# Wilcox's k- $\omega$ model

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#### 1 Wilcox's k- $\omega$ model

I have used the Finite Volume Method (FVM) approach. A non-uniform grid is generated using the DNS grid provided. For the given problem, full channel height (i.e. upto  $y_{max} = 2\delta$ ) is considered.

Initial guesses are:

 $k = 5 \times 10^{-3}$ ,  $\omega = 10^{-5}$ , u = randomly initialised between [0, 1] values

Boundary conditions applied are:

At wall: u = 0, k = 0,  $\omega_{y^{+} \le 3} = 6\nu/C_2 y^2$ 

The model constants are:

$$\beta^* = 0.09, \ \sigma_k = 2, \ \sigma_\omega = 2, \ C_1 = 5/9, \ C_2 = 3/40$$

Considering a 1D fully developed turbulent channel flow, the governing equations and its discretization is given below.

# 1.1 U equation

$$0 = -\frac{1}{\rho} \frac{\partial P}{\partial x} + \frac{\partial}{\partial y} \left[ (\nu + \nu_t) \frac{\partial U}{\partial y} \right]$$

Discretized equation is:

$$a_P U_P = a_N U_N + a_S U_S + b_U$$
 where  $a_N = \frac{(\nu + \nu_t)_n}{\delta y_n}$ ,  $a_S = \frac{(\nu + \nu_t)_s}{\delta y_s}$ , 
$$a_P = a_N + a_S, \ b_U = -\frac{1}{\rho} \frac{\partial P}{\partial x} \Delta y_P$$

# 1.2 k equation

$$0 = \frac{\partial}{\partial y} \left[ \left( \nu + \frac{\nu_t}{\sigma_k} \right) \frac{\partial k}{\partial y} \right] + P_k - \beta^* \omega k$$
where  $P_k = \nu_t \left( \frac{\partial U}{\partial y} \right)^2$ 

Discretized equation is:

$$a_P k_P = a_N k_N + a_S k_S + b_k$$

where 
$$a_N = \frac{\left(\nu + \frac{\nu_t}{\sigma_k}\right)_n}{\delta y_n}$$
,  $a_S = \frac{\left(\nu + \frac{\nu_t}{\sigma_k}\right)_s}{\delta y_s}$ ,  $a_P = a_N + a_S - S_P \Delta y_P$ ,  $S_P = -\beta^* \omega$ ,  $b_k = P_k \Delta y_P$ 

## 1.3 $\omega$ equation

$$0 = \frac{\partial}{\partial y} \left[ \left( \nu + \frac{\nu_t}{\sigma_\omega} \right) \frac{\partial \omega}{\partial y} \right] + \frac{\omega}{k} C_1 P_k - C_2 \omega^2$$

Discretized equation is:

$$a_P \omega_P = a_N \omega_N + a_S \omega_S + b_\omega$$
where  $a_N = \frac{\left(\nu + \frac{\nu_t}{\sigma_\omega}\right)_n}{\delta y_n}$ ,  $a_S = \frac{\left(\nu + \frac{\nu_t}{\sigma_\omega}\right)_s}{\delta y_s}$ ,
$$a_P = a_N + a_S - S_P \Delta y_P, \ S_P = -C_2 \omega^*, \ b_\omega = \frac{\omega^*}{k} C_1 P_k \Delta y_P$$

## 1.4 $\nu_t$ expression

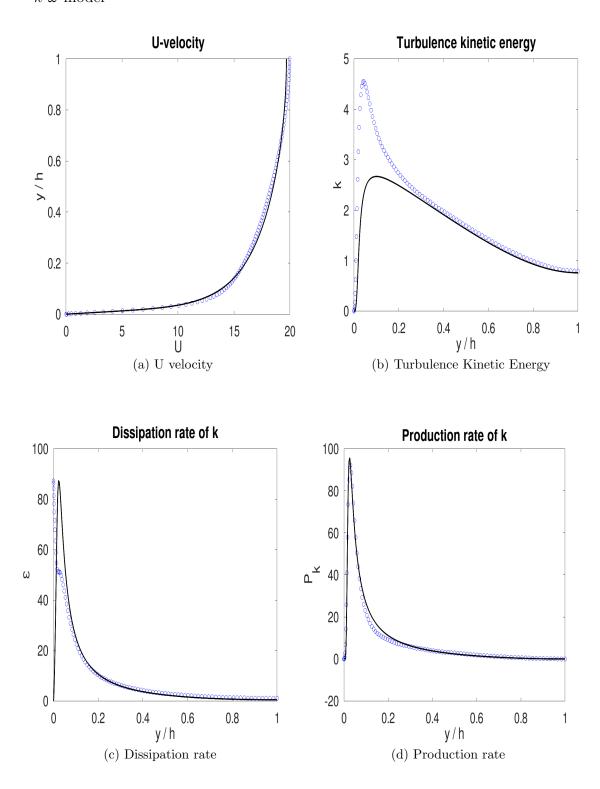
$$\nu_t = \begin{cases} l_m^2 \left| \frac{\partial U}{\partial y} \right| & \text{, for first 300 iterations} \\ \frac{\underline{k}}{\omega} & \text{, otherwise} \end{cases}$$

