

1.

The first equation $\frac{dx}{dt} = v \cos \varphi$ represents the velocity of the particle in the x-axis direction, which is equal to v times the cosine of the angle φ . The velocity of a particle in the x-axis direction is related to its current orientation φ .

The second equation $\frac{dy}{dt} = v \sin \varphi$ represents the velocity of the particle in the y-axis direction, which is equal to v times the sine of the angle φ . The velocity of a particle in the y-axis direction is related to its current orientation φ .

The third equation $\frac{d\varphi}{dt} = \sqrt{2D_R} \cdot \xi$ describes the change of the particle's angle φ with time. Here ξ is a random variable obeying Gaussian white noise, representing the randomness caused by rotational diffusion. The $\sqrt{2D_R}$ on the right side of the equation represents the intensity of the influence of the rotational diffusion coefficient D_R on the noise.

2.

The meaning of this equation is that by calculating the dot product of the motion direction at different time points, we can obtain the autocorrelation of the motion direction. The autocorrelation function takes a maximum value of 1 when $t=0$, indicating that the movement direction is exactly the same at the same time point. As the time interval t increases, the autocorrelation function gradually decreases. The exponential decay form $\exp(-|t|/\tau_R)$ expresses the memory of the motion direction, that is, the motion direction maintains relevance within a short time scale, but gradually loses relevance as time increases. This equation reveals that the direction of motion of a self-propelled particle is continuous. The longer the direction duration τ_R , the particle's direction of motion remains consistent for a longer period of time; while when τ is shorter, the particle's direction changes more frequently.

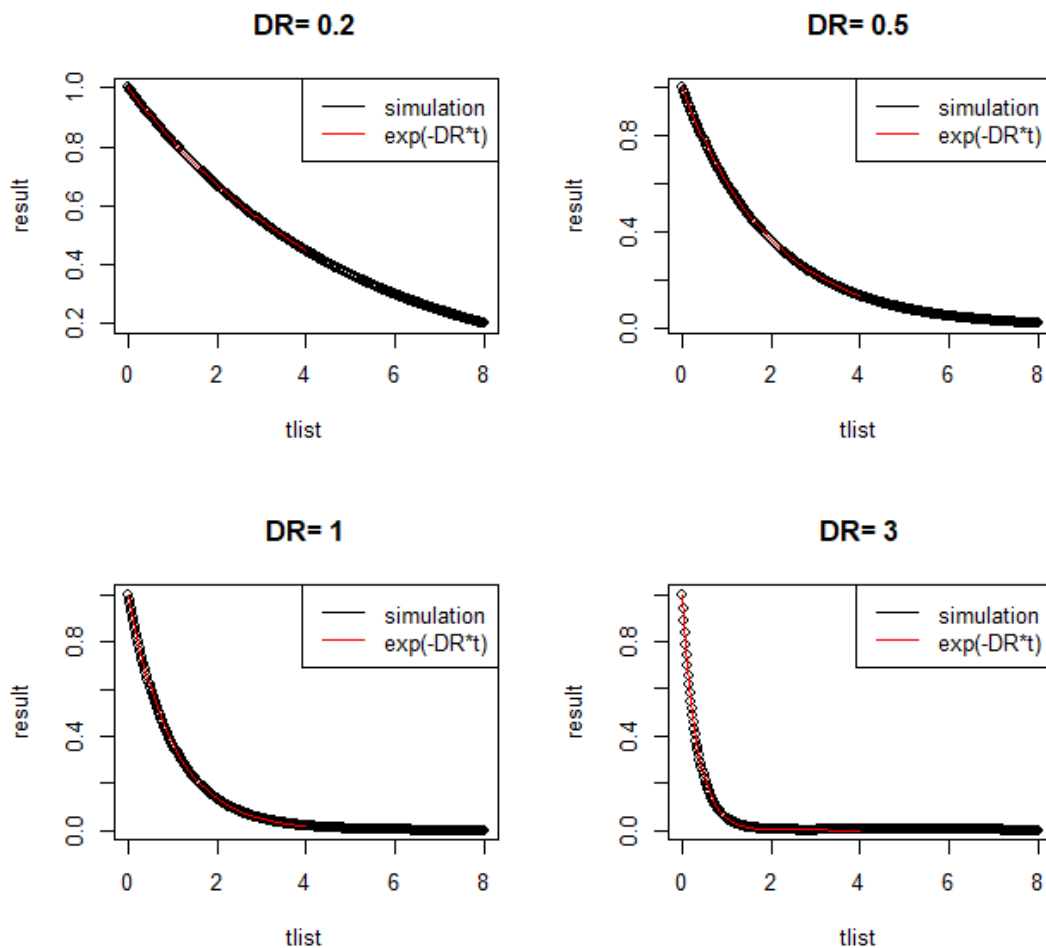
3.

