

Determination of the beam asymmetry Σ in η - and η' -photoproduction off the proton using Bayesian statistics

Master thesis for the CBELSA/TAPS collaboration

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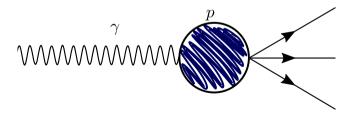
September 8/9 2022

Setting the scene

The Standard Model of Particle Physics

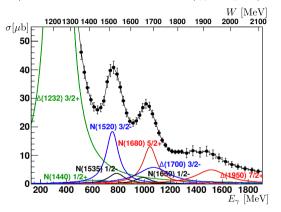
- ▶ matter consists of 12 (anti-)fermions
- ightharpoonup quarks interact via $strong\ interaction$
- ▶ form bound states: mesons $(q\bar{q})$ and baryons (qqq)

baryon spectroscopy (photoproduction) gives insight in strong interaction



Setting the scene

Observe resonances N^*/Δ^* in the cross sections $\sigma(\gamma p \to pM)$



Total cross section $\sigma(\gamma p \to p \pi^0)$ [Wunderlich et al. 2017]

→goal: (help to) identify contributing resonances as strong bound states!

- 1. Theoretical basics
- 2. Experimental Setup
- 3. Results

Determination of Σ_{η} using Bayesian statistics Determination of $\Sigma_{\eta'}$

4. Conclusion

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- ► resonances are broad, overlapping, require complicated partial-wave-analysis (PWA)
- ▶ constraints for the analysis can be derived from polarization observables
- ▶ ultimate goal: "complete experiment"; unambiguous, model-independent PWA solution → several single and double polarization observables needed

Beam-target polarization observables

| | target polarization | | | |
|----------------------|---------------------|---|----|----|
| photon | | x | y | z |
| unpolarized | σ_0 | - | T | - |
| linearly polarized | $-\Sigma$ | H | -P | -G |
| circularly polarized | - | F | - | -E |

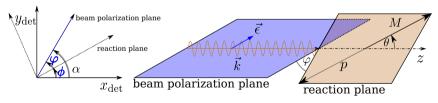
[Sandorfi et al. 2011]

Beam asymmetry Σ

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega}(E_{\gamma},\cos\theta,\varphi) = \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega_0}(E_{\gamma},\cos\theta) \cdot \left[1 - p_{\gamma}^{\mathrm{lin}}\Sigma\cos(2\varphi)\right]$$

polarization angle φ , polarization degree $p_{\gamma}^{\mathrm{lin}}$

[Sandorfi et al. 2011]



Definition of the polarization angle

- ▶ Polarization observables are input for further analysis
- ► Idea: increase amount of information gained from results using Bayesian inference

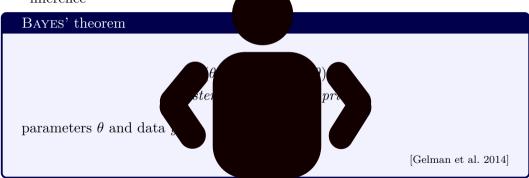
Bayes' theorem

$$p(\theta|y) \propto p(y|\theta) \cdot p(\theta)$$
$$posterior \propto likelihood \cdot prior$$

parameters θ and data y.

[Gelman et al. 2014]

- ▶ Polarization observables are input for further analysis
- ► Idea: increase amount of information gained from results using BAYESIAN inference



$$p(\theta|y) \propto p(y|\theta) \cdot p(\theta)$$

▶ prior $p(\theta)$ and likelihood $p(y|\theta)$ can easily be specified → gain distributions $p(\theta|y)$ instead of point estimates with error bars

