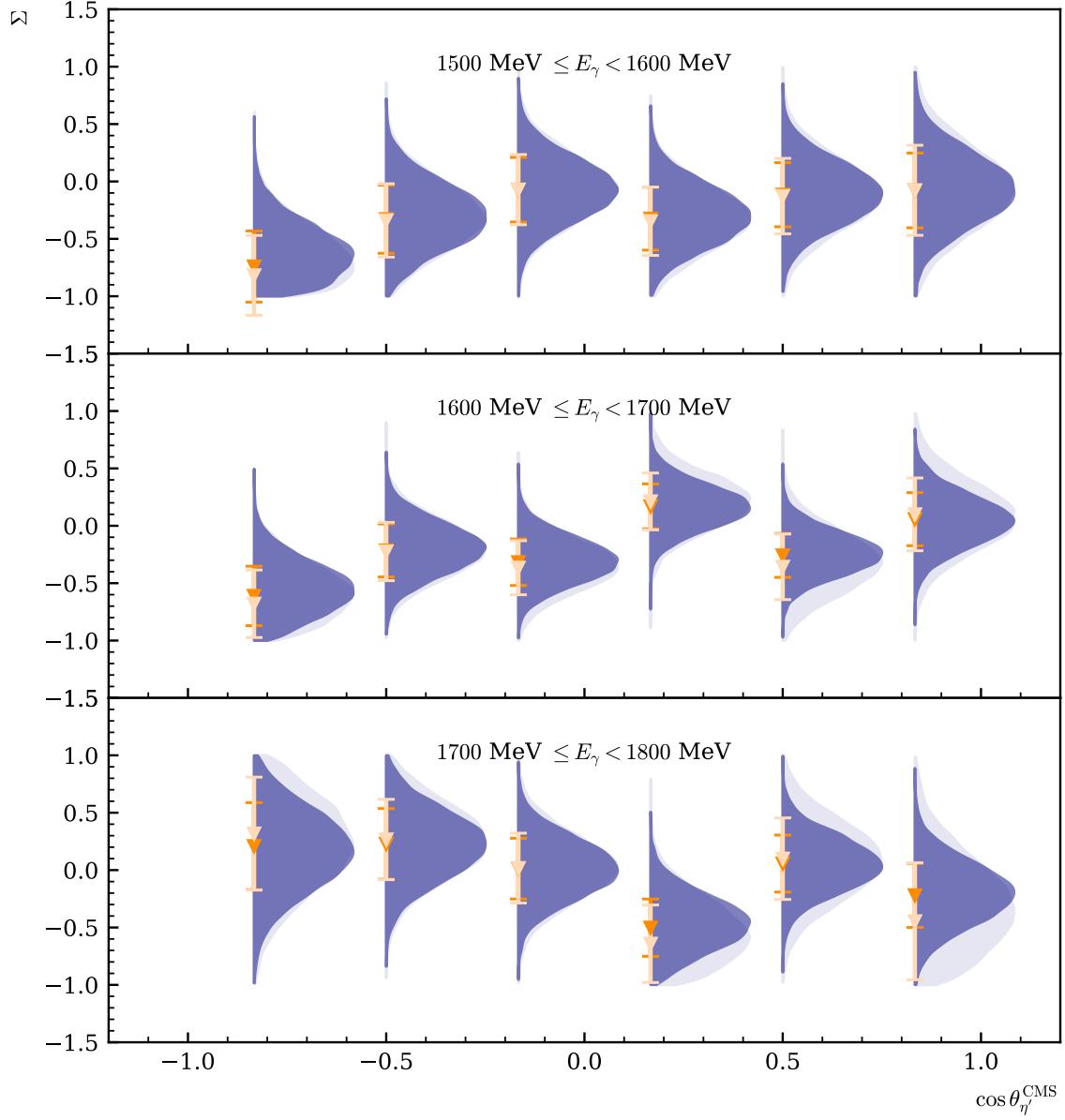
### 386 4.3.2 Application of event based fit to data

