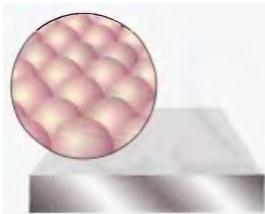


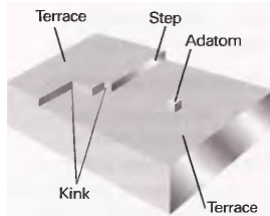
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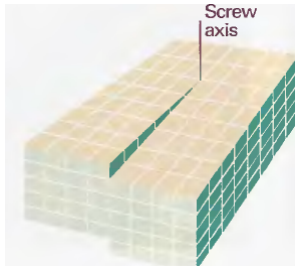


simple picture of a perfect crystal
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surface with defects





screw dislocation occurs where one region is pushed up through one or more unit cells relative to another region

The cut extends to the screw axis. As atoms lie along the step the dislocation rotates round the screw axis and is not annihilated

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1 m^2 of metal surface has $\approx 10^{19}$ atoms

\therefore each atom is struck about 10^8 times each second

Even if only a few collisions leave a molecule adsorbed to the surface, the time for which a freshly prepared surface remains clean is very short

Physisorption and chemisorption

Physisorption :

- long range, weak van der Waals interaction between adsorbate and substrate
- energy released on physisorption : $\mathcal{O}(\text{enthalpy of condensation})$
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insufficient to lead to bond breaking, so a physisorbed molecule retains its identity, although it is distorted on surface

Adsorbate	$\Delta_{ad} H^{\ominus}(\text{kJ/mol})$
CH ₄	-21
H ₂	-84
H ₂ O	-59
N ₂	-21

data at 298K

Chemisorption : molecules/atoms stick to surface by forming a chemical bond, and tend to find sites that maximize their coordination number with the substrate

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Adsorbate	Adsorbent		
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C ₂ H ₄	−427	−285	−243
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Except in special cases, chemisorption must be exothermic

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- These adsorbates also show order–disorder transitions when they are heated enough for thermal motion to overcome the particle–particle interactions, but not so much that they are desorbed

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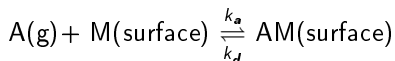
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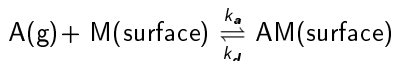
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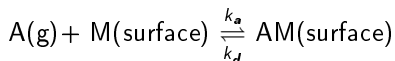
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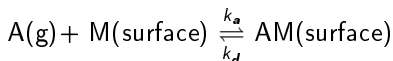
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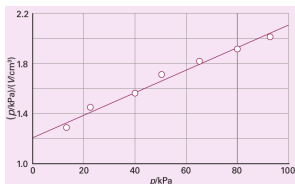
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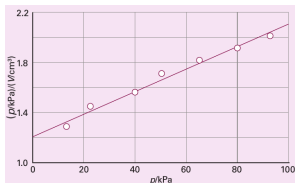
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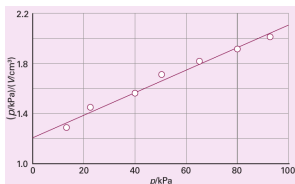
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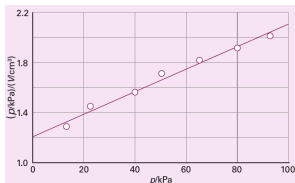
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write: $\theta = \frac{V}{V_{\infty}}$ and divide both sides by αV

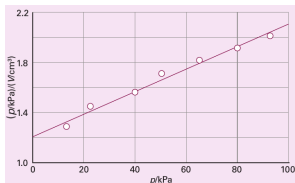
$$\text{then } \frac{p}{V} = \frac{p}{V_{\infty}} + \frac{1}{\alpha V_{\infty}}$$

p/kPa	13.3	26.7	40.0	53.3	66.7	80.0	93.3
$\frac{p}{V}$	1.30	1.44	1.57	1.69	1.81	1.92	2.02

slope=0.009 gives $V_{\infty} = 111\text{cm}^3$ (complete monolayer coverage)

intercept= 1.20

$$\text{gives } \alpha = \frac{1}{111 \times 1.20} = 7.51 \times 10^{-3} \text{ kPa}^{-1}$$



Ex. : adsorption of CO on charcoal at 273 K

p/kPa	13.3	26.7	40.0	53.3	66.7	80.0	93.3
V/cm ³	10.2	18.6	25.5	31.5	36.9	41.6	46.1