Bayesian parameter inference

For each parameter $\theta_n \in \theta$ we gain marginal posteriors

$$p(\theta_n|y) = \int d\theta_1 \cdots \int d\theta_{n-1} \int d\theta_{n-1} \cdots \int d\theta_N p(\theta_1 \dots \theta_N|y).$$

usually approximated using Markov-Chain Monte Carlo (MCMC) draws $\theta^{(s)}$

[Sivia and Skilling 2005]

$$p(\theta|y) \propto p(y|\theta)$$
:

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Bayesian parameter inference

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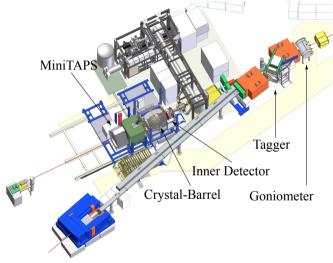
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CBELSA/TAPS experiment

- $\begin{tabular}{l} \hline & & & & \\ \hline & &$
- ▶ photon beam impinges on liquid hydrogen target: $\gamma p \rightarrow pM \rightarrow pX$
- ► measure decay products X of different final states: $M = \pi^0/\eta/\eta'/\ldots$
- ► data set: July-October 2013, 1065 h beam time



Overview of the experimental area, adapted from [Walther 2021]

- 1. Theoretical basics
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Determination of $\Sigma_{\eta'}$

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- ▶ Polarization observables are needed for different final states $(\pi^0, \eta, \eta', ...)$
- ▶ high precision measurement of beam asymmetry for η production recently published [F. Afzal et al. 2020]
- ▶ goal: confirm results using Bayesian fitting methods

Event selection (η)

analysis performed in 11x12 bins of $(E_{\gamma}, \cos \theta)$ for $\gamma p \to p \eta \to p \gamma \gamma$ by [F. Afzal et al. 2020]

Methods

Remember:
$$\frac{d\sigma}{d\Omega}(E_{\gamma},\cos\theta,\varphi) = \frac{d\sigma}{d\Omega_0}(E_{\gamma},\cos\theta) \cdot \left[1 - p_{\gamma}^{\ln}\Sigma\cos(2\varphi)\right]$$

Binned fit to event yield asymmetries

fit to event yield asymmetries $A(E_{\gamma}, \theta, \phi)$

$$= \frac{N^{\perp}(E_{\gamma}, \theta, \phi) - N^{\parallel}(E_{\gamma}, \theta, \phi)}{p_{\gamma}^{\parallel} N^{\perp}(E_{\gamma}, \theta, \phi) + p_{\gamma}^{\perp} N^{\parallel}(E_{\gamma}, \theta, \phi)} = \Sigma(E_{\gamma}, \theta) \cos\left(2\left(\alpha^{\parallel} - \phi\right)\right)$$

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Methods

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Unbinned maximum likelihood fit

Consider likelihood of each individual event

$$\tilde{p}(\phi, \Sigma) = \frac{\left(1 + p_{\gamma} \Sigma \cos\left(2\left(\alpha^{\parallel} - \phi\right)\right)\right) \cdot \epsilon\left(\phi\right)}{C}$$

Applying Bayesian approach to event yield asymmetries:

▶ assume Gaussian errors, i.e.

$$A(\phi) = \Sigma \cos \left(2\left(\alpha^{\parallel} - \phi\right)\right) + \epsilon$$

where $\epsilon \sim \mathcal{N}(0, \sigma)$

▶ likelihood $p(A|\Sigma)$ of each datapoint given by

$$y \sim \mathcal{N}\left(\Sigma \cos\left(2\left(\alpha^{\parallel} - \phi\right)\right), \sigma\right)$$

▶ prior:

$$p(\Sigma) \sim \mathcal{N}(0,1)_{[-1,1]}$$

Sample from posterior $p(\Sigma|A) \propto p(A|\Sigma) \cdot p(\Sigma)$!