387 Investigation of toy Monte Carlo experiments showed that the presence of polarized background  
 388 events cause a systematic shift in the extracted beam asymmetry. The shift may be compensated by a  
 389 modified likelihood function for the BAYESIAN fit or by reversing the shift *after* obtaining intermediate  
 390 results for  $\Sigma^{\text{meas}}$  with an unbinned maximum likelihood fit. By studying Monte Carlo simulations of  
 391 various mesonic final states it has been found that the main background contributions in the analysis of  
 392 the final state  $\gamma p \rightarrow p\eta' \rightarrow p\gamma\gamma$  is given by the false reconstruction of events  $\gamma p \rightarrow p2\pi^0 \rightarrow p4\gamma$ ,  
 393 see chapter 3. Conveniently, this includes the determination of the fraction of background events  $\delta$ , see  
 394 Figure 3.10. Furthermore, preliminary results for the beam asymmetry in double pion photoproduction  
 395 at the CBELSA/TAPS experiment using the same binning as this analysis [Mah22] were available so  
 396 that the previously discussed corrections can in fact be applied in each kinematic bin. It is hereby  
 397 assumed that the *entire* background is realized by  $2\pi^0$  production and the fraction  $\delta$  is determined by  
 398 an according fit of Monte Carlo spectra. This holds only approximately for each kinematic bin and is  
 399 thus a source of systematic error, which is discussed in the next subsection 4.3.3.

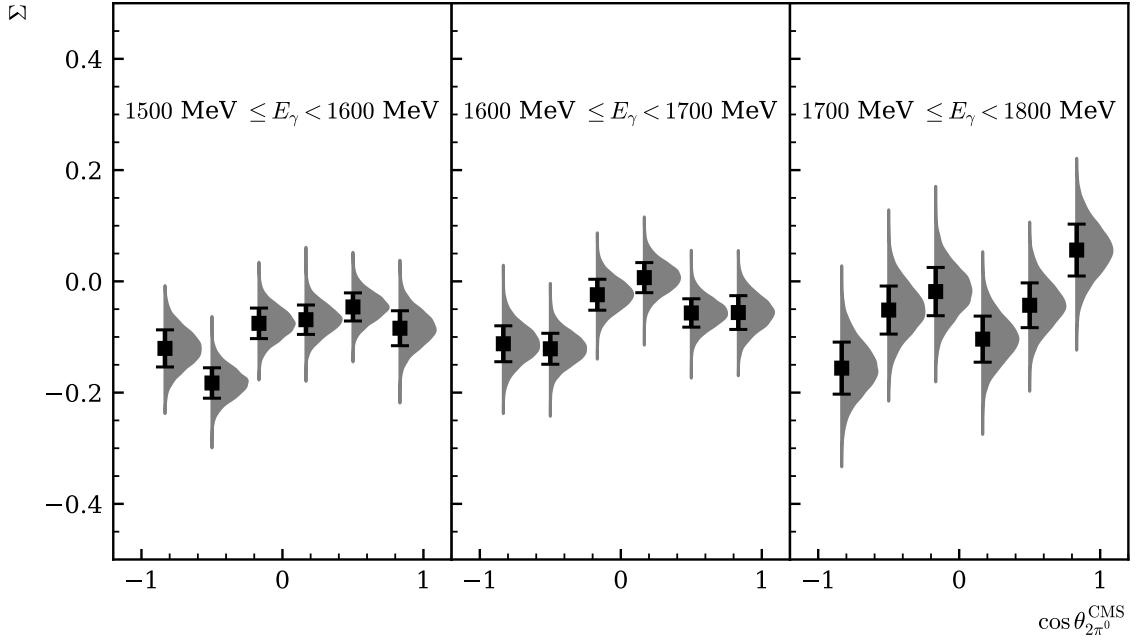
400 Figure 4.23 shows the final results for the beam asymmetry in the reaction  $\gamma p \rightarrow p\eta' \rightarrow p\gamma\gamma$   
 401 obtained from data taken at the CBELSA/TAPS experiment. Shown are two sets of results; the dark  
 402 blue distributions are the results of a unbinned BAYESIAN without any modifications to the original  
 403 model whatsoever. The dark orange data points are the corresponding results of an unbinned maximum  
 404 likelihood fit  $\Sigma^{\text{meas}}$ . The light blue distributions however depict the results that were obtained from  
 405 the modified BAYESIAN model, i.e. with the additional fit parameter  $\Sigma_2^{\text{true}}$ , the light orange data points  
 406 correspond to the corrected point estimates  $\Sigma^{\text{true}}$  according to Equation (4.42). In addition to the  
 407 beam asymmetry  $\Sigma_1 = \Sigma_{\eta'}$  the fit also determines marginal posteriors for  $\Sigma_2^{\text{true}} = \Sigma_{2\pi^0}$ . To check  
 408 consistency of the fit, these posterior distributions are compared to the results  $(\Sigma_2, \tau)$  that were used  
 409 as input [Mah22], they are shown in Figure 4.24.

