Experimental studies of the $\Lambda(1405)$

physics654 – Seminar on exotic multi-quark states

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Motivation

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vollständige sätze verme

What is special about the $\Lambda(1405)$?

- ~~its~~ mass does not fit well into constituent quark models ~~which do predict baryon masses well for other baryons~~
- invariant mass distribution (line shape) differs significantly from usual Breit-Wigner shape~~s~~
- candidate for an exotic multiquark state (bound system of $\overline{K}N$) since its mass lies just below production threshold

There are ~~(very)~~ many different theoretical approaches to explain this behavior

$\rightarrow$ There is need for more experimental data!

~~some plots/pictures?~~

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4. Conclusion

Continuous Electron Beam Accelerator Facility (CEBAF)

Wo?

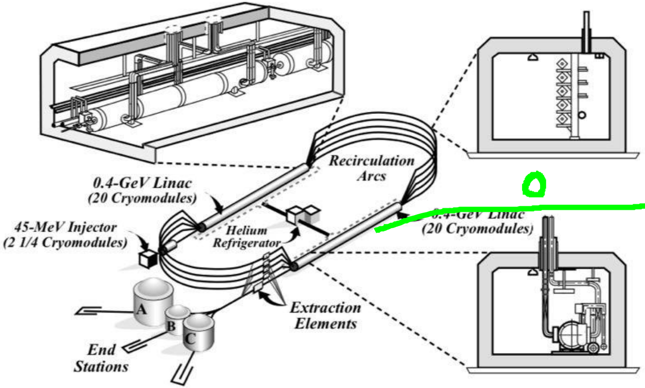


Figure 1: CEBAF layout at Jefferson Lab, [Mecking et al. 2003]

Das wird nicht in beamer gemacht(kein sin für referenzen)

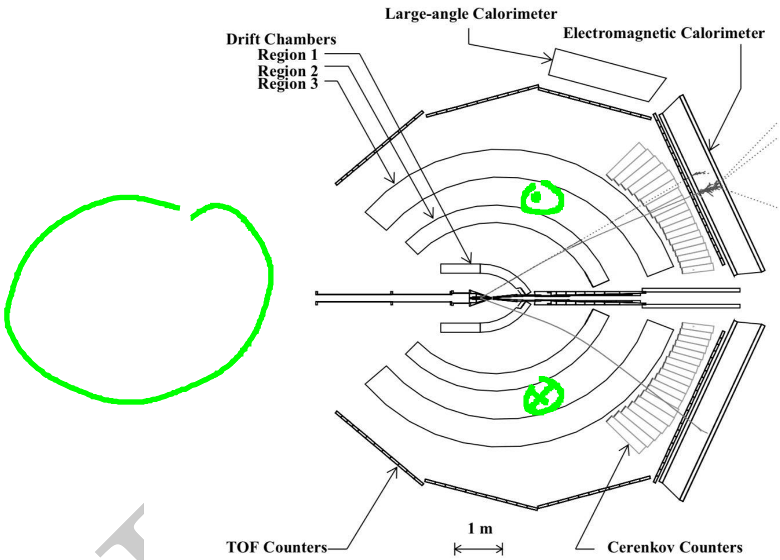
Continuous Electron Beam Accelerator Facility (CEBAF)

How can we access $\Lambda(1405)$ with this setup?

- employ radiator target for the electrons
- generate high energy photons using the *bremsstrahlung* process
- shoot high energy photons on proton target (LH2)
- then we can observe $\gamma p \to K^+ \Lambda(1405)$ while knowing p_γ, p_p

~~Wise zeigts du den Tagger nicht statdessen?~~

CEBAF Large Acceptance Spectrometer (CLAS)



Magnetfeld ı

Figure 2: CLAS layout at Jefferson Lab, [Mecking et al. 2003]

