

# Lecture 13

- ▶ Properties of Continuous-time Fourier Series
- ▶ Examples

# Properties of Continuous-time Fourier Series

**Notations** The pairing of a periodic signal with its Fourier series coefficients  $a_k$ s is denoted by the following notation:

$$x(t) \xleftrightarrow{\mathcal{FS}} a_k$$

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \quad a_k = \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt$$

► **Linearity:** If

$$x(t) \xleftrightarrow{\mathcal{FS}} a_k$$

$$y(t) \xleftrightarrow{\mathcal{FS}} b_k,$$

then

$$z(t) = Ax(t) + By(t) \xleftrightarrow{\mathcal{FS}} c_k = Aa_k + Bb_k.$$

**Proof:** By direct application of the definition of Fourier series.

► Time Shifting: If

$$x(t) \xleftrightarrow{\mathcal{FS}} a_k,$$

then

$$x(t - t_0) \xleftrightarrow{\mathcal{FS}} e^{-jk\omega_0 t_0} a_k = e^{-jk(2\pi/T)t_0} a_k.$$

**Proof:** The Fourier series coefficients  $b_k$  of  $y(t) = x(t - t_0)$  may be expressed as

$$b_k = \frac{1}{T} \int_T x(t - t_0) e^{-jk\omega_0 t} dt$$

Putting  $\tau = t - t_0$ , we get

$$b_k = \frac{1}{T} \int_T x(\tau) e^{-jk\omega_0(\tau+t_0)} d\tau = e^{-jk\omega_0 t_0} \underbrace{\left\{ \frac{1}{T} \int_T x(\tau) e^{-jk\omega_0 \tau} d\tau \right\}}_{a_k} = e^{-jk\omega_0 t_0} a_k.$$

Note that  $|b_k| = |a_k|$ .

► Time Reversal: If

$$x(t) \xleftrightarrow{\mathcal{FS}} a_k,$$

then

$$x(-t) \xleftrightarrow{\mathcal{FS}} a_{-k}.$$

Proof:

$$y(t) = x(-t) = \sum_{k=-\infty}^{\infty} a_k e^{-jk2\pi t/T} = \sum_{m=-\infty}^{\infty} a_{-m} e^{jm2\pi t/T} \quad [\text{Put } k = -m]$$

Therefore, the Fourier series coefficients  $b_k$  for  $y(t)$  are given by

$$b_k = a_{-k}.$$

Interesting facts:

- If  $x(t)$  is even i.e.,  $x(-t) = x(t)$ , then its Fourier series coefficients are also even, i.e.  $a_{-k} = a_k$ .
- Similarly, if  $x(t)$  is odd i.e.,  $x(-t) = -x(t)$ , then its Fourier series coefficients are also odd, i.e.  $a_{-k} = -a_k$ .

► Time Scaling: If

$$\underbrace{x(t)}_{\text{Period } T} \xleftrightarrow{\mathcal{FS}} a_k,$$

then

$$\underbrace{x(\alpha t)}_{\text{Period } T/\alpha} \xleftrightarrow{\mathcal{FS}} a_k.$$

Proof:

$$y(t) = x(\alpha t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 \alpha t} = \underbrace{\sum_{k=-\infty}^{\infty} a_k e^{jk(\alpha\omega_0)t}}_{\text{Fundamental frequency } \alpha\omega_0}$$

Therefore, the Fourier series coefficients  $b_k$  for  $y(t)$  are given by

$$b_k = a_k.$$

While the Fourier series coefficients have not changed, the Fourier series representation has changed due to the change of the fundamental frequency.

