

Quantum Spin

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April 2, 2015

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- Orbital momentum is momentum due to revolution, spin momentum is due to rotation
- Spin momentum is really just the sum of the orbital momentum of all the particles which make up the object

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- Instead we consider particles to have *intrinsic* angular momentum, **S**
- Elementary Particles each have a specific spin which never changes
- A few examples:

Particle	Spin
pi meson	0
e^{-}, p, n	$\frac{1}{2}$
photon	1
delta	$\frac{3}{2}$
graviton	2

Defining the Spin Operators

- We begin with the commutation relations:

$$[S_x, S_y] = i\hbar S_z \quad [S_y, S_z] = i\hbar S_x \quad [S_z, S_x] = i\hbar S_y$$

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and thus

$$S_{\pm} |sm\rangle = \hbar \sqrt{s(s+1) - m(m \pm 1)} |s(m \pm 1)\rangle \quad (2)$$

Here $S_{\pm} = S_x \pm iS_y$

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Here we have

$$\chi_+ = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

represents spin up and

$$\chi_- = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

represents spin down

- From equation 1, we have:

$$\mathbf{S}^2 \chi_+ = \frac{3}{4} \hbar^2 \chi_+$$

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- Solving this system we see that

$$\mathbf{S}^2 = \frac{3}{4} \hbar^2 \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

Spin 1/2

- From $\mathbf{S}_z\chi_+ = \frac{\hbar}{2}\chi_+$ and $\mathbf{S}_z\chi_- = -\frac{\hbar}{2}\chi_-$ it follows:

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- Referring back to equation 2, we have:

$$\mathbf{S}_+\chi_- = \hbar\chi_+ \quad \mathbf{S}_-\chi_+ = \hbar\chi_- \quad \mathbf{S}_+\chi_+ = \mathbf{S}_-\chi_- = 0$$

so

$$\mathbf{S}_+ = \hbar \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \quad \mathbf{S}_- = \hbar \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}$$

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- Recall $S_{\pm} = S_x \pm iS_y$. It follows that

$$\mathbf{S}_x = \frac{\hbar}{2} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad \mathbf{S}_y = \frac{\hbar}{2} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

Spin 1/2

- Since \mathbf{S}_x , \mathbf{S}_y and \mathbf{S}_z all have a factor of $\frac{\hbar}{2}$ we will let σ_x be the matrix in \mathbf{S}_x , and σ_y, σ_z defined similarly. Then $\mathbf{S} = \frac{\hbar}{2}\boldsymbol{\sigma}$ and

$$\sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad \sigma_y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \quad \sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

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- These are known as the **Pauli Spin Matrices**

Addition of Angular Momenta

- Let's consider two spin-1/2 particles: Electron and Proton
- There are four possible configurations

$$\uparrow\downarrow, \uparrow\uparrow, \downarrow\downarrow, \downarrow\uparrow$$

- How would we calculate the total angular momentum of the atom? We let

$$\mathbf{S} = \mathbf{S}^{(1)} + \mathbf{S}^{(2)}$$

- Each of the four states represent an eigenstate of S_z
- Therefore, the z components add

$$S_z \chi_1 \chi_2 = (S_z^1 + S_z^2) \chi_1 \chi_2 = (S_z^1 \chi_1) \chi_2 + \chi_1 (S_z^2 \chi_2)$$

Addition of Angular Momenta

- Continued...

$$\begin{aligned} S_z \chi_1 \chi_2 &= (S_z^1 + S_z^2) \chi_1 \chi_2 = (S_z^1 \chi_1) \chi_2 + \chi_1 (S_z^2 \chi_2) \\ &= (\hbar m_1 \chi_1) \chi_2 + \chi_1 (\hbar m_2 \chi_2) = \hbar(m_1 + m_2) \chi_1 \chi_2 \end{aligned}$$

- The quantum number m is just $m_1 + m_2$

$$\uparrow\uparrow : m = 1$$

$$\uparrow\downarrow : m = 0$$

$$\downarrow\uparrow : m = 0$$

$$\downarrow\downarrow : m = -1$$

Spin States

- We can apply the lowering operator to the state $\uparrow\uparrow$ to find the states of S :

$$\begin{aligned} S_-(\uparrow\uparrow) &= (S_-^{(1)} \uparrow) \uparrow + \uparrow (S_-^{(2)} \uparrow) \\ &= (\hbar \downarrow) \uparrow + \uparrow (\hbar \downarrow) = \hbar(\downarrow\uparrow + \uparrow\downarrow) \end{aligned}$$

- We can do this to obtain three states with $s = 1$
- In $|sm\rangle$ notation they are:

$$\left\{ \begin{array}{l} |1\ 1\rangle = \uparrow\uparrow \\ |1\ 0\rangle = \frac{1}{\sqrt{2}}(\uparrow\downarrow + \downarrow\uparrow) \\ |1\ -1\rangle = \downarrow\downarrow \end{array} \right\} \quad s = 1 \text{ (triplet)}.$$

- This makes $s = 1$ a triplet state

Singlet State

- On the contrary, if we look at $s = 0$ we have only

$$\left\{ |00\rangle = \frac{1}{\sqrt{2}}(\uparrow\downarrow - \downarrow\uparrow) \right\} \quad s = 0 \text{ (singlet)}.$$