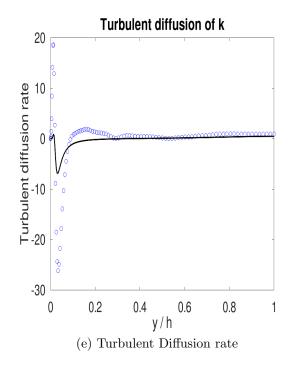
where 
$$l_m = \delta \left[ 0.14 - 0.08 \left( 1 - \frac{y}{\delta} \right)^2 - 0.06 \left( 1 - \frac{y}{\delta} \right)^4 \right]$$
,  $\delta = \text{channel half width}$ 

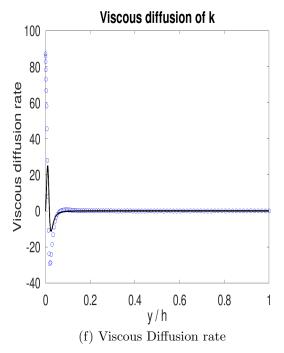
# 2 Results and Discussions

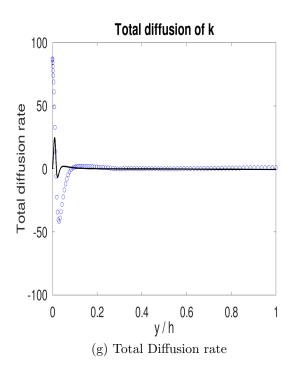
For all the plots below, the following legend style is adopted:

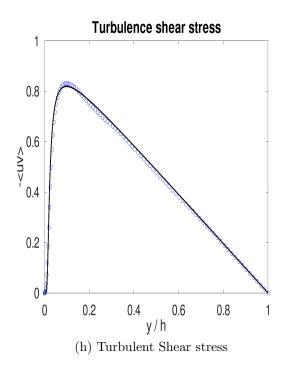
- o DNS data
- $-k-\omega$  model











## Wall Shear Stress

Analytical:  $\tau_w = 1.0$ 

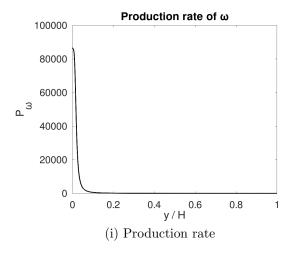
In k- $\omega$  model:  $\tau_w = 0.9999997$ 

#### **Near-Wall Region**

The k- $\omega$  model has produced a good result for U,  $P_k$ ,  $-\langle uv \rangle$ . However if we observe the dissipation rate of k (i.e.  $\varepsilon_k$ ), the model has calculated it as 0 while as per the DNS data, it has some maximum value at the wall (around 87). This is because the model calculates the  $\varepsilon$  using the expression:

$$\varepsilon = \omega \beta^* k$$

So at the wall, k=0 and thus, the model will calculate the  $\varepsilon$  as 0 at the wall. In the near-wall region, the turbulent diffusion rate and the viscous diffusion rate has not been accurately captured by the k- $\omega$  model. This has led to the overall under-prediction of TKE near the wall. We also neglected the pressure diffusion rate in the k modelled equation. As TKE has various transport rate terms (i.e. pressure diffusion rate, turbulent diffusion rate, and viscous diffusion rate), the under-predicted values of TKE has led to under-prediction of viscous diffusion rate ok k even though there is no modelling assumption made for this term.



From figure (i), the decoupling of k from  $\omega$  is observed as we move away from the wall. This also implies that the k- $\omega$  model performs better in near-wall regions. Away from the wall, the model does not perform satisfactorily (refer figure (a), (b)) because of the decoupling phenomenon of k and  $\omega$ . The need for implementing Low Reynolds Number formulation is eliminated which is not the case in the standard k- $\varepsilon$  model. However, implementing k- $\omega$  SST model (proposed by Menter, 1994) could improve the result in the regions away from the wall.