Applying Bayesian approach to unbinned fit:

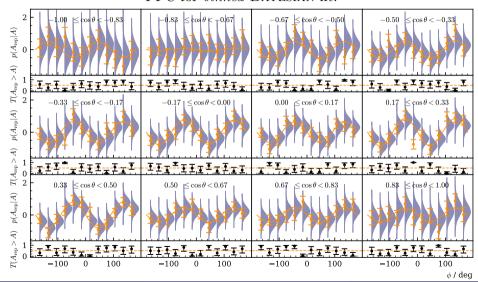
- ▶ event based likelihood given by product of all single-event likelihoods
- ▶ assign priors for all fit parameters (18 in total)
- \blacktriangleright truncate beam asymmetry to allowed region [-1,1]
- ▶ perform toy Monte Carlo experiments

Sample from posterior!

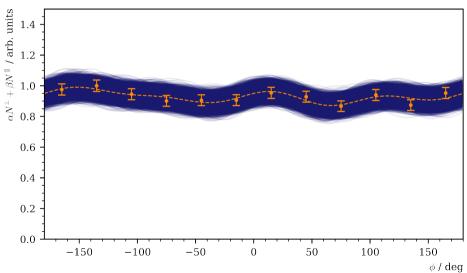
- ▶ All Bayesian fits performed using the *Python* frontend of *Stan*
- ► MCMC-sampling: adaptive Hamiltonian Monte-Carlo (HMC), i.e. No-U-Turn-Sampling (NUTS)
 - \blacktriangleright generate samples $\theta^{(1)}, \theta^{(2)}, \dots, \theta^{(S)}$ where each $\theta^{(t)}$ depends only on $\theta^{(t-1)}$
 - \triangleright simulate draws from the posterior by updating at point t such that the posterior increases (importance sampling)
- ▶ diagnosing convergence of MCMC:
 - potential scale reduction statistic $1.00 \lesssim \widehat{R} \lesssim 1.01$
 - ► Monte Carlo standard error (MCSE) 'small'
- ► Goodness of fit: posterior predictive checks (PPC)

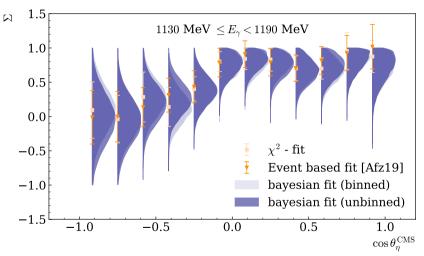


[Stan development team 2022; Hoffman and Gelman 2014]

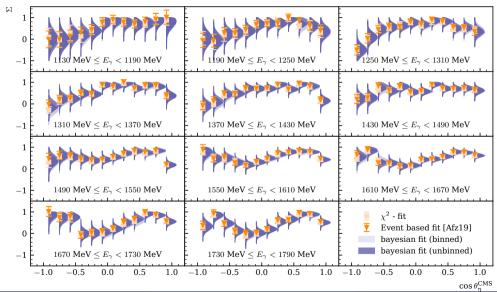


PPC for unbinned BAYESIAN fit:





Distinct advantage: sample only in physically allowed parameter space



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First, perform event selection regarding η' photoproduction

| η' | |
|--|---------------------------|
| Decay mode | Branching ratio |
| $\pi^+\pi^-\eta$ | 42.6% |
| $ ho^0 \gamma (o \pi^+ \pi^- \gamma)$ | 28.9%~(28.9%) |
| $\pi^0\pi^0\eta(o 6\gamma)$ | $22.8\% \ (8.8\%)$ |
| $\omega\gamma(\to\pi^+\pi^-\pi^0\gamma/\pi^0\gamma\gamma)$ | $2.52\% \ (2.2\%/0.21\%)$ |
| $\gamma\gamma$ | 2.3% |

[Workman et al. 2022]

