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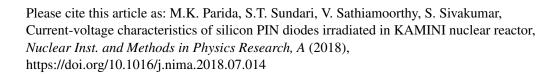
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Current-Voltage characteristics of silicon PIN diodes irradiated in KAMINI nuclear 1 2 Manoj Kumar Parida¹, S. Tripura Sundari¹*, V. Sathiamoorthy² and S. Sivakumar² 3 ¹Surface and Nanoscience Division, Materials Science Group, 4 5 ²Technical Services Division, Reactors Facilities Group 6 Indira Gandhi Centre for Atomic Research, 7 HBNI, Kalpakkam 603102, India 8 9 **Abstract** The present study reports on the investigation of current voltage (I-V) characteristics of 10 commercial planar Si-PIN diodes, irradiated in a typical thermal nuclear reactor – KAMINI for 11 neutron fluences ranging from 1×10^{14} to 1×10^{16} n/cm². The I-V Characteristics of the virgin and 12 neutron irradiated Si-PIN diodes are measured in ambient environment for the forward and 13 reverse biased conditions. In the forward bias condition, one of the consequences of increasing 14 the neutron irradiation fluence is the increase in knee voltage from ~ 0.5 V for virgin diode to 15 37.4 V for neutron irradiated diode with fluence 1x10¹⁶ n/cm². Further analysis of forward 16 characteristics, revealed increase in ideality factor from a typical value of ~2 for virgin diode to 17 an anomalous value of ~496 for the highest irradiated diode specimen. This increase is attributed 18 to the increasing neutron damage that the diodes undergo upon irradiation. Moreover, in the 19 reverse biased condition, the reverse leakage current increased by four orders magnitude from 20 10⁻⁹ to 10⁻⁵ ampere. A qualitative analysis of the forward and reverse I-V characteristics, showed 21 that the diodes change from a rectifying to ohmic behaviour with increase in fluence and this was 22 inferred from the decrease in 'gap' between the forward and reverse currents in the low voltage 23 regions. Quantitatively, the rectification ratio - ratio of the forward to reverse currents - was 24 calculated to be 10⁸ and 84 for the virgin and 1x10¹⁶n/cm² irradiated specimens, respectively. 25 The damage constant evaluated from the reverse bias I-V measurements conditions was found to 26 be $1.7683 \times 10^{-18} \text{ A/cm}$. 27

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- Keywords Si PIN diode, neutron damage, ideality factor, rectification ratio, electrical characterization
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Introduction

In recent years, semiconductor devices have found significant applications for high intensity radiation monitoring in research reactors, solar panels in outer space, lasers in military operations and in high energy physics as particle detectors [1,2]. Among the most widely used semiconductor devices, silicon (Si) PIN diodes have been employed for neutron dosimetry in radiation environment owing to their insensitivity to micro-phonics and amenability to miniaturization [3]. Other distinct inherent advantages of these diodes are fast signal response, good intrinsic detection efficiency, low voltage operation and better energy resolution in terms of pulse height spectra [4]. Moreover, standardization of planar technology process with respect to Si PIN diodes has led to the exploration of a choice of geometries and sizes for optimization of sensitivity of the diode for neutron dose measurement [3]. However, the performances of these devices are susceptible to radiation damage especially from neutrons [5]. It is therefore imperative to study their electrical characteristics from current voltage (I-V) measurements to understand their behaviour in the presence of neutron field. The effects of neutron irradiation on semiconductor Si PIN diodes have been studied in the context of LHC (Large Hardon Collider) too [6].

Several extensive studies on the effect of neutron irradiation on the electrical characteristics have also been conducted as part of ROSE (R&D for Silicon for future Experiments) project [7]. Beattie et al. reported increase in forward voltage of the PIN diode from 10 V to 100 V with increasing neutron fluence [8], when irradiated with neutrons of 1 MeV energy. A similar investigation on I-V characteristics of irradiated silicon detectors performed by Bosetti et al. [9] also, reports the change in the reverse and forward characteristics of diode in terms of rectification ratio, which changes drastically after critical fluence of irradiation. McPherson et al. [10] investigated the electrical characteristics of silicon PIN diodes subjected to 1 MeV neutrons, both prior to and after irradiation. They concluded that the radiation damage occurs only upto certain fluences and beyond that limit, the material becomes resistant to further damage. Additionally, commercial PIN diodes were characterized for their utilization as radiation monitors in LHC to cover 1 MeV equivalent neutrons by Ravotti et al. [11]. They also reported the shifting of the forward characteristics to higher voltages with increase in fluence. However, they obtained a 'thyristor' like behaviour for fluences more than $3x10^{13}$ n/cm².

Irradiation studies conducted on silicon detectors by Edwards et al. [5] indicate increase in the reverse leakage current of diode after exposure to neutrons. The increase in reverse leakage current is due to the formation of energy levels within the energy band gap [12], which aids in the thermal transition of carriers across the band gap. The I-V characteristics of non silicon based materials like GaAs and GaN exposed to thermal neutron irradiation for fluences upto $3x10^{14}$ n/cm² have been studied by Fauzi et al. [13]. They concluded that the performance of diode degraded owing to the displacement damage due to neutrons and gamma ionization. Another neutron irradiation studies on similar GaN materials at higher fluence of 10^{16} n/cm² have also been reported by P.Mulligan et al. [14]. They report that the devices showed insignificant change at higher neutron fluence of 10^{15} n/cm².

The above experiments have been conducted in controlled environments with respect to energy of the neutrons. There are not many studies on the electrical characterization of PIN diodes when subjected to typical reactor operation conditions also. Hasegawa et al. [15]. reported measurement of I-V characteristics of silicon junction detectors under typical reactor operation conditions. The fluences in their study, however, were limited to 1×10^{14} n/cm². The evaluated damage constant from their reverse current measurement was typically ~ 6.6×10^{-17} A/cm, which was two orders lesser than the other neutron irradiation experiments. However, no analysis was carried out to study the variation of the associated diode parameters such as rectification ratio and ideality factor in particular, as a function of neutron fluence, in reactor field conditions. It is therefore worthwhile and essential to examine the electrical behaviour of the Si PIN diodes in particular in real nuclear reactor field conditions so as to study their feasibility of employment in harsh radiation conditions. In the present work, a detailed evolution of forward and reverse I-V electrical characteristics have been investigated on silicon PIN diodes, irradiated in typical thermal research nuclear reactor - KAMINI, for neutron fluences varying from 1×10^{14} to 1×10^{16} n/cm².

Experimental

The experiments were conducted on commercial Si PIN diodes which were procured from M/s Bharat Electronics Limited with the following characteristics. The diode had an active area of 25 mm² and surrounded by a guard ring, which not only reduced the reverse leakage current, but also enabled the diode to be operated at higher voltages in reverse bias condition without

electrical breakdown [16]. The dark reverse leakage current of diodes were $< 10 \text{nA/cm}^2$ and having a maximum operating reverse bias voltage $\sim 100 \text{V}$. The window material was made of SiO₂ which is 0.5 µm thick and diodes had an operating temperature ranging from -65 °C to 125 °C.

