

## A new method for the measurement of series resistance of solar cells

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Received 23 September 1980, in final form 6 January 1981

**Abstract.** A method is proposed which uses the observed nonlinearity in the short-circuit current versus light intensity curve, at high intensities, to determine the series resistance of solar cells.

In this paper a new method for measuring the series resistance of a solar cell is proposed. The  $I$ – $V$  characteristics of a solar cell under illumination can be written in the form (Wolf and Rauschenbach 1963)

$$I = I_L - I_0 \left[ \exp \left( \frac{q(V + IR_s)}{AkT} \right) - 1 \right] \quad (1)$$

where  $I_L$  is the light generated current,  $R_s$  is the lumped series resistance of the solar cell and other symbols have their usual meaning. In the short-circuit configuration  $V=0$  and equation (1) reduces to

$$I_{sc} = I_L - I_0 \exp(qI_{sc}R_s/AkT) \quad (2)$$

where  $I_{sc}$  is the short-circuit current in the cell and the unity in the square brackets has been neglected. From equation (2) it can be seen that if  $I_{sc}R_s$  is negligible,  $I_{sc}$  varies linearly with intensity of illumination. This is the usual case at low intensities of illumination. However, as the intensity of illumination is increased  $I_{sc}$  increases, the term  $I_0 \exp(qI_{sc}R_s/AkT)$  is no longer negligible in comparison to  $I_L$ , and  $I_{sc}$  varies nonlinearly with intensity. The deviation from linearity is a manifestation of the series resistance and can be used to determine the value of series resistance.

Rearranging equation (2) we obtain

$$\ln [(I_L - I_{sc})/I_0] = qI_{sc}R_s/AkT. \quad (3)$$

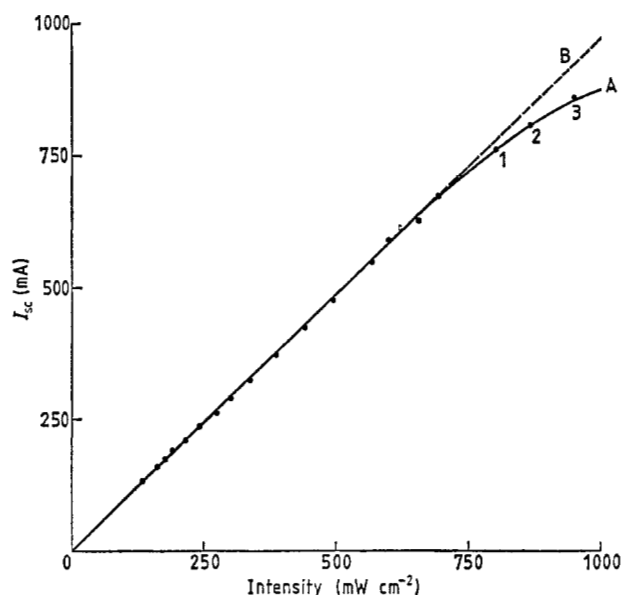
We can see from equation (3) that the plot of  $\ln (I_L - I_{sc})$  versus  $I_{sc}$  will be a straight line with a slope  $s = qR_s/AkT$ . If the diode factor  $A$  is known,  $R_s$  can be calculated from the slope  $s$ .

Using the method explained above, we have determined the value of  $R_s$  for space quality n+p (2 cm × 2 cm) Si solar cells fabricated at the Solid State Physics Laboratory (Delhi). The light source used was an AMO simulator from the Schoffel Corporation (USA). The intensity of illumination was varied by varying the distance between the

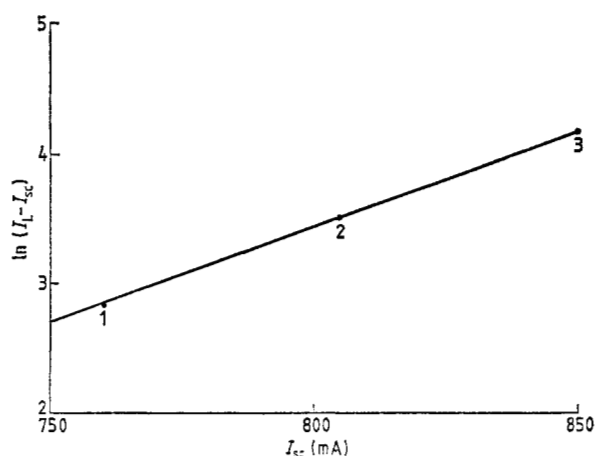
light source and the cell. The temperature of the cell was maintained constant with a cryostat. A compensated calibrated thermopile from Kipps and Zonen (Holland) was used, with the neutral density filters, to measure the intensity of light incident on the cell.

