

# QUANTUM ERGODICITY FOR POINT SCATTERERS ON ARITHMETIC TORI

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ABSTRACT. We prove an analogue of Shnirelman, Zelditch and Colin de Verdière's Quantum Ergodicity Theorems in a case where there is *no* underlying classical ergodicity. The system we consider is the Laplacian with a delta potential on the square torus. There are two types of wave functions: old eigenfunctions of the Laplacian, which are not affected by the scatterer, and new eigenfunctions which have a logarithmic singularity at the position of the scatterer. We prove that a full density subsequence of the new eigenfunctions *equidistribute in phase space*. Our estimates are uniform with respect to the coupling parameter, in particular the equidistribution holds for both the weak and strong coupling quantizations of the point scatterer.

## 1. INTRODUCTION

The point scatterer, namely the Laplacian with a delta potential, on a two-dimensional flat manifold is a popular model in the study of the transition between chaos and integrability in quantum systems. In 1990 Seba [14] considered this operator on a rectangle with irrational aspect ratio and Dirichlet boundary conditions and argued that the spectrum and eigenfunctions of the point scatterer display features such as level repulsion and a Gaussian value distribution, both of which are present in quantum systems with chaotic classical dynamics (cf. [4] and [2]), such as the quantization of the geodesic flow on hyperbolic manifolds or the flow in the Sinai billiard. In fact the point scatterer can be understood as a limit of the Sinai billiard where the radius shrinks to zero faster than the semiclassical wavelength.

The subject of this paper is a point scatterer on a flat torus. It has two types of eigenfunctions: first there are *old* eigenfunctions of the Laplacian, namely those which vanish at the position of the scatterer; the nonzero eigenvalues remain the same, though with multiplicities reduced by 1. Secondly, there are *new* eigenfunctions which diverge logarithmically near the position of the scatterer; the corresponding eigenvalues have multiplicity 1 and interlace with the old Laplace eigenvalues.

We shall only be concerned with the set of new eigenfunctions, i.e., the ones which are affected by the scatterer. In [11] it was proved that a full density subsequence<sup>1</sup> of the new eigenfunctions *equidistribute in position space* in the special case of a square torus. We extend the results of [11] and prove that a full density subsequence of these eigenfunctions in fact *equidistribute in phase space* — we thus establish an analogue of Shnirelman, Zelditch and Colin de Verdière's Quantum Ergodicity

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<sup>1</sup>See Section 3 for a precise definition of a “full density subsequence”.

Theorem in a case where there is no underlying chaotic dynamics and no classical ergodicity.

An analogue of this result for a cubic 3D torus was recently obtained by N. Yesha [18]. The situation for a square torus is very different from irrational tori, where the eigenfunctions are expected to localise in phase space on a finite number of momentum vectors [7, 1]. In the case where the aspect ratio is diophantine this can be proven rigorously for a full density subsequence of new eigenfunctions [8].

### 1.1. Spectrum of the point scatterer.

The formal operator

$$-\Delta + \alpha\delta_{x_0}, \quad \alpha \in \mathbb{R}$$

is realized using von Neumann's theory of self-adjoint extensions. We simply state the most important facts in this section in order to formulate the results of this paper. For a more detailed discussion of the self-adjoint realization of the point scatterer we refer the reader to the introduction and appendix of the paper [11].

Let  $\mathbb{T}^2 = \mathbb{R}^2 / 2\pi\mathbb{Z}^2$ . We consider the restriction of the positive Laplacian  $-\Delta$  to the domain

$$D_0 = C_c^\infty(\mathbb{T}^2 \setminus \{x_0\})$$

of functions which vanish near the position of the scatterer:

$$H = -\Delta|_{D_0}$$

The operator  $H$  is symmetric, but fails to be self-adjoint, in fact  $H$  has deficiency indices  $(1, 1)$ . Self-adjoint extension theory tells us that there exists a one-parameter family of self-adjoint extensions  $H_\varphi$ ,  $\varphi \in (-\pi, \pi]$ , which are restrictions of the adjoint  $H^*$  to the domain of functions  $f \in \text{Dom}(H^*)$  which satisfy the asymptotic

$$f(x) = C \left( \cos\left(\frac{\varphi}{2}\right) \frac{\log|x-x_0|}{2\pi} + \sin\left(\frac{\varphi}{2}\right) \right) + o(1), \quad x \rightarrow x_0$$

for some constant  $C \in \mathbb{C}$ . The case  $\varphi = \pi$  corresponds to  $\alpha = 0$ . In this paper we will study the operators  $H_\varphi$ ,  $\varphi \in (-\pi, \pi)$ .

The spectrum of the operator  $H_\varphi$  consists of two parts: “old” and “new” eigenvalues. Since  $H_\varphi$  is a self-adjoint realization of a rank one perturbation of the Laplacian, the effect is that each nonzero old Laplace eigenvalue appears, with multiplicity reduced by 1, in the spectrum of  $H_\varphi$ . Further, each old Laplace eigenvalue gives rise to a new eigenvalue with multiplicity 1. In fact, these new eigenvalues *interlace* with the multiplicity one sequence associated with the old Laplace eigenvalues.

There are two types of eigenfunctions of  $H_\varphi$  associated with the two parts of the spectrum:

- (A) “Old” eigenfunctions which vanish at  $x_0$  and therefore are not affected by the scatterer. These are simply eigenfunctions of the unperturbed Laplacian.
- (B) “New” eigenfunctions which feature a logarithmic singularity at  $x_0$ ; in fact they are given by Green’s functions  $G_\lambda = (\Delta + \lambda)^{-1}\delta_{x_0}$ .

We will study how eigenfunctions of type (B) are distributed in phase space as the eigenvalue tends to infinity. Denote by  $S$  the set of distinct eigenvalues of the Laplacian on  $\mathbb{T}^2$ , namely integers which can be represented as a sum of two squares:

$$S := \{n \in \mathbb{Z} : n = x^2 + y^2 \mid x, y \in \mathbb{Z}\}$$

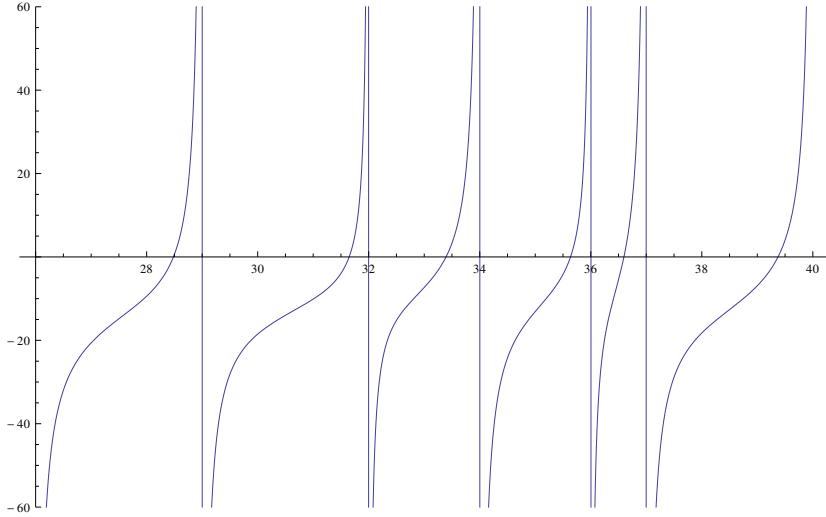


FIGURE 1. The picture shows a plot of the l. h. s. of equation (1.1) as a function of  $\lambda$ . The zeroes are the new eigenvalues corresponding to the self-adjoint extension with parameter  $\varphi = 0$ .

For given  $n \in S$  denote its multiplicity by

$$r_2(n) := \sum_{\substack{n=|\xi|^2 \\ \xi \in \mathbb{Z}^2}} 1,$$

i.e. the number of ways  $n$  can be written as a sum of two squares.

The eigenvalues of type (B) are solutions to the equation

$$(1.1) \quad \sum_{n \in S} r_2(n) \left( \frac{1}{n - \lambda} - \frac{n}{n^2 + 1} \right) = c_0 \tan \left( \frac{\varphi}{2} \right)$$

(see Figure 1 for a plot of the l. h. s.) where

$$(1.2) \quad c_0 = \sum_{n \in S} \frac{r_2(n)}{n^2 + 1}.$$

As mentioned earlier, they interlace with the distinct Laplace eigenvalues

$$S = \{0 < 1 < 2 < 4 < 5 < 8 < \dots\}$$

as follows

$$(1.3) \quad \lambda_{0,\varphi} < 0 < \lambda_{1,\varphi} < 1 < \lambda_{2,\varphi} < 2 < \lambda_{4,\varphi} < 4 < \lambda_{5,\varphi} < 5 < \lambda_{8,\varphi} < 8 < \dots$$

where the new eigenvalue associated with  $n \in S$  is denoted by  $\lambda_{n,\varphi}$ ; note that  $\lambda_{n,\varphi} < n$ .

**1.2. Strong coupling.** In the physics literature equation (1.1) is referred to as a “weak coupling” quantization. In fact, (cf. [12]) the new eigenvalues  $\lambda_{m,\varphi}$  “clump” with the Laplace eigenvalues  $m \in S$  in the sense that for a full density subsequence of  $S$ ,

$$0 < m - \lambda_{m,\varphi} \ll \frac{1}{(\log m)^{1-o(1)}}.$$

In particular the eigenvalue spacing distribution of the point scatterer coincides with that of the Laplacian, and the effect of the scatterer on the spectrum is quite weak in this quantization. (In a sense it corresponds to letting  $\alpha \rightarrow 0$  as  $\lambda \rightarrow \infty$ .)