The neutron irradiations were carried out using a pneumatic fast transfer system in the thermal research nuclear reactor – KAMINI [17], which is a reactor meant for irradiation experiments. The PIN diode specimens were packed in a polythene vial and transferred to the reactor core with a pneumatic fast transfer mechanism. The neutron flux and spectrum at irradiation location have been experimentally measured by the method of foil activation and is shown in figure 7 of ref. [18]. The corresponding neutron fluxes for various energy intervals are summarized in Table 1. The total neutron flux at the irradiation location was 6.01×10^{11} n/cm²·s, when the reactor power is 20 kW. The irradiation times were so chosen so as to result in neutron fluences ' ϕ ' ranging from 1×10^{14} to 1×10^{16} n/cm². The ambient temperature at the irradiation location was 300 K. Immediately after the specimens were extricated from the reactor, they had an activity in excess of 10 mR/h, which eventually was below 0.1 mR/h after about 4 hours. This activity was considered fit for handling the specimen and performing electrical measurements. It is to be noted that the diodes were also subjected to gamma dose rate (approx 3.14x10⁴ Gy/h). However, the contribution by gamma rays is neglected as the damage caused by it is orders of magnitude lesser than that caused by neutrons [19, 20].

For performing electrical measurements in forward and reverse bias conditions, the diodes were mounted in a Faraday chamber made of Aluminum metal. This chamber was specially designed to shield the diodes from external electromagnetic interference, apart from providing a light free dark environment. The measurements were carried out using Keysight B2912A current source meter. The measurement data were recorded in the PC through the universal serial bus (USB) port. The I-V characteristics in the forward and reverse bias conditions were measured for all the diodes prior to irradiation. For the forward bias condition, the measurements were carried out till 0.8V with a step size of 50 mV, while for the reverse bias conditions, the measurements were carried out from 0 V till 100 V, incrementing at the rate of 1 V. From literature [8], it is known that the forward knee voltage is increased upon irradiation. Therefore, for the measurements of forward currents in irradiated specimens, the upper limit on

- forward voltage was set at a value where a sharp increase in current was obtained. However, for
- the reverse bias conditions, the upper voltage (till 100 V) was kept the same as that in the virgin
- case. In order to estimate the depletion width of Si PIN diode, a reverse C-V measurement was
- performed using Agilent precision LCR meter (model E4980 A) from 0 V to 100V at constant
- frequency of 1 MHz.
- 128 Results and Discussion
- 129 a) Forward Current Analysis
- 130 1. Forward I-V characteristics of virgin PIN diode
- Figures 1(a) and 1(b) show the response of the forward current (I_f) of virgin PIN diode as a
- function of forward voltage (V_f) till 0.8V, in the linear and log-log scale, respectively. The
- forward I-V characteristics show normal rectification behaviour with a sharp increase in I_f close
- to a knee voltage (V_{knee}) of ~ 0.5 V.
- In general, for an ideal PIN diode [21], the forward current voltage (I_f-V_f) relation is given by
- the equation (1).

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$$I_f = I_0 \exp(\beta V_f) + I_{re} \exp(\beta V_f/2) \dots (1)$$

- where I_0 and I_{re} are the diffusion and recombination saturation currents, respectively and β
- $=q/\eta kT$, where q is the electronic charge, k is the Boltzmann constant, η is the ideality factor and
- 140 T is the temperature.
- From a linear fitting of I_f vs V_f to Eq. (1) for low and high voltage regions on the semilog scale,
- the saturation current (I_s) and ideality factor (η) , were evaluated from the intercepts and slopes,
- respectively. A point to note is that the knee voltage ' V_{knee} ' is taken as a voltage, above (below)
- which the total diode current is driven by diffusion (recombination) mechanism. Figure 2 shows
- the fitting of Eq. (1) for the virgin diode in the recombination and diffusion regions, which yields
- 146 $\eta = 2.1$ and $I_s = 58.3$ pA for the low voltage range, while for high voltage range, $\eta = 2.4$ and $I_s = 185$
- nA. The ideality factor, in general is ~ 1.2 -1.3 and ~2-3 for silicon and GaAs, respectively [22].
- For ideal p-n junctions devoid of any defects, the ideality factor is described by Sah-Noyce-
- Shockley theory [23], which yields $\eta = 1$ at low voltage range and $\eta = 2$ at high voltage ranges.
- Although the experimental ideality factor for the virgin PIN diode indicates a small deviation
- from the theoretical value $\sim 1(2)$ for low (high) voltage, the diode is still a good rectifier as seen

152	from the measured reverse leakage current (~ 1.69 nA at 100V) in the reverse bias condition as
153	discussed in later sections. The miniscule magnitude of reverse leakage current clearly indicates
154	that the rectification property of diode is still intact.
155	2. Forward I-V characteristics of neutron irradiated PIN diodes
156	Figures 3(a) and 3 (b) show the $I_{\rm f}$ response of the neutron irradiated Si PIN diodes as a function
157	of V_f , for neutron fluences ' ϕ ', ranging from $1x10^{14}$ n/cm² to $1x10^{16}$ n/cm² in the linear scale and
158	log-log scale, respectively.φ The plot for the virgin diode specimen is also shown for
159	comparison.
160	The forward I-V characteristics in these figures definitely show prominent deviations from
161	exponential behaviour of the virgin diode as shown in figures 1(a) and (b). In particular, as ϕ
162	increases, the value of V_{knee} at which the diode reaches limiting current also progressively
163	increases. Therefore, the I-V characteristics of the irradiated diodes could be measured for
164	increased bias voltages without an electrical breakdown of the diodes. The reason for the shift in
165	V_{knee} is due to the formation of traps caused by neutron damage of the Si PIN diode [3]. Higher
166	the neutron fluence, more the neutron damage, thereby leading to shifting V_{knee} to higher voltage
167	levels. As for the analysis, the I-V characteristics are fitted to Eq. (1) for the low and high
168	voltage regions to obtain the ideality factor and saturation currents. A representative fit of the
169	same is shown in figure 4.
170	The ideality factors and the saturation currents obtained for all the specimens from the fit are
171	shown in figures 5 and 6, respectively. The ideality factor $\boldsymbol{\eta}$ and the saturation currents consisting
172	of I_{o} and I_{re} , show distinct increase with increase in ϕ , which indicate considerable deviations
173	from the ideal diode behaviour. The Shockley-Read-Hall (SRH) recombination theory [24]
174	which assumes recombination via isolated point defect levels has been traditionally used to
175	explain the ideality factor (η) in diodes for the diffusion and recombination regions. In the
176	present experiment, η typically changes from ~2 for the virgin diodes to ~ 496 for the highest
177	irradiated fluence. Such large ideality factors have been observed in silicon solar cell devices
178	[25], BN/ZnO hetero structured rectifying diodes, phosphorus doped n-Ge/i-Ge/p-Si hetero-
179	structure diodes too [26]. The increase in ideality factor in the present studies is attributed to the
180	recombination currents being much higher than expected ones [26]. The evolution of increase in
181	η and associated saturation currents with fluence are due to the large defects that are created

thereby implying more recombination. Other factors such as tunnelling and metal semiconductor junctions also contribute to high η [26].

