

Density matrix or Density operator

Pure state : $\rho = |\psi\rangle\langle\psi|$

$$\langle A \rangle = \text{Tr}(\rho A)$$

General definition

$$\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$$

($|\psi_1\rangle, |\psi_2\rangle$)

$$\rho = p_1 |\psi_1\rangle\langle\psi_1| + p_2 |\psi_2\rangle\langle\psi_2|$$

$$p_1 + p_2 = 1$$

Matrix representation of ρ

$$\rho_{mn} = \langle m | \rho | n \rangle$$

$$\rho = |\psi\rangle\langle\psi|$$

$$\rho = \begin{pmatrix} \rho_{11} & \rho_{12} \\ \rho_{21} & \rho_{22} \end{pmatrix}$$

$$\rho_{12} = \langle 1 | \rho | 2 \rangle$$

$$\rho_{mn} = \langle m | \psi \rangle \langle \psi | n \rangle$$

$$\rho_{mm} = \langle m | \psi \rangle \langle \psi | m \rangle$$

$$= \underline{\underline{|\langle m | \psi \rangle|^2}}$$

$$\langle m | \psi \rangle$$

and $\langle \psi | m \rangle$
are complex
conjugates

$\{|m\rangle\}$ basis states

$$|\psi\rangle = \sum c_m |m\rangle$$

$$\langle m | \psi \rangle = c_m$$

$$\rho_{mm} = |c_m|^2$$

probability of find $|\psi\rangle$
in the state $|m\rangle$

Diagonal elements in the density matrix gives the probability of finding the state $|\psi\rangle$ in one the eigen states.

Off-diagonal elements $\underline{\underline{|m\rangle}}$ $m \neq n$

$$\rho_{mn} = \langle m | \rho | n \rangle = \langle m | \psi \rangle \langle \psi | n \rangle = c_m c_n^*$$

$$c_m = |c_m| e^{i\phi_m}$$

$$c_n = |c_n| e^{i\phi_n}$$

$$\rho_{mn} = |c_m| |c_n| e^{i(\phi_m - \phi_n)}$$

\Rightarrow off diagonal elements depend on the relative phases between the states $|m\rangle$ and $|n\rangle$

responsible for the phenomena of
quantum interference
coherence terms

$$|\psi\rangle = \frac{1}{\sqrt{2}} |0\rangle - \frac{i}{\sqrt{2}} |1\rangle \quad (\text{pure state})$$

$$= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

$$\rho = |\psi\rangle \langle \psi| = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & i \end{pmatrix}$$

$$= \frac{1}{2} \begin{pmatrix} 1 & i \\ -i & 1 \end{pmatrix}$$

$$\text{Tr}(\rho) = \frac{1}{2} + \frac{1}{2} = 1$$

$\boxed{\rho = \rho^+}$

$$\rho^+ = \left[\begin{matrix} \frac{1}{2} & \left(\begin{matrix} 1 & i \\ -i & 1 \end{matrix} \right) \\ \left(\begin{matrix} 1 & i \\ -i & 1 \end{matrix} \right)^T & \frac{1}{2} \end{matrix} \right]^C$$

$$= \frac{1}{2} \left(\begin{matrix} 1 & +i \\ -i & 1 \end{matrix} \right)$$

$$= \rho$$

Density operator
satisfies Liouville equation

$$i\hbar \frac{\partial \rho}{\partial t} = [H, \rho]$$

Mixed States

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

$$\underline{\underline{\text{Tr}(\rho) = 1}} \quad \underline{\underline{\text{Prove that}}}$$

1. $\text{Tr}(\rho) = 1$ (Both for pure and mixed states)

2.

