

**Physica Status Solidi A: Applications and Materials Science**  
**Influence of illumination spectrum on dissociation kinetic of iron-boron pairs in silicon**  
 --Manuscript Draft--

<b>Manuscript Number:</b>	
<b>Article Type:</b>	Research Article
<b>Corresponding Author:</b>	Oleg Olikh, Dr. Hab. Taras Shevchenko National University of Kyiv Kyiv, UKRAINE
<b>Corresponding Author E-Mail:</b>	olegolikh@knu.ua
<b>Order of Authors (with Contributor Roles):</b>	Oleg Olikh, Dr. Hab. (Conceptualization: Lead; Formal analysis: Lead; Investigation: Equal; Software: Lead; Validation: Equal; Visualization: Equal; Writing – original draft: Equal; Writing – review & editing: Equal)  Oleksandr Datsenko, PhD (Investigation: Equal; Validation: Equal; Writing – original draft: Equal; Writing – review & editing: Equal)  Serhiy Kondratenko (Project administration: Equal; Resources: Lead; Writing – original draft: Equal; Writing – review & editing: Equal)
<b>Section/Category:</b>	Gettering and Defect Engineering in Semiconductor Technology (GADEST 2024)
<b>Abstract:</b>	Herein, the results of an experimental study on the photo-dissociation kinetic of iron–boron (FeB) pairs in boron-doped Czochralski silicon using different light sources are reported. It was shown that the FeB dissociation rate depends not only on integrated light intensity and overall carrier generation rate, but also on spectral composition of illumination. The value of the material constant of dissociation K varies and has been determined to be within (1.5-3.8) fs. The investigation has revealed increase in the dissociation rate with an increase in photon energy. The results allowed us to conclude the dominant role of the recombination-enhanced defect reaction at the second stage of the pair dissociation.
<b>Suggested Reviewers:</b>	<p>Chang Sun Australian National University chang.sun@anu.edu.au His papers are devoted to defect in silicon (IEEE Journal of Photovoltaics, 2023, 13(4), pp. 524–534, Physica Status Solidi-RRL, 2021, 15(12), 2000520)</p> <p>Xuegong Yu Zhejiang University yuxuegong@zju.edu.cn He is skilled in properties of FeB pair in silicon {AIP ADVANCES 3, 082124 (2013), Journal of Crystal Growth, 2024, 630, 127602, International Journal of Photoenergy, 2015, 154574}</p> <p>Nabil Khelifati Centre de Recherche en Technologie des Semi-conducteurs pour l'Energetique khelifatinabil@crtse.dz He studies defects in silicon (Physica Status Solidi A, 2019, 216(17), 1900253, Materials Research Express, 2019, 6(5), 055907)</p> <p>Karsten Bothe Institute for Solar Energy Research Hamelin k.bothe@isfh.de He is a leading expert in the field of the influence of defects on the properties of silicon solar cells (Semiconductor Science and Technology, 2019, 34(3), 035030, IEEE Journal of Photovoltaics, 2021, 11(4), pp. 890–896, 9435415, Solar RRL, 2023, 7(13), 2300240,</p>

	<p>Solar RRL, 2024, 8(3), 2300873)</p> <p>Hele Savin Aalto University ville.vahanissi@aalto.fi</p> <p>His papers are devoted to iron and other defects in silicon solar cell (IEEE Journal of Photovoltaics, 2020, 10(6), pp. 1532–1537, 9157971, Physica Status Solidi A, 2019, 216(17), 1900253, )</p>
<b>Opposed Reviewers:</b>	
<b>Author Comments:</b>	
<b>Additional Information:</b>	
Question	Response
Please submit a plain text version of your cover letter here.	<p>Dear Editors,</p> <p>Enclosed with this letter you will find the electronic submission of manuscript entitled "Influence of illumination spectrum on dissociation kinetic of iron-boron pairs in silicon" by Oleg Olikh, Oleksandr Datsenko, and Serhiy Kondratenko.</p> <p>It is widely recognized, that defects significantly impact semiconductor properties. Therefore it's crucial to understand the parameters of defects and the mechanisms behind their alteration, which holds significant practical importance. Over several decades, extensive knowledge has been amassed regarding specific defects, with the iron-boron pair in silicon being a well-studied example. Therefore, the discovery of new findings or the resolution of contentious issues is particularly intriguing.</p> <p>It has been established that the rate of light-induced dissociation of FeB pairs is influenced by integrated illumination intensity, temperature, and the defect composition of the material. Our investigation found that the spectral composition of illumination is an additional important factor affecting dissociation efficiency. Specifically, we demonstrated that increased photon energy leads to higher photo-dissociation efficiency. The results indicate that the recombination-enhanced defect reaction is the more probable mechanism at the second stage of light-induced dissociation, as opposed to iron ion recharge. We are confident that this study, shedding new light on the long-studied FeB pair in silicon, will be of significant interest to your readers.</p> <p>The corresponding abstract has been accepted for poster presentation at the 20th Conference on Gettering and Defect Engineering in Semiconductor Technology (GADEST 2024, Reference ID 105868).</p> <p>This is an original paper which has not been simultaneously submitted as a whole or in</p>

	<p>part anywhere else.</p> <p>No elements of the work have been published in any form.</p> <p>No conflict of interest exists in the submission of this manuscript.</p>
	<p>We would very much appreciate if you would consider the manuscript for publication in the GADEST 2024 special issue in \emph{physica status solidi (a)}.</p>
	<p>Sincerely yours,</p> <p>Oleg Olikh and co-authors.</p> <p>Taras Shevchenko National University of Kyiv</p> <p>Kyiv 01601, Ukraine</p> <p>E-mail: <a href="mailto:olegolikh@knu.ua">olegolikh@knu.ua</a></p>
Do you or any of your co-authors have a conflict of interest to declare?	No. The authors declare no conflict of interest.
<b>Keywords:</b>	silicon; iron-boron pairs; light-induced dissociation; wavelength impact; dissociation rate

To: physica status solidi (a) Editorial Board  
Subject: Article Submit

Dear Editors,

Enclosed with this letter you will find the electronic submission of manuscript entitled “Influence of illumination spectrum on dissociation kinetic of iron-boron pairs in silicon” by Oleg Olikh, Oleksandr Datsenko, and Serhiy Kondratenko.

It is widely recognized, that defects significantly impact semiconductor properties. Therefore it's crucial to understand the parameters of defects and the mechanisms behind their alteration, which holds significant practical importance. Over several decades, extensive knowledge has been amassed regarding specific defects, with the iron-boron pair in silicon being a well-studied example. Therefore, the discovery of new findings or the resolution of contentious issues is particularly intriguing.

It has been established that the rate of light-induced dissociation of FeB pairs is influenced by integrated illumination intensity, temperature, and the defect composition of the material. Our investigation found that the spectral composition of illumination is an additional important factor affecting dissociation efficiency. Specifically, we demonstrated that increased photon energy leads to higher photo-dissociation efficiency. The results indicate that the recombination-enhanced defect reaction is the more probable mechanism at the second stage of light-induced dissociation, as opposed to iron ion recharge. We are confident that this study, shedding new light on the long-studied FeB pair in silicon, will be of significant interest to your readers.

The corresponding abstract has been accepted for poster presentation at the 20th Conference on Gettering and Defect Engineering in Semiconductor Technology (GADEST 2024, Reference ID 105868).

This is an original paper which has not been simultaneously submitted as a whole or in part anywhere else. No elements of the work have been published in any form. No conflict of interest exists in the submission of this manuscript.

We would very much appreciate if you would consider the manuscript for publication in the GADEST 2024 special issue in *physica status solidi (a)*.