Qualitatively, the forward I-V characteristics of all irradiated diodes become more and more shallow with increase in fluence ϕ (figures 3(a) and 3(b)). Moreover, for voltages below the knee voltage of virgin diode, I_f for all the irradiated diodes increases by three orders of magnitude. The large increase in I_f in this region for the lowest irradiated specimen $(1x10^{14})$ is indicative of the damage that the diode undergoes upon irradiation. This increase is attributed to the increased number of electron-hole pairs generated from the generation-recombination centres during the initial stages of irradiation. For higher fluences, the I_f decreases with increase in fluence as shown in figures 3(a) and 3(b). This decrease is due to the formation of traps which are formed by Frenkel defects [3]. The traps lead to increased scattering of carriers, resulting in decreased mobility of carriers. Consequently, the lifetime of the charge carriers [3] decreases thereby leading to increased resistance of the material upon irradiation. It is to be noted that even though the diodes were simultaneously exposed to the inevitable gamma dose rate $\sim 3.14 \times 10^4$ Gy/h, their consequential effect is neglected. This is due to the fact that the gamma ray damage in silicon is almost three orders of magnitude less than that caused by neutrons [19, 20].

In order to further understand the quantitative changes and trends in the evolution of I_f on irradiation, the V_f at which I_f is 17.5 mA, is monitored. This magnitude of I_f is chosen for two reasons. Firstly, the slope of I_f vs V_f is very large at I_f =17.5mA and does not change appreciably with further increase in V_f . Secondly, the current was limited by heating effects to avoid damage to the junction of the diode. The V_f corresponding to 17.5 mA is taken as knee voltage and is plotted in Figure 7 as a function of neutron fluence on a linear scale. It is clear from the graph that, the increase in V_{knee} is maximum in the initial stages of irradiation. The extent of increase in magnitude of V_{knee} with neutron fluence has also been reported in literature by Swartz et al. [27]. With the increase in neutron fluence, trap density is known to increase, which leads to the recombination of generated carriers, thereby decreasing the concentration of the intrinsic charge carriers. Consequently, there is a degradation of carrier lifetime and an increased resistivity upon irradiation. Incidentally, as seen from the fit shown in figure 7, the V_{knee} rises rapidly for low fluences and nearly saturates at a value 36.3 V for a fluence of $5x10^{15}$ n/cm². Such saturation behaviour is not reported in ref. [10], probably due to the fact that the I_f

212	was measured only till ~10-12 volts. The behaviour observed in the present study is almost		
213	similar to that reported in Swartz et al. [27] and indicates that there is little change in damage		
214	beyond a fluence of $4x10^{15}$ n/cm ² . This 'flattening off' is construed to be a consequence of the		
215	occupation of displaced atoms and impurities by the Frenkel pairs that have been created by the		
216	irradiation, thereby resulting in the material becoming radiation hard or radiation resistant.		
0.4.7			
217	b) Reverse current analysis		
218	The I-V characteristics in the reverse bias conditions for the virgin diode along with neutron		
219	irradiated diodes are shown in figures 8 and 9 in linear and semilog scale, respectively. The		
220	measurements were performed up to a maximum reverse bias voltage ' V_r ' of 100 V i.e., the limit		
221	to which the diodes can be operated in reverse biased condition. Since the interest in the present		
222	studies is on the change in characteristics upon neutron irradiation, the I_r - V_r profile of the virgin		
223	diode is taken as reference. The reverse leakage current (I _r) due to minority carriers is typically		
224	in the nanoampere (10-9A) range for the virgin diode. Apart from the minority carrier		
225	contributing to I _r , its magnitude also depends on the active volume of the diode, which in turn is		
226	determined by the depletion thickness and lateral extension [28].		
227	As for the neutron irradiated diodes, I _r increases by four orders of magnitude for the diode with		
228	lowest irradiated fluence of $1x10^{14} \text{ n/cm}^2$. The large change in I_r upon neutron irradiation is due		
229	to the fact that a large number of Frenkel pairs are created and unoccupied. Moreover, defect		
230	states are introduced upon neutron irradiation and they act as generation-recombination (g-r)		
231	centres [10]. These provide e-h pairs that are drawn by the applied reverse field contributing to		
232	the increased I_r [29]. It is pointed out that increase in I_r due to the minority charge carriers is a		
233	consequence of defects induced by the irradiated neutrons both at the surface and in the bulk of		
234	the diode. The increase in I_r is gradual, reaching a maximum value of ~ $5.7x10^{-4}$ A, for the diode		
235	irradiated to a fluence of 1x10 ¹⁶ n/cm ² .		
236	The behaviour of I_r at full depletion voltage which is ~ 100V, in the present experiment is shown		
237	in figure 10 as a function of ϕ . This behaviour is comparable to that reported by S.Moloi et al.		
238	[30] but does not follow a linear behaviour as reported by Hasegowa et al. [15] and Lemeilleur et		
239	al. [29]. Incidentally, a plot of I_r for lower V_r (20V), shows the same trend as V_r of 100V. The		

behaviour shown in figure 10 implies that change in minority carriers is very large for the lowest

fluence and thereafter the change is minimal. The curve tends to flatten indicating that the 241 material is already radiation hard [30]. 242 A qualitative behaviour of the effect of exposure of the diodes to neutrons can be inferred from 243 the gap between the I_f and I_r for both virgin and irradiated diodes. The progressive convergence 244 of 'gap' between the forward and reverse profiles (I_r-I_f) with increase in φ is shown in figures 11 245 (a-c) for a few representative specimens, namely virgin, 1×10^{14} n/cm² and 5×10^{15} n/cm². The 246 progressive decrease in the gap shown in figures 11, is indicative of the change in behaviour 247 from a diode like to ohmic behaviour and compares well with the reported behaviour [30]. The 248 249 forward and reverse currents are almost equal and vary linearly with applied voltage for high fluences. 250 The variations in the I-V characteristics are readily brought out by 'rate of change' in the V_{knee} 251 and reverse current I_r obtained from figures 7 and 10, respectively. These rates of change are 252 253 shown in log-log scale. The reference current and voltage for forward and reverse bias measurements were taken to be 17.5 mA and 100 V, respectively. For the evaluation of slope of 254 V_{knee} with fluence, the data was limited between virgin to 1x10¹⁵ n/cm², as beyond this fluence, 255 the V_{knee} saturates. The 'rate of change' of V_{knee} and I_r with fluence are 0.6 and 0.7 respectively, 256 as shown in figures 12 and 13, respectively. The magnitudes of the slopes in the present 257 258 experiment are nearly comparable and differ from that reported by Moloi et al. [30], where the slope of the reverse bias is reported to be twice that in the forward bias case. The difference in 259 the slopes from ref. [30] is owing to the fact that in the present work the slopes were estimated 260 till low neutron fluence as saturation in V_{knee} was observed at higher fluence. 261 **Damage Constant** 262 The damage constant ' α ' or volume reverse leakage current I_r/V_l , in the reverse bias conditions 263 is taken as a measure of the damage caused by irradiation [31] and I_r is used to evaluate the 264 radiation hardness property of the diode [31,32]. The constant ' α ' is defined through the relation, 265 $\Delta I = \alpha V_l$, where, ΔI is the difference in the reverse leakage current for the irradiated and virgin 266 diode for the full depletion voltage of 100V and V₁ is the volume of the depletion region in the 267 diode. The lateral dimension of the PIN diodes used in the experiments is 5mm x 5mm and the 268