410 Several remarks can be made regarding these results:

- 411 1. As has been found studying toy Monte Carlo simulations, the modified BAYESIAN fit again  
 412 successfully reconstructs two separate beam asymmetries. Figure 4.24 impressively shows the  
 413 good agreement of input data points  $(\Sigma_2, \tau)$  and posterior of  $\Sigma_2^{\text{true}}$ . On average  $\Sigma_2 \pm \tau$  covers  
 414 68% of the posterior distributions, corresponding to a  $1\sigma$  region, proving that no systematic  
 415 error is made by introducing the additional fit parameter  $\Sigma_2^{\text{true}}$ .
- 416 2. There is very good agreement between point estimates and corresponding posterior distributions.  
 417 The point estimates give mean and standard deviation of the marginal posteriors to good  
 418 approximation. If the point estimates are corrected according to Equation (4.42) their statistical  
 419 error increases. Similarly, the distributions from the modified BAYESIAN fit are wider, agreeing  
 420 with the corrected point estimates and errors.
- 421 3. Errors and distributions widths are rather large, which could be expected due to the small  
 422 amount of events that could be used to extract the results. However, distributions are rarely  
 423 truncated because they are mostly centered around  $\Sigma \approx 0$
- 424 4. Shifting the final results of the beam asymmetry depending on the amount of background  
 425 has marginal impact on the absolute scale. Although background contributions are far from  
 426 negligible they exhibit only asymmetries close to 0 (see Figure 4.24). Thus, the shifted values  
 427 still agree within statistical error bars or widths with the original values.



**Figure 4.23:** Final results for the beam asymmetry  $\Sigma$  in  $\eta'$  photoproduction. Two sets of results are shown: The dark blue distributions and orange data points with errorbars are obtained with an unbinned fit that does not consider any background contributions. The light blue distributions and data points are obtained with the modified BAYESIAN fit and by correcting the point estimates according to Equation (4.42), respectively. All errors are statistical errors only.



**Figure 4.24:** Results for the additionally fitted  $\Sigma_2^{\text{true}}$  (distributions) compared with the underlying data points [Mah22] with statistical errors. The error bars on average cover  $1\sigma$  of the distributions, indicating a successful fit. All errors are statistical errors only.