Shigehara [15] and later Bogomolny, Gerland and Schmit [3], with the intent of finding a model exhibiting level repulsion, considered another quantization sometimes referred to as a “strong coupling” quantization. There are various ways to arrive at this quantization condition from equation (1.1). For example, one may truncate the summation outside an energy window of size  $O(\lambda^\delta)$  where  $\delta > 0$  is fixed, and the new eigenvalues of the strong coupling quantization are then defined to be the solutions to the equation

$$(1.4) \quad \sum_{\substack{n \in S \\ |n - n_+(\lambda)| \leq n_+(\lambda)^\delta}} r_2(n) \left( \frac{1}{n - \lambda} - \frac{n}{n^2 + 1} \right) = c_0 \tan \left( \frac{\varphi}{2} \right),$$

where  $n_+(\lambda)$  denotes the smallest element of  $S$  which is larger than  $\lambda$ . (With  $\lambda$  denoting such a solution, the corresponding “new” eigenfunction is defined as a certain Green’s function  $G_\lambda$ , cf. Section 1.3.) A summation by parts argument (see for instance Lemma 3.1 in [16]) shows that

$$\sum_{\substack{n \in S \\ |n - n_+(\lambda)| > n_+(\lambda)^\delta}} r_2(n) \left( \frac{1}{n - \lambda} - \frac{n}{n^2 + 1} \right) = -\pi \log \lambda + O_\delta(1)$$

and hence the truncation given by (1.4) is equivalent to a logarithmic renormalisation of the r. h. s. of (1.1), namely, as  $\lambda \rightarrow \infty$ ,

$$(1.5) \quad \sum_{n \in S} r_2(n) \left( \frac{1}{n - \lambda} - \frac{n}{n^2 + 1} \right) = -\pi(1 + o_\delta(1)) \log \lambda = c_0 \tan \left( \frac{\varphi_\lambda}{2} \right),$$

if we allow  $\varphi_\lambda$  to depend on  $\lambda$  appropriately, and where the  $o_\delta(1)$  error term depends on the exponent  $\delta$ . Since the error term depends on  $\delta$  we note that there is no unique choice of strong coupling quantization; the key point is matching the leading order logarithmic term. The renormalization in (1.5) can be viewed as letting a boundary condition vary with the energy. Consequently  $D_{\varphi_\lambda}$ , the domain of the operator  $H_{\varphi_\lambda}$  is varying; this setting is reminiscent of problems in semiclassical analysis where boundary conditions are allowed to depend on the semiclassical parameter  $\hbar$ .

**Remark 1.** *In the weak coupling quantization the lowest new eigenvalue is always negative, but for the strong coupling quantization the lowest new eigenvalue may be either positive or negative. In the case of a positive lowest new eigenvalue, this eigenvalue would be denoted  $\lambda_1$  to keep our notation consistent, in particular ensuring that  $\lambda_n < n$  for  $n \in S$  always holds.*

We remark that in the statement of our main result, Theorem 1.1, the sequence  $\Lambda = \{\lambda_n\}$  will denote **any increasing sequence of numbers which interlace with  $S$** . In particular, it applies to the eigenvalues of the weak, as well as the strong, coupling quantizations.

**1.3. Semiclassical Measures.** Let  $a \in C^\infty(S^*\mathbb{T}^2)$ . Denote by  $\text{Op}(a)$  a zero-order pseudo-differential operator associated with  $a$  (see subsection 2.1 for more details.)

Let  $g_\lambda = G_\lambda / \|G_\lambda\|_2$ ,  $\lambda \notin S$ , where we recall that  $S$  denotes the set of Laplace eigenvalues and that  $G_\lambda = (\Delta + \lambda)^{-1}\delta_{x_0}$ . We are interested in weak limits of

measures  $d\mu_\lambda$  defined by the identity

$$(1.6) \quad \langle \text{Op}(a)g_\lambda, g_\lambda \rangle = \int_{S^*\mathbb{T}^2} ad\mu_\lambda.$$

**1.4. Main Result.** The following theorem holds generally for the  $L^2$ -normalized Green's functions  $g_\lambda$ . It states that the measures  $d\mu_\lambda$  defined by (1.6) converge weakly to Liouville measure as  $\lambda \rightarrow \infty$  along a full density subsequence of *any* increasing sequence  $\Lambda$  which interlaces with  $S$ . (Recall that  $S$  denotes set of unperturbed Laplace eigenvalues, namely the set of integers which can be represented as a sum of two squares.)

**Theorem 1.1.** *Let  $\Lambda$  be an increasing sequence which interlaces with  $S$ . For  $m \in S$ , denote by  $\lambda_m$  the largest element of  $\Lambda$  which is smaller than  $m \in S$ . There exists a full density subsequence  $S' \subset S$ , that does not depend on  $\Lambda$ , such that for all  $a \in C^\infty(S^*\mathbb{T}^2)$ ,*

$$(1.7) \quad \lim_{\substack{m \rightarrow \infty \\ m \in S'}} \langle \text{Op}(a)g_{\lambda_m}, g_{\lambda_m} \rangle = \int_{S^*\mathbb{T}^2} a(x, \varphi) \frac{dx d\varphi}{\text{vol}(S^*\mathbb{T}^2)}.$$

As already noted, the theorem holds in particular for the new eigenvalues of the weak and strong coupling quantizations of a point scatterer. Hence we have the following corollary of Theorem 1.1.

**Corollary 1.2.** *Quantum Ergodicity holds for the new eigenfunctions of weakly, as well as strongly, coupled point scatterers on  $\mathbb{T}^2$ .*

**Remark 2.** *Recall that the new eigenvalues in the strong coupling limit are given by the set of solutions  $\{\lambda_m\}_m$  to (1.5) (or alternatively, solutions to (1.4)), with corresponding new eigenfunctions given by the Green's functions  $G_{\lambda_m}$ . Although these Green's functions are eigenfunctions of different operators  $\{H_{\varphi_{\lambda_m}}\}_m$  (in fact, the domains of the operators change), it is natural to say that quantum ergodicity holds in the strong coupling limit if a full density subset of the collection of new eigenfunctions equidistribute.*

We further note that the counting function, or Weyl's law, for the set of new eigenvalues (cf. Theorem 3.3) satisfies

$$|\{n : \lambda_n \leq x\}| \ll \frac{x}{\sqrt{\log x}} = o(x),$$

while the counting function for the full set of eigenvalues (new and old, with multiplicity) is the same as for the unperturbed Laplacian, hence  $\gg x$ . Consequently, the sequence of new eigenvalues is of density zero within the full spectrum, and the approach of proving Quantum Ergodicity for the set of new eigenfunctions by computing first or second moments of matrix coefficients (e.g., see [19]) with respect to the full set of eigenfunctions seems unlikely to succeed.

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## 2. THE MATRIX ELEMENTS

**2.1. Quantization of phase space observables.** Consider a classical symbol  $a \in C^\infty(S^*\mathbb{T}^2)$ , where  $S^*\mathbb{T}^2 \simeq \mathbb{T}^2 \times S^1$  denotes the unit cotangent bundle of  $\mathbb{T}^2$ . We may expand  $a$  in the Fourier series (note that the Fourier coefficients decay rapidly since  $a$  is smooth)

$$(2.1) \quad a(x, \phi) = \sum_{\zeta \in \mathbb{Z}^2, k \in \mathbb{Z}} \hat{a}(\zeta, k) e^{i\langle \zeta, x \rangle + ik\phi}.$$

We choose a complex realization of the unit cotangent bundle  $S^*\mathbb{T}^2$  and parametrise the unit circle  $S^1$  at position  $x \in \mathbb{T}^2$  by the complex exponential map  $\varphi \mapsto e^{i\varphi}$ . We now want to associate with  $a$  a pseudodifferential operator  $\text{Op}(a) : C^\infty(\mathbb{T}^2) \rightarrow C^\infty(\mathbb{T}^2)$ . We choose the following symbol (we associate with  $\xi = (\xi_1, \xi_2)$  the complex number  $\tilde{\xi} := \xi_1 + i\xi_2$  and note that  $e^{ik\arg\tilde{\xi}} = (\tilde{\xi}/|\tilde{\xi}|)^k$ )

$$(2.2) \quad \sigma_a(x, \xi) = \begin{cases} \sum_{\zeta \in \mathbb{Z}^2, k \in \mathbb{Z}} \hat{a}(\zeta, k) \left( \frac{\tilde{\xi}}{|\tilde{\xi}|} \right)^k e^{i\langle \zeta, x \rangle}, & \xi \neq 0 \\ \sum_{\zeta \in \mathbb{Z}^2, k \in \mathbb{Z}} \hat{a}(\zeta, k) e^{i\langle \zeta, x \rangle}, & \xi = 0. \end{cases}$$

Claim: The symbol  $\sigma_a$ , as defined above, belongs to the class of toroidal symbols  $S_{1,0}^0(\mathbb{T}^2 \times \mathbb{Z}^2)$  as defined in [13], Part II, Section 4.1.2, Defn. 4.1.7, p. 344. To see this, define the difference operators

$$\Delta_{\xi_j} f(\xi) = f(\xi + e_j) - f(\xi),$$

where  $e_1 = (1, 0)$ ,  $e_2 = (0, 1)$ . By the mean value theorem for repeated differences, for  $f : \mathbb{R}^2 \rightarrow \mathbb{R}$  a smooth function,

$$\Delta_{\xi_1}^{\beta_1} \Delta_{\xi_2}^{\beta_2} f(\xi) = \partial_{\xi_1}^{\beta_1} \partial_{\xi_2}^{\beta_2} f(\xi) \Big|_{\xi=\xi'},$$

for some  $\xi' = \xi + (\beta'_1, \beta'_2)$  with  $(\beta'_1, \beta'_2) \in [0, \beta_1] \times [0, \beta_2]$ . With  $f_k(\xi)$  denoting the real, or imaginary, part of  $(\tilde{\xi}/|\tilde{\xi}|)^k$ , a quick calculation then gives that for integers  $\alpha_1, \alpha_2, \beta_1, \beta_2 \geq 0$ ,

$$\left| \partial_{x_1}^{\alpha_1} \partial_{x_2}^{\alpha_2} \Delta_{\xi_1}^{\beta_1} \Delta_{\xi_2}^{\beta_2} f_k(\xi) e^{i\langle \zeta, x \rangle} \right| \ll_{\alpha_1, \alpha_2, \beta_1, \beta_2} k^{\beta_1 + \beta_2} |\zeta_1|^{\alpha_1} |\zeta_2|^{\alpha_2} (1 + |\xi|)^{-\beta_1 - \beta_2}$$

This bound, together with the rapid decay of Fourier coefficients of  $a$ , implies that

$$|\partial_{x_1}^{\alpha_1} \partial_{x_2}^{\alpha_2} \Delta_{\xi_1}^{\beta_1} \Delta_{\xi_2}^{\beta_2} \sigma_a(x, \xi)| \leq C_{a, \alpha_1, \alpha_2, \beta_1, \beta_2} (1 + |\xi|)^{-\beta_1 - \beta_2},$$

thus confirming the claim.