Handwriting prove:

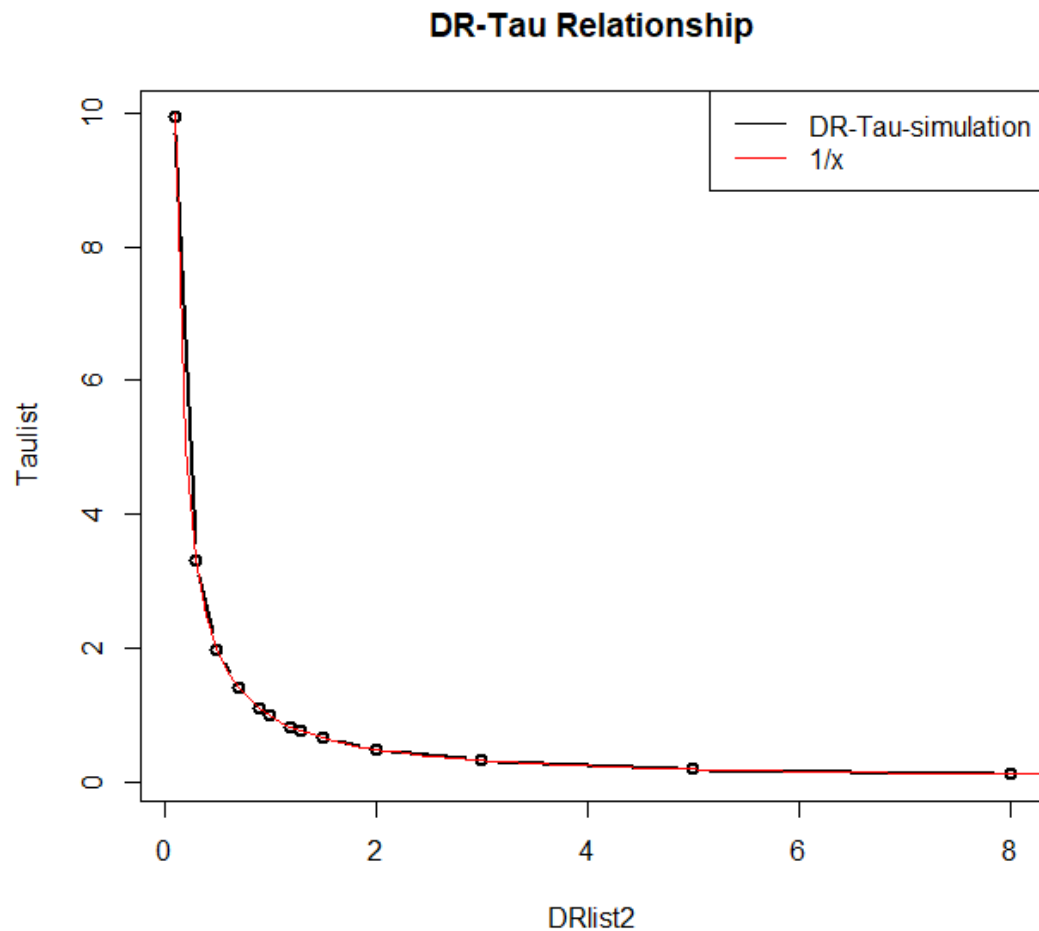
$$\begin{aligned}
 \mathbf{n}(t) &= (\cos \varphi(t), \sin \varphi(t)) \cdot (1, 0) + (\cos \varphi(t), \sin \varphi(t)) \cdot (0, 1) ; \varphi'(t) = \sqrt{2D_R} \cdot \xi \quad (0, \delta^2) \\
 \text{And } \langle \mathbf{n}(s+t) \cdot \mathbf{n}(s) \rangle &= \lim_{T \rightarrow \infty} \left[\frac{1}{T} \int_0^T \mathbf{n}(s+t) \cdot \mathbf{n}(s) ds \right] \rightarrow AC. \\
 \therefore \mathbf{n}(s+t) \cdot \mathbf{n}(s) &= (\cos \varphi(s+t), \sin \varphi(s+t)) \cdot (\cos \varphi(s), \sin \varphi(s)) \\
 &= \cos[\varphi(s+t) - \varphi(s)] \quad \text{take } \sqrt{2D_R} \text{ as } A. \\
 \varphi(s+t) - \varphi(s) &= \int_s^{s+t} A \xi(u) du \\
 \therefore AC &= \lim_{T \rightarrow \infty} \left[\frac{1}{T} \int_0^T \cos \left[\int_s^{s+t} A \xi(u) du \right] ds \right] \\
 &= \lim_{T \rightarrow \infty} \left[\frac{1}{T} \int_0^T \cos \left[A \cdot \int_s^{s+t} \xi(u) du \right] ds \right] \\
 \langle \xi(u) \rangle &= 0. \quad \xi(u) \text{ 只有 } \Delta u = 0 \text{ 时才 } \neq 0. \\
 \int_s^{s+t} \xi(u) du &\text{ 在 } t \rightarrow 0 \text{ 时可视为 } \sum_{i=1}^{n \rightarrow \infty} \Delta t \cdot \xi(s+i\Delta t), \text{ } n \text{ 个独立同分布.} \\
 E &\rightarrow 0; \text{ Var } \rightarrow \sum \Delta t = t. \quad \xrightarrow{\text{大数定律}} \sim \mathcal{N}(0, t)
 \end{aligned}$$