- This makes $s = 0$ a singlet state
- If we apply the raising or lowering operator to this we obtain zero
- Then, the total spin that two spin-1/2 particles can carry is either 1 or 0
- This is true if the triplet states are eigenvectors of S^2 with eigen value $2\hbar^2$
- The singlet state must also be an eigenvector of S^2 with eigenvalue of 0

Spin State Eigenvectors

- Let's check:

$$\text{spinsquared } S^2 = (S^{(1)} + S^{(2)}) \cdot (S^{(1)} + S^{(2)}) = (S^{(1)})^2 + (S^{(2)})^2 + 2S^{(1)} \cdot S^{(2)}. \quad (3)$$

- From previous equations we have

$$\begin{aligned} S^{(1)} \cdot S^{(2)}(\uparrow\downarrow) &= (S_x^{(1)} \uparrow)(S_x^{(2)} \downarrow) + (S_y^{(1)} \uparrow)(S_y^{(2)} \downarrow) + (S_z^{(1)} \uparrow)(S_z^{(2)} \downarrow) \\ &= \left(\frac{\hbar}{2} \downarrow\right)\left(\frac{\hbar}{2} \uparrow\right) + \left(\frac{i\hbar}{2} \downarrow\right)\left(\frac{-i\hbar}{2} \uparrow\right) + \left(\frac{\hbar}{2} \uparrow\right)\left(\frac{-\hbar}{2} \downarrow\right) \\ &= \frac{\hbar^2}{4}(2 \downarrow\uparrow - \uparrow\downarrow) \end{aligned}$$

- As we might expect, the opposite orientation gives:

$$S^{(1)} \cdot S^{(2)}(\downarrow\uparrow) = \frac{\hbar^2}{4}(2 \uparrow\downarrow - \downarrow\uparrow)$$

Proof Continued

- If we check:

$$S^{(1)} \cdot S^{(2)}|10\rangle = \frac{\hbar^2}{4} \frac{1}{\sqrt{2}} (2 \downarrow\uparrow - \uparrow\downarrow + 2 \uparrow\downarrow - \downarrow\uparrow) = \frac{\hbar^2}{4} |10\rangle$$

and

$$S^{(1)} \cdot S^{(2)}|00\rangle = \frac{\hbar^2}{4} \frac{1}{\sqrt{2}} (2 \downarrow\uparrow - \uparrow\downarrow - 2 \uparrow\downarrow + \downarrow\uparrow) = -\frac{3\hbar^2}{4} |00\rangle$$

- Then from ?? we conclude that

$$S^2|10\rangle = \left(\frac{3\hbar^2}{4} + \frac{3\hbar^2}{4} + 2\frac{\hbar^2}{4} \right) |10\rangle = 2\hbar^2 |10\rangle$$

and

$$S^2|00\rangle = \left(\frac{3\hbar^2}{4} + \frac{3\hbar^2}{4} + 2\frac{3\hbar^2}{4} \right) |00\rangle = 0$$

Spin State Eigenvectors

- We have just shown that if you combine spin $1/2$ with spin $1/2$ you will obtain spin 1 and 0 states.
- Spin 1 state is an eigenvector of S^2 with eigenvalue $2\hbar^2$ and spin 0 state is an eigenvector of S^2 with eigenvalue 0 .
- If we are to expand this to the general case we see that by combining spin s_1 with spin s_2 , we get a total spin of

$$s = (s_1 + s_2), (s_1 + s_2 - 1), (s_1 + s_2 - 2), \dots, |s_1 - s_2|$$

- This gives rise to the combined state of total spin s and z -component m as a linear combination of composite states:

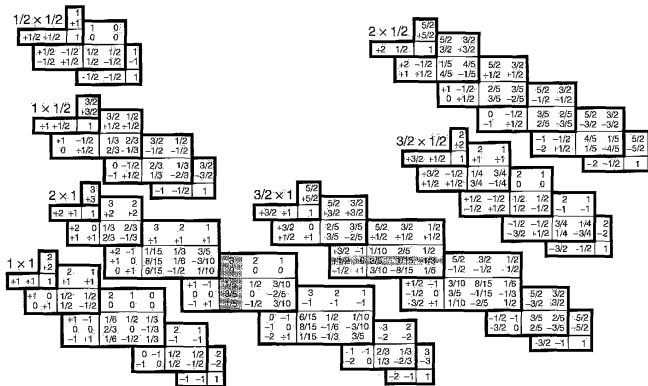
$$|sm\rangle = \sum_{m_1+m_2=m} C_{m_1 m_2 m}^{s_1 s_2 s} |s_1 m_1\rangle |s_2 m_2\rangle$$

Clebsch-Gordan Table

- The Clebsch-Gordan coefficients (a square root sign is present for every entry):

The table displays Clebsch-Gordan coefficients for the decomposition of the tensor product of two SU(2) representations into an SU(4) representation. The rows are labeled with the initial state $J_1 \times J_2$ and the final state J . The coefficients are arranged in a grid that is roughly 10 columns wide and 10 rows high, with some cells empty to maintain the staircase structure.

$$|30\rangle = \frac{1}{\sqrt{5}}|21\rangle|1-1\rangle + \sqrt{\frac{3}{5}}|20\rangle|10\rangle + \frac{1}{\sqrt{5}}|2-1\rangle|11\rangle$$



The End

Thanks for listening!