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Dipoles and defects caused by CO₂ plasma improve carrier transport of silicon solar cells

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Abstract

Carrier-selective contact is a fundamental issue for solar cells. For silicon heterojunction (SHJ) solar cells, it is important to improve hole transport because of the low doping efficiency of boron in amorphous silicon and the barrier stemming from valence band offset. Here, we develop a carbon dioxide (CO₂) plasma treatment (PT) process to form dipoles and defect states. We find a dipole moment caused by longitudinal distribution of H and O atoms. It improves hole transport and blocks electron transport and thus suppresses carrier recombination. In the meantime, the CO₂ PT process also results in defect states, which reduce passivation performance but improve hole hopping in the intrinsic amorphous layer. As a balance, an appropriate CO₂ PT process at the i/p interface increases fill factor and power conversion efficiency of SHJ solar cells. We emphasize, based on sufficient evidences, this work finds a distinct role of the CO₂ plasma in SHJ solar cells opposed to reported mechanisms.

KEYWORDS

carrier transport, CO₂ plasma, defect states, dipole moments, SHJ solar cells

1 | INTRODUCTION

Silicon heterojunction (SHJ) solar cells have attracted continuous attention due to low-temperature processing, high open-circuit voltage (V_{OC}), high power conversion efficiency (PCE), low temperature coefficient, and light-induced performance increase.^{1–4} In particular, a world-record PCE of 26.8% for SHJ solar cells was achieved in 2022 on a 274.4 cm² monocrystalline-silicon (c-Si) wafer.^{5,6}

For solar cells, carrier selectivity plays a key role in the device performance. One of the efficient approaches is to modify the work function (WF) of materials by forming a dipole. The interfacial dipole causes changes in the band structure and thus regulates carrier transport. Incorporating dipole interlayer has been an important interfacial engineering strategy in organic, polymer, and perovskite solar cells. Feng et al.⁷ applied water-alcohol-soluble polyelectrolyte dipole interlayers in single-junction organic solar cells and achieved a PCE of

17.6%. Choi et al.⁸ reported that the WF of ITO increased from 4.77 to 5.02 eV by a dipole interlayer, which contributed to the smooth hole transport of polymer solar cells. Ok et al.⁹ investigated the merits of transition dipole moment to the performances of perovskite solar cells. Canil et al.¹⁰ found electron acceptor or donor molecules resulted in the positive or negative WF shifted up to several hundreds of milielectron volts. Furthermore, the interfacial dipole is also reported in the research field of SHJ solar cells.^{11–13}

Plasma treatment (PT), such as hydrogen (H₂) PT, has been widely studied to improve performance of SHJ solar cells. H₂ PT is generally considered to improve chemical passivation of SHJ solar cell. However, due to the good passivation of the hydrogenated amorphous silicon (a-Si:H) layers, the improvement by this strategy is quite limited, as shown in Table 1.^{14–17} In this study, we report that carbon dioxide (CO₂) PT at the i/p interface after the deposition of intrinsic hydrogenated amorphous silicon (i-a-Si:H) injects oxygen (O) atoms to the a-Si:H network and getters hydrogen (H) atoms from the plasma of the subsequent deposition process. We find H atoms are further away from the i-a-Si:H surface in comparison to O atoms. This leads to an interfacial dipole whose positive end pointing towards the p-type hydrogenated amorphous silicon (p-a-Si:H) layer. As a result, the surface WF (W_{surface}) of i-a-Si:H film significantly decreases by about 0.44 eV, and the energy band bends upward. Consequently, the dipole reduces the hole hopping distance and blocks electron transport. In addition, CO₂ PT also causes defect states that improve hole transition¹⁸ but has a slightly negative impact on chemical passivation.¹⁹ We optimized balance CO₂ PT duration for the i/p interface of SHJ solar cells, which boosts their fill factor (FF) from 77.61% to 82.37% and PCE from 22.16% to 23.59%.

2 | EXPERIMENTAL SECTION

2.1 | Materials

In this study, a-Si:H film was deposited in a plasma enhanced chemical vapor deposition (PECVD) system (ULVAC-CME400), and three different types of a-Si:H film were deposited in independent chambers. The PT process was taken place in the chamber in which doped a-Si:H film was deposited. Each stack would not be exposed to air before production was completed. For Secondary Ion Mass Spectroscopy

TABLE 1 Improvements in V_{OC} and FF through the H₂ PT reported by researchers.

	i-p (a) ¹⁴	i-p (b) ¹⁵	i-n ¹⁶	Both sides ¹⁷
V_{OC} w/o H ₂ PT (mV)	724	~721	722	734.5
V_{OC} with H ₂ PT (mV)	729	~722	728	736.1
ΔV_{OC} (mV)	5	~1	6	1.6
FF w/o H ₂ PT (%)	70.2	~72	78.5	77.23
FF with H ₂ PT (%)	71.3	~73	79.4	77.32
ΔFF (%)	1.1	~1	0.9	0.09