$\therefore \int_s^{s+t} \xi(u) du$ 可视为在 $\xi(s)$ 处采样, $t \rightarrow 0$. LHS $\approx \xi(s) \cdot \sqrt{t}$.
 $\therefore AC = \lim_{T \rightarrow \infty} \left[\frac{1}{T} \int_0^T \cos(A\sqrt{t} \cdot \xi(s)) ds \right] \approx E(\cos(K \xi(s)))$
 $\cos x = \frac{e^{ix} + e^{-ix}}{2}$ 展开 $\rightarrow K$.
 $\Rightarrow AC = \lim_{T \rightarrow \infty} \left[\frac{1}{T} \int_0^T \frac{1}{2} [e^{iK \xi(s)} + e^{-iK \xi(s)}] ds \right]$
 高维特征函数, $X = \xi(s)$, $\varphi_X(m) = E(e^{imX}) = e^{-\frac{m^2 \sigma^2}{2}}$ (σ^2)
 $Y = K \xi(s) \Rightarrow \varphi_Y(m) = e^{-\frac{K^2 \sigma^2}{2}}$. 代入 AC
 $\Rightarrow AC = \frac{1}{2} [E(e^{iK \xi(s)}) + E(e^{-iK \xi(s)})] = \frac{1}{2} [\varphi_Y(\frac{1}{\sqrt{t}}) + \varphi_Y(-\frac{1}{\sqrt{t}})] = e^{-\frac{K^2 \sigma^2}{2}}$
 $K^2 = A^2 t = 2DRt$, if $\sigma^2 = 1 \Rightarrow AC = e^{-DRt}$ ($t > 0$)
 if $DR = \frac{1}{\tau_R} \Rightarrow AC = e^{-\frac{t}{\tau_R}}$ ($t > 0$)

Simulation result:



Find τ and compare the relationship between τ and D_R :