width of the depletion region was computed from a measurement of capacitance in reverse bias

conditions(C-V). Figure 14 shows the reverse C-V measurement of virgin diode $(1x10^{16} \text{ n/cm}^2)$ in which the capacitance (C) decreases with increase in reverse voltage and becomes constant beyond 60 V. The initial decrease in capacitance is attributed to the increase in depletion width of diode with voltage, while the invariance above a certain voltage ($\sim 60\text{V}$) indicates a complete depletion of the diode. The depletion region width estimated from the C-V measurement is $\sim 300\mu\text{m}$ [33].

The difference in the I_r per unit volume estimated at the reverse bias voltage of 100V for all the five irradiated diodes is shown in figure 15. From a linear fitting of ΔI vs ϕ , the damage constant α is estimated to be 1.7683 x 10⁻¹⁸(A/cm). The damage constant ' α ' in this experiment is lesser compared to that reported in ref. [15]. This is due to use of a wide spectrum of incident neutrons energy rather than a monoenergetic neutron beam.

Rectification ratio (RR)

It is well known that the I-V characteristics, which demonstrate the rectifying behaviour, are crucial for device applications. In the context of diodes, rectification is monitored and quantified in terms of rectifying ratio 'RR', which expresses the ratio of the $I_{\rm f}$ to the $I_{\rm r}$. Mathematically, RR is

$$RR = \frac{I_{fV_{f,max}}}{I_{rV_{f,max}}}....(2)$$

In the above expression, $V_{f,max}$ is the maximum forward voltage of the diode which can be applied within the permissible limits of operation and is used as the reference voltage for calculation of the RR. In the present experiment, as seen from figure 3(a), V_{fmax} varies with the neutron irradiation fluence. As seen from the graph (see figure 16), the virgin diode has the highest rectification ratio RR of ~10⁸, with the current conduction in one direction only. This high value of RR is owing to the fact that the I_f is in the mA range, while the leakage current in the reverse bias condition is in the nA range. Large RR has also been reported in all printed organic diodes [34] and metal induced lateral crystallization in a-Si: H in solar cell application [35].

In the present experiment, RR decreases drastically to $\sim 10^3$ and ~ 84 for diodes exposed to ϕ of 1×10^{14} n/cm² and 1×10^{16} n/cm², respectively as shown in figure 16. This decreasing trend is

consistent with that observation by Bosetti et al. [9], although the magnitudes vary. The decrease in RR clearly indicates that the neutron irradiated diodes progressively lose their rectifying behaviour with increasing ϕ . The reason for the loss of rectifying behaviour is the decrease (increase) in the concentration of majority (minority) charge carriers. It is well known that irradiation induces defects and leads to loss of crystallinity, thereby resulting in decreased lifetime and mobility [3] .Such defect centres are known to increase the reverse leakage current and consequently resulting in loss of rectification [9].The variation of RR vs ϕ has been studied in literature and a critical fluence was observed [9] in which the RR is equal to zero. However, in the present experiment, it is found that the RR is \sim 84 even for a fluence of 1×10^{16} n/cm², which is higher than that reported in ref [9].

The quantities, RR and η are inverse to each other and the same is shown in figure 17. A large RR results in nearly ideal diodes where the ideality factor is ~1 in the diffusion region. It is pertinent to note that the ideality factor of an ideal diode is 1(2) in the diffusion (recombination) region with the concomitant RR being very large (typically 10 8 or more). Figure 17 indicates a clear change in the behaviour of the diode from a rectifying one to ohmic one. This is a consequence of deceasing (increasing) forward (reverse) current with increasing fluence. Such a correlation has also been reported for nanocrystalline pn diodes [35].

Conclusion

This paper discusses the effect of neutron irradiation on the current voltage characteristics of silicon PIN diodes in a typical thermal research reactor for fluences ranging from $1x10^{14}$ to $1x10^{16}$ n/cm². The measured forward I-V characteristics, is increasingly becoming shallow upon increased irradiation fluence. Moreover, the knee voltage increases from 0.5V for the virgin diode to ~ 36.3 V for diode irradiated to a fluence of ~ $5x10^{15}$ n/cm² and thereafter saturates. The shift and increase in knee voltage is attributed to increased trap density [36], which in turn leads to the recombination of the majority carriers, thereby decreasing the concentration of the intrinsic charge carriers. The increase in knee voltage indicates that diodes can be operated at higher voltages [37]. The ideality factors extracted from the fit of the forward I-V characteristics to diode equations in the low/high voltage regions varied from a typical value of ~ 2 for the virgin case to ~ 496 for the highest irradiated fluence. The large ideality factors observed in this study are caused by the defects that are created upon irradiation.

With increase in neutron fluence, the reverse leakage currents also increased by almost four
orders of magnitude owing to the production of a large number of defect states that act as
generation-recombination (g-r) centres upon neutron irradiation. Another important consequence
of the increase in fluence of irradiation is the transformation of diode properties from a rectifying
to ohmic one, which was inferred from the progressive decrease in gap between the forward and
reverse currents. Further, quantitatively, the rectification ratio decreased from 10^8 for virgin
diode to 84 for the maximum neutron irradiated fluence of 1x10 ¹⁶ n/cm ² . The damage constant
evaluated from the reverse leakage current indicates that the diode hardness property is still in
intact after the irradiation. The quantitative behaviour of diodes in terms of rectification ratio
shows that its rectification property has been degraded drastically with the increase in neutron
fluence.