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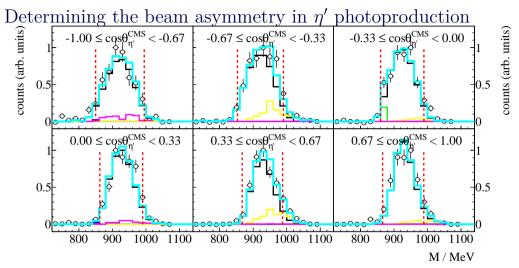
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[Workman et al. 2022]

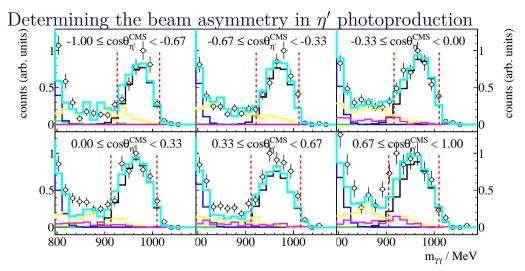
Event selection (η')

analysis performed in 3x6 bins of
$$(E_{\gamma}, \cos \theta)$$
 for $\gamma p \to p \eta' \to p \gamma \gamma$, $E_{\gamma} \in [1500, 1800]$ MeV
$$p_{\gamma} + p_{p} = p_{\eta'} + p_{\text{recoil}} = \underbrace{p_{\gamma_{1}} + p_{\gamma_{2}}}_{=p_{\eta'}} + p_{\text{recoil}}$$
$$p_{\gamma} + p_{p} = p_{\eta'} + p_{X} = \underbrace{p_{\gamma_{1}} + p_{\gamma_{2}}}_{=p_{\pi'}} + p_{X}$$

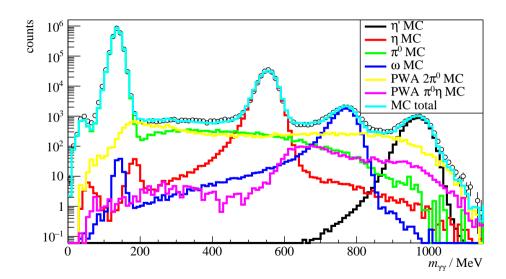
- \blacktriangleright one charged, two uncharged detector hits in coincidence with beam γ
- ► Coplanarity $\Delta \phi = \phi_{n'} \phi_{\text{recoil}} \stackrel{!}{=} 180^{\circ}$
- ightharpoonup Polar angle $\theta_X \stackrel{!}{=} \theta_{\text{recoil}}$
- ightharpoonup Missing mass $m_X \stackrel{!}{=} m_p$
- ▶ Invariant mass $m_{\gamma\gamma} \stackrel{!}{=} m_{\eta'}$



data points: m_X , turquoise: total MC, black: η' MC, yellow: $2\pi^0$ MC, magenta: $\pi^0\eta$ MC

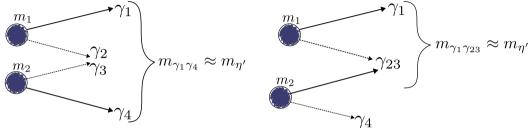


data points: $m_{\gamma\gamma}$, turquoise: total MC, black: η' MC, yellow: $2\pi^0$ MC, magenta: $\pi^0\eta$ MC, blue: ω MC

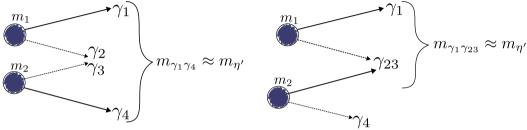


Determining the beam asymmetry in η' photoproduction Why these contributions from 4γ final states $2\pi^0$ (and $\pi^0\eta$)??