► **Multiplication:** If

$$x(t) \xleftrightarrow{\mathcal{FS}} a_k$$

$$y(t) \xleftrightarrow{\mathcal{FS}} b_k,$$

then

$$x(t)y(t) \xleftrightarrow{\mathcal{FS}} h_k = \sum_{l=-\infty}^{\infty} a_l b_{k-l}.$$

**Proof:**

$$\begin{aligned} x(t)y(t) &= \sum_{l=-\infty}^{\infty} a_l e^{jl\omega_0 t} \sum_{n=-\infty}^{\infty} b_n e^{jn\omega_0 t} = \sum_{l=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} a_l b_n e^{j(l+n)\omega_0 t} \\ &= \sum_{l=-\infty}^{\infty} \sum_{k=-\infty}^{\infty} a_l b_{k-l} e^{jk\omega_0 t} \quad [\text{By putting } k = l + n] \\ &= \sum_{k=-\infty}^{\infty} \underbrace{\left( \sum_{l=-\infty}^{\infty} a_l b_{k-l} \right)}_{h_k} e^{jk\omega_0 t} \end{aligned}$$

The Fourier series coefficients  $h_k$  is the discrete-time convolution of the sequences representing Fourier coefficients of  $x(t)$  and  $y(t)$ .

► Conjugation and Conjugate Symmetry: If

$$x(t) \xleftrightarrow{\mathcal{FS}} a_k$$

then

$$x^*(t) \xleftrightarrow{\mathcal{FS}} a_{-k}^*$$

Proof:

$$x^*(t) = \sum_{k=-\infty}^{\infty} a_k^* e^{-jk\omega_0 t} = \sum_{k=-\infty}^{\infty} a_{-k}^* e^{jk\omega_0 t}$$

Special cases:

- If  $x(t)$  is real i.e.,  $x(t) = x^*(t)$ , then the Fourier series coefficients will be *conjugate symmetric*:

$$a_{-k} = a_k^*.$$

- If  $x(t)$  is real, then  $a_0$  is real and  $|a_k| = |a_{-k}|$ .
- If  $x(t)$  is real and even, then the Fourier coefficients are also real and even.
- If  $x(t)$  is real and odd, then its Fourier coefficients are purely imaginary and odd.  
In specific,  $a_0 = 0$  if  $x(t)$  is real and odd.

## ► Parseval's Relation for Continuous-Time Periodic Signals:

If

$$x(t) \xleftrightarrow{\mathcal{FS}} a_k$$

then

$$\frac{1}{T} \int_T |x(t)|^2 dt = \sum_{k=-\infty}^{\infty} |a_k|^2.$$

Proof:

$$\begin{aligned} \frac{1}{T} \int_T |x(t)|^2 dt &= \frac{1}{T} \int_T x(t) x^*(t) dt = \frac{1}{T} \int_T \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} x^*(t) dt \\ &= \frac{1}{T} \sum_{k=-\infty}^{\infty} a_k \int_T x^*(t) e^{jk\omega_0 t} dt = \sum_{k=-\infty}^{\infty} a_k \underbrace{\left\{ \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt \right\}^*}_{a_k} = \sum_{k=-\infty}^{\infty} |a_k|^2 \end{aligned}$$

Note that  $|a_k|^2$  is the average power in the  $k$ th harmonic since:

$$\frac{1}{T} \int_T |a_k e^{jk\omega_0 t}|^2 dt = \frac{1}{T} \int_T |a_k|^2 dt = |a_k|^2.$$

Thus, the total average power in a periodic signal equals the sum of the average powers in all of its harmonic components.



