### A Tale of a Boundary-Layer Problem

### Martin Kunz

Faculty of Nuclear Sciences and Physical Engineering, Czech Technical University in Prague

May 5, 2023

### **Abstract**

I will demonstrate the solution with all its steps (and missteps) to the boundary layer problem

$$\varepsilon y'' = y \cdot y' - y,\tag{1}$$

such that

$$y(0)=1$$
 and  $y(1)=-1$  and where  $\varepsilon\ll 1.$  (2)

The journey towards the solution will take us on an excurse through some basics of *perturbation theory* and *asymptotic approximation matching*. I will demonstrate visualizations that will give you insight into the reasoning behind the steps taken towards the solution serving as an exhibition for the leverage in problem solving accessible to us through the computational power at our fingertips.

### Outline

- 1 Motivation
- 2 Theory
- 3 Solution of the DE
- 4 Interactive Solution

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Have we encountered other non-linear DEs?



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Newtons gravitational equation



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How did we manage to solve these equations?

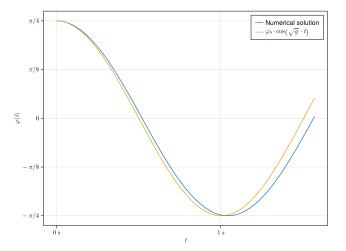


Figure: Comparison of the numerical solution and the small angle approximation of the simple pendulum DE for the initial angle  $\varphi(0)=\pi/4$  and  $\dot{\varphi}(0)=0$ .

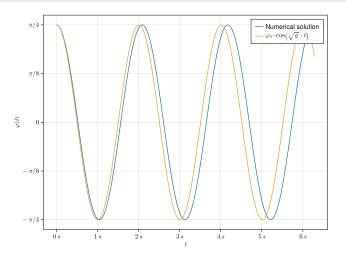


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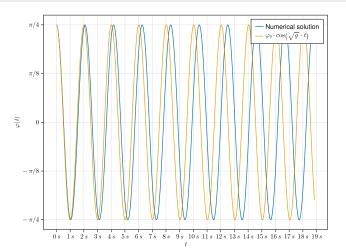


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ordinary single unknown function second-order contains a second order derivative first-degree highest degree of the unknown function is 1 homogeneous does not contain a free term in the x variable non-linear contains a mixed term (boundary-value) constrained by values on the boundaries singularly-perturbed contains one small term

How good is our numerical solution?

### Note (Analytical solution - Salas et al. 2014)

There is an analytical solution to the pendulum DE in the form

$$\varphi(t) = 2 \cdot \arcsin \left( \sin(\varphi_0/2) \cdot \operatorname{cd}(t \cdot \sqrt{g/L}, \, \sin^2(\varphi_0/2)) \right), \text{ where }$$

$$\operatorname{cd}(u, m) = \operatorname{cn}(u, m) / \operatorname{dn}(u, m)$$

$$\operatorname{cn}(u, m) = \cos(\operatorname{am}(u, m)) \quad \operatorname{dn}(u, m) = \frac{\operatorname{d}}{\operatorname{d}u}(\operatorname{am}(u, m))$$

and  $am(u, m) = \varphi$  is the Jacobi amplitude which is the upper-bound of the integral such that

$$u = \int_0^{\varphi} \frac{\mathrm{d}\vartheta}{\sqrt{1 - m\sin^2(\vartheta)}}.$$

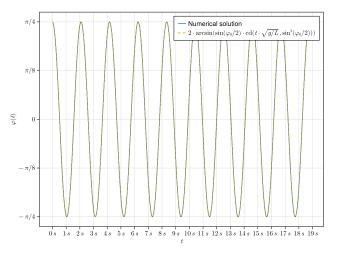


Figure: Comparison of the analytical solution and the numerical solution of the simple pendulum DE for the initial angle  $\varphi(0)=\pi/4$  and  $\dot{\varphi}(0)=0$ .

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- We can prove some good qualities about the solution (uniform convergence, etc.)
- Numerical solution is not guaranteed to be correct
- Size of the latent space for the solution! (we cannot search the whole space and find all the solutions numerically)

Numerical Methods

How precise is a perturbation solution?

Numerical Methods

#### Note (Perturbation solution - Beléndez et al. 2006)

There is an approximation to the simple pendulum problem based on a homotopy perturbation method. It is uniformly convergent and is precise from the first few terms of the expansion.

#### Numerical Methods

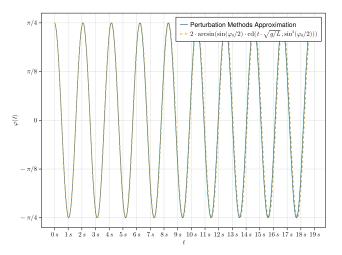


Figure: Comparison of the perturbation solution and the analytical solution of the simple pendulum DE for the initial angle  $\varphi(0)=\pi/4$  and  $\dot{\varphi}(0)=0$ .