Sincerely yours,  
Oleg Olikh and co-authors.  
Taras Shevchenko National University of Kyiv  
Kyiv 01601, Ukraine  
E-mail: olegolikh@knu.ua

# <sup>1</sup> Influence of illumination spectrum on dissociation kinetic of iron- <sup>2</sup> boron pairs in silicon

<sup>5</sup> Oleg Olikh\* Oleksandr Datsenko Serhiy Kondratenko

<sup>9</sup> Prof. O. Olikh, Dr. O. Datsenko, Prof. S. Kondratenko

<sup>10</sup> Taras Shevchenko National University of Kyiv, 64/13, Volodymyrska Street, 01601, Kyiv, Ukraine

<sup>11</sup> Email Address: olegolikh@knu.ua

<sup>13</sup> Keywords: *silicon, iron-boron pairs, light-induced dissociation, wavelength impact, dissociation rate*

<sup>15</sup> Herein, the results of an experimental study on the photo-dissociation kinetic of iron–boron (FeB) pairs in boron-doped Czochralski  
<sup>16</sup> silicon using different light sources are reported. It was shown that the FeB dissociation rate depends not only on integrated light  
<sup>17</sup> intensity and overall carrier generation rate, but also on spectral composition of illumination. The value of the material constant of  
<sup>18</sup> dissociation  $K$  varies and has been determined to be within  $(1.5 - 3.8) \times 10^{-15}$  s. The investigation has revealed increase in the disso-  
<sup>19</sup> ciation rate with an increase in photon energy. The results allowed us to conclude the dominant role of the recombination-enhanced  
<sup>20</sup> defect reaction at the second stage of the pair dissociation.

21

22

## <sup>23</sup> 1 Introduction

<sup>24</sup> Defects significantly impact semiconductor properties. Although minimizing device dimensions to nanome-  
<sup>25</sup> ters shifts some focus from extended to point defects, physical properties still rely heavily on the pres-  
<sup>26</sup> ence and distribution of these irregularities. Hence, many strategies for enhancing semiconductor struc-  
<sup>27</sup> tures, including radiation and temperature treatments or certain fabrication conditions, strive to de-  
<sup>28</sup> crease the defect concentration or neutralize its effects [1, 2, 3]. For instance, in the case of photovoltaic  
<sup>29</sup> devices, we must understand and optimize the carrier properties tied to defects and impurities [1]. Such  
<sup>30</sup> controlled alteration methods of the defective subsystem have been generalized under the term “defect  
<sup>31</sup> engineering” and are extremely important from a practical standpoint.

<sup>32</sup> Successful defect engineering hinges on an in-depth understanding of defect properties. Key factors are  
<sup>33</sup> defect formation energy, transition energy levels, self-compensating effects, nonradiative recombination  
<sup>34</sup> caused by defects, and the mechanism of reconstruction and diffusion [1]. Considering the extraordinary  
<sup>35</sup> diversity of possible intrinsic and impurity defects, complete information on all of them is lacking even  
<sup>36</sup> for silicon, which is the most studied semiconductor. Nevertheless, it must be noted that considerable  
<sup>37</sup> data have been amassed on silicon, and have a solid understanding of some defects [4].

<sup>38</sup> For instance, such defects are iron impurity, a common, detrimental, and often unavoidable contaminant  
<sup>39</sup> in photovoltaic silicon [3, 5], and iron-boron pair. Specifically, iron atoms are known to be at the inter-  
<sup>40</sup> stitial sites, and  $\text{Fe}_i^+$  are highly efficient recombination centers [6]. In p-type Si at room temperature,  
<sup>41</sup> iron atoms are almost predominantly bound into complexes with dopants (B, Ga, Al, In). This defect  
<sup>42</sup> demonstrates bistable behavior: the stable state is defined by the configuration in which the Fe occupies  
<sup>43</sup> the first nearest tetrahedral interstitial site closest to the substituent atom, whereas, in the metastable  
<sup>44</sup> configuration, Fe is at the second  $T_d$  interstitial site [7]. The energy levels associated with iron and its  
<sup>45</sup> complexes, as well as the respective carriers capture cross-sections, are well-established [4, 8]. Among the  
<sup>46</sup> acceptor-iron pairs, the complex FeB is the most thoroughly investigated, primarily due to the widespread  
<sup>47</sup> use of Si:B in the fabrication of various devices, such as solar cells. However, it is worth mentioning that  
<sup>48</sup> gallium is gaining increasing attention as an acceptor dopant whose incorporation, for instance, can help  
<sup>49</sup> to mitigate the light and elevated temperature-induced degradation [9].

<sup>50</sup> The dynamics of FeB pairs are also examined. It's established that FeB pairs can be dissociated through  
<sup>51</sup> illumination, minority carrier injection, and thermal treatment at 200 °C [10]. In the context of illumina-  
<sup>52</sup> tion, the dissociation rate  $R_d$  is influenced by the overall carrier generation rate  $G$  [11, 10, 12, 13]:

$$\text{62} \quad R_d = K \left( \frac{G}{N_{\text{FeB}}} \right)^2, \quad (1)$$

63

64

65

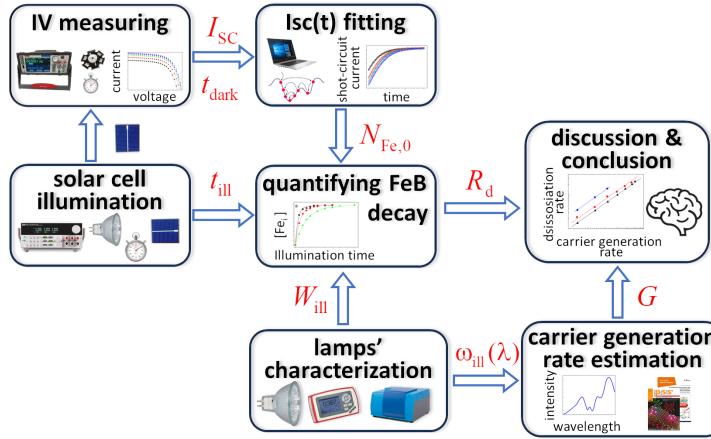


Figure 1: Investigation framework

where  $N_{\text{FeB}}$  is the pair concentration,  $K$  is the constant of material. To achieve almost complete dissociation of the FeB pairs, it is necessary for the illumination power to exceed  $0.1 \text{ W cm}^{-2}$  [14]. The dissociation process of FeB pairs by electron capture unfolds in two stages [15, 10]: the initial stage involves the neutralization of Fe and the elimination of the Coulombic attraction between the pair components. The mechanism of the second stage is contentious; it may involve either the recharge of the iron ion or the recombination-enhanced defect reaction (REDR) triggered by electron-hole recombination.

It should be noted that despite the extensive data on the properties of iron-related defects in silicon, intensive research persists. In particular, efforts focus on analyzing the impact of high-intensive illumination [16] or dopant compensation [17], alongside clarifying the second-stage mechanism of dissociation [5] or reassessing recombination parameters [18].

This study aims to investigate the effect of the light spectrum on the dissociation kinetics of FeB pairs in silicon. While pair dissociation is typically carried out using a halogen lamp [11, 5] or 904 nm laser [16, 10, 19], there is limited understanding of how the light source influences this process. By studying the impact of different illumination spectra on FeB dissociation, we aim to provide valuable insights for defect engineering and the efficient transformation of detrimental impurity iron atoms into a highly mobile interstitial state within the active region of a silicon device. Besides, such information, in our opinion, can help make the right choice between existing options for the second stage of pair decay.

