Derivation of Euler-Lagrange Equation for ROF denoising algorithm

Kisung You kyoustat@gmail.com

June 25, 2019

1 Rudin-Osher-Fatemi TV-L2 Denoising Model

Since this article is not meant for comprehensive introduction to the problem, I will not thoroughly review theoretic cues. Depending on which viewpoint one takes upon the problem, a common framework of energy formulation on desirable denoised image u given a noisy picture u_0 is

$$E[u|u_0] = \lambda ||u - u_0||_2^2 + ||\nabla u||_1 \tag{1}$$

where the first term measures closeness of desirable solution to the noisy picture and the second is smoothness penalty - also known as regularization - that penalizes non-smoothness on pursued solution. λ , which originates from Lagrange multiplier, plays a balancing role between noise model and image's smoothness requirement. In this formulation, large λ value weighs closeness and small one focuses on smoothing effect. First-order condition for a minimizer u^* is that its derivative $\nabla E = 0$ and this Gâteaux derivative is achieved via Euler-Lagrange equation. A classical solution is to proceed with time-marching method,

$$u_t = -\nabla E[u|u_0] \tag{2}$$

to reach at minimum energy level. For reference, see Rudin, Osher, Fatemi (1992) "Nonlinear total variation based image removal algorithm" [1].

2 Basic Calculus

First we need is divergence theorem. Suppose Ω is a compact subset of \mathbb{R}^n with a piecewise smooth boundary denoted by $\partial\Omega$. If **F** is a continuously differentiable vector field which is defined on Ω , we have

$$\int_{\Omega} (\nabla \cdot \mathbf{F}) dV = \int_{\partial \Omega} (\mathbf{F} \cdot \mathbf{n}) dS, \tag{3}$$

where left part is volume integral over Ω which expresses, in some sense, change within the entire volume and right side is surface integral over the boundary $\partial\Omega$ of Ω , which somehow measures outwards flux on its surface.

In addition, we need following equality

$$\int_{\Omega} \nabla \cdot (u \nabla v) = \int_{\Omega} \nabla u \cdot \nabla v + \int_{\Omega} u \nabla \cdot (\nabla v) = \int_{\Omega} \nabla u \cdot \nabla v + \int_{\Omega} u (\nabla^{2} v). \tag{4}$$

3 Derivation of Euler-Lagrange equation

A standard Euler-Lagrange equation pertains to the concept of directional derivative in that we can achieve in via following

$$0 = \nabla \mathbf{E} = \frac{d}{dt}\Big|_{t=0} \mathbf{E}[u + tv] \tag{5}$$

and calculation is as follows.

$$\begin{split} \frac{d}{dt}\mathbf{E}[u+tv] &= \frac{d}{dt}\left(\lambda\int_{\Omega}(u+tv-u_0)^2 + \int_{\Omega}|\nabla(u+tv)|\right) \\ &= \lambda\frac{d}{dt}\left(\int_{\Omega}u^2 + t^2v^2 + u_0^2 + 2tuv - 2tu_0v - 2uu_0\right) \\ &+ \frac{d}{dt}\left(\int_{\Omega}|\nabla(u+tv)|\right) \\ &= \lambda\int_{\Omega}(2tv^2 + 2uv - 2u_0v) + \frac{d}{dt}\left(\int_{\Omega}|\nabla(u+tv)|\right). \end{split}$$

For the second term above, using chain rule, we have

$$\frac{d}{dt} \left(\int_{\Omega} |\nabla(u+tv)| \right) = \int_{\Omega} \frac{\nabla(u+tv)}{|\nabla(u+tv)|} \cdot \nabla v$$

since

$$\frac{d}{dx}|x| = \frac{x}{|x|} = sgn(x) = \begin{cases} 1 & \text{if } x > 0\\ -1 & x < 0\\ (-1, 1) & \text{otherwise.} \end{cases}$$

Combining these two and inserting t = 0, we have

$$0 = \nabla \mathbf{E} = 2\lambda \int_{\Omega} ((u - u_0)v) + \int_{\Omega} \left(\frac{\nabla u}{|\nabla u|} \cdot \nabla v \right)$$
 (6)

and we need to further apply basic calculus results from the above in that

$$\begin{split} \int_{\Omega} \left(\frac{\nabla u}{|\nabla u|} \cdot \nabla v \right) &= -\int_{\Omega} \left(\nabla \cdot \frac{\nabla u}{|\nabla u|} \right) v + \int_{\Omega} \nabla \cdot v \frac{\nabla u}{|\nabla u|} \\ &= -\int_{\Omega} \operatorname{div} \left(\frac{\nabla u}{|\nabla u|} \right) v + \int_{\Omega} \operatorname{div} \left(v \frac{\nabla u}{|\nabla u|} \right) \\ &= -\int_{\Omega} \operatorname{div} \left(\frac{\nabla u}{|\nabla u|} \right) v + \int_{\partial \Omega} \frac{v}{|\nabla u|} \left(\nabla u \cdot n \right). \end{split}$$

Therefore, if we set $\nabla u \cdot n = 0$, we have

$$0 = \int_{\Omega} \left((2\lambda(u - u_0)) - \operatorname{div}\left(\frac{\nabla u}{|\nabla u|}\right) \right) v \tag{7}$$

and since it should be satisfied with arbitrary direction v, we finally have

$$\begin{cases} (u - u_0) - \frac{1}{2\lambda} \operatorname{div}\left(\frac{\nabla u}{|\nabla u|}\right) = 0 & \text{in } \Omega\\ \frac{\partial u}{\partial n} = 0 & \text{on } \partial\Omega. \end{cases}$$
(8)

Then, time-marching methods can be written as following differential equation with Neumann boundary condition,

$$\begin{cases} \frac{\partial u}{\partial t} = (u_0 - u) + \frac{1}{2\lambda} \operatorname{div}\left(\frac{\nabla u}{|\nabla u|}\right) & \text{in } \Omega\\ \frac{\partial u}{\partial n} = 0 & \text{on } \partial\Omega. \end{cases}$$

References

[1] Leonid I. Rudin, Stanley Osher, and Emad Fatemi. Nonlinear total variation based noise removal algorithms. *Physica D: Nonlinear Phenomena*, 60(1-4):259–268, November 1992.