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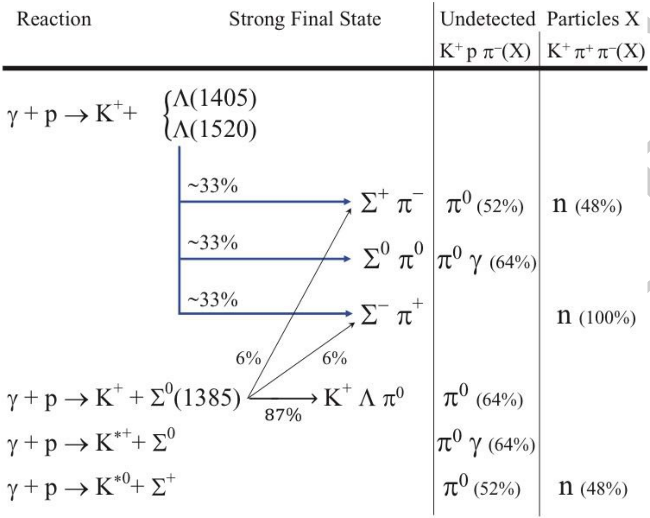
Reaction kinematics



$$\Sigma^+ \rightarrow p\pi^0$$
$$\Sigma^+ \rightarrow n\pi^+$$
$$\Sigma^0 \rightarrow \gamma\Lambda \qquad \rightarrow \gamma p\pi^-$$
$$\Sigma^- \rightarrow n\pi^-$$
$$\Lambda \rightarrow p\pi^-$$

~~**Figure 3**~~: Possible and studied reactions in the analysis of the lineshapes of $\Lambda(1405)$, ~~taken from~~ [Moriya, Schumacher, Adhikari et al. 2013]

Event selection

~~There are~~ two sets of reactions that the detector sees

1. $K^+ p \pi^-$
2. $K^+ \pi^+ \pi^-$

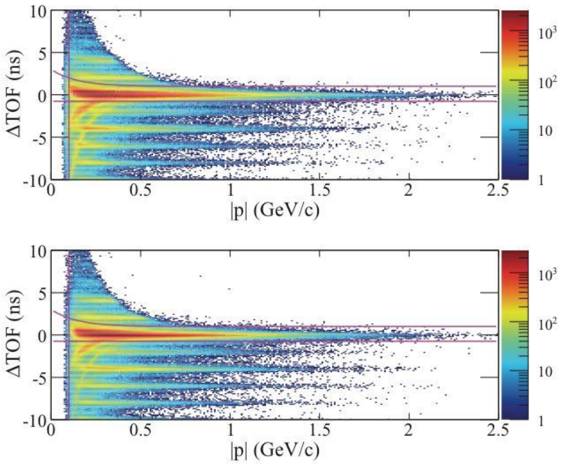
~~There are~~ many cuts that can be made that apply to both

Initial selection of particles

- ~~Initial~~ final state Kaon ~~selection~~
- fiducial cuts
- remove false K^+ due to π^+ or p
- Loose ΔTOF cut ($\Delta \mathrm{TOF} = t_{\text{meas}} - t_{\text{calc}} = t_{\text{meas}} - \frac{l\sqrt{p^2+m_0^2}}{cp}$)
- vertex z cut
- minimum $|\mathbf{p}|$ cuts
- Precise ΔTOF cuts

- p, beta particle idententifc

Event selection



~~Figure 4~~: ΔTOF for $\pi^{\pm}$ @ $2.35 < W < 2.45\,$GeV, applied cuts are shown in magenta. [Moriya, Schumacher, Adhikari et al. 2013]

Beide sehen zimlich gleich aus ich würde nur eins z

Event selection

Data divided into bins of 100 MeV \sqrt(s) and 0.1 cos(\theta^{K^+}_CMS) s

~~In all channels, the data was divided into 10 bins of energy spanning 100 MeV in the CMS energy $W = \sqrt{s}$ and 20 angle bins in the CMS kaon angle.~~
$\rightarrow$ the following analysis was performed independently in every bin of energy and angle!

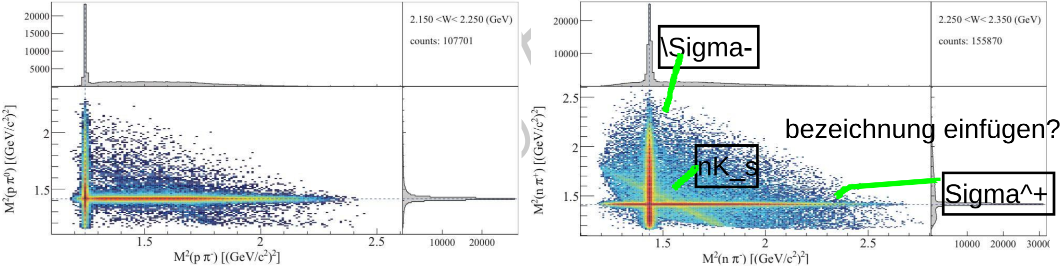
Event selection

extracting $\Lambda\pi^0$ and $\Sigma^+\pi^-$

- reminder: $\Lambda \to p\pi^-, \Sigma^+ \to p\pi^0$
- final state particles: $K^+p\pi^-(\pi^0)$
- determine p_π via missing mass fit
- apply cuts based on fits to the invariant masses $M_{p\pi^-}$ and $M_{p\pi^0}$

extracting $\Sigma^+\pi^-$ and $\Sigma^-\pi^+$

- reminder: $\Sigma^\pm \to n\pi^\pm$
- final state particles: $K^+\pi^+\pi^-(n)$
- determine p_n via missing mass fit
- apply cuts based on fits to the invariant masses $M_{n\pi^\pm}$