In **Figure 1**, the main stages of the research are illustrated. First step was determination of the dissociation rate of FeB pairs under illumination with different integral intensities. Three light sources from different manufacturers were used (further details are described in Section 4). To measure number of interstitial iron atoms formed over fixed time under strong illumination the kinetics of short-circuit current was used. The result is presented in Subsection 2.1. Subsection 2.2 deals with estimating the carrier generation rate using spectra of sample illumination and considering the effects of light reflection, absorption by free carriers, and effective absorption depths. The obtained results showed that the efficiency of light-induced dissociation increases with decreasing photon wavelength — see Subsection 2.3. Finally, we conclude this paper in Section 3.

## 2 Results and Discussion

### 2.1 Dissociation rate determination

The equilibrium between free  $\text{Fe}_i$  and  $\text{Fe}_i\text{B}_{\text{Si}}$  is known to be determined by the following equations [20, 5, 10]



where  $R_a$  is the association rate. As a result, the concentration of unpaired interstitial iron atoms  $N_{\text{Fe}_i}$  depending on illumination time  $t_{\text{ill}}$  during light-induced dissociation can be described as follows [11, 12,

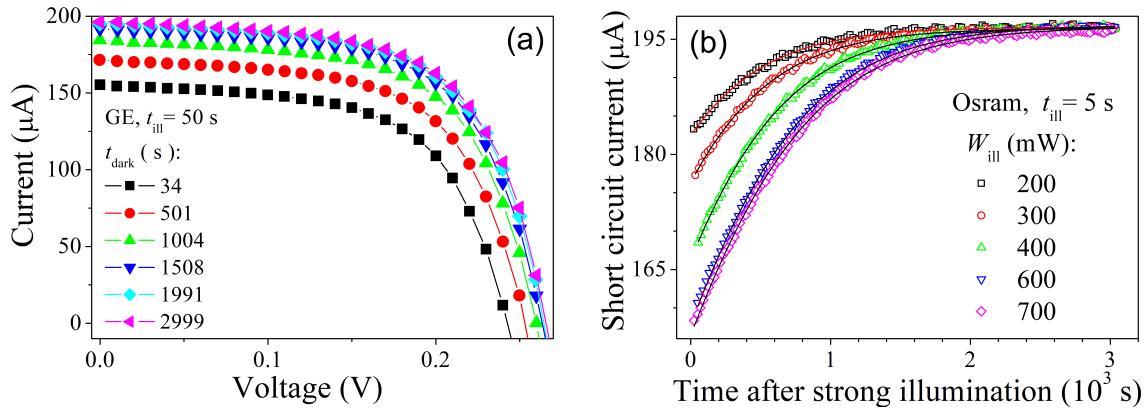


Figure 2: Typical current-voltage characteristics measured under low-intensive LED illumination at 940 nm across different periods following exposure to strong light (halogen lamp) (panel a) and short circuit current plotted as a function of the time after high-intensive illumination (panel b). The marks are the experimental results and the lines on panel b are the curves fitted according to [23, 21]. Light sources: GE (a), Osram (b).  $t_{\text{ill}}$ , s: 50 (a), 5(b).  $W_{\text{ill}} = 400$  mW (a).  $T = 340$  K.

$$N_{\text{Fe}_i}(t_{\text{ill}}) = \left( N_{\text{Fe},\text{eq}} - N_{\text{Fe},\text{tot}} \frac{R_d}{R_d + R_a} \right) \exp[-(R_d + R_a)t_{\text{ill}}] + N_{\text{Fe},\text{tot}} \frac{R_d}{R_d + R_a}, \quad (3)$$

where  $N_{\text{Fe},\text{tot}}$  is the total concentration of the impurity iron,  $N_{\text{Fe},\text{eq}}$  represents the concentration of unpaired interstitial iron atoms in the equilibrium state (in darkness,  $N_{\text{Fe},\text{eq}} = N_{\text{Fe}_i}(t_{\text{ill}} \leq 0)$ ). It's important to highlight that  $N_{\text{Fe},\text{eq}}$  is significantly influenced by temperature and the Fermi level location [20]. Specifically, in the case of p-type Si with a hole concentration of  $1.36 \times 10^{15}$  cm $^{-3}$  (which corresponds to the base of the structure under investigation), at a temperature of  $T = 300$  K,  $N_{\text{Fe},\text{eq}}$  constitutes merely about 1% of  $N_{\text{Fe},\text{tot}}$ , rendering it negligible for practical considerations. However, when the temperature rises to 340 K, the proportion of  $N_{\text{Fe},\text{eq}}$  increases to approximately 14.5%.

After the cessation of illumination, only the process of association occurs, and the time dependence of  $\text{Fe}_i$  concentration can be expressed as follows [20, 22]:

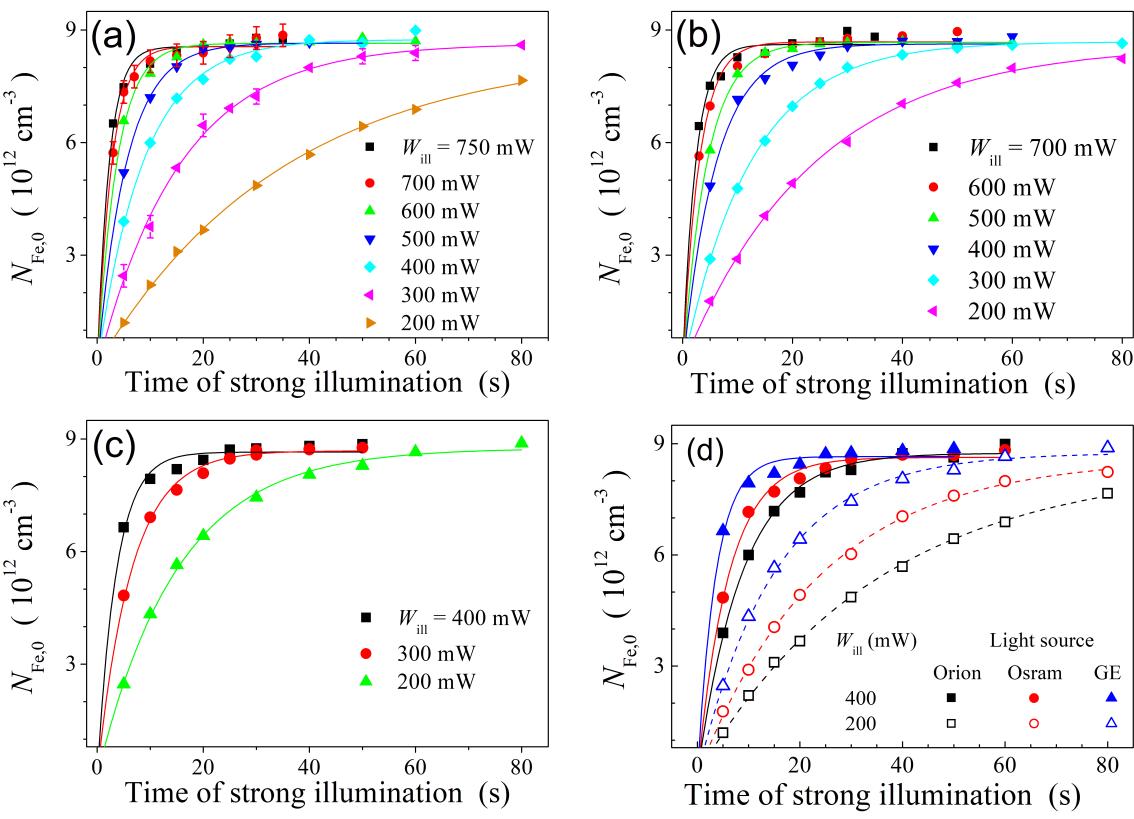
$$N_{\text{Fe}_i}(t_{\text{dark}}) = (N_{\text{Fe},0} - N_{\text{Fe},\text{eq}}) \times \exp(-R_a t_{\text{dark}}) + N_{\text{Fe},\text{eq}}, \quad (4)$$

where  $t_{\text{dark}}$  is the time after stopping of strong illumination,  $N_{\text{Fe},0}$  is the concentration of interstitial iron atoms formed after illumination,  $N_{\text{Fe},0} = N_{\text{Fe}_i}(t_{\text{dark}} = 0) = N_{\text{Fe}_i}(t_{\text{ill}})$ .