The action of the pseudodifferential operator  $\text{Op}(a)$  is then defined by multiplication on the Fourier side, analogously to Defn. 4.1.9 in [13]. Therefore, we have

$$(2.3) \quad \begin{aligned} (\text{Op}(a)f)(x) &= \sum_{\xi \in \mathbb{Z}^2} \sigma_a(x, \xi) \hat{f}(\xi) e^{i\langle \xi, x \rangle} \\ &= \sum_{\xi \in \mathbb{Z}^2 \setminus \{0\}} \sum_{\zeta \in \mathbb{Z}^2, k \in \mathbb{Z}} \hat{a}(\zeta, k) \left( \frac{\tilde{\xi}}{|\tilde{\xi}|} \right)^k \hat{f}(\xi) e^{i\langle \xi + \zeta, x \rangle} \\ &\quad + \sum_{\zeta \in \mathbb{Z}^2, k \in \mathbb{Z}} \hat{a}(\zeta, k) \hat{f}(0) e^{i\langle \zeta, x \rangle} \end{aligned}$$

We can now read off the action of  $\text{Op}(a)$  on the Fourier coefficients:

$$(2.4) \quad (\widehat{\text{Op}(a)f})(\xi) = \sum_{\zeta \in \mathbb{Z}^2, k \in \mathbb{Z}} \hat{a}(\zeta, k) \left( \frac{\tilde{\xi} - \tilde{\zeta}}{|\tilde{\xi} - \tilde{\zeta}|} \right)^k \hat{f}(\xi - \zeta), \quad \xi \neq \zeta$$

(recall that  $\tilde{\xi} := \xi_1 + i\xi_2$  and that the Fourier coefficients  $\hat{a}(\zeta, k)$  decay rapidly) and

$$(2.5) \quad (\widehat{\text{Op}(a)f})(\zeta) = \sum_{\zeta \in \mathbb{Z}^2, k \in \mathbb{Z}} \hat{a}(\zeta, k) \hat{f}(0).$$

In terms of the Fourier coefficients the matrix elements of  $\text{Op}(a)$  can be written as

$$(2.6) \quad \langle \text{Op}(a)f, f \rangle = \sum_{\xi \in \mathbb{Z}^2} (\widehat{\text{Op}(a)f})(\xi) \overline{\hat{f}(\xi)}$$

In particular, for the observable  $e_{\zeta, k}(x, \phi) = e^{i\langle \zeta, x \rangle + ik\phi}$ , we have

$$(2.7) \quad \langle \text{Op}(e_{\zeta, k})f, f \rangle = \sum_{\xi \in \mathbb{Z}^2 \setminus \{\zeta\}} \left( \frac{\tilde{\xi} - \tilde{\zeta}}{|\tilde{\xi} - \tilde{\zeta}|} \right)^k \overline{\hat{f}(\xi)} \hat{f}(\xi - \zeta) + \overline{\hat{f}(\zeta)} \hat{f}(0).$$

**2.2. Mixed modes.** If  $\zeta \neq 0$ , we have the bound

$$(2.8) \quad |\langle \text{Op}(e_{\zeta, k})f, f \rangle| \leq \sum_{\xi \in \mathbb{Z}^2} |\hat{f}(\xi)| |\hat{f}(\xi - \zeta)|.$$

In the case  $f = g_\lambda = G_\lambda / \|G_\lambda\|_2$  we have the  $L^2$ -expansion

$$G_\lambda(x, x_0) = \frac{1}{4\pi^2} \sum_{\xi \in \mathbb{Z}^2} c(\xi) e^{i\langle x, \xi \rangle}$$

where  $c(\xi) = \frac{1}{|\xi|^2 - \lambda}$ . We obtain

$$(2.9) \quad |\langle \text{Op}(e_{\zeta, k})g_\lambda, g_\lambda \rangle| \leq \frac{\sum_{\xi \in \mathbb{Z}^2} |c(\xi)| |c(\xi - \zeta)|}{\sum_{\xi \in \mathbb{Z}^2} |c(\xi)|^2}.$$

In [11] it was proved that there exists a full density subsequence  $S' \subset S$  such that for any nonzero lattice vector  $\zeta \in \mathbb{Z}^2$  the matrix elements of  $\text{Op}(e_{\zeta, k})$  vanish as  $n \rightarrow \infty$  along  $S'$ . The following result was obtained.

**Theorem 2.1. (Rudnick-U., 2012)** *Let  $\Lambda$  be an increasing sequence which interlaces with  $S$ . Denote by  $\lambda_m$  the largest element of  $\Lambda$  which is smaller than  $m \in S$ . There exists a subsequence  $S' \subset S$  of full density such that for any  $\zeta \in \mathbb{Z}^2$ ,  $\zeta \neq 0$ ,  $k \in \mathbb{Z}$*

$$(2.10) \quad \lim_{\substack{n \rightarrow \infty \\ n \in S'}} \langle \text{Op}(e_{\zeta, k})g_{\lambda_n}, g_{\lambda_n} \rangle = 0.$$

**Remark 3.** *The above result is only stated for the weak coupling quantization in [11]. However, the proof in fact works for any interlacing sequence, in particular for the strong coupling quantization. To see this, we briefly recall the key steps of the proof in [11]. The first step is to show that the Green's functions  $G_\lambda$  can be approximated by truncated Green's functions  $G_{\lambda, L}$ , where  $L = \lambda^\delta$  for a specific choice of  $\delta > 0$  and the truncation drops all lattice vectors  $\xi$  outside an annulus  $A(\lambda, L) = \{\xi \in \mathbb{Z}^2 \mid ||\xi|^2 - \lambda| \leq L\}$ . The subsequence  $S' \subset S$  is chosen in such a way to ensure that the lattice points inside the annulus  $A(\lambda, L)$  are sufficiently well-spaced; this then implies that  $c(\xi - \zeta) \ll 1/L$  for  $\zeta \in \mathbb{Z}^2$  fixed and  $\xi \in A(\lambda, L)$ .*

A second condition requires that the neighboring Laplace eigenvalues are not too far apart in order for the lower bound  $\|G_\lambda\|_2 \gg 1/\lambda^{o(1)}$  to hold. These two key properties only depend on the arithmetic properties of the neighboring Laplace eigenvalues, and not on the location of the new eigenvalue itself.

**2.3. Pure momentum modes.** Let us consider the case  $\zeta = 0$ . We rewrite the matrix elements as (cf. eq. (2.4))

$$\begin{aligned} \langle \text{Op}(e_{0,k})g_\lambda, g_\lambda \rangle &= \frac{\sum_{\xi \in \mathbb{Z}^2 \setminus \{0\}} (\tilde{\xi}/|\tilde{\xi}|)^k |c(\xi)|^2 + |c(0)|^2}{\sum_{\xi \in \mathbb{Z}^2} |c(\xi)|^2} \\ (2.11) \quad &= \frac{\frac{1}{\lambda^2} + \sum_{n \in S \setminus \{0\}} \frac{w_k(n)}{(n-\lambda)^2}}{\frac{1}{\lambda^2} + \sum_{n \in S \setminus \{0\}} \frac{r_2(n)}{(n-\lambda)^2}} \end{aligned}$$

where  $w_k(n)$ , for  $n \in S$ , is a certain exponential sum defined as follows: with

$$\Lambda_n := \{z = x + iy \in \mathbb{Z}[i] : |z|^2 = n\},$$

denoting the set Gaussian integers of norm  $n$  (we can interpret these as lattice points lying on a circle of radius  $\sqrt{n}$ ), we define

$$(2.12) \quad w_k(n) := \sum_{z \in \Lambda_n} \left( \frac{z}{|z|} \right)^k.$$

### 3. PURE MOMENTUM OBSERVABLES ON THE SQUARE TORUS

We begin by introducing some convenient notation. Given a set  $S \subset \mathbb{Z}$ , let  $S(x) := S \cap [1, x]$ . We say that a subset  $S_1 \subset S$  is of *full density* if  $|S_1(x)| = (1 + o(1)) \cdot |S(x)|$  as  $x \rightarrow \infty$ . In what follows,  $S$  will always denote the set of integers that can be represented as sums of two integer squares.

For  $k \neq 0$ , we can now construct a full density subsequence  $S'_k \subset S$  such that  $\langle \text{Op}(e_{0,k})g_{\lambda_n}, g_{\lambda_n} \rangle \rightarrow 0$  as  $\lambda_n \rightarrow \infty$  along  $n \in S'_k$ . (Recall that  $\lambda_n$  denotes the perturbed eigenvalue associated with the Laplace eigenvalue  $n \in S$ .)

**Proposition 3.1.** *For a given integer  $k \neq 0$ , there exists a subsequence  $S'_k \subset S$ , of full density, such that for  $n \in S'_k$*

$$(3.1) \quad |\langle \text{Op}(e_{0,k})g_{\lambda_n}, g_{\lambda_n} \rangle| \ll (\log \lambda_n)^{1/4 - \log 2/2 + o(1)}.$$

We note that  $1/4 - \log 2/2 = -0.09657 \dots < 0$ .

**3.1. Preliminary Results.** Before we can give the proof of Proposition 3.1 we state a number of necessary results whose proof can be found in the number theory literature, or in Section 5. We first recall Rieger's bound on pair correlation type sums for integers that are sums of two squares.

**Theorem 3.2** ([10], Satz 2). *Let  $f(n)$  denote the characteristic function of  $S$ , the set integers representable as sums of two integer squares. If  $0 < |h| \ll x$  then*

$$\sum_{n \leq x} f(n)f(n+h) \ll \frac{c(h)x}{\log x}$$

where  $c(h) := \prod_{\substack{p|h \\ p \equiv 3 \pmod{4}}} (1 + 1/p)$ .