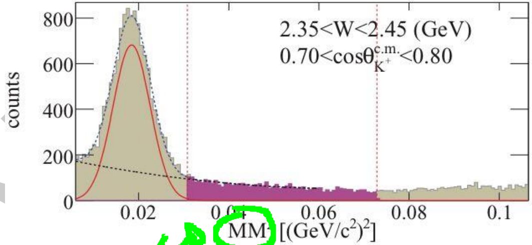


~~**Figure 5:**~~ Dalitz-like plots of the above mentioned invariant masses, ~~taken from~~ [Moriya, Schumacher, Adhikari et al. 2013]

Event selection

extracting $\Sigma^0\pi^0$

- reminder: $\Sigma^0 \to \gamma\Lambda \to \gamma p\pi^-$
- final state particles: $K^+p\pi^-(\gamma\pi^0)$ - missing mass fit is not applicable here: demand the missing mass is sufficiently greater than m_π
- make cuts based on the invariant mass $M_{p\pi^-}$
- now the missing mass ($\gamma p \to K^+X$) gives the $\Sigma^0\pi^0$ lineshape



~~**Figure 6:**~~ Invariant mass of the $\gamma\pi$ system, selection range in magenta. ~~Taken from~~ [Moriya, Schumacher, Adhikari et al. 2013]

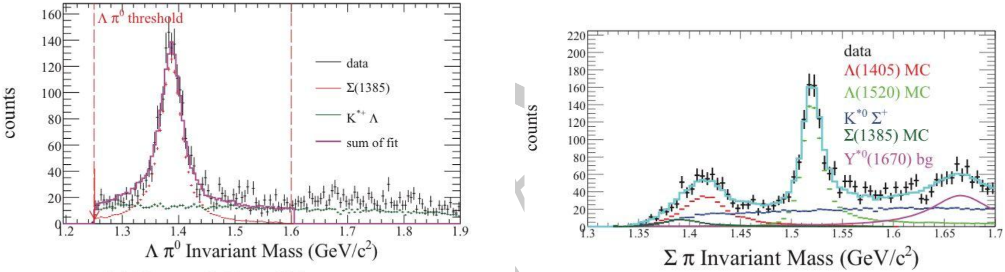
Measurements and analysis

- Now the signal regions have been established
- The true lineshape of the $\Lambda(1405)$ has to be extracted from the vast of reactions $\rightarrow$ any other contributions have to be ~~excluded~~ subtracted
- Strategy: use of MONTE-CARLO fits to the data, simulating the contribution of other resonances using the PDG widths and masses

~~I will not really go into detail here...(??)~~

Measurements and analysis

Some fit results:



~~Figure 7~~: Sample fit results of invariant mass spectra for a single bin in energy and angle, ~~taken from~~ [Moriya, Schumacher, Adhikari et al. 2013]

Interpretation of the results

Having substracted all unwanted reactions, one can obtain the true $\Lambda(1405)$ lineshapes:

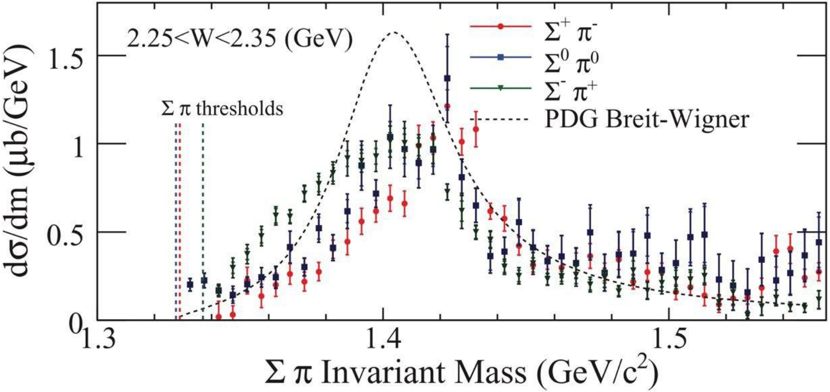
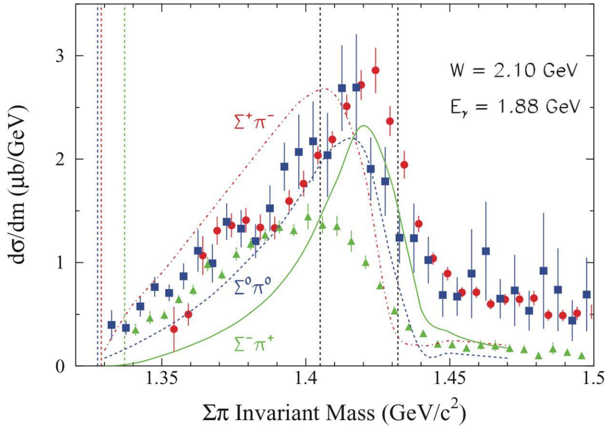


Figure 8: Lineshapes of the $\Lambda(1405)$ for the 3 different decay channels and the PDG Breit-Wigner. The data were summed over all angles for better statistics. [Moriya, Schumacher, Adhikari et al. 2013]

Interpretation of the results

There are in fact predictions of the lineshapes differing by decay channel [Nacher et al. 1999].
→ main idea: not one amplitude, but two due to isospin decomposition.



~~Figure 9:~~ Lineshapes of the Λ(1405) for the 3 different decay channels and the prediction of [Nacher et al. 1999], ~~taken from~~ [Moriya, Schumacher, Adhikari et al. 2013]

Interpretation of the results

MORIYA ET AL. saw two main reasons for the lineshapes differing from a simple BREIT-WIGNER:

1. Isospin decomposition
2. channel coupling between the detected $\Sigma\pi$ and $N\overline{K}$ final states

Isospin decomposition

let

$$|t_I|^2 = |\langle I, 0|T^{(I)}|\gamma p\rangle|^2,$$

then we can write (neglecting $I = 2$ using CGK)

$$|T_{\pi^-\Sigma^+}|^2 = \frac{1}{3}|t_0|^2 + \frac{1}{2}|t_1|^2 - \frac{2}{\sqrt{6}} \boxed{\text{t0t1}} \cos\phi_{01}$$

$$|T_{\pi^0\Sigma^0}|^2 = \frac{1}{3}|t_0|^2$$

$$|T_{\pi^+\Sigma^-}|^2 = \frac{1}{3}|t_0|^2 + \frac{1}{2}|t_1|^2 + \frac{2}{\sqrt{6}} \cos\phi_{01}$$

Channel coupling

the t_I are described by one or two BREIT-WIGNER amplitudes with mass dependent widths Γ

$\rightarrow$ modify the amplitude preserving analyticity [Flatté 1976] including the $N\overline{K}$ decay channel available at threshold $m_K + m_p \approx 1434\,\text{MeV}$

Interpretation of results

Fits with two $I = 1$ and one $I = 0$ amplitudes lead to best agreement with measured data also show $\Sigma^-\pi^+, \Sigma^0\pi^0$?

ja zeig ein drittel von allen

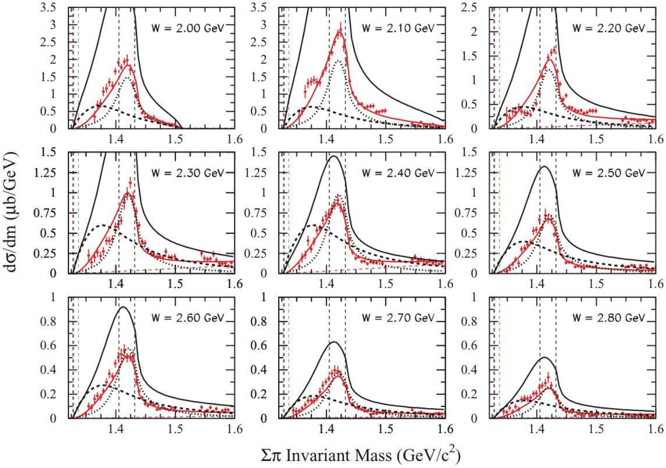


Figure 10: Data and fits for $\Sigma^+\pi^-$ for different bins in W. $I = 0$ (solid black), narrow $I = 1$ (dotted black) and wide $I = 1$ (dashed black). Background in dashed red. [Moriya, Schumacher, Adhikari et al. 2013]