The study examined the dependence of  $N_{\text{Fe},0}$  in silicon solar cells on illumination time  $t_{\text{ill}}$  using different integral illumination intensities  $W_{\text{ill}}$  (200 – 750 mW) and light sources (three halogen lamps, labeled as Orion, Osram, and GE, and described in detail in Section 4). The experiments were conducted at a temperature of 340 K. The values of  $N_{\text{Fe},0}$  were determined using a methodology [23, 21] based on the fitting of the kinetics of short-circuit current  $I_{SC}$  under low-intensity monochromatic illumination. Specifically, after strong illumination with a duration of  $t_{\text{ill}}$ , the current-voltage characteristic ( $I$ - $V$ ) of the solar cell was measured every 21 seconds over a time  $t_{\text{dark}}$  interval of approximately 3000 seconds.

**Figure 2a** shows some typical  $I$ - $V$  curves. It can be seen that upon cessation of illumination, there exists a gradual augmentation in both the short-circuit current and the open-circuit voltage. This phenomenon is indicative of a decrease in the recombination activity of the defective subsystem, which is a result of the transition of interstitial iron to a bound state with an acceptor. Moreover, at the end of the measurement interval, the minute changes in the  $I$ - $V$  curves denote that the selected interval of 50 minutes is sufficient to complete the association.

**Figure 2b** illustrates the dependencies  $I_{SC}(t_{\text{dark}})$  after illumination with different intensities. As previously shown [21], the magnitude of the change in  $I_{SC}$  after the dark recovery period inherently correlates with the concentration of  $\text{Fe}_i$  formed as a result of light-induced dissociation of FeB pairs. From examining the presented data, it is evident that escalating  $W_{\text{ill}}$  leads to an augmentation in the dissociation efficiency. Concurrently, the recovery time remains insensitive to the illumination parameters, which conforms to expectations, given that the latter is determined by  $R_a$  — see Equation (4).



29 Figure 3: The relationships between the concentration of FeB pairs following intense illuminations of varying intensities  
30 and the illumination duration. Light source: Orion (a), Osram (b), GE (c). Panel d highlights variations in the dissociation  
31 of pairs induced by different light sources. The marks are the experimental results, the lines are the fitted curves using  
32 Equation (6).  $T = 340$  K.

33  
34 It should be noted that besides  $N_{Fe,0}$  values, the fitting of short-circuit current [23, 21] allows for the  
35 estimation of the energy of  $Fe_i$  migration  $E_m$  and bulk lifetime  $\tau_{other}$ , which arises from recombination  
36 channels other than Fe-related defects and intrinsic recombination. The obtained value  $E_m = (0.650 \pm$   
37  $0.005)$  eV coincides with that wellknown value [12, 24, 25]. This coincidence confirms that the investi-  
38 gated processes are indeed associated with rebuilding, as described by Equation (2). In turn, the value  
39 of  $E_m$  allows for the estimation of the recombination rate [10, 12, 24]:

$$43 R_a^{-1} = 5.7 \times 10^5 \frac{s}{K \text{cm}^3} \times \frac{T}{p} \exp\left(\frac{E_m}{kT}\right). \quad (5)$$

44 Thus, in our case,  $R_a = (1.68 \pm 0.03) \times 10^{-3} \text{ s}^{-1}$ .

45 Regarding the value of  $\tau_{other}$ , it was found to significantly exceed the lifetime associated with Shockley-  
46 Read-Hall (SRH) recombination on Fe-related defects (about  $2.2 \mu\text{m}$ ). Notably, according Möller *et al.*  
47 [10], such a condition is essential for the accurate determination of the constant  $K$ , which is included in  
48 Equation (1).

49 The dependencies of the concentration of interstitial atoms on illumination time are shown in **Figure 3**.  
50 From the data, it's evident that the pair dissociation rate is significantly influenced by the illumination  
51 intensity. This effect is consistent across all utilized light sources. Nonetheless, the  $W_{ill}$  value is not the  
52 exclusive determining factor for the pair dissociation rate, as demonstrated in Figure 3d. For instance,  
53 when using the GE source, pair dissociation occurs most efficiently. With Osram, the process unfolds  
54 more slowly, and illumination with Orion, under otherwise identical conditions, proves to be the least ef-  
55 fective in terms of altering the state of FeB pairs.

56 The experimentally obtained dependencies  $N_{Fe,0}(t_{ill})$  were fitted using the equation

$$57 N_{Fe,0}(t_{ill}) = A \exp(-t_{ill}/\tau_{dis}) + N_{Fe,fit}, \quad (6)$$

## 2.2 Carrier generation rate estimation

1 Table 1: Fitting results of experimental dependencies  $N_{\text{Fe},0}(t_{\text{ill}})$  using Equation (6) and defect parameter estimation using  
 2 Equations (7-8).

$W_{\text{ill}}$ [mW]	Light source	fitting parameters			defect parameters	
		$\tau_{\text{dis}}$ [s]	$N_{\text{Fe},\text{fit}}$ [ $10^{12} \text{ cm}^{-3}$ ]	$R^2$	$R_d$ [ $10^{-3} \text{ s}^{-1}$ ]	$N_{\text{Fe},\text{tot}}$ [ $10^{12} \text{ cm}^{-3}$ ]
750	Orion	$2.2 \pm 0.2$	$8.6 \pm 0.1$	0.993	450	8.6
700	Orion	$2.7 \pm 0.2$	$8.7 \pm 0.1$	0.995	370	8.7
	Osram	$2.4 \pm 0.2$	$8.6 \pm 0.1$	0.992	410	8.6
600	Orion	$3.7 \pm 0.2$	$8.65 \pm 0.06$	0.998	270	8.7
	Osram	$3.0 \pm 0.2$	$8.69 \pm 0.08$	0.995	330	8.7
500	Orion	$5.5 \pm 0.2$	$8.65 \pm 0.04$	0.999	180	8.7
	Osram	$4.5 \pm 0.1$	$8.7 \pm 0.1$	0.998	220	8.8
400	Orion	$8.8 \pm 0.3$	$8.74 \pm 0.06$	0.998	110	8.8
	Osram	$6.1 \pm 0.2$	$8.63 \pm 0.08$	0.997	160	8.7
	GE	$3.6 \pm 0.3$	$8.7 \pm 0.1$	0.996	280	8.7
300	Orion	$15.7 \pm 0.6$	$8.6 \pm 0.1$	0.998	62	8.8
	Osram	$12.4 \pm 0.1$	$8.69 \pm 0.02$	0.999	79	8.8
	GE	$6.5 \pm 0.2$	$8.69 \pm 0.05$	0.998	150	8.8
200	Orion	$35 \pm 3$	$8.5 \pm 0.3$	0.998	27	8.8
	Osram	$24 \pm 1$	$8.6 \pm 0.1$	0.999	40	8.9
	GE	$15.1 \pm 0.5$	$8.7 \pm 0.1$	0.999	65	8.8

23  
 24 where  $\tau_{\text{dis}}$  is the characteristic dissociation time, and  $N_{\text{Fe},\text{fit}}$  concentration of dissociated pairs at sat-  
 25 uration. The fitting results, shown in Figure 3 as lines and listed in the **Table 1**, include the coefficients  
 26 of determination  $R^2$ . The high values of  $R^2$  (greater than 0.99) confirm the suitability of the chosen ap-  
 27 proximation formula.