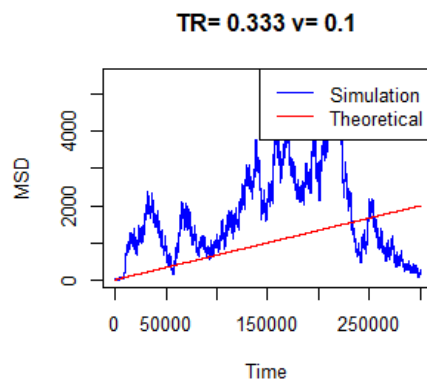
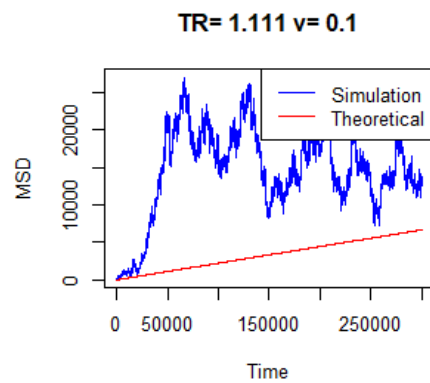
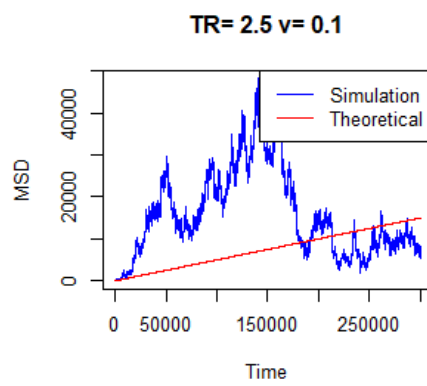
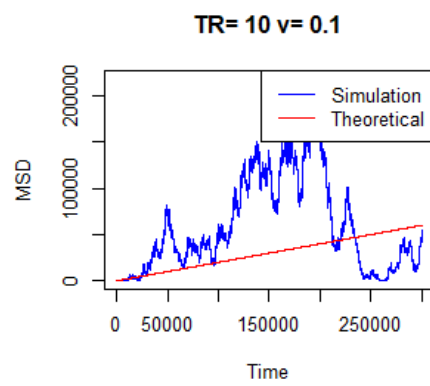
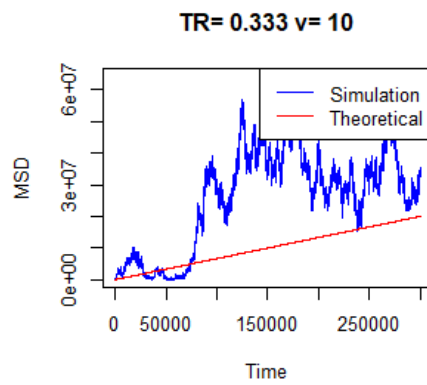
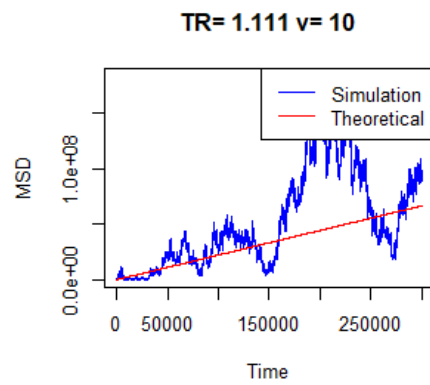
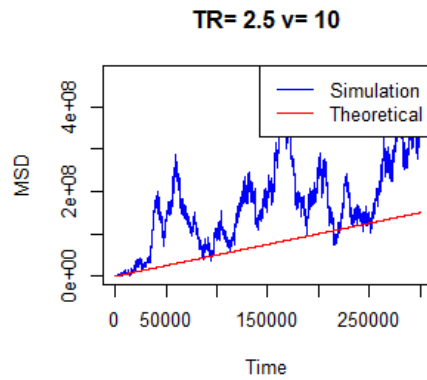
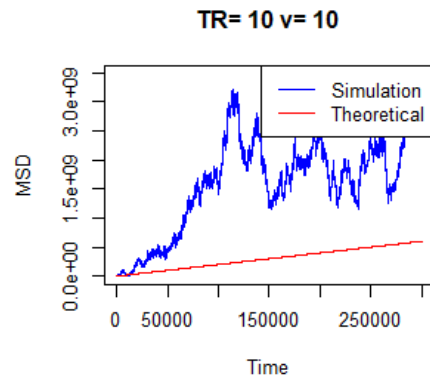


4.