**Remark 4.** Rieger's result is stated for  $h > 0$  and summing over  $n \leq x + h$ , but since  $f(n) = 0$  for  $n < 0$  and we assume  $|h| \ll x$ , the above formulation follows immediately (albeit possibly with a worse absolute constant.) Moreover, since  $c(h) \leq \sum_{d|h} 1/d$ , it easily follows that

$$(3.2) \quad \sum_{0 < |h| \leq T} c(h) \ll T, \quad T \rightarrow \infty.$$

We shall also need to recall a fundamental fact about the size of  $S(x)$ .

**Theorem 3.3** (Landau, see [9], §183.). *There exists  $c > 0$  such that*

$$(3.3) \quad |S(x)| = \frac{c \cdot x}{\sqrt{\log x}} (1 + O(1/\log x))$$

as  $x \rightarrow \infty$ .

Given  $n \in S$ , let  $\omega_1(n)$  denote the number of prime divisors of  $n$  that are congruent to one modulo four, i.e.,

$$\omega_1(n) := \sum_{p|n, p \equiv 1 \pmod{4}} 1$$

We shall use Erdős-Kac type techniques to prove (see Section 5.2) the following structure result about the factorizations of “typical” integers in the set  $S$ .

**Proposition 3.4.** *We have*

$$(3.4) \quad \frac{1}{|S(x)|} \sum_{n \in S(x)} \omega_1(n) = \frac{1}{2} \log \log x + O(\log \log \log x)$$

and

$$(3.5) \quad \frac{1}{|S(x)|} \sum_{n \in S(x)} \omega_1(n)^2 = \frac{1}{4} (\log \log x)^2 + O((\log \log x) \cdot \log \log \log x)$$

This, together with Chebychev's inequality, immediately gives the following normal order result on  $\omega_1(n)$ .

**Corollary 3.5.** *Fix  $\epsilon > 0$ . Then, as  $x \rightarrow \infty$ ,*

$$(3.6) \quad |\{n \in S(x) : |\omega_1(n) - \frac{1}{2} \log \log n| < (\log \log n)^{1/2+\epsilon}\}| = |S(x)| \cdot (1 + o_\epsilon(1)).$$

From the corollary, we deduce (see Section 5.2 for details) the following weak analog of a normal order result for  $r_2(n)$ .

**Corollary 3.6.** *As  $x \rightarrow \infty$ ,*

$$(3.7) \quad |\{n \in S(x) : r_2(n) = (\log n)^{(\log 2)/2 \pm o(1)}\}| = |S(x)| \cdot (1 + o(1)).$$

We shall also need the following  $L^2$ -bound on the exponential sums  $w_k(n)$ ,

**Proposition 3.7.** *If  $k \neq 0$ , then*

$$(3.8) \quad \sum_{n \in S(x)} |w_k(n)|^2 \ll_k x$$

In particular, by Chebychev's inequality, the number of  $n \in S(x)$  for which  $|w_k(n)| > T$  is at most  $x/T^2$ , and we find that  $|w_k(n)| \leq (\log n)^{1/4+\epsilon}$  holds for almost all  $n \in S(x)$ .

The result readily follows from a Halberstam-Richert type inequality, see section 5.1 for more details.

**3.2. Proof of Proposition 3.1.** We begin by noting that since the set  $\Lambda_n$  is invariant under multiplication by  $i$ ,  $w_k(n) = 0$  unless  $4|k$ . Hence the case  $k \not\equiv 0 \pmod{4}$  is essentially trivial on recalling (2.11). Thus, in what follows we will always assume that  $4|k$ , and  $k \neq 0$ .

We next introduce some further notation. Given  $m \in S$ , let  $m_+, m_- \in S$  denote the nearest neighbor (in  $S$ ) to the right, respectively left, and similarly, let  $m_{++}, m_{--}$  denote the second nearest neighbors to the right, respectively left.

Define  $S_1 \subset S$  by successively removing a zero density subset of elements for which the following properties do not hold. Namely, let  $S_1$  consist of those  $m \in S$  for which the following properties, as  $m \rightarrow \infty$ , all hold:

- (1) Multiplicities are near their (logarithmic) normal order in the following sense:

$$r_2(m) = (\log m)^{(\log 2)/2 \pm o(1)}, \quad r_2(m_-) = (\log m)^{(\log 2)/2 \pm o(1)}.$$

- (2) There is nearly square root cancellation in exponential sums:

$$|w_k(m)| \leq (\log m)^{1/4+o(1)}, \quad |w_k(m_-)| \leq (\log m)^{1/4+o(1)}.$$

- (3) There are no near neighbors:  $m_+ - m \geq (\log m)^{1/2-o(1)}$ , and  $m - m_- \geq (\log m)^{1/2-o(1)}$ .
- (4) There are no near second neighbors:  $m_{++} - m_+ \geq (\log m)^{1/2-o(1)}$ , and  $m_- - m_{--} \geq (\log m)^{1/2-o(1)}$ .
- (5) Neighbors are not too far away:  $m_+ - m \leq (\log m)^{1/2+o(1)}$ ,  $m - m_- \leq (\log m)^{1/2+o(1)}$ , and  $m_- - m_{--} \leq (\log m)^{1/2+o(1)}$ .
- (6) There are not too many “close” neighbors in the following sense: for  $T \ll m$ ,

$$|\{n \in S : |n - m| \leq T\}| \ll \frac{T(\log T)^2}{(\log m)^{1/2-o(1)}}$$

- (7) For  $W \in [(\log m)^{1/4} \cdot (\log \log m)^2, (\log m)^2]$  there are

$$\gg W^2 / ((\log m)^{1/2} (\log \log m) (\log W)^2)$$

elements in  $S$  that lie between  $m$  and  $n$  if  $|w_k(n)| \geq W$ .

- (8) For  $W \geq (\log m)^2$  there are  $\gg W^{3/2} / \log W$  elements in  $S$  that lie between  $m$  and  $n$  if  $|w_k(n)| \geq W$ .
- (9) For  $\epsilon > 0$  and  $G \in [2, m^{1-\epsilon}]$ ,

$$H_G(m) := \sum_{\substack{n \in S, n \neq m \\ |m - n| \geq G}} \frac{1}{|m - n|^2} \ll_\epsilon \frac{(\log G)^2}{G(\log m)^{1/2-o(1)}}.$$

**Remark 5.** We tacitly assume that  $o(1)$  is chosen so that  $(\log m)^{o(1)} \rightarrow \infty$  as  $m \in S$  tends to infinity; we will also use the convention that the sign of  $o(1)$  is important, in particular  $(\log m)^{-o(1)} \rightarrow 0$ .

We defer the proof that  $S_1$  has full density inside  $S$  to Section 3.3.

**Remark 6.** Here, and what follows we will without comment make use of the fact that  $\log m \sim \log m_- \sim \log m_+ \sim \log m_{++}$  etc. To see this, any crude bound on  $|m_+ - m|$ ,  $|m_{++} - m|$  etc suffices, e.g. the trivial bound  $|m_+ - m| \ll m^{1/2}$  which follows from bounding the distance to the nearest square. Moreover, we also use the fact that for almost all  $m \in S(x)$ , e.g.  $m \in [x/\log x, x]$ , we have  $\log m \sim \log x$ .

The following will be used to show that the numerator in (2.11) is essentially given by two terms.

**Lemma 3.8.** *If  $m \in S_1$ , then*

$$(3.9) \quad \sum_{n \in S, n \neq m, m_-} \frac{|w_k(n)|}{|m - n|^2} \ll \frac{1}{(\log m)^{3/4-o(1)}}$$

*Proof.* Fix  $m \in S_1(x)$ . To simplify the notation, let  $L = \log m$ . To bound the sum

$$\sum_{n \in S, n \neq m, m_-} \frac{|w_k(n)|}{|m - n|^2}$$

we split it into parts according to the size of  $|w_k(n)|$ .

*Small  $|w_k(n)|$ :*  $|w_k(n)| \leq L^{1/4}(\log L)^2$ . Since  $m \in S_1$ , its nearest neighbors, by property 3 are of distance at least  $L^{1/2-o(1)}$  away from  $m$ . Thus, the contribution from  $n$  for which  $|w_k(n)| \leq L^{1/4}(\log L)^2$  is, by property 9,

$$\begin{aligned} &\ll \sum_{\substack{n \in S \\ |n-m| \geq L^{1/2-o(1)}}} \frac{L^{1/4}(\log L)^2}{|m - n|^2} = L^{1/4} \cdot (\log L)^2 \cdot H_{L^{1/2-o(1)}}(m) \\ &\ll \frac{L^{1/4}(\log L)^4}{L^{1/2-o(1)} L^{1/2-o(1)}} = \frac{1}{L^{3/4-o(1)}}. \end{aligned}$$

*Medium  $|w_k(n)|$ :*  $|w_k(n)| \in [L^{1/4}(\log L)^2, L^2]$ . For terms in the sum for which  $n \geq 2m$ , we use the crude bound  $|w_k(n)| \leq r_2(n) \ll \sqrt{n}$  and find that the total contribution is  $\ll \sum_{n \geq m} n^{-3/2} \ll m^{-1/2}$ , and hence it is enough to consider terms for which  $n < 2m$ .