The measured values of  $I_{sc}$  have been plotted as a function of the incident light intensity in figure 1. We see from figure 1 that  $I_{sc}$  is linear with  $I_L$  up to about  $500 \text{ mW cm}^{-2}$ . However, in figure 1, the departure from linearity becomes apparent only beyond about  $700 \text{ mW cm}^{-2}$ .

As we have discussed earlier, in the linear region  $I_{sc} \simeq I_L$ . Since  $I_L$  always varies linearly with intensity (Wolf and Rauschenbach 1963), the extrapolated linear region



**Figure 1.** Measured values of short-circuit current  $I_{sc}$  (curve A) plotted as a function of intensity of illumination. The dashed curve (curve B) corresponds to  $I_L$  extrapolated from the linear region of  $I_{sc}$ . Points 1-3 refer to figure 2.



**Figure 2.**  $\ln(I_L - I_{sc})$  as obtained from figure 1 plotted as a function of  $I_{sc}$ . The points correspond to the numbered points in figure 1.

represented by the dashed curve in figure 1 gives  $I_L$ . From these two curves  $\ln(I_L - I_{sc})$  is calculated, and is plotted as a function of  $I_{sc}$  in figure 2. We see that the plot is a straight line and agrees with the theory. The value of  $A$  is determined (Neugroschel *et al* 1977) from dark-forward characteristics of the cell and was found to be equal to 1. The value of  $R_s$  determined from the slope of the plot in figure 2 was found to be  $0.38 \Omega$ . This value was found to agree with the value obtained from other methods (Imamura and Portscheller 1970).

We shall now discuss the relative merits of our method in comparison to other methods for measuring  $R_s$ . We consider the following three well known methods: (i) the illuminated curve method (Wolf and Rauchenbach 1963); (ii) the p-n junction-dark-forward characteristic method (Imamura and Portscheller 1970); and (iii) the method due to Rajkanan and Shewchun (1979). In (i) the whole of the I-V characteristic has to be measured at two different intensities. In (ii)  $I_{sc}$  and  $V_{oc}$  have to be measured at different light intensities along with the dark-forward characteristics corresponding to each  $V_{oc}$ . This method actually yields  $R_s$  in the dark mode of solar cell operation. In (iii)  $V_{oc}$ ,  $I_{sc}$  and the dark diode voltage corresponding to each  $I_{sc}$  have to be measured at different light intensities.

In contrast to these methods, our method requires the measurement of a single parameter  $I_{sc}$  as a function of light intensity and is therefore quicker and more accurate. We do need to know the value of the diode factor  $A$ , which is usually available from independent measurements or can be obtained quite accurately by measuring the dark-forward characteristic of the solar cell. It may be remarked that in our method it is not necessary to measure the absolute values of the light intensity: it would suffice to vary the incident light intensity by a known ratio, such as by using the neutral density filters.

In all the other three methods given above, it has been assumed that  $I_{sc} = I_L$ . This assumption is valid only at low light intensities, as is well known and is apparent from figure 1. Our method, on the other hand, depends upon the departure of  $I_{sc}$  from  $I_L$  at high intensities. Since the effect of  $R_s$  on the performance of a solar cell is important only at high intensities, our method has the advantage of giving  $R_s$  in such operating conditions. Our method of course cannot be used to give  $R_s$  at low intensities. This method can therefore be used to supplement the information about  $R_s$  obtained by other methods at low intensities.

The chief disadvantage of our method is that measurements have to be made over a wide range of intensities from low intensities where  $I_{sc} = I_L$  to high intensities where the departure of  $I_{sc}$  from linearity becomes substantial. If the value of  $A$  over this range is not constant, it would introduce nonlinear behaviour in the  $\ln(I_L - I_{sc})$  versus  $I_{sc}$  plot. The nonlinearity in this plot may also be caused by the dependence of  $R_s$  on the light intensity. Thus our method will not be reliable in a region of intensity where  $\ln(I_L - I_{sc})$  has a nonlinear dependence on  $I_{sc}$ , but has the advantage of immediately displaying the intensity dependence of  $R_s$  and/or  $A$ . In the present case, as shown in figure 2,  $\ln(I_L - I_{sc})$  is found to be linear with respect to  $I_{sc}$  until at least  $1000 \text{ mW cm}^{-2}$ , and therefore the measured value of  $R_s$  should be quite reliable.

However, at high illumination levels the depletion layer width changes due to voltage drop ( $I_{sc}R_s$ ) across the p-n junction and  $I_{sc}$  may vary nonlinearly with  $I_L$ . We have made detailed numerical calculations using equations (1), (2) and (3) of Agarwala *et al* (1980). These calculations show that even at the highest illumination levels used in our experiments, the change in total  $I_{sc}$  due to the change in the width of the depletion layer is less than 0.2% and so can be neglected.

**References**

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