28 The comparison of Equations (3) and (6) reveals a relationship between the fitting parameters and de-  
 29 fect characteristics, specifically:

$$\tau_{\text{dis}}^{-1} = R_a + R_d, \quad (7)$$

$$N_{\text{Fe},\text{fit}} = N_{\text{Fe},\text{tot}} \frac{R_d}{R_d + R_a}. \quad (8)$$

30 The fitting parameters and recombination rate of  $1.68 \times 10^{-3} \text{ s}$  obtained were used to calculate the  $N_{\text{Fe},\text{tot}}$   
 31 and  $R_d$  values listed in the Table 1. Data suggest that regardless of the light source used and  $W_{\text{ill}}$ , the  
 32 observed total concentration of impurity iron atom is consistently  $N_{\text{Fe},\text{tot}} = (8.7 \pm 0.1) \times 10^{12} \text{ cm}^{-3}$ . This  
 33 stability supports the accuracy of the analysis. However, the FeB dissociation rate may vary significantly  
 34 for the same intensity value depending on the light source used.

35 According to Wijaranakula [20], at the specified value of  $N_{\text{Fe},\text{tot}}$ , the equilibrium (in darkness) concen-  
 36 trations of interstitial iron atoms  $N_{\text{Fe},\text{eq}}$  and FeB pairs  $N_{\text{FeB}}$  at  $T = 340 \text{ K}$  are  $1.3 \times 10^{12} \text{ cm}^{-3}$  and  
 37  $7.4 \times 10^{12} \text{ cm}^{-3}$ , respectively. The values of  $N_{\text{Fe},\text{eq}}$  and  $N_{\text{FeB}}$  were used to estimate the minority carrier  
 38 diffusion length  $L_n$  in the base of the used solar cell. It was assumed that the dominant recombination  
 39 processes are SRH recombination at  $\text{Fe}_i$  and FeB and intrinsic recombination. The required electron mo-  
 40 bility  $\mu_n$  value were taken from Klaassen [26], the capture cross sections and energy levels for  $\text{Fe}_i$  and  
 41 FeB from Rougieux *et al.* [8], the coefficients of band-to-band radiation recombination and Auger recom-  
 42 bination from Niewelt *et al.* [27] and Black & Macdonald [28], respectively. The calculated value was  
 43 found to be  $L_n = 80 \mu\text{m}$ , which is remarkably close to the value of  $86 \mu\text{m}$  obtained from the study of  
 44 temperature dependencies of short-circuit current — see Supplementary materials.

56  
 57 **2.2 Carrier generation rate estimation**

58 The dissociation rate of FeB pairs during light-induced decay is well known to be dependent on the car-  
 59 rier generation rate — see Equation (1). Our subsequent objective involved determining the values of  
 60  $G$  for various light sources. The measured dependencies of spectral intensity  $w_{\text{ill}}$  of illumination inci-  
 61 dent on the sample under diverse conditions are shown in **Figure 4**. It is crucial to highlight that our  
 62

## 2.2 Carrier generation rate estimation

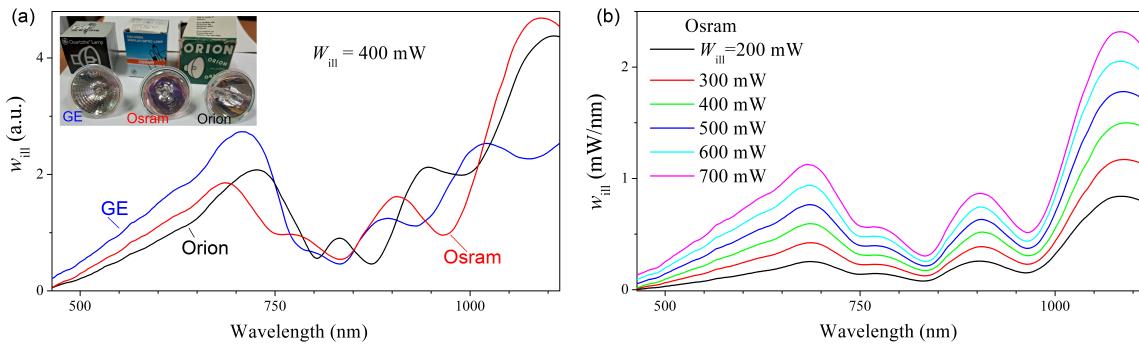


Figure 4: The spectra of sample illumination in the case of using different light sources with the same integral intensity  $W_{\text{ill}} = 400 \text{ mW}$  (panel a) and a single source (Osram) at various  $W_{\text{ill}}$  values (panel b). The inset shows photos of light sources.

focus is specifically on the light reaching the sample; hence, the spectrum is altered not only by the infrared transparency of the lamp reflector but also by absorption in the fiber utilized to transmit the light flux to the solar cell. Analogous modifications to the illumination spectra have been observed previously [29]. Figure 4a displays discrepancies in the illumination spectra obtained from different light sources, attributed to variations in the operational temperatures of the halogen lamps and differences in reflectors (photos of the lamps are in the inset of Figure 4a). It is important to note that the upper limit of the spectra in Fig 4 (1120 nm) is limited by the silicon bandgap, which, according to Passler [30], corresponds to 1.11 eV at 340 K. Furthermore, Figure 4b demonstrates the change in the Osram spectrum with integral intensity increasing. Notably, in addition to the expected increase in the curve's area, a minor spectrum shift towards shorter wavelengths is observed. These behaviour is typical for all used light sources.

Carrier generation rate was estimated as follows:

$$G = \int g(\lambda) d\lambda, \quad (9)$$

where spectral carrier generation rate  $g$

$$g = \frac{w_{\text{ill}} \lambda}{hc} \frac{(1 - R) A_{\text{bb}}}{S d_{\text{eff}}}, \quad (10)$$

where  $n_{\text{ph}} = \frac{w_{\text{ill}} \lambda}{hc}$  is the spectral photon flux,  $R$  is the reflectance,  $A_{\text{bb}}$  is the fraction of the band-to-band transitions,  $S$  is the illuminated area of sample,  $d_{\text{eff}}$  is the effective width of carrier generation.

In calculating the value of  $R$ , we employed an approach [31], which accounted for the presence of antireflective and passivating layers on the front surface of the sample, as well as the effects of multiple reflections. The resulting spectral dependence of  $R$  is shown in Figure S3 of the Supplementary materials.

The expression for the e-h pair generating fraction of the Lambertian absorptance in a solar cell can be written as [32]:

$$A_{\text{bb}}(\lambda) = \frac{\alpha_{\text{bb}}}{\alpha_{\text{bb}} + \alpha_{\text{fca}}} \frac{(1 - T_r)(1 + T_r)n_r^2}{n_r^2 - (n_r^2 - 1)T_r^2}, \quad (11)$$

with

$$\begin{aligned} T_r &= (1 - x) \exp(-x) + x^2 E_1(x), \\ x &= (\alpha_{\text{bb}} + \alpha_{\text{fca}})d, \\ E_1(x) &= \int_x^\infty t^{-1} \exp(-t) dt, \end{aligned}$$

where  $\alpha_{\text{bb}}$  is the absorption coefficient due to e-h pair generation by band-to-band transitions;  $\alpha_{\text{fca}}$  is the absorption coefficient due to free carrier absorption;  $n_r$  is the refractive index;  $d$  is the width of the device.