$\text{Tr}(\rho^2) < 1$ for mixed state

Proof

1. (a) Pure state:

$$\rho = |\psi\rangle \langle \psi|$$

$$|\psi\rangle = \sum c_m |m\rangle$$

$$\begin{aligned}
 \text{Tr}(\rho) &= \sum_n \langle n | \rho | n \rangle \\
 &= \sum_n \langle n | \psi \rangle \langle \psi | n \rangle \\
 &= \sum_n |\langle n | \psi \rangle|^2 \\
 &= \sum_n |\langle n | m \rangle|^2 |c_m|^2 \quad \text{($\psi = \sum c_m |m\rangle$)} \\
 &= \sum_n |c_n|^2 \\
 &= 1
 \end{aligned}$$

(6) Mixed state

$$\begin{aligned}
 \rho &= \sum_i p_i |\psi_i\rangle \langle \psi_i| \\
 \text{Tr}(\rho) &= \sum_i \sum_m p_i \langle m | \psi_i \rangle \langle \psi_i | m \rangle \\
 &= \sum_i \sum_m p_i |\langle m | \psi_i \rangle|^2 \quad \boxed{|\psi_i\rangle = \sum_n c_n |n\rangle} \\
 &= \sum_i \sum_m \sum_n p_i |\langle m | n \rangle|^2 |c_m|^2 \\
 &= \sum_i \sum_n p_i |c_n|^2 \\
 &= \sum_i p_i \sum_n |c_n|^2 \\
 &= \sum_i p_i \quad \therefore \sum_n |c_n|^2 = 1 \\
 &= 1
 \end{aligned}$$

$$2. \quad \rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$$

$$\rho^2 = \sum_i \sum_j p_i p_j |\psi_i\rangle\langle\psi_i| |\psi_j\rangle\langle\psi_j|$$

$$\text{Tr}(\rho^2) = \sum_n \langle n | \rho^2 | n \rangle$$

$$= \sum_n \sum_i \sum_j p_i p_j \langle n | \psi_i \rangle \langle \psi_i | \psi_j \rangle \langle \psi_j | n \rangle$$

$$= \sum_n \sum_i \sum_j p_i p_j \langle \psi_i | \psi_j \rangle \langle \psi_j | n \rangle \langle n | \psi_i \rangle$$

$$= \sum_i \sum_j p_i p_j \langle \psi_i | \psi_j \rangle \langle \psi_j | \underbrace{\left(\sum_n |n\rangle \langle n| \right)}_{I} |\psi_i \rangle$$

$$= \sum_i \sum_j p_i p_j \langle \psi_i | \psi_j \rangle \langle \psi_j | \psi_i \rangle$$

$$= \sum_i \sum_j p_i p_j |\langle \psi_i | \psi_j \rangle|^2$$

$$\leq \left[\sum_i p_i \right]^2$$

$$\leq 1$$

The equality holds only if $|\langle \psi_i | \psi_j \rangle|^2 = 1$
for every pair of states $|\psi_i\rangle$ and $|\psi_j\rangle$.

Thus

$$\text{Tr } \rho^2 = 1 \quad \text{for pure state}$$

$$\text{Tr } \rho^2 < 1 \quad \text{for mixed state}$$

A mixed state cannot be represented by a state vector but it is described by density operator

① Consider an ensemble 50% is in state $|0\rangle^{\otimes \infty}$ and 50% is in state $\frac{|0\rangle + |1\rangle}{\sqrt{2}}^{\otimes \infty}$

$$\rho = \sum p_i |\psi_i\rangle \langle \psi_i|$$

$$\begin{aligned} \rho &= \frac{1}{2} |0\rangle \langle 0| + \frac{1}{2} \left(\frac{|0\rangle + |1\rangle}{\sqrt{2}} \right) \otimes \left(\frac{|0\rangle + |1\rangle}{\sqrt{2}} \right) \\ &= \frac{1}{2} |0\rangle \langle 0| + \frac{1}{4} \left[|0\rangle \langle 0| + |0\rangle \langle 1| + |1\rangle \langle 0| + |1\rangle \langle 1| \right] \end{aligned}$$

$$= \underbrace{\frac{3}{4} |0\rangle \langle 0|}_{\text{Partially mixed}} + \underbrace{\frac{1}{4} |1\rangle \langle 1|}_{\text{Partially mixed}} + \underbrace{\frac{1}{4} \left[|0\rangle \langle 1| + |1\rangle \langle 0| \right]}_{\text{Partially mixed}}$$

$$\rho = \begin{pmatrix} 3/4 & 1/4 \\ 1/4 & 1/4 \end{pmatrix} \quad \text{(Partially mixed)}$$

1 }
2 }