(SIMS) analyses, c-Si/i-a-Si:H/p-a-Si:H stacks with and without a 20 s CO₂ PT process after the deposition of i-a-Si:H film were prepared. Time of Flight Secondary Ion Mass Spectroscopy (TOF-SIMS) analyses were carried out using a TOF-SIMS 5 spectrometer (ION TOF, GmbH-Muenster, Germany). The analysis chamber was maintained at a pressure of 1×10^{-9} mbar. A pulsed 30 keV Bi⁺ primary ion source was employed for analysis at a current of 1.0 pA. Depth profiles were acquired using a Cs⁺ sputter beam with 500 eV for H⁺, O⁺, Si⁺, and B⁺ negative ions under the dual beam mode. Kelvin Probe Force Microscope (KPFM, KEYSIGHT Technologies 7500) measurements were performed to obtain surface potential of c-Si/i-a-Si:H stacks with and without a 20 s CO₂ PT process and a 30 s H₂ PT process. Details of the CO₂ PT process and the H₂ PT process are shown in Table S1. The surface potential is associated with the WF, which is determined by contact potential difference (CPD) between a Pt-coated conductive cantilever probe and the samples. The WF of Pt-coated tip and sample can be calculated using Equation (1).

$$CPD = (\Phi_{\text{tip}} - \Phi_{\text{sample}})/e, \quad (1)$$

where Φ_{tip} is the WF of the tip, Φ_{sample} is the WF of the sample surface, and e is the elementary charge of electron. Herein, the WF of the Pt/Cr-coated tip was calibrated with the HOPG sample ($\Phi_{\text{HOPG}} = 4.6$ eV, $CPD = 420$ mV). Therefore, the Φ_{tip} is 5.02 eV. The Φ_{sample} can be calculated using Equation (2):

$$\Phi_{\text{sample}} = \Phi_{\text{tip}} - e \times CPD. \quad (2)$$

For ultraviolet photoelectron spectroscopy (UPS, Thermo Fisher, ESCALAB 250Xi) measurements, c-Si/i-a-Si:H stacks without PT (REF), with a CO₂ PT process, and with a H₂ PT process after the CO₂ PT process were prepared, respectively. For Fourier-Transform Infrared spectroscopy (FTIR, Perkin Elmer, Spectrum 100) measurements, i-a-Si:H and p-a-Si:H thin films (with and without a 20 s CO₂ PT) were deposited on >3,000 Ωcm float-zone c-Si substrates, and their infrared absorptions were characterized using the transmission mode. The crystallinity of p-a-Si:H film was probed by Raman spectroscopy (WITEC alpha 300R) with a wavelength of 532 nm. Moreover, the dark conductivity was measured through Keithley 6487 at 25°C. The minority carrier lifetime (τ_{eff}) and implied V_{OC} (iV_{OC}) of semi-cells (SHJ solar cells without finger and bus bars) were conducted with a Sinton WCT-120. Schematics of samples for measurements are shown in Figure S1.

2.2 | Devices

The architecture of 6-in. rear-emitter SHJ solar cells are shown in Figure 4B. N-type c-Si (100) wafers (150 μm , $156 \times 156 \text{ mm}^2$, resistivity: 0.3–2.1 Ωcm) were textured with KOH, followed by the standard Radio Corporation of America (RCA) cleaning process, and subsequently dipped in 2% diluted HF for 1 min. Intrinsic and doped a-Si:H films were deposited on both sides of the wafers, and some

samples underwent a CO₂ PT process after the deposition of i-a-Si:H film. After the deposition of doped a-Si:H film, the tungsten-doped indium oxide (IWO) layers were deposited by the reactive plasma deposition (RPD) technique. Subsequently, the Ag electrodes were screen-printed with conventional busbar and finger patterns on the IWO layers. The completed solar cells were characterized by standard 1-Sun light current–voltage (I–V test in-house built setup with a Yamashita class AAA light simulator at 1.5 air mass). Pseudo fill factor (pFF) was measured by the Sinton WCT-120 with Suns-Voc. Electrochemical impedance spectroscopy (EIS) measurements were carried out by Bio-Logic SAS VSP-300 in the dark. Dark J–V curves were scanned from –1 to 1 V. Nyquist plots were obtained at a bias of 0.25 V with an amplitude of 5 mV and frequency range from 0.1 Hz to 1 kHz.

3 | RESULTS AND DISCUSSION

3.1 | Dipoles and defects induced by CO₂ PT

By SIMS measurements, Figure 1A shows the H and O distributions in c-Si/i-a-Si:H/p-a-Si:H stacks with (EXP) and without (REF) a 20 s CO₂ PT process after the i-a-Si:H deposition. The CO₂ discharge in PECVD is known to show intense vacuum-ultraviolet irradiation, producing active oxygen species radicals. These radicals oxidize the surface of the sample to form O-rich thin film, so there are more O atoms near the i/p interface. As shown in Figure S2, the SiO stretching mode at 1,044 cm^{–1} (Fourier-transform infrared spectroscopy; FTIR) of the EXP sample confirms the injection of O atoms.²⁰ Interestingly, H intensity also increases substantially near the i/p interface. The presence of strained Si–O bonds in the O-rich thin film gives rise to an additional channel of interaction of H atoms with film networks.²¹ These extra H atoms are possibly from the plasma of SiH₄ and H₂ during the p-a-Si:H deposition.²² The position of the half peak width of the O and H intensity of the EXP sample is marked in Figure 1A, and it is found that the hydrogen peak is further away from the interface compared to the oxygen peak. Since O is more electronegative than H, local charges are shifted towards O atoms, which causes large dipoles pointing from the i-a-Si:H layer towards the p-a-Si:H layer, as schematic in Figure 1B.