## 2.3 Effect of illumination spectrum on FeB pair decay

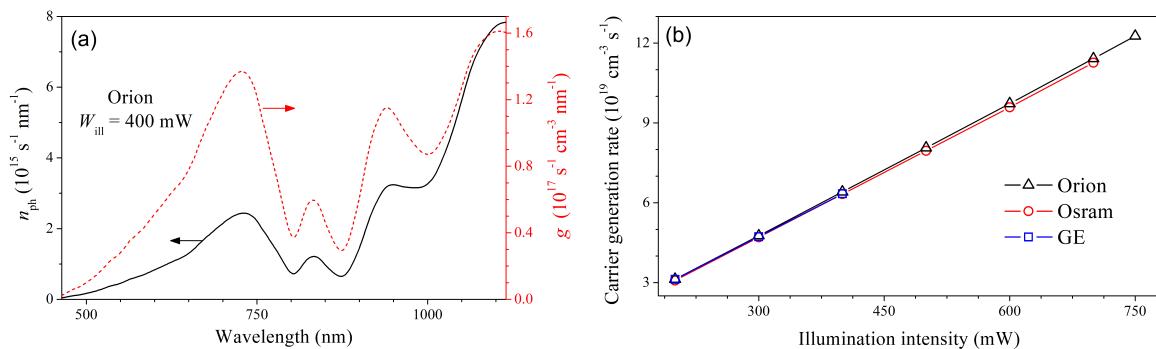


Figure 5: (a) Photon flux spectral density (left axis, solid line) and carrier generate rate spectral density (right axis, dashed line). Orion light source,  $W_{\text{ill}} = 400 \text{ mW}$ . (b) Dependencies of carrier generation rate on illumination intensity for different light sources.

In our calculations of  $A_{\text{bb}}$  by using Equations (11), we took  $\alpha_{\text{bb}}$  and  $n_r$  from Green[33],  $\alpha_{\text{fca}}$  from Baker-Finch *et al.* [34]. The spectral dependence of the fraction of the band-to-band transitions can be found in Supplementary materials (Figure S5).

When determining the carrier generation volume, we applied Bowden&Sinton approach [35] to thick silicon wafers, where the diffusion length or light absorption depth is significantly less than the sample thickness. In these cases, the distribution of carriers is heavily skewed toward the illuminated surface, rendering the use of the arithmetic mean of carrier concentration unsuitable. Consequently, the average values are computed using carrier concentration as a weighting function, and effective generation region width is determined as follows [35]:

$$d_{\text{eff}}(\lambda) = \frac{\left( \int_0^d \Delta n dx \right)^2}{\int_0^d \Delta n^2 dx}, \quad (12)$$

where  $\Delta n$  is the increase in minority carrier density due to illumination

$$\Delta n(x) = \frac{\alpha_{\text{bb}} n_{\text{ph}} L_n^2 q}{(\alpha_{\text{bb}}^2 L_n^2 - 1) k T \mu_n} \left[ \exp \left( -\frac{x}{L_n} \right) - \exp(-\alpha_{\text{bb}} x) \right]. \quad (13)$$

In calculations, we used  $L_n$  value, determined in Section 2.1. Some dependencies for different  $L_n$  values are shown in Figure S4 (Supplementary materials).

The consideration of dependencies  $R(\lambda)$ ,  $A_{\text{bb}}(\lambda)$ , and  $d_{\text{eff}}(\lambda)$  alters the spectral carrier generation rate compared to the spectral photon flux, leading to an increased contribution to e-h pairs generating from shorter wavelength light, as illustrated in Figure 5a. In Figure 5b, variations in the total carrier generation rate are depicted with increasing light intensity for different light sources. It is evident that differences exist between the light sources, yet the absolute contrasts in the magnitude of  $G$  for various light sources under  $W_{\text{ill}} = \text{const}$  conditions do not surpass 2 percent, with Orion registering the highest carrier generation rate value. Notably, without considering  $R$ ,  $A_{\text{bb}}$ , and  $d_{\text{eff}}$  and calculating  $G$  as the ratio of incident photons number to the total sample volume, the qualitative outcomes would remain quite similar. The main difference would lie in the absolute values of the carrier generation rate.

Thus, the discrepancies previously (Section 2.1) noticed in the value of  $R_d$  under identical illumination intensity levels cannot be attributed to variations in the carrier generation rate among different light sources, even when factoring in the quadratic dependency of the dissociation rate on  $G$ . Hence, there must be another underlying cause for these differences.

### 2.3 Effect of illumination spectrum on FeB pair decay

The dependencies  $R_d(G)$  in logarithmic scale are presented in Figure 6a. The linear patterns observed in the data indicate a clear power law relationship between the dissociation rates of FeB pairs and the

## 2.3 Effect of illumination spectrum on FeB pair decay

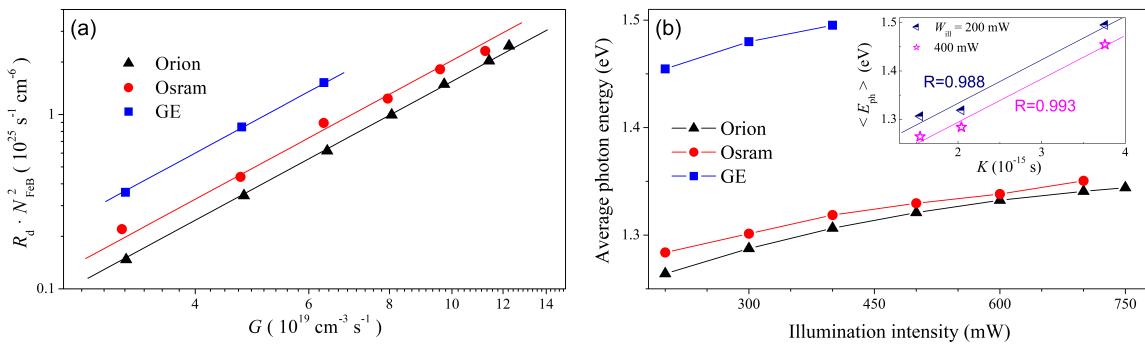


Figure 6: (a) FeB pair dissociation rate plotted as  $R_d \cdot N_{FeB}^2$  over the light induced generation rate. The solid lines show the quadratic dependence according to Equation (1). (b) Dependencies of average photon energy on illumination intensity for different light sources. The inset shows prefactor  $K$  vs average photon energy for the different light sources and illumination intensities. The lines are linear fitting curves. Coefficients of correlation are shown as well.

carrier generation rate. The figure also exhibits the fit results using Equation (1). High correlation coefficients exceeding 0.998 validate the applicability of the quadratic dependence expressed by Equation (1) to our data. It should be noted that Khelifati *et al.* [12] stipulate the inclusion of  $R_d(1 + \tau_{FeB}/\tau_{other})^2$  (where  $\tau_{FeB}$  is the lifetime associated with recombination on FeB pairs) on the left-hand side of Equation (1) rather than  $R_d$ . However, in our cases  $\tau_{other} \gg \tau_{FeB}$ , as previously acknowledged in Section 2.1, this additional multiplier can be disregarded.

The prefactor  $K$  values determined from the fitting are  $3.8 \times 10^{-15} \text{ s}$  for the GE light source,  $2.0 \times 10^{-15} \text{ s}$  for the Osram, and  $1.5 \times 10^{-15} \text{ s}$  for the Orion source.  $K$  serves as an important parameter linked to the phenomenon of FeB pairs' dissociation by illumination [12], and the values obtained in this study are compared with those ( $4.2 \times 10^{-17} - 5 \times 10^{-15} \text{ s}$ ) previously presented [11, 10, 12]. It is essential to note that in prior research, disparities in the constant  $K$  value for diverse samples were ascribed to variations in defect composition and the presence of alternative recombination channels apart from iron-related defects [11, 10]. Early, the dependence of  $R_d$  on temperature was reported as well [36]. However, in our case, distinct  $K$  values were obtained for the same structure under identical conditions, such as temperature and integrated light intensity, with differences in the illumination spectra only.