Let  $W_i = 2^i L^{1/4}(\log L)^2$  for integer  $i \geq 0$  such that  $2^i L^{1/4}(\log L)^2 \leq L^2$  and consider  $n$  such that  $|w_k(n)| \in [W_i, W_{i+1}]$ . By property 7, the number of elements in  $S$  between  $n$  and  $m$  is  $\gg W_i^2/(L^{1/2+o(1)}(\log W_i)^2)$ . Thus, using the bound on the number of close neighbors, i.e., take  $T = |n - m|$  in property 6 (note that  $T \ll m$  when  $n < 2m$ ), we must have  $\frac{T(\log T)^2}{(\log m)^{1/2-o(1)}} \gg W_i^2/(L^{1/2+o(1)}(\log W_i)^2)$ , which implies that  $|n - m| \gg W_i^{2-o(1)}$ . Thus, by property 9 (take  $G = W_i^{2-o(1)}$ ),

$$\begin{aligned} &\sum_{\substack{n \in S \\ |w_k(n)| \in [W_i, W_{i+1}]} } \frac{|w_k(n)|}{|n - m|^2} \ll \sum_{n \in S: |m-n| \gg W_i^{2-o(1)}} \frac{W_i}{|n - m|^2} \ll \frac{W_i(\log W_i)^2}{L^{1/2-o(1)} W_i^{2-o(1)}} \\ &= \frac{1}{L^{1/2-o(1)} \cdot W_i^{1-o(1)}} = \frac{1}{L^{1/2-o(1)} \cdot (2^i L^{1/4}(\log L)^2)^{1-o(1)}} \ll \frac{1}{L^{3/4-o(1)}(3/2)^i}. \end{aligned}$$

Summing over relevant  $i \geq 0$ , we find that the total contribution is  $\ll \frac{1}{L^{3/4-o(1)}}$ .

*Large  $|w_k(n)|$ :*  $|w_k(n)| \geq L^2$ . Let  $W_i = 2^i L^2$  and consider  $n$  such that  $|w_k(n)| \in [W_i, 2W_i]$ . By property 8, we must then have  $|n - m| \gg W_i^{3/2}/\log W_i$ , and hence the contribution is (using the bound  $\sum_{k \geq A} 1/k^2 \ll 1/A$ )

$$\ll \sum_{\substack{n \in S \\ |n-m| \gg W_i^{3/2}/\log W_i}} \frac{W_i}{|n - m|^2} \ll \frac{W_i \cdot \log W_i}{W_i^{3/2}} \ll \frac{1}{W_i^{1/2-o(1)}} = \frac{1}{(2^i L^2)^{1/2-o(1)}}.$$

Summing over  $i \geq 0$ , the total contribution is

$$\ll \frac{1}{L^{1-o(1)}} \sum_{i \geq 0} 2^{-(1/2-o(1))i} \ll \frac{1}{L^{1-o(1)}}.$$

□

*Proof of Proposition 3.1 using Lemma 3.8.* Recalling that  $m_- < \lambda_m < m$ , we note that that  $|\lambda_m - n| \geq |m - n|$  if  $n > m$ . Moreover, for  $n \leq m_-$  the minimum of  $|\lambda_m - n|/|m - n| = 1 - \frac{m-\lambda_m}{m-n}$  (as  $n \leq m_-$  ranges over elements in  $S$ ) is attained for  $n = m_-$ , and consequently

$$|\lambda_m - n|/|m - n| \geq |\lambda_m - m_-|/|m - m_-| \geq |m_- - m_-|/|m - m_-|$$

which, by properties 4, and 5, is  $\gg (\log m)^{1/2-o(1)}/(\log m)^{1/2+o(1)} = 1/(\log m)^{o(1)}$ . Hence  $|\lambda_m - n| \gg (\log m)^{-o(1)}|m - n|$  holds for  $n \neq m, m_-$ , and thus

$$\sum_{n \in S, n \neq m, m_-} \frac{|w_k(n)|}{|\lambda_m - n|^2} \leq (\log m)^{o(1)} \cdot \sum_{n \in S, n \neq m, m_-} \frac{|w_k(n)|}{|m - n|^2}$$

Let  $M = \min(|\lambda_m - m|^2, |\lambda_m - m_-|^2)$ . Trivially  $\lambda_m \gg m^{1/2}$ , and by property 1, Lemma 3.8 implies that

$$\frac{1/\lambda_m^2 + \sum_{n \in S} \frac{|w_k(n)|}{|\lambda_m - n|^2}}{1/\lambda_m^2 + \sum_{n \in S} \frac{r_2(n)}{|\lambda_m - n|^2}} \ll \frac{O(1/m) + (|w_k(m)| + |w_k(m_-)|)/M + \frac{1}{(\log m)^{3/4-o(1)}}}{\frac{O(1/m) + (\log m)^{(\log 2)/2-o(1)}}{M}}$$

which, by property 2, is

$$(3.10) \quad \ll \frac{(\log m)^{1/4+o(1)} + \frac{M}{(\log m)^{3/4-o(1)}}}{(\log m)^{(\log 2)/2-o(1)}}.$$

Recalling that  $\lambda_m \in [m_-, m]$ , property 5 implies that  $M \ll (\log m)^{1+o(1)}$  and we thus find that (3.10) is

$$\ll \frac{(\log m)^{1/4+o(1)} + \frac{(\log m)^{1+o(1)}}{(\log m)^{3/4-o(1)}}}{(\log m)^{(\log 2)/2-o(1)}} = \frac{1}{(\log m)^{(\log 2)/2-1/4-o(1)}} = o(1)$$

as  $m \rightarrow \infty$ , since  $(\log 2)/2 - 1/4 = 0.09657 \dots$ . Recalling that  $\log m \gg \log \lambda_m$  (cf. Remark 6) the proof is concluded.

□

### 3.3. Proof that $S_1$ has full density.

3.3.1. *Property (1).* That  $r_2(m) = (\log m)^{(\log 2)/2+o(1)}$  holds for almost all  $m \in S$  follows from corollary 3.6. To ensure that  $r_2(m_-) = (\log m)^{(\log 2)/2+o(1)}$  also holds, we remove the right neighbor of those  $m$  for which  $r_2(m) = (\log m)^{(\log 2)/2+o(1)}$  is not true; this removes another zero density set. (By Remark 6,  $\log m_+ = (1 + o(1)) \log m$ .)

3.3.2. *Property (2).* By Proposition 3.7

$$\sum_{n \in S(x)} |w_k(n)|^2 \ll x$$

and Chebychev's inequality, together with  $|S(x)| \sim cx/\sqrt{\log x}$ , then gives that  $|w_k(m)| \leq (\log m)^{1/4+o(1)}$  holds for almost all  $m \in S(x)$ . Removing right neighbors, as in the proof of Property (1), the same holds for  $|w_k(m_-)|$ .

3.3.3. *Property (3).* Let  $f$  denote the characteristic function of  $S$ . By Theorem 3.2,

$$\begin{aligned} \sum_{m \leq x} \sum_{h:0<|h|\leq(\log m)^{1/2-o(1)}} f(m)f(m+h) &\ll \frac{x}{\log x} \sum_{h:0<|h|\leq(\log x)^{1/2-o(1)}} c(h) \\ &\ll \frac{x}{\log x} \cdot (\log x)^{1/2-o(1)}. \end{aligned}$$

Thus, by Chebychev's inequality,

$$\sum_{h:0<|h|\leq(\log m)^{1/2-o(1)}} f(m)f(m+h) < 1$$

holds for almost all  $m$  in  $S(x)$ . Consequently, almost all  $m \in S$  have no nearby neighbors.

3.3.4. *Property (4).* We use the same proof as the one used for showing that Property (3) holds.

3.3.5. *Property (5).* Let  $n_1 < n_2 \dots < n_I \leq x$  denote ordered representatives of the elements in  $S(x)$ , and let  $s_i = n_{i+1} - n_i$ . Since  $\sum_{i < I} s_i \leq x$ , Chebychev's inequality implies that  $s_i \leq (\log n_i)^{1/2+o(1)}$  holds for almost all  $n_i \in S$ ; consequently  $n_+ - n \leq (\log n_i)^{1/2+o(1)}$  for almost all  $n \in S$ .

A similar argument shows that  $n_i - n_{i-2} \leq (\log n_i)^{1/2+o(1)}$  also holds for almost all  $n_i$ . Hence both  $n - n_- \leq (\log n_i)^{1/2+o(1)}$  and  $n_- - n_{--} \leq (\log n_i)^{1/2+o(1)}$  holds for almost all  $n \in S$ .

3.3.6. *Property (6).* The argument is similar to the one used to prove property 3: again let  $f$  denote the characteristic function of  $S$ . Then, as  $T \ll m \leq x$ , Theorem 3.2 gives that

$$\sum_{m \leq x} \sum_{h:0<|h|\leq T} f(m)f(m+h) \ll \frac{x}{\log x} \sum_{h:0<|h|\leq T} c(h) \ll \frac{xT}{\log x}$$

and Chebychev's inequality implies that

$$\sum_{h:0<|h|\leq T} f(m)f(m+h) \geq \frac{T(\log T)^2}{(\log x)^{1/2-o(1)}}$$

holds for at most  $\frac{x}{(\log x)^{1/2+o(1)}(\log T)^2}$  exceptional elements  $m \in S(x)$ . Taking  $T_i = 2^i$ , removing the exceptional elements, and summing over  $i$  we find that we have removed

$$\ll \frac{x}{(\log x)^{1/2+o(1)}} \sum_{i \geq 0} 1/i^2 = o(|S(x)|)$$

elements.

Thus, property 6 holds for  $T$  being a power of two. To see that it holds for all  $T \ll m$ , take  $i$  to be the smallest integer such that  $T_i = 2^i \geq T$  and note that  $T \in [T_i/2, T_i]$ .

**3.3.7. Property (7).** Given  $n \in S$  such that

$$w = |w_k(n)| \in [(\log n)^{1/4}(\log \log n), (\log n)^2],$$

remove  $2 \cdot w^2 / ((\log n)^{1/2}(\log \log n)(\log w)^2)$  neighbors to the left of  $n$ , and  $2 \cdot w^2 / ((\log n)^{1/2}(\log \log n)(\log w)^2)$  neighbors to the right, and let  $R_n$  denote the set of such removed elements.

Fix  $x$  and consider the number of removed elements in  $[1, x]$ . We claim that if  $l \in S(x)$  has been removed, then  $l \in R_n$  for some  $n \leq 2x$ . To show this, we note that given an integer  $t$ , we can always find  $l_1, l_2 \in S$  such that  $l_1 < t < l_2$ , and  $l_2 - l_1 \ll \sqrt{t}$  (just take nearby squares), and since  $|w_k(n)| \leq r_2(n) \leq n^{o(1)}$ , any  $R_n$  will be contained in an interval of length  $\ll n^{1/2+o(1)}$ .