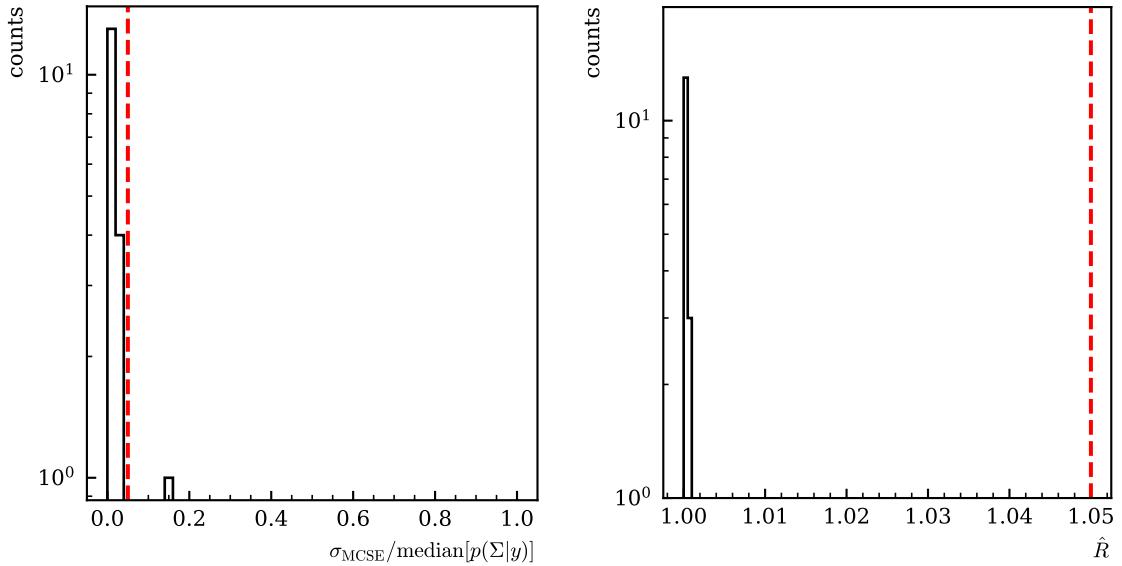
428 To confirm between- and in-chain convergence for all fits, the  $\widehat{R}$ -values of all fit parameters are  
 429 checked to be  $1 \lesssim \widehat{R} < 1.05$ . Due to the large number of samples this criterion is fulfilled without  
 430 complications. Furthermore the relative MCSE with respect to the median of the posterior distributions  
 431 is checked to confirm the accuracy of all chains. It is  $\sigma_{\text{mcse}}/\text{median}(p(\Sigma|y)) \leq 0.05$  for all but one  
 432 fit where the beam asymmetry is very close to zero. Figure 4.25 shows both quantities for all fits.

433 A last consistency check of the posterior predictive distributions of the efficiency function reveals  
 434 no results aside from the expectations. Figure 4.26 shows all posterior predictive draws of the detector  
 435 coefficients  $a^{(s)}, b^{(s)}$  plugged into the efficiency function  $\epsilon(\phi)$ . The mean of all draws describes the  
 436 polarization weighted sum of event yields within statistical errors. The observed efficiency is less  
 437 peaked than the simulated one (Figure 4.22), confirming the results that have been found in section 4.2.  
 438 Yet, more inefficiencies are observed compared to the  $p\eta$  final state due to the decreased statistics.