This equation shows that when t is much larger than τ_R , the mean square displacement is proportional to time t . The growth rate of $MSD(t)$ is linear and proportional to time t , as a particle moves with a constant velocity v over a longer time scale, its MSD also increases in a linear manner. The coefficient represents the magnitude of the displacement. When t is much larger than τ_R , this product can be approximately regarded as a constant.

5.

Verify MSD:



Appendix

#3

```
DRlist=c(0.2,0.5,1,3)
```

```
tlist=seq(0,8,0.02)
```

```
integration=function(DR,t,T=10000){  
  set.seed(1)  
  samples=rnorm(n=T,0,1)  
  integrand=cos(sqrt(DR*2)*sqrt(t)*samples)  
  integral=sum(integrand)/T  
  return(integral)  
}
```

```
par(mfrow=c(2,2))
```

```
T=100000
```

```
result=c()
```

```
for (DR in DRlist){  
  for (i in 1:length(tlist)){  
    result[i]=integration(DR,tlist[i],T)  
  }  
}
```

```
plot(tlist,result,type="p",col="black",main=paste("DR=",DR))  
curve(exp(-DR*x),from=0,to=4,add=TRUE,col="red",lwd=1.5)  
legend("topright", legend = c("simulation", "exp(-DR*t)"), col = c("black", "red"), lty = 1)  
}
```

```
par(mfrow=c(1,1))
```

#求 Tau 并对比 Tau 和 DR 的关系

```
model=function(t, Tau) exp(-t / Tau)
```

```
init_params=list(Tau=1)
```

```
DRlist2=c(0.1,0.3,0.5,0.7,0.9,1,1.2,1.3,1.5,2,3,5,8)
```

```
Taulist=c()
```

```
for (DR in DRlist2){  
  for (i in 1:length(tlist)){  
    result[i]=integration(DR,tlist[i],T)  
  }  
}
```

```
fit=nls(result ~ model(tlist,Tau), start =init_params)
```

```

    Tau=coef(fit)["Tau"]
    Taulist=c(Taulist,Tau)
}

plot(DRlist2,Taulist,type="b",col="black",lwd=2,main="DR-Tau Relationship")
curve(1/x,from=0,to=10,add=TRUE,col="red",lwd=1)
legend("topright", legend = c("DR-Tau-simulation", "1/x"), col = c("black", "red"), lty = 1)

```

```

#5
par(mfrow=c(2,2))
# 设置参数值
vlist=c(0.1,10) # 速度
DRlist3=c(0.1,0.4,0.9,3) # 扩散系数

# v=2 # 速度
# DR=0.5 # 扩散系数

# 设置模拟参数
nsteps=1000000 # 模拟步数
dt=0.3 # 时间步长

# 初始化变量
x=c(0) # 初始位置 x
y=c(0) # 初始位置 y
phi=c(0) # 初始角度  $\varphi$ 

# 模拟运动轨迹
for (v in vlist) {
  for (DR in DRlist3) {
    TR=1/DR
    set.seed(100*v)
    for (i in 2:nsteps) {
      xi=rnorm(1) # 生成高斯白噪声
      dx_dt=v*cos(phi[i-1]) # 计算 x 轴速度
      dy_dt=v*sin(phi[i-1]) # 计算 y 轴速度
      dphi_dt=sqrt(2*DR)*xi # 计算角速度

      x[i]=x[i-1]+dx_dt*dt # 更新 x 坐标
      y[i]=y[i-1]+dy_dt*dt # 更新 y 坐标
      phi[i]=phi[i-1]+dphi_dt*dt # 更新角度
    }

    # 计算模拟结果的 MSD

```

```

t=dt*(1:nsteps) # 时间向量
MSD=x^2+y^2 # MSD = x^2 + y^2

# 计算理论结果的 MSD
MSD_theoretical=(2*v^2/DR)*t

# 绘制模拟结果和理论结果
plot(t, MSD, type="l", col="blue", xlab="Time", ylab="MSD", main=paste("TR=",round(TR,
3),"v=",v))
lines(t, MSD_theoretical, type="l", col="red")
legend("topright", legend=c("Simulation", "Theoretical"), col=c("blue", "red"), lty=1)
}
}

```