Hence it suffices to bound the union of  $R_n$  for  $n \leq 2x$ . The removed contribution from  $n$  for which  $n \leq x/(\log x)^{10}$  is at most  $\frac{x \cdot (\log x)^4}{(\log x)^{10}} = o(|S(x)|)$  (here we use the assumption  $w \leq (\log n)^2$ ).

On the other hand, for  $n \in [x/(\log x)^{10}, 2x]$ , we have  $\log n = (1 + o(1)) \log x$ . Let  $W_i = 2^i(\log x)^{1/2}(\log \log x)$ , and consider  $n \in S(x)$  such that  $|w_k(n)| \in [W_i, 2W_i]$ . By Proposition 3.7 and Chebychev's inequality, the number of such  $n$  is  $\ll \frac{x}{W_i^2}$ , and the total number of removed elements is thus

$$\ll \frac{x}{W_i^2} \cdot \frac{W_i^2}{(\log x)^{1/2}(\log \log x)(\log W_i)^2} \ll \frac{x}{(\log x)^{1/2}(\log \log x) \cdot i^2}$$

Summing over  $i \geq 0$  we find that the total number of removed elements is

$$\ll \frac{x}{(\log x)^{1/2}(\log \log x) \cdot i^2} = o(|S(x)|).$$

**3.3.8. Property (8).** Arguing as before, if  $|w_k(n)| \geq (\log n)^2$  let  $w = |w_k(n)|$  and remove the nearest  $2w^{3/2}/\log w$  neighbors to the right and left of  $n$ ; let  $R_n$  denote the set of removed neighbors.

Fix  $x$  and consider the number of removed elements in  $[1, x]$ . We first note that  $|w_k(n)| \leq r_2(n) \ll n^{1/100}$  holds for all  $n \in S$ . Consequently  $R_n$ , if non-empty, contains at most  $n^{3/200}$  neighbors of  $n$  which (since  $|S(2y)| - |S(y)| \gg y/\sqrt{\log y}$  for all  $y$  by Landau) implies that if  $l \in S(x)$  and  $l$  belongs to some  $R_n$ , then  $n \leq 2x$ .

Consider first the removed contribution coming from  $R_n$  for which  $n \leq \sqrt{x}$ . Since  $|w_k(n)| \leq r_2(n) \ll n^{1/100}$ , the total contribution is

$$\ll \sqrt{x} \cdot (x^{1/100})^{3/2} = o(|S(x)|).$$

If  $n \in [\sqrt{x}, 2x]$  and  $|w_k(n)| \geq (\log n)^2$ , we have

$$|w_k(n)| \geq (\log x)^2/100.$$

Define  $W_i = 2^i \cdot (\log x)^2/100$  and consider the removed contribution from  $R_n$  for which  $|w_k(n)| \in [W_i, 2W_i]$ . By Proposition 3.7 and Chebychev's inequality, the number of such  $n \in S(2x)$  is  $\ll \frac{x}{W_i^2}$  and the associated removed contribution is

$$\ll \frac{x \cdot (W_i^{3/2}/\log W_i)}{W_i^2} \ll \frac{x}{W_i^{1/2} \log W_i} \ll \frac{x}{(2^i(\log x)^2)^{1/2}} = \frac{x}{2^{i/2} \log x}.$$

Summing over  $i \geq 0$  we find that the total contribution is

$$\ll \sum_{i \geq 0} \frac{x}{2^{i/2} \log x} = o(|S(x)|).$$

3.3.9. *Property (9).* The final property is an immediate consequence of the following Lemma.

**Lemma 3.9.** *If  $\epsilon > 0$  then for almost all  $m \in S(x)$ , we have, for any  $T \in [2, x^{1-\epsilon}]$ ,*

$$\sum_{\substack{n \in S \\ |n-m| \geq T \\ n \geq 2x}} \frac{1}{(m-n)^2} \ll \frac{(\log T)^2}{T(\log x)^{1/2-o(1)}}$$

*Proof.* We first bound the sum over  $n \in S \setminus S(2x)$ , i.e., those  $n$  for which  $n \geq 2x$ :

$$\sum_{\substack{n \in S \\ |n-m| \geq T \\ n \geq 2x}} \frac{1}{(m-n)^2} \leq \sum_{k \geq x} 1/k^2 \ll 1/x = o\left(\frac{(\log T)^2}{T(\log x)^{1/2-o(1)}}\right).$$

(Recall that  $m \leq x$  since  $m \in S(x)$ , and that  $T \leq x^{1-\epsilon}$ .)

Next we note that

$$\sum_{\substack{m,n \in S(2x) \\ |n-m| \geq T}} \frac{1}{(m-n)^2} = \sum_{k \geq T} \frac{|\{m, n \in S(2x) : |m-n|=k\}|}{k^2}$$

By Theorem 3.2,

$$|\{m, n \in S(2x) : |m-n|=h\}| \ll \frac{x \cdot c(h)}{\log x}$$

and, by partial summation and using that  $c(h)$  is bounded on average (cf. (3.2)),

$$\sum_{\substack{m,n \in S(2x) \\ |n-m| \geq T}} \frac{1}{(m-n)^2} \ll \frac{x}{\log x} \sum_{h \geq T} \frac{c(h)}{h^2} \ll \frac{x}{T \log x}.$$

By Chebychev's inequality, the number of  $m \in S(2x)$  for which  $\sum_{\substack{n \in S(2x) \\ |n-m| \geq T}} \frac{1}{(m-n)^2} \geq \frac{(\log T)^2}{T(\log x)^{1/2-o(1)}}$  holds is thus

$$\ll \frac{x}{T \log x} \sqrt{\frac{(\log T)^2}{T(\log x)^{1/2-o(1)}}} = \frac{x}{(\log x)^{1/2+o(1)} \cdot (\log T)^2}$$

Taking  $T_i = 2^i$ , summing over  $i \ll \log x$ , and recalling that  $|S(x)| \sim x/\sqrt{\log x}$  we find that the property holds in the special case of  $T$  being a power of two. The result for general  $T$  follows by taking the largest  $i$  such that  $T_i = 2^i \leq T$  and noting that  $T_i \in [T/2, T]$ .  $\square$

#### 4. PROOF OF THEOREM 1.1

Let  $a \in C^\infty(S^*\mathbb{T}^2)$  be a smooth observable with rapidly decaying Fourier expansion

$$a(x, \phi) = \sum_{\zeta \in \mathbb{Z}^2, k \in \mathbb{Z}} \hat{a}(\zeta, k) e^{i\langle \zeta, x \rangle + ik\phi}.$$

Since  $\text{Op}(e_{\zeta,k})$  is unitary (cf. (2.4)) for all  $\zeta, k$ , we have  $|\langle \text{Op}(e_{\zeta,k})g_\lambda, g_\lambda \rangle| = |\langle g_\lambda, g_\lambda \rangle|$ , and the rapid decay of Fourier coefficients then shows that given  $\epsilon > 0$ , there exists  $J$  such that

$$(4.1) \quad |\langle (\text{Op}(a) - \text{Op}(P_J))g_\lambda, g_\lambda \rangle| \leq \sum_{|\zeta|, |k| > J} |\hat{a}(\zeta, k)| |\langle \text{Op}(e_{\zeta,k})g_\lambda, g_\lambda \rangle| \leq \epsilon$$

holds *uniformly* in  $\lambda$ , where  $P_J(x, \phi)$  is the trigonometric polynomial

$$(4.2) \quad P_J(x, \phi) = \sum_{\substack{\zeta \in \mathbb{Z}^2, k \in \mathbb{Z} \\ |\zeta|, |k| \leq J}} \hat{a}(\zeta, k) e^{i\langle \zeta, x \rangle + ik\phi}$$

obtained by truncating the Fourier expansion of  $a$ . Hence it is enough to show that for any fixed  $J \geq 1$ ,

$$(4.3) \quad \langle \text{Op}(P_J)g_\lambda, g_\lambda \rangle \rightarrow \frac{1}{\text{vol}(S^*\mathbb{T}^2)} \int_{S^*\mathbb{T}^2} a d\mu = \hat{a}(0, 0)$$

as  $\lambda \rightarrow \infty$  along a full density subsequence of  $S$ .

Now, given  $J \geq 1$ , let

$$\tilde{S}_J := \bigcap_{|k| \leq J} (S'_k \cap S')$$

where  $S' \subset S$  denotes the full density sequence of Theorem 2.1. (Since  $S'$  and  $S'_k$  have full densities for all  $k \neq 0$ , so does  $\tilde{S}_J$  for all  $J$ .) It follows from the previous two sections that

$$(4.4) \quad \langle \text{Op}(P_J)g_\lambda, g_\lambda \rangle \rightarrow \frac{1}{\text{vol}(S^*\mathbb{T}^2)} \int_{S^*\mathbb{T}^2} P_J d\mu = \hat{a}(0, 0)$$

as  $\lambda \in \tilde{S}_J \rightarrow \infty$ .

In order to construct the full density sequence of Theorem 1.1 we use a standard diagonalisation argument (see for instance [4]) to extract such a sequence from the list of sequences  $\{\tilde{S}_J\}_J$ . By construction  $\tilde{S}_{J+1} \subset \tilde{S}_J$ . Choose  $M_J$  such that for all  $X > M_J$

$$(4.5) \quad \frac{\#\{\lambda \in \tilde{S}_J \mid \lambda \leq X\}}{\#\{\lambda \in S \mid \lambda \leq X\}} \geq 1 - \frac{1}{2^J}$$

and let  $S'_\infty$  be such that  $S'_\infty \cap [M_J, M_{J+1}] = \tilde{S}_J \cap [M_J, M_{J+1}]$  for all  $J$ . Then  $S'_\infty \cap [0, M_{J+1}]$  contains  $\tilde{S}_J \cap [0, M_{J+1}]$  and therefore  $S'_\infty$  is of full density in  $S$ . Moreover, for any  $J \geq 1$ , we have

$$(4.6) \quad \langle \text{Op}(P_J)g_\lambda, g_\lambda \rangle = \int_{S^*\mathbb{T}^2} P_J d\mu_\lambda \rightarrow \frac{1}{\text{vol}(S^*\mathbb{T}^2)} \int_{S^*\mathbb{T}^2} P_J d\mu = \hat{a}(0, 0)$$

as  $\lambda \rightarrow \infty$  along  $S'_\infty$  since  $S'_\infty \cap (M_{J+1}, \infty) \subset \tilde{S}_J \cap (M_{J+1}, \infty)$ .

