Institute for Theoretical Physics

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Nonlinear dynamics

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## **Solutions to Assignment 11**

## Exercise 1 [6 points]: Diffusion + chemistry: Turing instability

Consider the brusselator model

$$\partial_t u = D_u \partial_x^2 u + a - (b+1)u + u^2 v,$$
  
$$\partial_t v = D_u \partial_x^2 v + bu - u^2 v,$$

where u(x,t), v(x,t) are concentration fields.

- 1. Determine the homogeneous stationary base state (with  $u \neq 0, v \neq 0$ ).
- 2. Calculate the onset of the finite wavelength (Turing) instability, as discussed for the general case in the lecture.

HINT: In the eigenvalue problem, it is again enough to study  $\sigma=0$ . Solve for the control parameter b(q), minimize with respect to q and show that you get  $b_c^{\text{Turing}}=(1+a\sqrt{D_u/D_v})^2$  and  $q_c=\sqrt{\frac{a}{\sqrt{D_u}D_v}}$ .

3. To observe the finite wavelength pattern, the control parameter b needs to be beyond  $b_c^{\rm Turing}$  but below the threshold of the other instability occuring in the system, the oscillatory Hopf instability having  $b_c^{\rm Hopf}=1+a^2$  (see lecture). What restriction do  $b_c^{\rm Turing}$  and  $b_c^{\rm Hopf}$  hence imply for the diffusion coefficients of the two species?

## SOLUTION:

1. For the stationary *homogeneous* state, we can put all derivatives to zero and have to solve the following two equations simultaneously:

$$0 = f_u = a - (b+1)u + u^2v (1)$$

$$0 = f_v = bu - u^2 v. (2)$$

The only solution we find is  $\vec{w_0} = \begin{pmatrix} u_0 \\ v_0 \end{pmatrix} = \begin{pmatrix} a \\ b/a \end{pmatrix}$ .

2. The perturbation of the state is given by

$$\delta \dot{\vec{w}} = \begin{pmatrix} \delta \dot{u} \\ \delta \dot{v} \end{pmatrix} = \begin{pmatrix} \frac{\partial f_u}{\partial u} - D_u q^2 & \frac{\partial f_u}{\partial v} \\ \frac{\partial f_v}{\partial u} & \frac{\partial f_v}{\partial v} - D_v q^2 \end{pmatrix} \begin{pmatrix} \delta u \\ \delta v \end{pmatrix} = A \begin{pmatrix} \delta u \\ \delta v \end{pmatrix}, \tag{3}$$

with  $\delta \vec{w} = \vec{w_q} e^{iqx + \sigma(q)t}$ . Hence, a solution for  $\sigma$  can be found via the eigenvalue problem  $\sigma \vec{w_q} = A\vec{w_q}$  and hence the equation for  $\sigma$  is:

$$\sigma_{1,2}(q) = \frac{1}{2}TrA \pm \frac{1}{2}\sqrt{(TrA)^2 - 4detA}.$$
 (4)

Evaluating the homoegenous part of matrix A at  $\vec{w}_0$  yields:

$$A|_{\vec{w_0}} = \begin{pmatrix} -(b+1) + 2u_0v_0 & u_0^2 \\ b - 2u_0v_0 & -u_0^2 \end{pmatrix} \Big|_{\vec{w_0}} = \begin{pmatrix} b - 1 & a^2 \\ -b & -a^2 \end{pmatrix}, \tag{5}$$

At the onset of a possible finite  $q_c$  instability,  $\sigma = 0$ . Therefore, we just have to solve det A = 0 for the full system including the spacial derivatives. This leads to:

$$(b-1-D_uq^2)(-a^2-D_vq^2)+a^2b=0. (6)$$

Solving this for b(q) gives:

$$b(q) = \frac{a^2}{D_v q^2} + \frac{D_v a^2 D_u}{D_v} + D_u q^2. \tag{7}$$

Minimizing the function gives:

$$0 = \partial_q b(q) = -2\frac{a^2}{D_v q^3} + 2D_u q \tag{8}$$

$$q_c = \sqrt{\frac{a}{\sqrt{D_u D_v}}} \tag{9}$$

and plugging this back into the equation for b gives:

(10)

$$b_c = \left(1 + a\sqrt{\frac{D_u}{D_v}}\right)^2. \tag{11}$$

Note that for both instabilities b > 1 such that  $A_{uu} = b - 1 > 0$  and  $A_{vv} = -a^2 < 0$ . This means that u is "activated" while v is inhibited around the stationary homogeneous solution. One can also say that due to the bu-term in the equation for v, u activates v if it is small (and hence  $-u^2v$  is small).

3. b must be below the Hopf instability, i.e. one needs:

$$b < b_H = 1 + a^2. (12)$$

For Turing one needs:

$$b > b_T = \left(1 + a\sqrt{\frac{D_u}{D_v}}\right)^2. \tag{13}$$
 which leads to:

Hence we need  $b_T < b < b_H$  which leads to:

$$\left(1 + a\sqrt{\frac{D_u}{D_v}}\right)^2 < 1 + a^2 \rightarrow \sqrt{\frac{D_u}{D_v}} < \frac{\sqrt{1 + a^2 - 1}}{a}.$$
 (14)

The r.h.s. is of order a, hence finite and rather large, hence  $D_u \ll D_v$ , implying again short range activation and long range inhibition.

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