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Theoretical basics $\boxed{\text{Spin}}$

The $\Lambda(1405)$ is so far (mostly) assumed to have $J^P = \frac{1}{2}^-$, but this has not been determined experimentally

Measuring spin

- consider the strong decay $Y^* \to Y\pi$, with J^P the spin and parity of Y^*
- the $Y\pi$ angular distribution will only depend on J

$$I(\theta_Y) = \text{const.} \qquad J = 1/2$$

$$I(\theta_Y) \propto 1 + \frac{3(1-2p)}{2p+1}\cos^2\theta_Y \qquad J = 3/2,$$

  where θ_Y is the polar angle of the decay direction of Y in the Y^* rest frame, p describes the fraction of spin projections along the z axis
- uniform decay pattern is best evidence for spin $J = 1/2$

[Moriya, Schumacher, Aghasyan et al. 2014]

Theoretical basics

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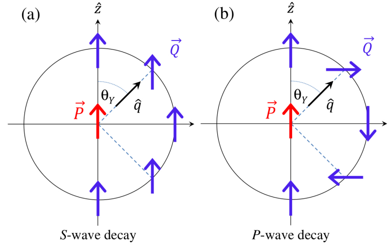
Measuring parity

- the key to accessing the parity lies in determining the Polarization transfer to the decay product Y which we will denote $\mathbf{Q}$
- the angular distribution of $\mathbf{Q}$ will only depend on $\mathbf{P}$

$$\mathbf{Q}(\theta_Y) = \text{const.} \qquad J^P = 1/2^-$$
$$\mathbf{Q}(\theta_Y) = -\mathbf{P} + 2(\mathbf{P}\cdot\mathbf{q})\mathbf{q} \qquad J^P = 1/2^+$$

- $\mathbf{Q}$ can be measured from weak decay angular distribution of Y



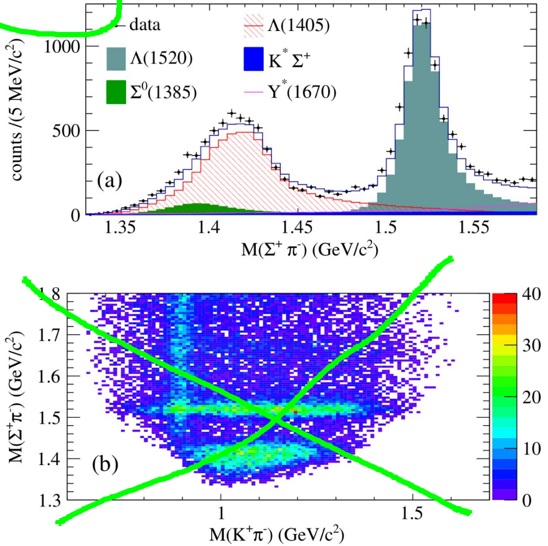
~~**Figure 11:**~~ Polarization transfer in the strong decay $Y^* \to Y\pi$, ~~taken from~~ [Moriya, Schumacher, Aghasyan et al. 2014]

[Moriya, Schumacher, Aghasyan et al. 2014 and Ref. therein]

Measurements and analysis $\boxed{\text{BackUp}}$

Event selection

- select kinematic region where the $\Sigma\pi$ invariant mass is dominated by the $\Lambda(1405) \rightarrow M_{\Sigma\pi} \in$ 1.30 GeV to 1.45 GeV
- inspect nine bins in energy and angle, namely with CM energy at 2.6, 2.7 and 2.8 GeV and the three forwardmost kaon angle bins each



~~**Figure 12:**~~ $\Sigma\pi$ ~~and $K\pi$~~ invariant mass in the vicinity of the $\Lambda(1405)$, ~~taken from~~ [Moriya, Schumacher, Aghasyan et al. 2014]