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References

350		
351 352 353	[1]	N.I.C. Omar, N.K.A.M. Rashid, J.A. Karim, J. Abdullah, N.F. Hasbullah, Effects of neutron on reverse bias characteristics of commercially available Si and GaAs diodes, Aust. J. Basic Appl. Sci. 6 (2012) 211–216.
354	[2]	G.Lutz, Semiconductor and electronic devices, Springer, 1999.
355		
356	[3]	I. Anokhin, O. Zinets, A. Rosenfeld, M. Lerch, M. Yudelev, V. Perevertaylo, M.
357		Reinhard, M. Petasecca, Studies of the Characteristics of a Silicon Neutron Sensor,
358		IEEE Trans. Nucl. Sci. 56 (2009) 2290–2293. doi:10.1109/TNS.2009.2024150.
359		
360	[4]	S.J. Bates, D.J. Munday, M.A. Parker, F. Anghinolfi, A. Chilingarov, A. Ciasnohova,
361		M. Glaser, E. Heijne, P. Jarron, F. Lemeilleur, J.C. Santiard, R. Bonino, A.G. Clark,
362		H. Kambara, C. Gössling, B. Lisowski, A. Rolf, S. Pilath, H. Feick, E. Fretwurst, G.
363		Lindström, T. Schulz, R.A. Bardos, G.W. Gorfine, G.F. Moorhead, G.N. Taylor, S.N.
364		Tovey, Recent results of radiation damage studies in silicon, Nucl. Instrum. Methods
365		Phys. Res. Sect. Accel. Spectrometers Detect. Assoc. Equip. 344 (1994) 228–236.
366		doi:10.1016/0168-9002(94)90675-0.
367		
368 369 370	[5]	M. Edwards, G. Hall, S. Sotthibandhu, Neutron radiation damage studies of silicon detectors, Nucl. Instrum. Methods Phys. Res. Sect. Accel. Spectrometers Detect. Assoc. Equip. 310 (1991) 283–286. doi:10.1016/0168-9002(91)91044-V.
371 372 373	[6]	M. Moll, Development of radiation hard sensors for very high luminosity colliders—CERN-RD50 project, Nucl. Instrum. Methods Phys. Res. Sect. Accel. Spectrometers Detect. Assoc. Equip. 511 (2003) 97–105. doi:10.1016/S0168-9002(03)01772-8.
374	[7]	G. Lindström, M. Ahmed, S. Albergo, P. Allport, D. Anderson, L. Andricek, M.M.
375	r. 1	Angarano, V. Augelli, N. Bacchetta, P. Bartalini, R. Bates, U. Biggeri, G.M. Bilei, D.
376		Bisello, D. Boemi, E. Borchi, T. Botila, T.J. Brodbeck, M. Bruzzi, T. Budzynski, P.
377		Burger, F. Campabadal, G. Casse, E. Catacchini, A. Chilingarov, P. Ciampolini, V.
378 379		Cindro, M.J. Costa, D. Creanza, P. Clauws, C. Da Via, G. Davies, W. De Boer, R. Dell'Orso, M. De Palma, B. Dezillie, V. Eremin, O. Evrard, G. Fallica, G. Fanourakis,
380		H. Feick, E. Focardi, L. Fonseca, E. Fretwurst, J. Fuster, K. Gabathuler, M. Glaser, P.
381		Grabiec, E. Grigoriev, G. Hall, M. Hanlon, F. Hauler, S. Heising, A. Holmes-Siedle,
382		R Horisberger G Hughes M Huhtinen I Ilvashenko A Ivanov BK Jones L

383		Jungermann, A. Kaminsky, Z. Kohout, G. Kramberger, M. Kuhnke, S. Kwan, F.
384		Lemeilleur, C. Leroy, M. Letheren, Z. Li, T. Ligonzo, V. Linhart, P. Litovchenko, D.
385		Loukas, M. Lozano, Z. Luczynski, G. Lutz, B. MacEvoy, S. Manolopoulos, A.
386		Markou, C. Martinez, A. Messineo, M. Mikuž, M. Moll, E. Nossarzewska, G.
387 388		Ottaviani, V. Oshea, G. Parrini, D. Passeri, D. Petre, A. Pickford, I. Pintilie, L. Pintilie, S. Pospisil, R. Potenza, C. Raine, J.M. Rafi, P.N. Ratoff, R.H. Richter, P. Riedler, S.
389		Roe, P. Roy, A. Ruzin, A.I. Ryazanov, A. Santocchia, L. Schiavulli, P. Sicho, I. Siotis,
390		T. Sloan, W. Slysz, K. Smith, M. Solanky, B. Sopko, K. Stolze, B. Sundby Avset, B.
391		Svensson, C. Tivarus, G. Tonelli, A. Tricomi, S. Tzamarias, G. Valvo, A. Vasilescu,
392		A. Vayaki, E. Verbitskaya, P. Verdini, V. Vrba, S. Watts, E.R. Weber, M. Wegrzecki,
393		I. Wegrzecka, P. Weilhammer, R. Wheadon, C. Wilburn, I. Wilhelm, R. Wunstorf, J.
394		Wüstenfeld, J. Wyss, K. Zankel, P. Zabierowski, D. Žontar, Radiation hard silicon
395		detectors—developments by the RD48 (ROSE) collaboration, Nucl. Instrum. Methods
396		Phys. Res. Sect. Accel. Spectrometers Detect. Assoc. Equip. 466 (2001) 308–326.
397		doi:10.1016/S0168-9002(01)00560-5.
398		
399	[8]	L.J. Beattie, A. Chilingarov, T. Sloan, Forward-bias operation of Si detectors:: a way
400		to work in high-radiation environment, Nucl. Instrum. Methods Phys. Res. Sect.
401		Accel. Spectrometers Detect. Assoc. Equip. 439 (2000) 293-302. doi:10.1016/S0168-
402		9002(99)00840-2.
403		
404	[9]	M. Bosetti, N. Croitoru, C. Furetta, C. Leroy, S. Pensotti, P. Rancoita, M. Rattaggi, M.
405		Redaelli, A. Seidman, Study of current-voltage characteristics of irradiated silicon
406		detectors, Nucl. Instrum. Methods Phys. Res. Sect. B Beam Interact. Mater. At. 95
407		(1995) 219–224. doi:10.1016/0168-583X(94)00439-0.
408		
409	[10]	M. McPherson, B.K. Jones, T. Sloan, Effects of radiation damage in silicon p - i - n
410		photodiodes, Semicond. Sci. Technol. 12 (1997) 1187. doi:10.1088/0268-
411		1242/12/10/003.
412		
413	[11]	F. Ravotti, M. Glaser, M. Moll, F. Saigne, BPW34 Commercial p-i-n Diodes for High-
414		Level 1-MeV Neutron Equivalent Fluence Monitoring, IEEE Trans. Nucl. Sci. 55
415		(2008) 2133–2140. doi:10.1109/TNS.2008.2000765.