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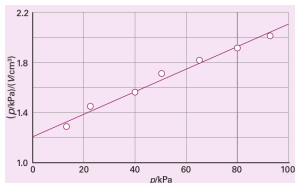
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volumes in definition of θ are those of the free gas measured under same conditions of temperature and pressure, not the volume the adsorbed gas occupies when attached to surface



adsorption with dissociation: A_2 adsorbs as $2A$

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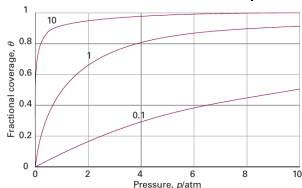
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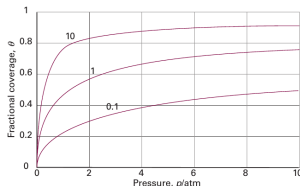
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$$\text{Equilibrium: } \theta = \frac{\sqrt{\alpha p}}{1 + \sqrt{\alpha p}},$$

surface coverage depends more weakly on pressure than for non-dissociative adsorption



dissociative



non-dissociative

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reduce the pressure and reduce the number of impacts on the surface

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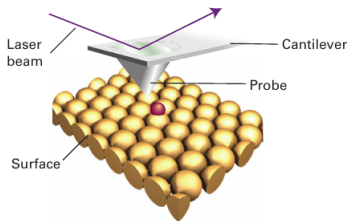
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\Rightarrow each surface atom hit once every $10^5 - 10^6$ or \sim once a day

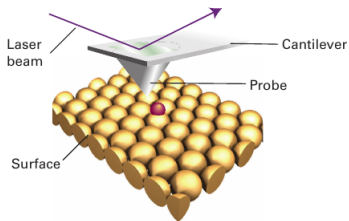
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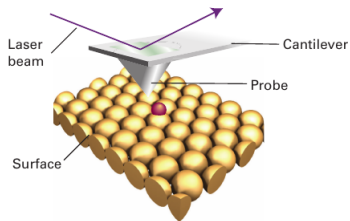


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sharpened tip attached to a cantilever is scanned across surface

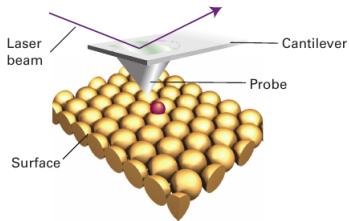


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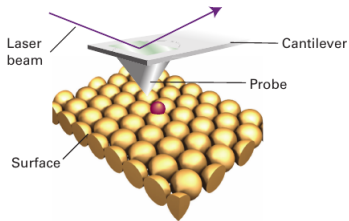
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 \therefore no current needs to pass between sample and probe, the technique can be applied to non-conducting surfaces and to liquid samples

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In 'non-contact', or 'tapping mode', the tip bounces up and down with specified frequency and never quite touches the surface

Amplitude of the tip's oscillation changes when it passes over species adsorbed on surface

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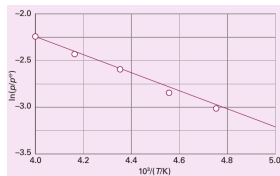
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$-\ln \frac{p}{p^\ominus}$	3.22	3.01	2.81	2.63	2.47	2.32



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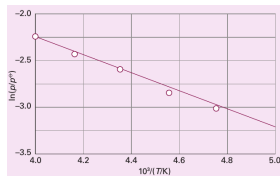
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$$\text{slope} = -0.904, \Delta_{ad} H^\ominus = - (0.904 \times 10^3 \text{ K}) \times R = -7.52 \text{ kJ mol}^{-1}$$



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- number of isotherms deal with cases where deviations from Langmuir isotherm are important

BET isotherm: accounts for multilayer adsorption

$$\frac{V}{V_{\text{mon}}} = \frac{cz}{(1-z)\{1-(1-c)z\}} \text{ with } z = \frac{p}{p^*}; c = e^{(\Delta_{\text{des}} H^\ominus - \Delta_{\text{vap}} H^\ominus)}$$

p^* = vapour pressure above a layer of adsorbate that is more than one molecule thick and which resembles pure bulk liquid

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Freundlich isotherm : accounts for substrate-substrate interactions