Measurements and analysis

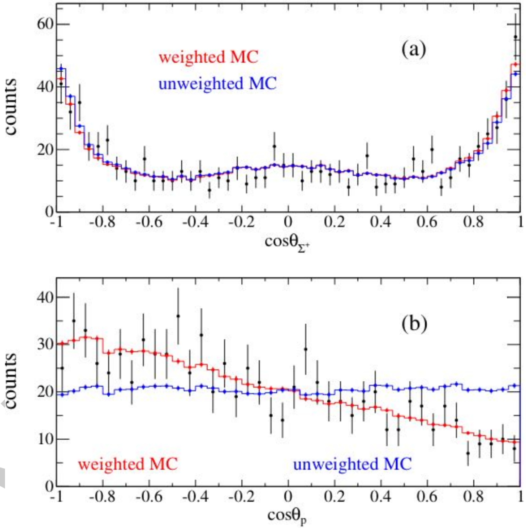
aufspalten + extra bild

Analysis procedure

- 1 plot the angular distribution of the projections $\cos\theta_\Sigma$ and $\cos\theta_p$ for each bin
- 1 test each spin hypothesis using Monte-Carlo maximum likelihood fits, which employ angular decay probability distributions according to each hypothesis for $\Sigma\pi$ and $p\pi$. From the fit Q_z will be determined
- 2 test parity hypotheses by determining $Q_z(\cos\theta_\Sigma)$
- 1 compare each hypothesis by calculating a χ^2 probability

Result: data is consistent with $J^P = 1/2^-$ but does in principle not rule out $J^P = 3/2^+$. $1/2^+$ and $3/2^-$ can be discarded.

Measurements and analysis



~~Figure 13:~~ Distributions of the projections of (a)$\cos\theta_\Sigma$ and (b) $\cos\theta_p$ @ $2.65 < W < 2.75\,\mathrm{GeV}$ and $0.70 < \cos\theta < 0.80$, ~~taken from~~ [Moriya, Schumacher, Aghasyan et al. 2014]

Measurements and analysis

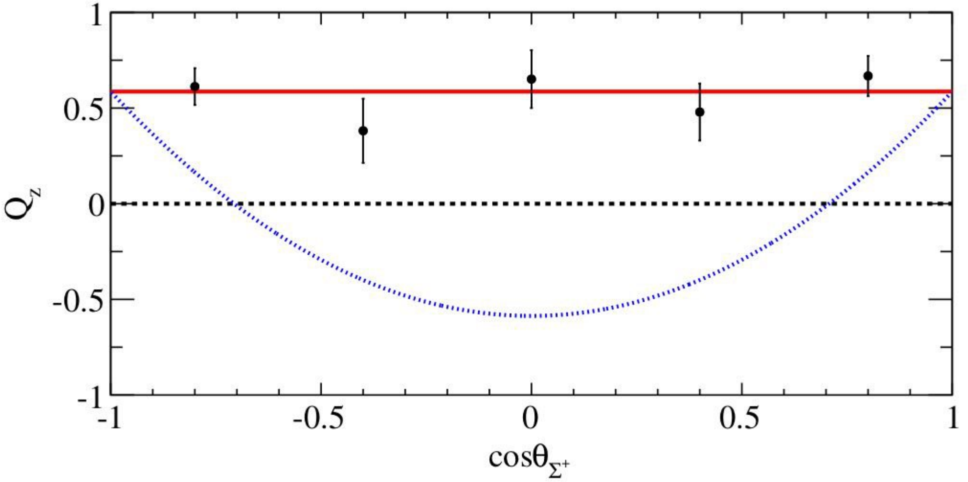


Figure 14: angular distribution of the polarization Q_z @ $2.65 < W < 2.75\,\mathrm{GeV}$ and $0.70 < \cos\theta < 0.80$. Red: average, blue: expectation for P-wave decay. Taken from [Moriya, Schumacher, Aghasyan et al. 2014]

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Conclusion

Lineshape measurement

- the CLAS detector was used to study $\gamma p \to K^+\Lambda(1405)$
- after selecting the correct events the true lineshape was extracted using Monte-Carlo sim. of the yield of other resonances
- the lineshapes in different decay channels differ from each other and from a simple Breit-Wigner
- a phenomenological isospin decomposition model was able to describe the data

Spin parity measurement

- the angular distribution of the decay $\Lambda(1405) \to \Sigma^+\pi^-$ was studied
- the angular distributions were tested against various J^P hypotheses
- the data is consistent with $J^P = 1/2^-$ but does not exclude $J^P = 3/2^+$

References

- Flatté, S.M. (1976). 'Coupled-channel analysis of the $\pi\eta$ and $K\overline{K}$ systems near $K\overline{K}$ threshold'. In: *Physics Letters B* 63.2, pp. 224–227. ISSN: 0370-2693. DOI: https://doi.org/10.1016/0370-2693(76)90654-7. URL: https://www.sciencedirect.com/science/article/pii/0370269376906547.
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