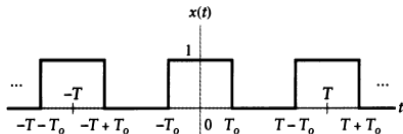
Property	Section	Periodic Signal	Fourier Series Coefficients
		$\left. \begin{array}{l} x(t) \\ y(t) \end{array} \right\} \begin{array}{l} \text{Periodic with period } T \text{ and} \\ \text{fundamental frequency } \omega_0 = 2\pi/T \end{array}$	$\begin{array}{l} a_k \\ b_k \end{array}$
Linearity	3.5.1	$Ax(t) + By(t)$	$Aa_k + Bb_k$
Time Shifting	3.5.2	$x(t - t_0)$	$a_k e^{-jk\omega_0 t_0} = a_k e^{-jk(2\pi/T)t_0}$
Frequency Shifting		$e^{jM\omega_0 t} x(t) = e^{jM(2\pi/T)t} x(t)$	$a_{k-M}$
Conjugation	3.5.6	$x^*(t)$	$a_{-k}^*$
Time Reversal	3.5.3	$x(-t)$	$a_{-k}$
Time Scaling	3.5.4	$x(\alpha t), \alpha > 0$ (periodic with period $T/\alpha$ )	$a_k$
Periodic Convolution		$\int_T x(\tau)y(t - \tau)d\tau$	$T a_k b_k$
Multiplication	3.5.5	$x(t)y(t)$	$\sum_{l=-\infty}^{+\infty} a_l b_{k-l}$
Differentiation		$\frac{dx(t)}{dt}$	$jk\omega_0 a_k = jk \frac{2\pi}{T} a_k$
Integration		$\int_{-\infty}^t x(\tau) d\tau$ (finite valued and periodic only if $a_0 = 0$ )	$\left(\frac{1}{jk\omega_0}\right) a_k = \left(\frac{1}{jk(2\pi/T)}\right) a_k$
Conjugate Symmetry for Real Signals	3.5.6	$x(t)$ real	$\begin{cases} a_k = a_{-k}^* \\ \Re\{a_k\} = \Re\{a_{-k}\} \\ \Im\{a_k\} = -\Im\{a_{-k}\} \\  a_k  =  a_{-k}  \\ \angle a_k = -\angle a_{-k} \end{cases}$
Real and Even Signals	3.5.6	$x(t)$ real and even	$a_k$ real and even
Real and Odd Signals	3.5.6	$x(t)$ real and odd	$a_k$ purely imaginary and odd
Even-Odd Decomposition of Real Signals		$\begin{cases} x_e(t) = \mathcal{E}\{x(t)\} & [x(t) \text{ real}] \\ x_o(t) = \mathcal{O}\{x(t)\} & [x(t) \text{ real}] \end{cases}$	$\begin{cases} \Re\{a_k\} \\ j\Im\{a_k\} \end{cases}$

Parseval's Relation for Periodic Signals

$$\frac{1}{T} \int_T |x(t)|^2 dt = \sum_{k=-\infty}^{+\infty} |a_k|^2$$

# Examples

**Example 1:** Determine the Fourier series representation of the square wave depicted in the following figure:

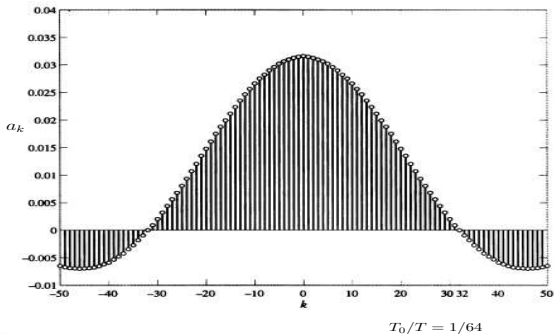
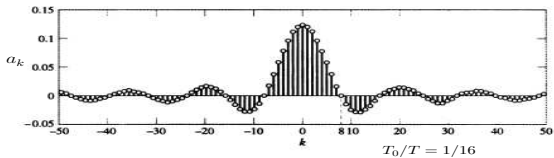
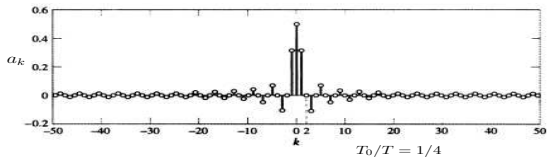


**Solution:**

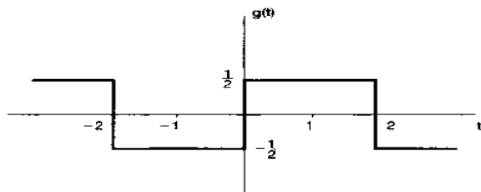
$$\begin{aligned} a_k &= \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\omega_0 t} dt = \frac{1}{T} \int_{-T_0}^{T_0} e^{-jk\omega_0 t} dt = -\frac{1}{Tjk\omega_0} e^{-jk\omega_0 t} \Big|_{-T_0}^{T_0} \\ &= \frac{2}{Tk\omega_0} \left( \frac{e^{jk\omega_0 T_0} - e^{-jk\omega_0 T_0}}{2j} \right), \quad k \neq 0 \\ &= \frac{2 \sin(k\omega_0 T_0)}{Tk\omega_0} = \frac{2T_0}{T} \operatorname{sinc} \left( k \frac{2T_0}{T} \right), \quad k \neq 0 \quad [\operatorname{sinc}(u) = \frac{\sin(\pi u)}{\pi u}] \end{aligned}$$