In other words, the obtained data indicate that in the analysis of light-induced dissociation of FeB pairs, it is necessary to consider not only the quantity of photo-generated excess charge carriers but also the energies of the photons that lead to their appearance. For such an energy characterization of light sources, we used the average photon energy  $\langle E_{ph} \rangle$ :

$$\langle E_{ph} \rangle = \frac{\int \frac{hc}{\lambda} n_{ph}(\lambda) d\lambda}{\int n_{ph}(\lambda) d\lambda}. \quad (14)$$

The summary of the results concerning the  $\langle E_{ph} \rangle$  values is shown in Figure 6b. In particular, it demonstrates the shift of the emission spectrum of light sources towards the short-wavelength region with an increase in the  $W_{ill}$  value, as illustrated in Figure 4a. And a clear conclusion can be drawn upon comparing the data in Figure 3,5b, 6a, 6b and Table 1: with rising average photon energy, the light-induced dissociation of FeB pairs becomes more pronounced. Specifically, the constant  $K$  increases, the dissociation rate  $R_d$  escalates, and correspondingly, the illumination time necessary for a complete complex decay decreases. Consequently, for the dissociation of FeB pairs, the energy expended during the thermalization of non-equilibrium carriers also holds significance.

The obtained results offer some conclusions about the mechanism of FeB dissociation. As discussed in the literature and previously mentioned, two possible ways of the second decay stage are typically considered: REDR and the recharge of the iron ion. REDR arises from strong electron-lattice coupling at the defect site and involves the utilization of local vibrational energy to promote pair dissociation [10, 5, 14]. The observed correlation between dissociation rate and photon energy in this study supports the REDR process. Specifically, as photon energy increases, the production of non-equilibrium phonons during thermalization also rises. Furthermore, the increase in  $R_d$  value, as found in the experiment, signifies

1 the active involvement of these quasi-particles in the dissociation of FeB pairs. Notably, recent research  
 2 [5] focusing on a detailed analysis of the dissociation and association reactions of the iron-boron pairs  
 3 similarly concluded the predominant role of REDR processes.  
 4

5

### 6 3 Conclusion

7

8 The effect of the illumination spectra on the dissociation of FeB pairs in p-Si was investigated in this pa-  
 9 per. We reported the results of a experimental systematic study on FeB dissociation rate in solar cell  
 10 based on Cz-Si, which were carried out using different light source and illumination intensities.  
 11

12 The findings showed that the time required for a total dissociation of the FeB pairs not only becomes  
 13 shorter with the increase of illumination intensity but also significantly depends on light source. As a re-  
 14 sult, the value of material constant  $K$  varies and has been determined to be  $3.8 \times 10^{-15}$  s,  $2.0 \times 10^{-15}$  s,  
 15 and  $1.5 \times 10^{-15}$  s for three different light sources. The study of the illumination spectrum led to the con-  
 16 clusion that the efficiency of FeB photo-dissociation increases with a decrease in light wavelength. The  
 17 observation of an increase in the dissociation rate with photon energy indicates that the REDR effect  
 18 prevails as the dominant factor during the second stage of light-induced dissociation of FeB pairs. Fur-  
 19 thermore, the obtained results could help to develop defect engineering procedures for effectively con-  
 20 verting iron impurities in silicon into high-mobility states, which could significantly impact semiconduc-  
 21 tor technology.  
 22

23

24

### 25 4 Experimental Section

26

27 The  $n^+$ - $p$ - $p^+$ -Si samples were used in the experiment. The structure was fabricated from a  $380\text{ }\mu\text{m}$  thick  
 28  $p$ -type boron-doped Czochralski silicon (100) wafer with hole concentration  $p = 1.36 \times 10^{15}\text{ cm}^{-3}$ . The  
 29  $n^+$  emitter with sheet resistance of about  $20 - 30\text{ }\Omega/\square$  and thickness of  $0.7\text{ }\mu\text{m}$  was formed by phos-  
 30 phorus diffusion. The anti-recombination isotype barrier was created by using a  $p^+$  layer ( $10 - 20\text{ }\Omega/\square$ ,  
 31  $0.6\text{ }\mu\text{m}$ ) formed by boron diffusion. On the front surface, the antireflective and passivating  $\text{SiO}_2$  (40 nm)  
 32 and  $\text{Si}_3\text{N}_4$  (30 nm) layers were formed. The solid and grid Al contacts were formed by magnetron sput-  
 33 tering on the rear and front surfaces respectively.  
 34

35 Three powerful halogen lamps from different manufacturers were used for sample illumination and were  
 36 employed for the light-induced dissociation of FeB pairs:  
 37

38

- 39 • Orion Haltlichtspiegel 52240.0, 24 V, 200 W (labeled Orion in the paper);  
 40
- 41 • Osram 64653 HLX ELC, 24 V, 250 W (Osram);  
 42
- 43 • General Electric 43537 H271, 20 V, 150 W (GE);  
 44
- 45

46 The light sources were powered by the DC Power Supply ITECH IT6332B, which allowed setting the  
 47 current passing through the lamp with an accuracy of up to 1 mA.  
 48

49 The illumination originating from the sources was transmitted to the sample via fiber. The sources' emis-  
 50 sion underwent calibration at the fiber output, thereby enabling a direct assessment of the light flux in-  
 51 cident on the sample. The relationships between the integral illumination intensity and the current con-  
 52 sumed by the sources were measured using an optical power and energy meter Thorlabs PM100D and a  
 53 high-resolution sensor S401C. The illumination spectra measurements were performed using spectrome-  
 54 ter IKC-6 with germanium photodiode. When analyzing the spectra, the spectral sensitivity of the pho-  
 55 todiode was considered.  
 56

57 The current-voltage characteristics were measured using a Keithley 2450 source meter and low-intensity  
 58 monochromatic light source (light-emitting diode SN-HPIR940nm-1W with light wavelength 940 nm  
 59 and intensity of about  $400\text{ }\mu\text{W}$ ). The LED radiation intensity was stabilized by utilizing the W1209 ther-  
 60 mostat and a power supply regulated by a circuit incorporating positive feedback and digital control.  
 61

62

63

64

65

1 The measurements were carried out at a temperature of 340 K. The sample temperature was driven by a  
2 thermoelectric cooler, controlled by an STS-21 sensor, and maintained constant through a PID algorithm  
3 embedded in the software that serves the experimental setup.  
4

## 5 Supporting Information

6 Supporting Information is available from the Wiley Online Library or from the author.  
7

## 8 Acknowledgements

9 The authors are grateful for the help with calculating the reflectance by solar cells to Prof. Vitaliy Kosty-  
10 lyov.  
11

## 12 Conflict of Interest

13 The authors declare no conflict of interest.  
14

## 15 Data Availability Statement

16 The data that support the findings of this study are available from the corresponding author upon rea-  
17 sonable request.  
18