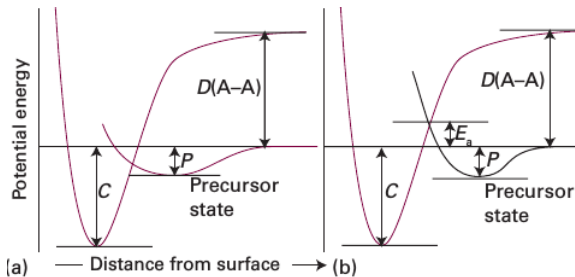
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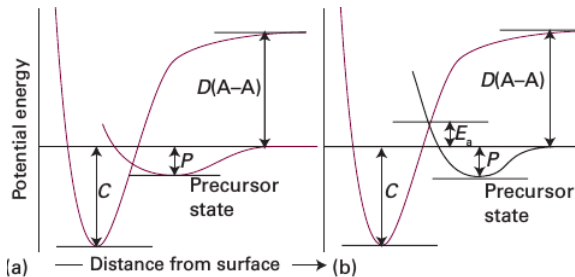
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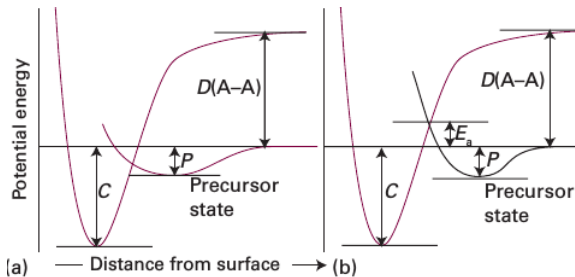


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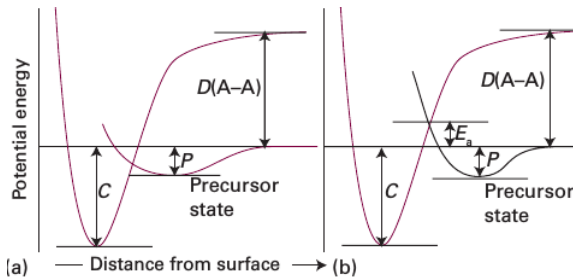
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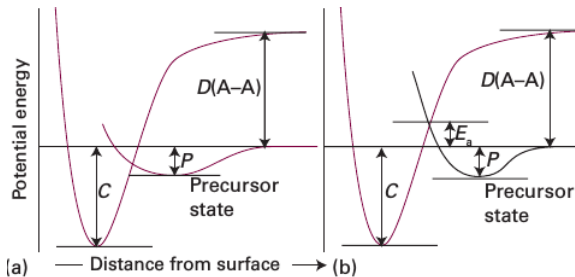
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Physisorption is usually fast, but can appear slow if adsn. takes place on porous medium

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ratio of rates of chemisorption on equal areas of the two faces at 250 K is

$$\frac{\text{Rate}(1)}{\text{Rate}(2)} = \frac{Ae^{-\frac{E_{act}(1)}{RT}}}{Ae^{-\frac{E_{act}(2)}{RT}}} = 11$$

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but for N_2 on Rhenium, $s < 10^{-2}$

Simple assumption : $s = (1 - \theta) s_0$,

where s_0 = sticking probability on a perfectly clean surface

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$$t_{\frac{1}{2}} = \frac{\ln 2}{k_d} = \tau_0 e^{\frac{E_d}{RT}}; \tau_0 = \frac{\ln 2}{A}$$

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$$t_{\frac{1}{2}} = \frac{\ln 2}{k_d} = \tau_0 e^{\frac{E_d}{RT}}; \tau_0 = \frac{\ln 2}{A}$$

If $\frac{1}{\tau_0} \approx$ vibrational frequency of weak particle-surface bond ($\approx 10^{12}\text{Hz}$) and $E_d \approx 25\text{kJ mol}^{-1}$, then residence half-lives $\approx 10\text{ns}$ at room temp.

Rate of desorption: Desorption is always activated because particles have to be lifted from the foot of a potential well

physisorbed particle vibrates in its shallow potential well, and might shake itself off the surface after a short time

The temperature dependence of first-order rate of departure can be expected to be Arrhenius-like,

with an activation energy for desorption, $E_d \approx$ enthalpy of physisorption

$$k_d = Ae^{-\frac{E_d}{RT}};$$

half-life for remaining on surface depends on T

$$t_{\frac{1}{2}} = \frac{\ln 2}{k_d} = \tau_0 e^{\frac{E_d}{RT}}; \tau_0 = \frac{\ln 2}{A}$$

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then residence half-lives $\approx 10\text{ns}$ at room temp.

Lifetimes close to 1 s are obtained by lowering T to $\approx 100\text{K}$

For chemisorption, with $E_d \approx 100\text{kJ mol}^{-1}$, and $\tau_0 = 10^{-14}\text{s}$, we get residence half-lives $\approx 3000\text{s}$ at room temp and 1 s at 350 K