We have  $a_0 = \frac{1}{T} \int_{-T_0}^{T_0} 1 dt = \frac{2T_0}{T}$ . This can also be verified from

$$a_0 = \lim_{k \rightarrow 0} a_k = \lim_{k \rightarrow 0} \frac{2 \sin(k\omega_0 T_0)}{Tk\omega_0} = \frac{2T_0}{T}$$



**Example 2:** Determine the Fourier series representation of the square wave depicted in the following figure:



**Solution:** Compare this square wave with that of Example 1, i.e.,  $x(t)$ . Note that here we have period  $T = 4$  and  $T_0 = 1$ . The signal  $g(t)$  is given by  $g(t) = x(t - 1) - 1/2$ . The Fourier coefficients of  $x(t)$  are given by

$$a_k = \frac{2T_0}{T} \text{sinc}\left(k \frac{2T_0}{T}\right) = \frac{1}{2} \text{sinc}(k/2)$$

By the time-shifting property, the Fourier coefficients of  $x(t - 1)$  are given by

$$b_k = a_k e^{-jk\pi/2}.$$

The Fourier coefficients of the dc offset of  $g(t)$ , i.e.,  $-1/2$  is given by

$$c_k = \begin{cases} 0, & \text{for } k \neq 0 \\ -\frac{1}{2}, & \text{for } k = 0 \end{cases}$$

Invoking the linearity property, the Fourier coefficients of  $g(t)$  are given by

$$d_k = \begin{cases} \frac{1}{2}\text{sinc}(k/2)e^{-jk\pi/2}, & \text{for } k \neq 0 \\ 0, & \text{for } k = 0 \end{cases}$$

**Example 3:** Evaluate the Fourier series representation of the following signal:

$$x(t) = \sin(3\pi t) + \cos(4\pi t).$$

**Solution:** The time period of  $\sin(3\pi t)$  is  $T_1 = 2/3$  and the time period of  $\cos(4\pi t)$  is  $T_2 = 1/2$ . Therefore, the time period of  $x(t)$  is  $T = \text{lcm}(T_1, T_2) = 2$ ,  
 $\omega_0 = 2\pi/T = \pi$ .

$$\begin{aligned} x(t) &= \sin(3\pi t) + \cos(4\pi t) \\ &= \frac{1}{2j}e^{j(3)\pi t} - \frac{1}{2j}e^{j(-3)\pi t} + \frac{1}{2}e^{j(4)\pi t} + \frac{1}{2}e^{j(-4)\pi t} \end{aligned}$$

Therefore, the Fourier series coefficients are given by:

$$a_k = \begin{cases} \frac{1}{2} & k = \pm 4 \\ \frac{1}{2j} & k = 3 \\ -\frac{1}{2j} & k = -3 \\ 0 & \text{otherwise} \end{cases}$$

**Example 4:** Determine the time-domain signal represented by the following Fourier series coefficients:

$$a_k = \left(-\frac{1}{3}\right)^{|k|}, \omega_0 = 1.$$

**Solution:**

$$\begin{aligned} x(t) &= \sum_{m=-\infty}^{\infty} \left(-\frac{1}{3}\right)^{|k|} e^{jkt} = \sum_{m=0}^{\infty} \left(-\frac{1}{3}e^{jt}\right)^k + \sum_{m=1}^{\infty} \left(-\frac{1}{3}e^{-jt}\right)^k \\ &= \frac{1}{1 + \frac{1}{3}e^{jt}} - \frac{\frac{1}{3}e^{-jt}}{1 + \frac{1}{3}e^{-jt}} \\ &= \frac{8}{10 + 6 \cos(t)} \end{aligned}$$