- 2 Theory
  - Mathematical Basics
  - Perturbation Theory
  - Boundary Layers

Mathematical Basics

#### Theorem (Taylor's theorem)

Given a function  $f(\varepsilon)$ , suppose its (n+1)st derivative  $f^{(n+1)}$  is continuous for  $\varepsilon_a < \varepsilon < \varepsilon_b$ . Suppose that the points  $\varepsilon, \varepsilon_0 \in (\varepsilon_a, \varepsilon_b)$ . Then

$$f(\varepsilon) = f(\varepsilon_0) + (\varepsilon - \varepsilon_0)f'(\varepsilon_0) + \ldots + \frac{1}{n!}(\varepsilon - \varepsilon_0)^n f^{(n)}(\varepsilon_0) + R_{n+1},$$

where the error term  $R_{n+1}$  is given by

$$R_{n+1} = \frac{1}{(n+1)!} (\varepsilon - \varepsilon_0)^{n+1} f^{(n+1)}(\xi),$$

and  $\xi \in ]\varepsilon, \varepsilon_0[$ 

Mathematical Basics

Theory

### Definition (Bachmann-Landau Symbols - Holmes 2012)

**1**  $f = O(\varphi)$  as  $\varepsilon \to \varepsilon_0^+$  if there is a constant k and  $\tilde{\varepsilon}$  so that

$$|f(\varepsilon)| \le k |\varphi(\varepsilon)|$$
 for  $\varepsilon_0 < \varepsilon < \tilde{\varepsilon}$ 

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 $f = o(\varphi)$  as  $\varepsilon \to \varepsilon_0^+$  if  $\forall \delta > 0$  there is a  $\tilde{\varepsilon}$  so that

$$|f(\varepsilon)| \le \delta |\varphi(\varepsilon)|$$
 for  $\varepsilon_0 < \varepsilon < \tilde{\varepsilon}$ 

#### Definition (Asymptotic approximation - Holmes 2012)

Given a function  $f(\varepsilon)$  and  $\varphi(\varepsilon)$  we say that  $\varphi(\varepsilon)$  is an asymptotic approximation to  $f(\varepsilon)$  as  $\varepsilon \to \varepsilon_0^+$  whenever  $f = \varphi + o(\varphi)$  as  $\varepsilon \to \varepsilon_0^+$ . We write  $f \sim \varphi$  as  $\varepsilon \to \varepsilon_0^+$ .

#### Definition (Asymptotic expansion - Holmes 2012)

1 The sequence of functions  $\varphi_1(\varepsilon), \varphi_2(\varepsilon), \ldots$  form an asymptotic sequence (are well ordered) as  $\varepsilon \to \varepsilon_0^+$  if  $\varphi_{m+1} = o(\varphi_m)$  as  $\varepsilon \to \varepsilon_0^+$  for  $\forall n$ .

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Theory

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- 2 Let  $\varphi_1(\varepsilon), \varphi_2(\varepsilon), \ldots$  be basis functions. Then  $f(\varepsilon)$  has an asymptotic expansion to n terms for this sequence if

$$f = \sum_{k=1}^{m} a_k \varphi_k + o(\varphi_m)$$
 for  $m \in \hat{n}$  as  $\varepsilon \to \varepsilon_0^+$ ,

where  $a_k$  are independent of  $\varepsilon$ . We can write

$$f \sim a_1 \varphi_1(\varepsilon) + a_2 \varphi_2(\varepsilon) + \ldots + a_n \varphi_n(\varepsilon)$$
 as  $\varepsilon \to \varepsilon_0^+$ .

## Asymptotic Solution of the Projectile Problem

$$\frac{\mathrm{d}^2 x}{\mathrm{d}t^2} = -\frac{g \cdot R^2}{(x+R)^2} \quad \text{for} \quad 0 < t.$$

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$$\tau = t \cdot g/v_0$$
$$y(\tau) = x(t) \cdot g/v_0^2$$

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Doing so transforms the problem into

$$\frac{\mathrm{d}^2 y}{\mathrm{d}\tau^2} = -\frac{1}{(1+\varepsilon y)^2} \quad \text{for} \quad 0 < \tau,$$

where

$$y(0) = 0$$
 and  $y'(0) = 1$  and  $\varepsilon = v_0^2/Rg$ .

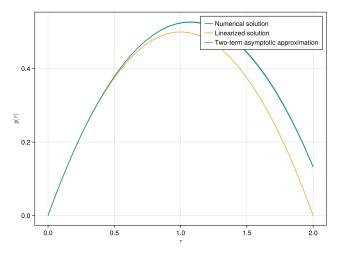


Figure: Comparison of the two term approximation, the linearised problem (one-term approximation) and the numerical solution for  $\varepsilon=1/10$ 

Boundary Layers

## Boundary Layer

A boundary layer in lay terms is an interval of the solution of a DE where the derivative is large compared with surrounding times.