416	[12]	A. Ruzin, Recent results from the RD-48 (ROSE) Collaboration, Nucl. Instrum.
417		Methods Phys. Res. Sect. Accel. Spectrometers Detect. Assoc. Equip. 447 (2000) 116-
418		125. doi:10.1016/S0168-9002(00)00179-0.
419		
420	[13]	D.A. Fauzi, N.K.A.M. Rashid, J.A. Karim, M.R.M. Zin, N.F. Hasbullah, O.A.S.
421		Fareed, Electrical performances of commercial GaN and GaAs based optoelectronics
422		under neutron irradiation, IOP Conf. Ser. Mater. Sci. Eng. 53 (2013) 012029.
423		doi:10.1088/1757-899X/53/1/012029.
424		
425	[14]	P. Mulligan, J. Qiu, J. Wang, L.R. Cao, Study of GaN radiation sensor after in-core
426		neutron irradiation, in: 2013 3rd International Conference on Advancements in
427		Nuclear Instrumentation, Measurement Methods and Their Applications (ANIMMA),
428		2013: pp. 1–5. doi:10.1109/ANIMMA.2013.6727935.
429		
430	[15]	M. Hasegawa, S. Mori, T. Ohsugi, H. Kojima, A. Taketani, T. Kondo, M. Noguchi,
431		Radiation damage of silicon junction detectors by neutron irradiation, Nucl. Instrum.
432		Methods Phys. Res. Sect. Accel. Spectrometers Detect. Assoc. Equip. 277 (1989) 395-
433		400. doi:10.1016/0168-9002(89)90768-7.
434		
435	[16]	V.Mishra, V.D.Srivastava, S.K.Kataria, Role of guard rings in improving the
436		performance of silicon detectors, Pramana-Journal of Physics. 65 (2005) 259-272.
437		
438	[17]	D.K. Mohapatra, P. Mohanakrishnan, Measurement and prediction of neutron spectra
439		in the Kalpakkam mini reactor (KAMINI), Appl. Radiat. Isot. Data Instrum. Methods
440		Use Agric. Ind. Med. 57 (2002) 25–33.
441		
442	[18]	G.V.S.A. Kumar, S. Sen, E. Radha, J.S.B. Rao, R. Acharya, R. Kumar, C.R.
443		Venkatasubramani, A.V.R. Reddy, M. Joseph, Studies on neutron spectrum
444		characterization for the Pneumatic Fast Transfer System (PFTS) of KAMINI reactor,
445		Appl. Radiat. Isot. 124 (2017) 49–55. doi:10.1016/j.apradiso.2017.03.009.

446	[19]	J.R. Srour, D.M. Long, D.G. Millward, R.L. Fitzwilson and W.L. Chadsey, Radiation
447		damage of silicon junction detectors Enhancement of Electronic Materials, (Noyes
448		Publications), 1984.
449		
450 451 452	[20]	J.P. Raymond, E.L. Petersen, Comparison of Neutron, Proton and Gamma Ray Effects in Semiconductor Devices, IEEE Transactions on Nuclear Science. 34 (1987) 1621–1628. doi:10.1109/TNS.1987.4337526.
453	[21]	A. Bar-Lev, Semiconductor and electronic devices, Prentice-Hall, 1993.
454		
455	[22]	V.R.V. Pillai, S.K. Khamari, V.K. Dixit, T. Ganguli, S. Kher, S.M. Oak, Effect of γ-
456		ray irradiation on breakdown voltage, ideality factor, dark current and series resistance
457		of GaAs p-i-n diode, Nucl. Instrum. Methods Phys. Res. Sect. Accel. Spectrometers
458		Detect. Assoc. Equip. 685 (2012) 41–45. doi:10.1016/j.nima.2012.05.062.
459		
460	[23]	C.Sah,R.N. Noyce,W. Shockley, Carrier Generation and Recombination in p-n
461		Junctions and p-n Junction Characteristics. Proceedings of the IRE 45 (1957) 1228-
462		1243. http://dx.doi.org/10.1109/JRPROC.1957.278528
463		
464	[24]	W. Shockley, W.T. Read, Statistics of the Recombinations of Holes and Electrons,
465		Phys. Rev. 87 (1952) 835–842. doi:10.1103/PhysRev.87.835.
466		
467 468 469	[25]	O. Breitenstein, P. Altermatt, K. Ramspeck, A. Schenk, The Origin of Ideality Factors N> 2 of Shunts and Surfaces in the Dark I-V Curves of Si Solar Cells, n.d.,21 st European Photovoltaic Solar Energy Conference., 2006.
470	[26]	M. Brötzmann, U. Vetter, H. Hofsäss, BN/ZnO heterojunction diodes with apparently
471		giant ideality factors, J. Appl. Phys. 106 (2009) 063704. doi:10.1063/1.3212987.
472		
473	[27]	J.M. Swartz, M.O. Thurston, Analysis of the Effect of Fast-Neutron Bombardment on
474		the Current-Voltage Characteristic of a Conductivity-Modulated p-i-n Diode, J. Appl.
475		Phys. 37 (1966) 745–755. doi:10.1063/1.1708249.
476		

477	[28]	Y. Murakami, T. Shingyouji, Separation and analysis of diffusion and generation
478		components of pn junction leakage current in various silicon wafers, J. Appl. Phys. 75
479		(1994) 3548–3552. doi:10.1063/1.356091.
480		
481	[29]	F. Lemeilleur, M. Glaser, E.H.M. Heijne, P. Jarron ,E. Occelli, Neutron-induced
482		radiation damage in silicon detectors, IEEE Transactions on Nuclear Science, 39, 1992
483		551-557.
484		
485	[30]	S.J. Moloi, M. McPherson, The current and capacitance response of radiation-
486		damaged silicon PIN diodes, Phys. B Condens. Matter. 404 (2009) 3922-3929.
487		doi:10.1016/j.physb.2009.07.123.
488		
489 490	[31]	G. Lindström, M. Moll, E. Fretwurst, Radiation hardness of silicon detectors – a challenge from high-energy physics, Nucl. Instrum. Methods Phys. Res. Sect. Accel.
491 492		Spectrometers Detect. Assoc. Equip. 426 (1999) 1–15. doi:10.1016/S0168-9002(98)01462-4.
493	[32]	F.Honniger, Radiation Damage in Silicon - Defect Analysis and Detector Properties,
494		Ph.D dissertation submitted in 2007.
495		
496 497	[33]	V.Mishra, Study of Silicon Detectors, Ph.D dissertation submitted in 2002 , Mumbai University.
498	[34]	S. Ali, J. Bae, C.H. Lee, Organic diode with high rectification ratio made of
499		electrohydrodynamic printed organic layers, Electron. Mater. Lett. 12 (2016) 270-275.
500		doi:10.1007/s13391-015-5202-y.
501		
502	[35]	J.D. Hwang, K.S. Lee, A High Rectification Ratio Nanocrystalline p-n Junction Diode
503		Prepared by Metal-Induced Lateral Crystallization for Solar Cell Applications, J.
504		Electrochem. Soc. 155 (2008) H259-H262. doi:10.1149/1.2840618.
505		
506	[36]	V. Sopko, B. Sopko, D. Chren, J. Dammer, Study of PIN diode energy traps created by
507		neutrons, J. Inst. 8 (2013) C03014. doi:10.1088/1748-0221/8/03/C03014.
508		