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acceptance for $2\pi^0$ events is almost vanishing $A(E_{\gamma}, \cos \theta) < 2 \cdot 10^{-3}$, yet

$$R = \frac{\sigma_{2\pi^0} \cdot \text{BR}_{2\pi^0 \to \gamma\gamma} \cdot \tilde{A}_{2\pi^0}}{\sigma_{n'} \cdot \text{BR}_{n' \to \gamma\gamma} \cdot \tilde{A}_{n'}} = \frac{5 \,\text{µb} \cdot 0.9765 \cdot 2 \cdot 10^{-3}}{1 \,\text{µb} \cdot 0.023 \cdot 0.61} \approx 0.7!$$

explains high background contributions of up to 45%

[Workman et al. 2022; Crede et al. 2009; Dieterle et al. 2020]

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- ► simple: we don't
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- ▶ simple: we don't
- ► real photons mimic a two-photon final state, no sensible additional cuts have been found
- ▶ main background from $2\pi^0$ photoproduction
- ▶ beam asymmetry for this reaction determined by [Mahlberg 2022]
- \triangleright correct estimates for Σ according to the amount of background

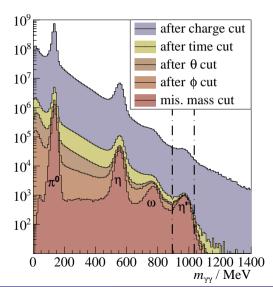
$$\Sigma^{\text{meas}} = (1 - \delta) \cdot \Sigma_{\eta'} + \delta \Sigma_{2\pi^0}$$

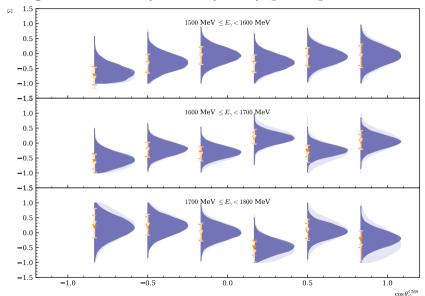
total: $\sim 8000 \ \eta' \rightarrow \gamma \gamma$ events

- ▶ perform unbinned fit as maximum likelihood fit and BAYESIAN fit
- ► Bayesian fit: modify likelihood

$$\tilde{p}\left(\phi, \Sigma\right) \to \tilde{p}\left(\phi, (1 - \delta)\Sigma_{1} + \delta\Sigma_{2}^{true}\right)$$
$$\Sigma_{2}^{true} \sim \mathcal{N}(\Sigma_{2}^{meas}, \tau)$$

▶ unbinned maximum likelihood fit: shift point estimates after fit





Systematic error:
$$\Delta \Sigma_{\eta'}^{\text{sys}} = \sqrt{\left(\frac{\Delta p_{\gamma}}{p_{\gamma}} \Sigma_{\eta'}\right)^2 + \left(\Delta \Sigma_{\eta'}\right)^2}$$
.

▶ polarization degree

$$\frac{\Delta p_{\gamma}}{p_{\gamma}} = \begin{cases} 0.05 & E_{\gamma} < 1600 \,\text{MeV}, \\ 0.08 & \text{otherwise} \end{cases}$$

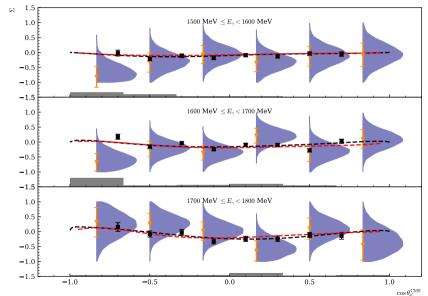
▶ background contributions

$$\Sigma^{\text{meas}} = (1 - \delta_1 - \delta_2) \cdot \Sigma_{n'} + \delta_1 \Sigma_{2\pi^0} + \delta_2 \Sigma^{\text{r bkg}},$$

thus:

$$\Delta\Sigma_{\eta'} = \max \left[\left| \frac{\Sigma^{\text{meas}} - \delta_1 \cdot \Sigma_{2\pi^0} \pm \delta_2 \cdot 1}{1 - \delta_1 - \delta_2} - \frac{\Sigma^{\text{meas}} - \delta \cdot \Sigma_{2\pi^0}}{1 - \delta} \right| \right]$$