19

20

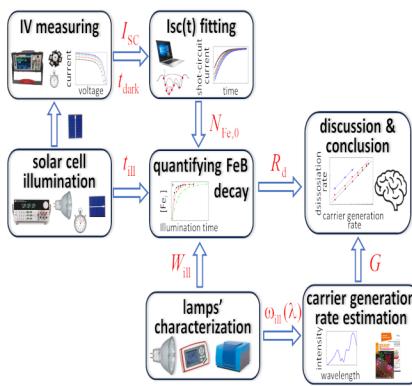
## 21 References

- 22
- 23 [1] X. Cai, S.-H. Wei, *J. Appl. Phys.* **2023**, *134*, 22 220901.  
24
- 25 [2] J. Vobecky, *Phys. Status Solidi A* **2021**, *218*, 23 2100169.  
26
- 27 [3] J. Frascaroli, P. Monge Roffarello, I. Mica, *Phys. Status Solidi A* **2021**, *218*, 23 2100206.  
28
- 29 [4] M. K. Juhl, F. D. Heinz, G. Coletti, D. Macdonald, F. E. Rougiewux, F. Schindle, T. Niewelt, M. C.  
30 Schubert, In *2018 IEEE 7th World Conference on Photovoltaic Energy Conversion (WCPEC) (A*  
31 *Joint Conference of 45th IEEE PVSC, 28th PVSEC & 34th EU PVSEC)*. **2018** 0328–0332.  
32
- 33 [5] C. Sun, Y. Zhu, M. Juhl, W. Yang, F. Rougiewux, Z. Hameiri, D. Macdonald, *Phys. Status Solidi*  
34 *RRL* **2021**, *15*, 12 2000520.  
35
- 36 [6] E. Weber, *Appl. Phys. A* **1983**, *30*, 1 1.  
37
- 38 [7] H. Nakashima, T. Sadoh, T. Tsurushima, *Phys. Rev. B* **1994**, *49*, 24 16983.  
39
- 40 [8] F. E. Rougiewux, C. Sun, D. Macdonald, *Sol. Energy Mater. Sol. Cells* **2018**, *187* 263 .  
41
- 42 [9] L. Ning, L. Song, J. Zhang, *J. Alloys Compd.* **2022**, *912* 165120.  
43
- 44 [10] C. Möller, T. Bartel, F. Gibaja, K. Lauer, *J. Appl. Phys.* **2014**, *116*, 2 024503.  
45
- 46 [11] L. J. Geerligs, D. Macdonald, *Appl. Phys. Lett.* **2004**, *85*, 22 5227.  
47
- 48 [12] N. Khelifati, H. S. Laine, V. Vähänissi, H. Savin, F. Z. Bouamama, D. Bouhafs, *Phys Status Solidi*  
49 *A* **2019**, *216*, 17 1900253.  
50
- 51 [13] S. Herlufsen, D. Macdonald, K. Bothe, J. Schmidt, *physica status solidi (RRL) – Rapid Research*  
52 *Letters* **2012**, *6*, 1 1.  
53
- 54 [14] D. H. Macdonald, L. J. Geerligs, A. Azzizi, *J. Appl. Phys.* **2004**, *95*, 3 1021.  
55
- 56 [15] L. Kimerling, J. Benton, *Physica B+C* **1983**, *116*, 1 297.  
57
- 58 [16] X. Zhu, D. Yang, X. Yu, J. He, Y. Wu, J. Vanhellemont, D. Que, *AIP Adv.* **2013**, *3*, 8 082124.  
59
- 60 [17] X. Zhu, X. Yu, P. Chen, Y. Liu, J. Vanhellemont, D. Yang, *Int. J. Photoenergy* **2015**, *2015* 154574.  
61
- 62 [18] T. T. Le, Z. Zhou, A. Chen, Z. Yang, F. Rougiewux, D. Macdonald, A. Liu, *J. Appl. Phys.* **2024**,  
63 *135*, 13 133107.  
64
- 65

- 1 [19] K. Lauer, C. Möller, D. Debbih, M. Auge, D. Schulze, In *Gettering and Defect Engineering in Semi-*  
2 *conductor Technology XVI*, volume 242 of *Solid State Phenomena*. Trans Tech Publications Ltd,  
3 **2016** 230–235.
- 4
- 5 [20] W. Wijaranakula, *J. Electrochem. Soc.* **1993**, *140*, 1 275.
- 6
- 7 [21] O. Olikh, V. Kostylyov, V. Vlasiuk, R. Korkishko, Y. Olikh, R. Chupryna, *J. Appl. Phys.* **2021**,  
8 *130*, 23 235703.
- 9
- 10 [22] J. D. Murphy, K. Bothe, M. Olmo, V. V. Voronkov, R. J. Falster, *J. Appl. Phys.* **2011**, *110*, 5  
11 053713.
- 12
- 13 [23] O. Olikh, V. Kostylyov, V. Vlasiuk, R. Korkishko, R. Chupryna, *J. Mater. Sci.: Mater. Electron.*  
14 **2022**, *33*, 16 13133.
- 15
- 16 [24] J. Tan, D. Macdonald, F. Rougier, A. Cuevas, *Semicond Sci. Technol.* **2011**, *26*, 5 055019.
- 17
- 18 [25] D. Macdonald, A. Cuevas, L. J. Geerligs, *Appl. Phys. Lett.* **2008**, *92*, 20 202119.
- 19
- 20 [26] D. Klaassen, *Solid-State Electron.* **1992**, *35*, 7 953.
- 21
- 22 [27] T. Niewelt, B. Steinhauser, A. Richter, B. Veith-Wolf, A. Fell, B. Hammann, N. Grant, L. Black,  
23 J. Tan, A. Youssef, J. Murphy, J. Schmidt, M. Schubert, S. Glunz, *Sol. Energ. Mat. Sol.* **2022**, *235*  
24 111467.
- 25
- 26
- 27 [28] L. E. Black, D. H. Macdonald, *Sol. Energ. Mat. Sol.* **2022**, *234* 111428.
- 28
- 29 [29] M. Libra, V. Poulek, P. Kourim, *Research in Agricultural Engineering* **2017**, *63*, 1 10.
- 30
- 31 [30] R. Pässler, *Phys. Rev. B* **2002**, *66* 085201.
- 32
- 33 [31] N. Klyui, V. Kostylyov, A. Rozhin, V. Gorbulik, V. Litovchenko, M. Voronkin, N. Zaika, *Opto-*  
34 *Electr. Rev.* **2000**, *8*, 4 402.
- 35
- 36 [32] S. Schäfer, R. Brendel, *IEEE J. Photovolt.* **2018**, *8*, 4 1156.
- 37
- 38 [33] M. A. Green, *Prog. Photovoltaics Res. Appl.* **2022**, *30*, 2 164.
- 39
- 40 [34] S. C. Baker-Finch, K. R. McIntosh, D. Yan, K. C. Fong, T. C. Kho, *J. Appl. Phys.* **2014**, *116*, 6  
41 063106.
- 42
- 43 [35] S. Bowden, R. A. Sinton, *J. Appl. Phys.* **2007**, *102*, 12 124501.
- 44
- 45 [36] J. Lagowski, P. Edelman, A. M. Kontkiewicz, O. Milic, W. Henley, M. Dexter, L. Jastrzebski, A. M.  
46 Hoff, *Appl. Phys. Lett.* **1993**, *63*, 22 3043.
- 47
- 48
- 49
- 50
- 51
- 52
- 53
- 54
- 55
- 56
- 57
- 58
- 59
- 60
- 61
- 62
- 63
- 64
- 65

1  
2  
3  
4  
5  
6  
7  
8  
9  
10  
11  
12  
13  
14  
15  
16  
17  
18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60  
61  
62  
63  
64  
65

## Table of Contents



The results of an study on the kinetic of dissociation of FeB pairs in Cz-Si:B using different light sources are reported. Dissociation rate was shown to depend not only on intensity but also on spectral composition of illumination. The investigation has revealed increase in the dissociation efficiency with a decrease in wavelength and dominant role of the recombination-enhanced defect reaction



Click here to access/download  
**Supporting Information**  
SupIMat.docx