## 5. NUMBER THEORETIC BACKGROUND

**5.1. Bounding mean values of multiplicative functions.** We recall that  $r_2(n)/4$  is a *multiplicative function*, i.e.,  $r_2(mn)/4 = r_2(m)/4 \cdot r_2(n)/4$  if  $(m, n) = 1$ , and similarly  $w_k(n)/4$  is also multiplicative (e.g., see the proof of Proposition 6 in [5].)

In particular, both functions are determined by their values at prime powers, and we have

$$\frac{r_2(p^e)}{4} = \begin{cases} e+1 & \text{if } p \equiv 1 \pmod{4}, \\ 1 & \text{if } p \equiv 3 \pmod{4} \text{ and } e \text{ is even, or if } p=2, \\ 0 & \text{if } p \equiv 3 \pmod{4} \text{ and } e \text{ is odd.} \end{cases}$$

For  $p \equiv 1 \pmod{4}$ , define the angle  $\theta_p \in [0, \pi/4)$  by  $\cos \theta_p = x/\sqrt{x^2 + y^2}$ , where  $x^2 + y^2 = p$  for  $x, y \in \mathbb{Z}$  and  $0 \leq y \leq x$ . We then have (if  $4|k$ )

$$\frac{w_k(p^e)}{4} = \begin{cases} \sum_{l=0}^e e^{i \cdot \theta_p \cdot (e-2l)} & \text{if } p \equiv 1 \pmod{4}, \\ 1 & \text{if } p \equiv 3 \pmod{4} \text{ and } e \text{ is even,} \\ 0 & \text{if } p \equiv 3 \pmod{4} \text{ and } e \text{ is odd,} \\ \pm 1 & \text{if } p=2. \end{cases}$$

In particular (again for  $4|k$ ),  $w_k(2)/4 = (-1)^{k/4}$ , and for odd primes we have

$$(5.1) \quad \frac{w_k(p)}{4} = \begin{cases} 2 \cos(k\theta_p) & \text{for } p \equiv 1 \pmod{4}, \\ 0 & \text{for } p \equiv 3 \pmod{4}. \end{cases}$$

Now, let  $f$  be a non-negative multiplicative function such that for all prime powers  $f(p^k) \ll \gamma^k$  holds for some  $\gamma < 2$ , and

$$\sum_{p \leq x} f(p) = \frac{x}{\log x} \cdot (\tau + o(1)),$$

as  $x \rightarrow \infty$ , for some constant  $\tau$ . Satz 1 of Wirsing [17] then implies that

$$\sum_{n \leq x} f(n) \ll_{\tau} \frac{x}{\log x} \cdot \prod_{p \leq x} \left(1 + \frac{f(p)}{p} + \frac{f(p^2)}{p^2} + \dots\right).$$

**5.1.1. Proof of Proposition 3.7.** For  $k$  fixed, define a multiplicative function  $f(n) := (|w_k(n)|/4)^2$  (recall that  $|w_k(n)|/4$  is multiplicative.) By (5.1), we have

$$f(p) = \begin{cases} 1 & \text{for } p=2, \\ (2 \cos(k\theta_p))^2 & \text{for } p \equiv 1 \pmod{4}, \\ 0 & \text{for } p \equiv 3 \pmod{4}, \end{cases}$$

and we find that

$$\sum_{p \leq x} f(p) = 1 + \sum_{p \leq x, p \equiv 1 \pmod{4}} f(p) = 1 + \sum_{p \leq x, p \equiv 1 \pmod{4}} (2 \cos(2\pi k\theta_p))^2.$$

Thus Hecke's result on angular equidistribution of split Gaussian primes (see [6]) gives that

$$(5.2) \quad \sum_{p \leq x} f(p) = \frac{x}{\log x} \cdot \left(\frac{1}{2} \cdot \int_0^1 (2 \cos(2\pi k\theta))^2 d\theta + o(1)\right) = \frac{x}{\log x} \cdot (1 + o(1))$$

Hence Wirsing's Satz 1 applies (also note that  $f(p^k) \ll k^2$  for all  $p, k$ ), thus

$$\sum_{n \leq x} |w_k(n)|^2 \ll \sum_{n \leq x} f(n) \ll \frac{x}{\log x} \cdot \prod_{p \leq x} \left(1 + \frac{f(p)}{p} + \frac{f(p^2)}{p^2} + \dots\right).$$

Now, since  $\sum_{k=2}^{\infty} f(p^k)/p^k \leq \sum_{k=2}^{\infty} (k+1)^2/p^k \ll 1/p^2$ , we find that

$$\prod_{p \leq x} \left(1 + \frac{f(p)}{p} + \frac{f(p^2)}{p^2} + \dots\right) \ll \exp \left( \sum_{p \leq x} \frac{f(p)}{p} \right) = \exp(\log \log x + O(1)),$$

where the final equality follows from (5.2) and partial summation. Hence

$$\sum_{n \leq x} |w_k(n)|^2 \ll \frac{x}{\log x} \cdot \exp(\log \log x + O(1)) \ll x.$$

**5.2. Erdős-Kac Theory.** Let  $\omega(n)$  denote the number of distinct prime factors of an integer  $n$ . The celebrated Erdős-Kac theorem assert that the distribution of  $\left\{ \frac{\omega(n) - \log \log n}{\sqrt{\log \log n}} \right\}_{n \leq x}$  is given by the standard normal distribution as  $x \rightarrow \infty$ ; in particular, a typical integer of size  $x$  has about  $\log \log x$  prime factors. We shall need some analogous, but weaker, results for elements in  $S$ .

Recall that given  $n \in S$ ,  $\omega_1(n)$  denotes the number of prime factors, congruent to one modulo four, of  $n$ ; i.e., with  $\sum'_p$  denoting the sum over  $p \equiv 1 \pmod{4}$ ,

$$\omega_1(n) := \sum'_{p|n} 1.$$

**5.2.1. Proof of Proposition 3.4.** Using that at most four primes  $p \geq x^{1/4}$  can divide an integer  $n \leq x$ , together with  $\sum'_{p \leq x^{1/4}} |\{n \in S(x) : p|n\}| = \sum'_{p \leq x^{1/4}} |S(x/p)|$ , we find that

$$\begin{aligned} \sum_{n \in S(x)} \omega_1(n) &= \sum_{n \in S(x)} \sum'_{p|n} 1 = \sum_{n \in S(x)} \left( \sum'_{p|n, p \leq x^{1/4}} 1 + O(1) \right) \\ &= \sum'_{p \leq x^{1/4}} |S(x/p)| + O(|S(x)|) \end{aligned}$$

By Landau,  $|S(x/p)| = \frac{cx}{p\sqrt{\log(x/p)}} \cdot (1 + O(1/\log(x/p)))$ , and thus, what will be the main term, is given by

$$\sum'_{p \leq x^{1/4}} |S(x/p)| = \sum'_{p \leq x^{1/4}} \frac{cx}{p\sqrt{\log(x/p)}} \cdot (1 + O(1/\log(x/p)))$$

If  $p \in [x^{1/\log \log x}, x^{1/4}]$  then  $\log(x/p) \gg \log x$  and thus the contribution from such primes  $p$  is

$$\ll \frac{x}{\sqrt{\log x}} \sum'_{p \in [x^{1/\log \log x}, x^{1/4}]} 1/p \ll \frac{x}{\sqrt{\log x}} \cdot \log \left( \frac{\log x^{1/4}}{\log x^{1/\log \log x}} \right) \ll \frac{x \log \log \log x}{\sqrt{\log x}}$$

which is of the same order as the claimed error term in the first assertion of the Proposition.

Now, if  $p \leq x^{1/\log \log x}$  then  $\log p \leq \log x / \log \log x$  and thus

$$\sqrt{\log(x/p)} = \sqrt{\log x - \log p} = \sqrt{\log x} \left( 1 - O\left(\frac{1}{\log \log x}\right) \right)$$

Hence, by the analogue of Mertens' theorem for primes in progressions<sup>2</sup>, together with  $1/\log(x/p) \ll 1/\log x$  for  $p \leq x^{1/\log\log x}$ , we find that

$$\begin{aligned} & \sum'_{p \leq x^{1/\log\log x}} \frac{cx}{p\sqrt{\log(x/p)}} \cdot (1 + O(1/\log(x/p))) \\ &= \frac{cx}{\sqrt{\log x}} \cdot (1 + O(1/\log x)) \cdot (1 + O(1/\log\log x)) \sum'_{p \leq x^{1/\log\log x}} 1/p = \\ & \quad \frac{cx}{\sqrt{\log x}} \cdot (1 + O(1/\log\log x)) \cdot \left( \frac{1}{2} \log \left( \frac{\log x}{\log\log x} \right) + O(1) \right) = \\ & \quad \frac{cx}{\sqrt{\log x}} \cdot (1 + O(1/\log\log x)) \cdot \left( \frac{1}{2} \log\log x + O(\log\log\log x) \right) = \\ & \quad \frac{cx}{\sqrt{\log x}} \left( \frac{1}{2} \log\log x + O(\log\log\log x) \right) \end{aligned}$$

Dividing by  $|S(x)|$  and again using Landau's Theorem, the proof of the first assertion is concluded.