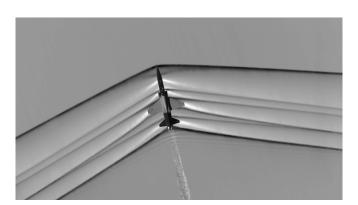


Figure: Shlieren image of the air waves of a supersonic airflow produced by a fighter jet [Source: NASA (2015)]

- 3 Solution of the DE
  - Locating the Boundary Layers
  - Symmetry of the Unknown
  - Phase Portrait
  - Perturbation Theory
    - Outer Solution
  - Inner Solution

Locating the Boundary Layers

# Plausibility Argument of no Border Boundary Layers - Holmes 2012

If we are looking for a boundary layer at x=0, then y'<0 and if we assume that y''>0, then  $\varepsilon y''>0$ , but  $y\cdot (y'-1)$  on the left hand side is both >0 on some interval and <0. Therefore we can rule out a boundary layers at the interval borders for y''>0.

Symmetry of the Unknown

Is there any symmetry in the DE?

$$(x,y) \rightarrow (1-x,-y)$$

Phase Portrait

#### Phase Portrait of the DE

$$\varepsilon y'' = y \cdot y' - y$$

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## Phase Portrait of the DE

$$y'(t) = z(t)$$
$$z'(t) = \frac{1}{\varepsilon} \cdot y(t) \cdot (z(t) - 1)$$

Phase Portrait

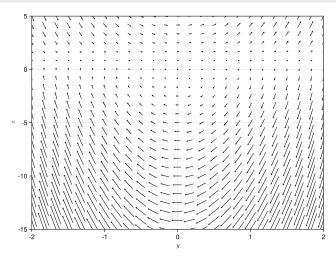


Figure: Arrows representing the vector field of our system of differential equations

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Phase Portrait

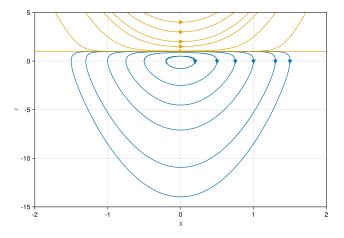


Figure: Trajectories of imaginary particles in the phase portrait of our system of differential equations for  $\varepsilon=0.1$ 

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Phase Portrait

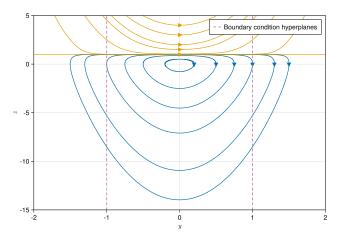


Figure: Trajectories of imaginary particles in the phase portrait of our system of differential equations for  $\varepsilon=1$  with the boundary condition hyperplanes

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Phase Portrait

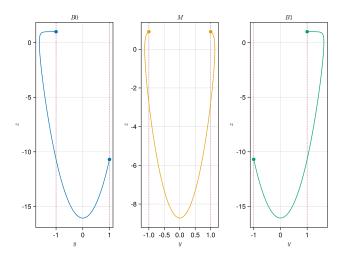


Figure: Solutions  $B_0$ , M and  $B_1$  in the phase plane of the DE

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$$y_0(x) = x + a$$
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$$y_0(x) = x + a$$
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$$y_0^L(x) = 1 + x$$
 and  $y_0^R(x) = x - 2$ 

Inner Solution

## Scaling

For the boundaries, we have to scale the solutions as

$$X \equiv \frac{x - x_0}{\delta},$$

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If  $\delta$  is chosen correctly, then Y and all its derivatives should be O(1) as  $\varepsilon \to 0^+$  and we can find the limit when  $\varepsilon/\delta^2 = 1/\delta$  or  $\delta = \varepsilon$ . The inner equation layer is

$$\frac{\mathrm{d}^2 Y}{\mathrm{d}X^2} = Y \frac{\mathrm{d}Y}{\mathrm{d}X} - \varepsilon Y.$$

#### Inner Solutions

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When we create these composite solutions, we get

$$\tilde{y}^{M} = x - \frac{1}{2} - \frac{3}{2} \cdot \tanh\left[\frac{3}{4\varepsilon}(x - \frac{1}{2})\right] + O(\varepsilon)$$

$$\tilde{y}^{B0} = x - 2 \cdot \tanh\left[\frac{x}{\varepsilon} - \tanh^{-1}(1/2)\right] + O(\varepsilon)$$

$$\tilde{y}^{B1} = x - 1 - 2 \cdot \tanh\left[\frac{x - 1}{\varepsilon} + \tanh^{-1}(1/2)\right] + O(\varepsilon)$$

4 Interactive Solution

#### References

- Beléndez, Augusto et al. (2006). "Application of the homotopy perturbation method to the non-linear pendulum". In: *European Journal of Physics* 28.1, p. 93.
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- Salas, Alvaro H et al. (2014). "Exact solution to Duffing equation and the pendulum equation". In: *Applied Mathematical Sciences* 8.176, pp. 8781–8789.