[F. N. Afzal 2019; Eberhardt 2012]



- ► measurement at CLAS [Collins et al. 2017]: (black points),
- etaMAID PWA [Tiator et al. 2018]: (black, dashed lines),
- BnGa PWA
 [Anisovich
 et al. 2018]:
 (red, dashed
 lines)

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Conclusion

Summary

- $ightharpoonup \Sigma$ extracted for η and η' final state
- \blacktriangleright η results obtained with BAYESIAN fit agree with previous results
- \blacktriangleright η' results agree with previous measurements

Outlook

- ► use posterior distributions for PWA calculations
- ▶ increase precision of existing data , i.e. investigate $\eta' \to \pi^0 \pi^0 \eta$
- \blacktriangleright measurement of observables near η' production threshold

BACKUP & REFERENCES

Full PDF for unbinned maximum likelihood fit

$$-\ln \mathcal{L} = \sum_{i=1}^{n} -\ln(p_{\text{prompt}}(\phi_i, p_{\gamma,i}, \Sigma, a_1 \dots a_4, b_1 \dots b_4)) + \sum_{i=1}^{m} -\ln\left(p_{\text{sideband}}(\phi_j, p_{\gamma,j}, \Sigma^{\text{bkg}}, a_1^{\text{bkg}} \dots a_4^{\text{bkg}}, b_1^{\text{bkg}} \dots b_4^{\text{bkg}})\right)$$

where

$$p_{\text{prompt}} = f_{\text{sig}} \cdot \tilde{p}(\phi, p_{\gamma}, \Sigma, a_{1} \dots a_{4}, b_{1} \dots b_{4})$$

$$+ (1 - f_{\text{sig}}) \cdot \tilde{p}(\phi, p_{\gamma}, \Sigma^{\text{bkg}}, a_{1}^{\text{bkg}} \dots a_{4}^{\text{bkg}}, b_{1}^{\text{bkg}} \dots b_{4}^{\text{bkg}})$$

$$p_{\text{sideband}} = \tilde{p}(\phi, p_{\gamma}, \Sigma^{\text{bkg}}, a_{1}^{\text{bkg}} \dots a_{4}^{\text{bkg}}, b_{1}^{\text{bkg}} \dots b_{4}^{\text{bkg}})$$

and

$$\tilde{p}(\phi, \Sigma) = \frac{\left(1 + p_{\gamma} \Sigma \cos\left(2\left(\alpha^{\parallel} - \phi\right)\right)\right) \cdot \left(\sum_{k=0}^{4} a_{k} \sin(k\phi) + b_{k} \cos(k\phi)\right)}{1 - \frac{1}{2} a_{2} p_{\gamma} \Sigma}$$

Additional theoretical basics

Unpolarized differential cross section

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \frac{1}{4}\rho \sum_{\mathrm{spins}} |\langle f|\mathcal{F}|i\rangle|^2,$$

where

$$\mathcal{F} = i(\vec{\sigma} \cdot \vec{\epsilon})F_1 + (\vec{\sigma} \cdot \hat{q})(\vec{\sigma} \cdot (\hat{k} \times \vec{\epsilon}))F_2 + i(\vec{\sigma} \cdot \hat{k})(\hat{q} \cdot \vec{\epsilon})F_3 + i(\vec{\sigma} \cdot \hat{q})(\hat{q} \cdot \vec{\epsilon})F_4$$

 F_i : complex CGLN Amplitudes

[Chew et al. 1957]

 $\frac{d\sigma}{d\Omega} \in \mathbb{R}$, not sufficient do determine \mathcal{F} unambiguously

 \rightarrow Polarization Observables can be related to F_i

Diagnostics of a BAYESIAN fit

- \triangleright \hat{R} : measure of convergence for chains
- ▶ Monte-Carlo-Standard-Error: measure for adequate sample size
- ▶ posterior predictive checks: "goodness of fit"

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