The variance estimate is similar: since  $n \leq x$  can have at most 4 prime divisors  $p \geq x^{1/4}$ , we have

$$\begin{aligned} \sum_{n \in S(x)} \omega_1(n)^2 &= \sum_{n \in S(x)} \left( \sum'_{p \leq x, p|n} 1 \right)^2 = \sum_{n \in S(x)} \left( \sum'_{p \leq x^{1/4}, p|n} 1 + O(1) \right)^2 \\ &= \sum_{n \in S(x)} \left( \left( \sum'_{p \leq x^{1/4}, p|n} 1 \right)^2 + 2 \sum'_{p \leq x^{1/4}, p|n} 1 + O(1) \right) \end{aligned}$$

The total contribution from the last two terms in the inner sum is, by our first assertion (regarding the mean value of  $\omega_1(n)$ ),

$$\ll \frac{x \log\log x}{\sqrt{\log x}} + |S(x)| \ll \frac{x \log\log x}{\sqrt{\log x}}.$$

With  $[a, b] = ab/(a, b)$  denoting the least common multiple of integers  $a, b$ , we have

$$\begin{aligned} \sum_{n \in S(x)} \left( \sum'_{p \leq x^{1/4}, p|n} 1 \right)^2 &= \sum'_{p_1, p_2 \leq x^{1/4}} |S(x/[p_1, p_2])|^2 \\ &= \sum'_{p_1, p_2 \leq x^{1/4}} |S(x/p_1 p_2)| + \sum'_{p \leq x^{1/4}} |S(x/p)| - \sum'_{p \leq x^{1/4}} |S(x/p^2)| \end{aligned}$$

The latter two terms are of lower order than the claimed main term — the argument used to estimate the mean of  $\omega_1(n)$  implies that

$$\sum'_{p \leq x^{1/4}} |S(x/p)| \ll \frac{x \log\log x}{\sqrt{\log x}}$$

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<sup>2</sup>We shall only need that  $\sum'_{p \leq x} 1/p = 1/2 \cdot \log\log x + O(1)$ , a simple consequence of the prime number theorem for arithmetic progressions.

and, again using Landau, we find that

$$\sum'_{p \leq x^{1/4}} |S(x/p^2)| \ll \frac{x}{\sqrt{\log x}} \sum'_{p \leq x^{1/4}} 1/p^2 \ll \frac{x}{\sqrt{\log x}}.$$

As for the double sum over small primes, again by Landau,

$$\begin{aligned} (5.3) \quad \sum'_{p_1, p_2 \leq x^{1/4}} |S(x/p_1 p_2)| &= \sum'_{p_1, p_2 \leq x^{1/4}} \frac{cx}{p_1 p_2 \sqrt{\log(x/(p_1 p_2))}} (1 + O(1/\log(x/(p_1 p_2)))) \\ &= \sum'_{p_1, p_2 < x^{1/\log \log x}} \dots + 2 \cdot \sum'_{\substack{p_1 \leq x^{1/\log \log x} \\ p_2 \in [x^{1/\log \log x}, x^{1/4}]}} \dots + \sum'_{p_1, p_2 \in [x^{1/\log \log x}, x^{1/4}]} \dots \end{aligned}$$

Again by the analogue of Mertens' Theorem for arithmetic progressions,

$$\sum'_{p \in [x^{1/\log \log x}, x^{1/4}]} 1/p \ll \log \log \log x$$

and

$$\sum'_{p \leq x^{1/\log \log x}} 1/p = \frac{1}{2} \log \log x + O(\log \log \log x)$$

hence the contribution from the latter two sums in (5.3) is

$$\ll \frac{x}{\sqrt{\log x}} ((\log \log \log x) \cdot \log \log x + (\log \log \log x)^2) \ll \frac{x \cdot (\log \log \log x) \cdot \log \log x}{\sqrt{\log x}}$$

which is of the same size as the claimed error term.

Finally, yet again by Landau, and that  $\log(x/(p_1 p_2)) = \log x(1 + O(1/\log \log x))$  for  $p_1, p_2 \leq x^{1/\log \log x}$ , we find that

$$\begin{aligned} \sum'_{p_1, p_2 < x^{1/\log \log x}} |S(x/p_1 p_2)| &= \sum'_{p_1, p_2 < x^{1/\log \log x}} \frac{cx}{p_1 p_2 \sqrt{\log(x/(p_1 p_2))}} (1 + O(1/\log x)) \\ &= \sum'_{p_1, p_2 < x^{1/\log \log x}} \frac{cx}{p_1 p_2 \sqrt{\log x}} (1 + O(1/\log x))(1 + O(1/\log \log x)) \\ &= \frac{cx}{\sqrt{\log x}} \left( \sum'_{p < x^{1/\log \log x}} 1/p \right)^2 (1 + O(1/\log \log x)) \end{aligned}$$

Again by Mertens's for primes in progressions, we find that the main term equals

$$\frac{cx}{\sqrt{\log x}} \left( \frac{1}{2} \log \log x - O(\log \log \log x) \right)^2 \cdot (1 + O(1/\log \log x))$$

Dividing by  $|S(x)|$  and using Landau again, the main term is thus

$$\frac{1}{4} (\log \log x)^2 + O((\log \log x) \cdot \log \log \log x).$$

5.2.2. *Proof of Corollary 3.6.* Define a multiplicative function

$$f(n) := \frac{r_2(n)}{4 \cdot 2^{\omega_1(n)}}$$

If  $p \equiv 1 \pmod{4}$ , then  $f(p^e) = (e+1)/2$ ; if  $p \equiv 3 \pmod{4}$  then  $f(p^{2e+1}) = 0$  whereas for even exponents  $f(p^{2e}) = 1$ . Using Wirsing's Satz 1 again, we find (recall that  $\sum'_{p \leq x}$  denotes the over primes  $p \equiv 1 \pmod{4}$ ) that

$$\begin{aligned} \sum_{n \in S(x)} f(n) &= \sum_{n \leq x} f(n) \ll \frac{x}{\log x} \exp \left( \sum_{p \leq x} f(p)/p \right) = \frac{x}{\log x} \exp \left( \sum'_{p \leq x} 1/p + O(1) \right) \\ &\ll \frac{x}{\log x} \exp \left( \frac{1}{2} \log \log x + O(1) \right) \ll \frac{x}{(\log x)^{1/2}} \ll |S(x)| \end{aligned}$$

(here we again have used Mertens' Theorem for arithmetic progressions.) Chebyshev's inequality then implies that the number of  $n \in S(x)$  for which  $f(n) \geq \log \log \log n$  is  $o(|S(x)|)$ . In particular, we find that

$$2^{\omega_1(n)} \leq r_2(n)/4 \leq 2^{\omega_1(n)} \cdot \log \log \log n$$

holds for almost all  $n \in S(x)$ . Now, since Corollary 3.5 implies that  $\omega_1(n) = (1/2 + o(1)) \log \log n$  for almost all  $n \in S(x)$ , we find that

$$r_2(n) = 2^{(1/2+o(1)) \log \log n} = (\log n)^{(\log 2)/2+o(1)}$$

holds for almost all  $n \in S(x)$ .

## REFERENCES

- [1] G. Berkolaiko, J. P. Keating and B. Winn, *No quantum ergodicity for star graphs*, Comm. Math. Phys. Vol. 250, 259–285, 2004.
- [2] O. Bohigas, M. J. Giannoni and C. Schmit, *Characterization of Chaotic Quantum Spectra and Universality of Level Fluctuation Laws*, Phys. Rev. Lett. 52, 1–4, 1984.
- [3] E. Bogomolny, U. Gerland and C. Schmit, *Singular Statistics*, Phys. Rev. E, Vol. 63, No. 3, 2001.
- [4] Y. Colin de Verdière, *Ergodicité et fonctions propres du laplacien*. Comm. Math. Phys. Vol. 102, No. 3, 497–502, 1985.
- [5] L. Fainsilber, P. Kurlberg, and B. Wennberg, *Lattice points on circles and discrete velocity models for the Boltzmann equation*, SIAM J. Math. Anal. 37, No. 6, 1903–1922 (electronic), 2006.
- [6] E. Hecke, *Eine neue Art von Zetafunktionen und ihre Beziehungen zur Verteilung der Primzahlen*. Math. Z. 6, No. 1-2, 11–51, 1920.
- [7] J. P. Keating, J. Marklof and B. Winn, *Localised eigenfunctions in Šeba billiards*. J. Math. Phys. 51, no. 062101.
- [8] P. Kurlberg and H. Ueberschär, *Localised semiclassical measures for point scatterers on diophantine tori*. in preparation
- [9] E. Landau, *Handbuch der Lehre von der Verteilung der Primzahlen. 2 Bände*. Chelsea Publishing Co., New York, 1953. 2nd ed., With an appendix by Paul T. Bateman.
- [10] G. J. Rieger, *Aufeinanderfolgende Zahlen als Summen von zwei Quadraten*. Nederl. Akad. Wetensch. Proc. Ser. A 68 = Indag. Math., 27:208–220, 1965.
- [11] Z. Rudnick, H. Ueberschär, *Statistics of wave functions for a point scatterer on the torus*. Comm. Math. Phys., Vol. 316, No. 3, 763–782, 2012.
- [12] Z. Rudnick, H. Ueberschär, *On the eigenvalue spacing distribution for a point scatterer on a flat torus*. Annales Henri Poincaré, to appear
- [13] M. Ruzhansky, V. Turunen, *Pseudo-Differential Operators and Symmetries: Background Analysis and Advanced Topics*, Birkhäuser (2010).
- [14] P. Šeba, *Wave chaos in a singular quantum billiard*. Phys. Rev. Lett. 64, 1855–1858, 1990.

- [15] T. Shigehara, *Conditions for the appearance of wave chaos in quantum singular systems with a pointlike scatterer*, Phys. Rev. E, Vol. 50, No. 6, 1994.
- [16] H. Ueberschär, *Quantum Chaos for point scatterers on flat tori*, Phil. Trans. R. Soc. A 372 (2014), 20120509.
- [17] E. Wirsing, *Das Asymptotische Verhalten von Summen über multiplikative Funktionen*, Math. Ann. 143, 75–102, 1961.
- [18] N. Yesha, *Uniform distribution for the perturbed eigenfunctions of a point scatterer on the three-dimensional torus*, in preparation
- [19] S. Zelditch, *Uniform distribution of eigenfunctions on compact hyperbolic surfaces*, Duke Math. J. 55, 919–941, 1987.

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