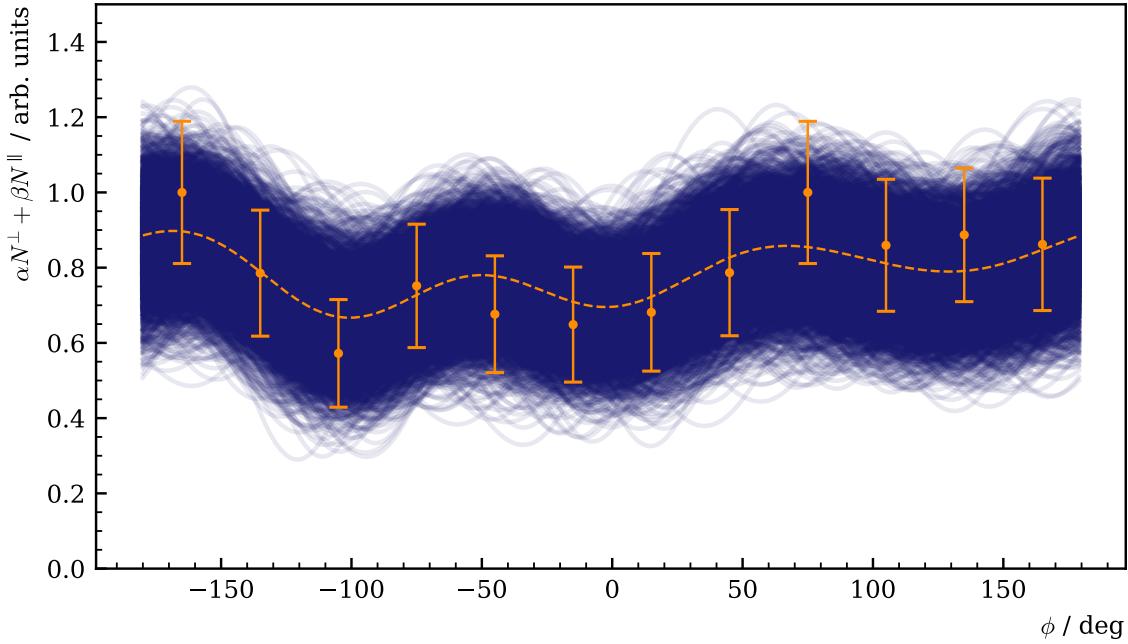
### 439 4.3.3 Systematic error

440 There are two major sources for systematic error in the discussed analysis. On one hand, the  
 441 determination of the polarization degree of the incident photon beam is not exactly accurate and thus  
 442 influences the fit results which heavily depend on the linear polarization degree. On the other hand,  
 443 neglecting polarized background events that have been falsely regarded as  $2\pi^0$  events when shifting  
 444 the intermediate results may create systematic errors.

The relative error of the polarization degree can be estimated during the determination process where analytically calculated bremsstrahlung spectra (ANB) are fitted to measured spectra [Afz19].



**Figure 4.25:** MCMC diagnostics for the event based BAYESIAN fit. Left: MCSE, Right:  $\hat{R}$ -value. The critical values not to be exceeded are marked by the dashed lines.



**Figure 4.26:** Posterior predictive checks of the kinematic bin  $1700 \text{ MeV} \leq E_\gamma < 1800 \text{ MeV}$ ,  $0.67 \leq \cos \theta < 1$  using the draws from the marginal posteriors of the detector coefficients  $a, b$  (opaque blue lines). The mean values are marked by the dashed line and follow the distribution of the data points which are the polarization weighted sum of event yields, using 12  $\phi$  bins.

This allows to give the relative error in the beam energy region that was analyzed as [Afz19]

$$\frac{\Delta p_\gamma}{p_\gamma} = \begin{cases} 0.05, & \text{if } E_\gamma < 1600 \text{ MeV} \\ 0.08, & \text{otherwise.} \end{cases} \quad (4.46)$$

It was assumed that

$$\Sigma^{\text{meas}} = (1 - \delta) \cdot \Sigma^{\text{true}} + \delta \Sigma^{\text{bkg}}, \quad (4.47)$$

with  $\Sigma^{\text{true}} = \Sigma_{\eta'}$  and  $\Sigma^{\text{bkg}} = \Sigma_{2\pi^0}$  and according modifications to the obtained results have been made. However regarding the entire background contributions to double pion photoproduction holds only approximately in each bin. Actually, the measured value for the beam asymmetry is given by

$$\Sigma^{\text{meas}} = (1 - \delta_1 - \delta_2) \cdot \Sigma_{\eta'} + \delta_1 \Sigma_{2\pi^0} + \delta_2 \Sigma^{\text{bkg}}, \quad (4.48)$$