509	[3/]	A. Chilingarov, 1. Sloan, Operation of heavily irradiated silicon detectors under
510		forward bias, Nuclear Instruments and Methods in Physics Research Section A
511		Accelerators, Spectrometers, Detectors and Associated Equipment. 399 (1997) 35-37
512		doi:10.1016/S0168-9002(97)00940-6.
513		
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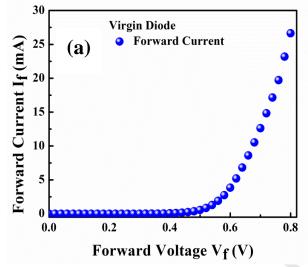
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541	Table Captions
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543	Table 1. Neutron spectrum and flux measured at neutron irradiation location
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571	Figure Capt	ions
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573	Figure 1	Forward I-V Characteristics of virgin diode in (a) linear and (b) log-log scale.
574 575 576	Figure 2	Fitting of diode equation (Eq. 1) for low and high voltage regions for virgin diode, in order to obtain ideality factor (η) and saturation current (I_s) in the respective region.
577 578	Figure 3	Forward I-V Characteristics of neutron irradiated diodes in (a) linear and (b) log-log scale.
579 580	Figure 4	Fitting of diode equation (Eq.1) to a typical neutron irradiated diode (ϕ =1 x10 ¹⁵ n/cm ²) for low and high voltage regions.
581 582	Figure 5	Variation of η with neutron fluence ϕ (n/cm ²), obtained from fit to Eq. 1 for low and high voltage ranges (Note: line is guide to the eye).
583 584	Figure 6	Variation of saturation current with neutron fluence ϕ (n/cm ²) obtained from fit to Eq.1 for low and high voltage range (Note: line is guide to the eye).
585 586	Figure 7	Variation of knee voltage (V_{knee}) with neutron fluence at constant forward current of 17.5 mA.
587 588	Figure 8	Reverse I-V characteristics of virgin and neutron irradiated diodes in linear scale.
589 590	Figure 9	Reverse I-V characteristics of virgin and neutron irradiated diodes in semi-log scale.
591 592	Figure 10	Reverse leakage current at bias voltage of 20V and 100 V as a function of neutron fluence.
593 594 595	Figure 11	Low voltage gap between I_f and I_r for (a) virgin (b) $5x10^{14}$ n/cm ² and (c) $5x10^{15}$ n/cm ² showing the progressive decrease in gap between I_f and I_r with increasing fluence.
596 597	Figure 12	'Rate of change' of knee voltage at forward current of 17.5mA with neutron fluence in log-log scale.

598	Figure 13	'Rate of change' of reverse current with neutron fluence in log-log scale.
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600	Figure 14	Reverse capacitance voltage characteristics of virgin diode.
601 602	Figure 15	Variation of current density with neutron fluence and slope showing the damage constant ' α '.
603 604	Figure 16	Variation of rectification ratio RR of virgin and neutron irradiated diodes as function of neutron fluence ϕ (n/cm ²).
605	Figure 17	Variation of rectification ratio RR and ideality factor (η) with neutron fluence φ
606		(n/cm^2) .
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622	Table.1		
623		Neutron Energy Interval	Neutron Flux (cm ⁻² ·s ⁻¹)
624		0 - 0.56 eV	3.45 x10 ¹⁰
625		0.56 ev – 0.5 MeV	4.30 x10 ¹¹
626		0.5 MeV- 2 MeV	1.36 x10 ¹¹
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Figure 1