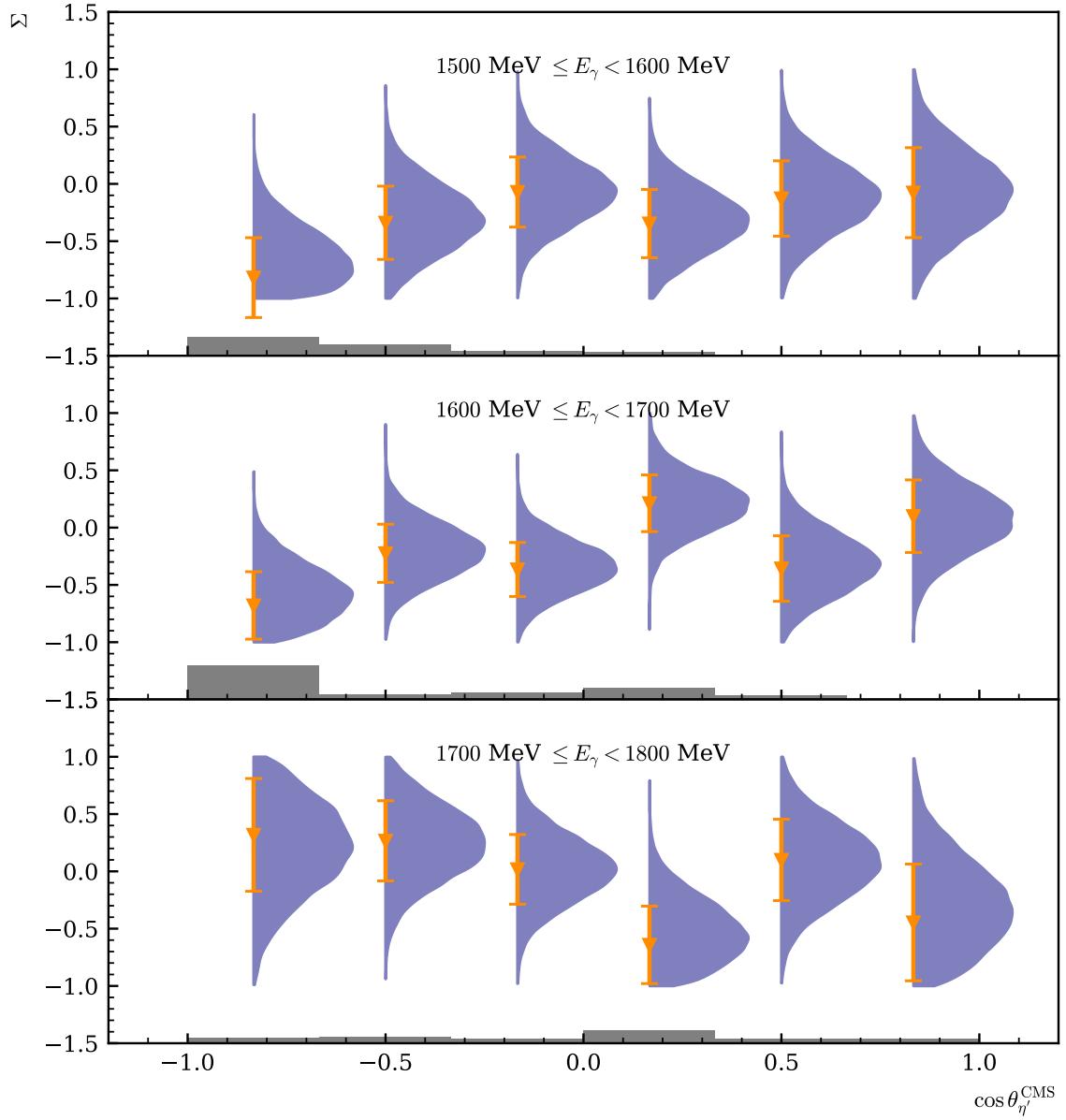
where  $\delta_1$  is the relative contribution of  $2\pi^0$  production events and  $\delta_2$  the relative contribution of other background events, e.g.  $\pi^0\eta$  production, that have been neglected before that may exhibit an asymmetry  $\Sigma^{\text{bkg}}$ . The systematic uncertainty of  $\Sigma_{\eta'}$  can then be determined by setting  $\Sigma^{\text{bkg}} = \pm 1$  and determining the maximum deviation from the previous results. The results from the unbinned maximum likelihood fit are used which to good approximation also give the median and mean of the posterior distributions.

$$\Delta \Sigma_{\eta'} = \max \left[ \left| \frac{\Sigma^{\text{meas}} - \delta_1 \cdot \Sigma_{2\pi^0} \pm \delta_2 \cdot 1}{1 - \delta_1 - \delta_2} - \frac{\Sigma^{\text{meas}} - \delta \cdot \Sigma_{2\pi^0}}{1 - \delta} \right| \right] \quad (4.49)$$

The total absolute systematic error is then given by

$$\Delta \Sigma_{\eta'}^{\text{sys}} = \sqrt{\left( \frac{\Delta p_\gamma}{p_\gamma} \Sigma_{\eta'} \right)^2 + (\Delta \Sigma_{\eta'})^2}. \quad (4.50)$$

<sup>445</sup> Figure ?? shows all results with systematic uncertainties given on the bottom of each plot.



**Figure 4.27:** Final results for the beam asymmetry  $\Sigma_{\eta'}$  for all energy and angular bins. Only the corrected results from the unbinned maximum likelihood fit and distributions from the modified BAYESIAN fit are shown. The bottom of each plot indicates the systematic error as gray bars. It was determined as previously discussed.



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