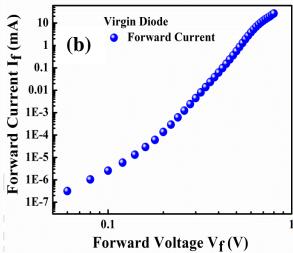


Figure 2

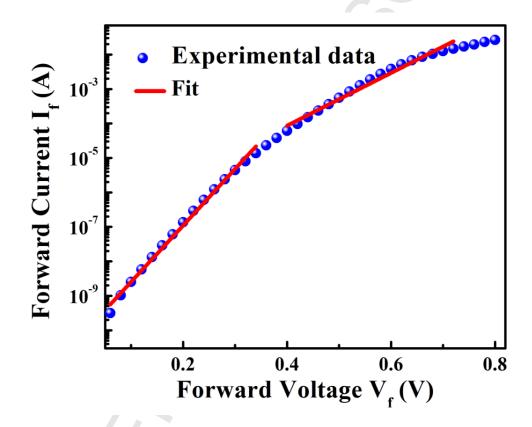


Figure 3

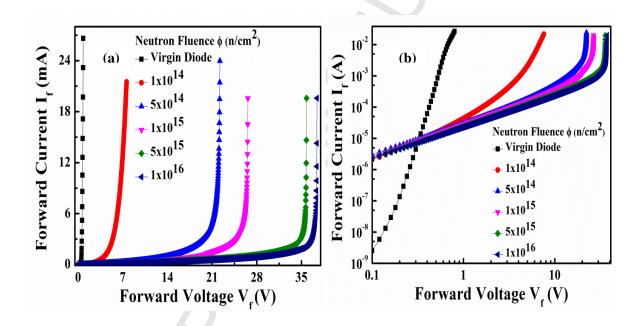


Figure 4

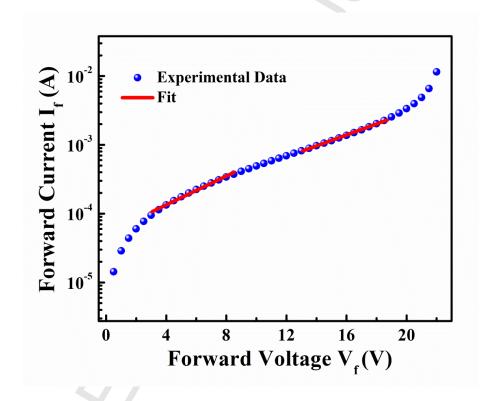


Figure 5

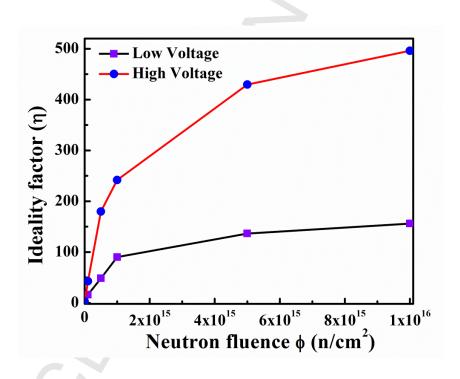


Figure 6

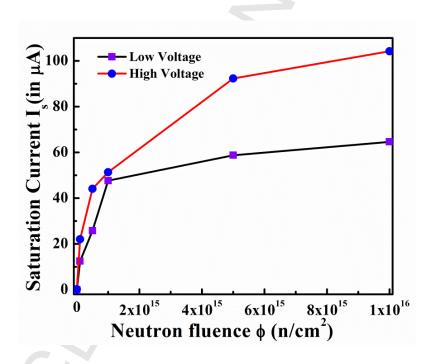


Figure 7

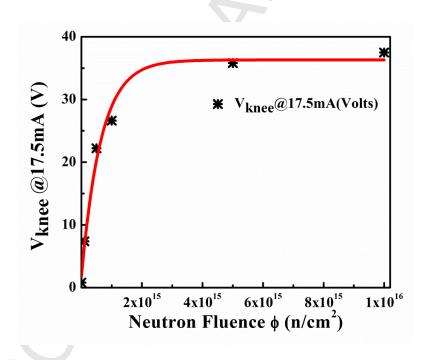


Figure 8

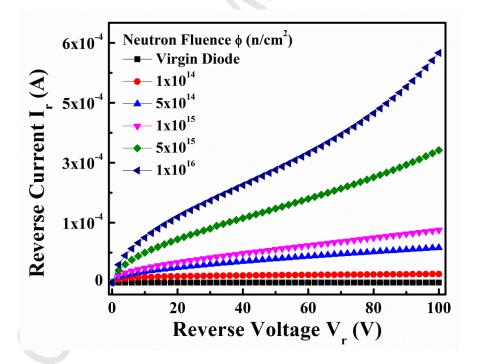


Figure 9

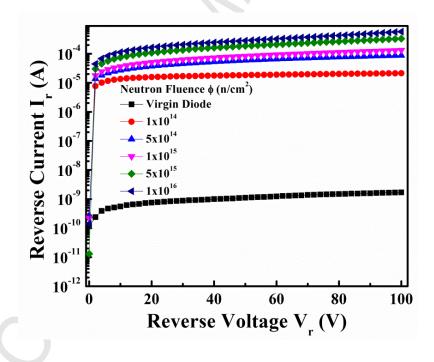


Figure 10

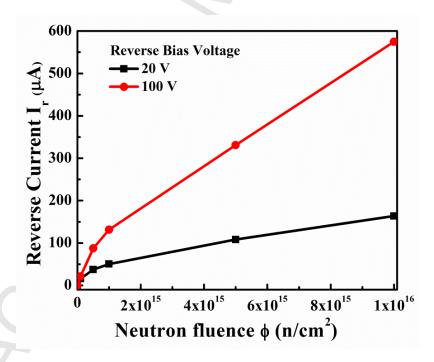
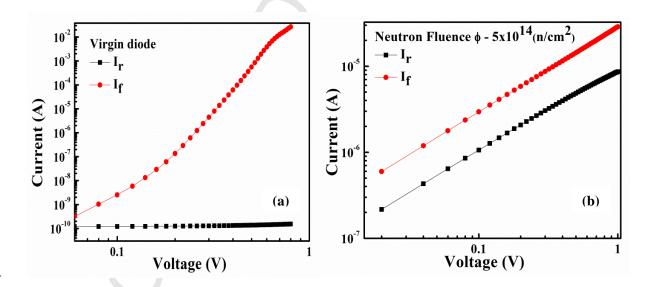


Figure 11



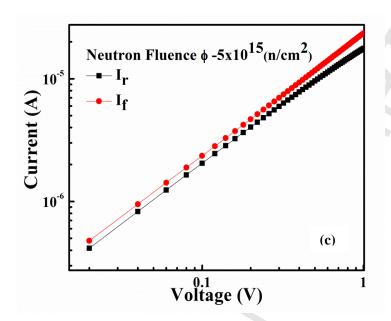
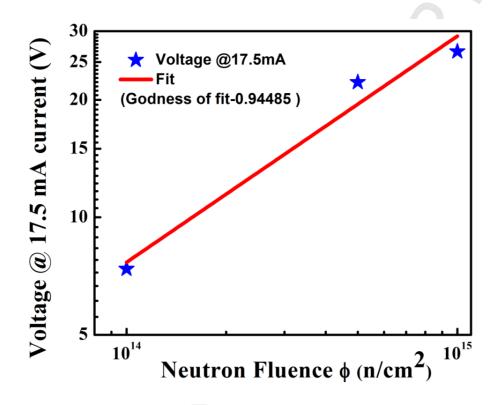


Figure 12



844 Figure 13

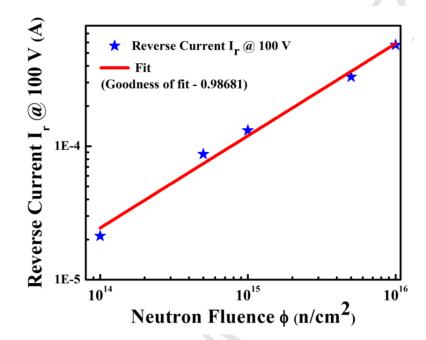


Figure 14

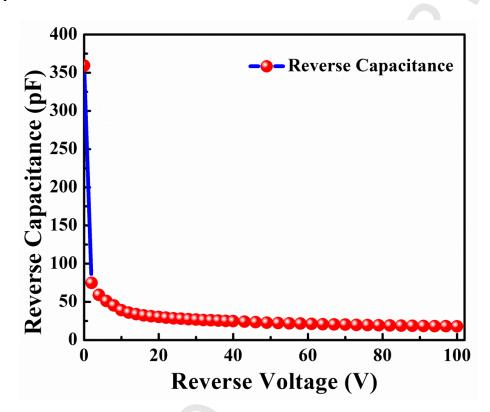


Figure 15

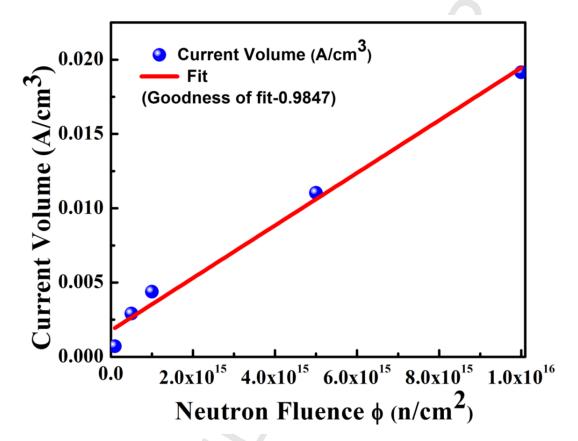


Figure 16

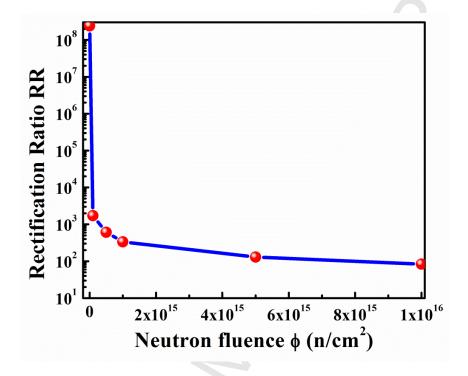


Figure 17