Figure S3 shows the effect of dipoles on the vacuum level and thus WF_{surface} . Dipoles (point outside with respect to the i-a-Si:H surface) will shift the local vacuum level downwards and therefore decrease WF_{surface} of film.^{23,24} To confirm the formation of the dipole on the i-a-Si:H surface, WF_{surface} of i-a-Si:H film with and without a CO₂ PT process and a H₂ PT process was measured by KPFM. The H₂ PT process after the CO₂ PT process was to mimic the H₂ PT that would occur during the p-a-Si:H deposition. As shown in Figure 1C, the CPD of the i-a-Si:H film increases significantly after the PT process, which indicates decrease of WF_{surface} . WF_{surface} of i-a-Si:H film before and after the PT process is calculated to be 4.60 and 4.16 eV, respectively, suggesting the PT process decreased WF_{surface} by 0.44 eV. Furthermore, UPS measurements were utilized to investigate

the surface electronic states of i-a-Si:H film. The valence band edge curves of the i-a-Si:H film without PT (REF), with a CO₂ PT process and with a H₂ PT process after the CO₂ PT process are shown in Figure S4. The energy difference between the valence band edge (E_V) and the Fermi level (E_F) is 1.30, 1.08, and 1.35 eV, respectively. After CO₂ PT, electronegative O atoms gather on the surface of film, the energy band bends upward since negative charges on the surface, and the $E_F - E_V$ value decreases. After CO₂ PT, a H₂ PT process was carried out. Then, H atoms with weak electronegativity gather on the surface of O-rich film, the energy band bends downward since positive charges on the surface and $E_F - E_V$ value increases. These results also confirm that the PT process causes the existence of dipoles. Figure 1D represents a schematic diagram of the energy bands in the hole-selective contact stack with and without the CO₂ PT treatment. Since the energy band of the i-a-Si:H bends upward due to dipoles pointing from the i-a-Si:H layer towards the p-a-Si:H layer (green lines), there is a higher barrier for electron transport that is beneficial to suppress the charge recombination. The implication on hole transport is discussed in the next section.

Mazzarella et al.²⁵ and Jiang²⁶ found that CO₂ PT on the i/p interface of the SHJ solar cells promoted nucleation and improved the conductivity of p-type emitter layer. To study the structural features of p-a-Si:H film, Raman spectra of samples without and with a 20 s CO₂ PT process are shown in Figure S5. Spectra exhibit no changes, demonstrating CO₂ PT does not promote the formation of nanocrystalline phase. Dark current–voltage curve of p-a-Si:H thin film is shown in Figure S6. The curve changes slightly, which indicates that the dark conductivity of p-a-Si:H thin film is hardly improved after CO₂ PT. Our observations are distinctly different from other researches, indicating that there are different mechanisms underlying the impact of CO₂ PT under low pressure on silicon solar cells in this study.

We treated tungsten-doped indium oxide (IWO)/n-type hydrogenated amorphous silicon(n-a-Si:H)/i-a-Si:H/n-c-Si/i-a-Si:H/p-a-Si:H/IWO samples with different CO₂ PT duration on the i-a-Si:H layer at the p-a-Si:H side. Figure 2A,B shows that both τ_{eff} and iV_{OC} decrease slightly with short CO₂ PT duration, τ_{eff} and iV_{OC} of the sample with a 30 s CO₂ PT process decrease significantly compared to that of the REF sample. The pseudo fill factor (pFF) characterizes losses due to injection-dependent recombination.²⁷ Figure 2C shows pFF of SHJ solar cells with different CO₂ PT duration before the deposition of the p-a-Si:H layer. Each plot is an average value from three pieces of 6-in. rear-emitter SHJ solar cells (~244 cm²). Along with the increase of CO₂ PT duration, pFF continues to decrease from 84.8% (REF) to 83.03% (30 s). This passivation degradation indicates an increase of defect states.¹⁹ Figure 2D shows the forward dark current density–voltage (J–V) characteristics of SHJ solar cells with different CO₂ PT duration. The forward current in SHJ solar cell is dominated by the diffusion current and the tunneling current at higher bias voltage and lower bias voltage, respectively.²⁸ The tunneling current of cells with a 10 s or 20 s CO₂ PT process is slightly larger than that of the REF cell, but the tunneling current increases significantly when CO₂ PT duration is increased to 30 s. The increase of the tunneling current is caused by the increase of the depletion region recombination and

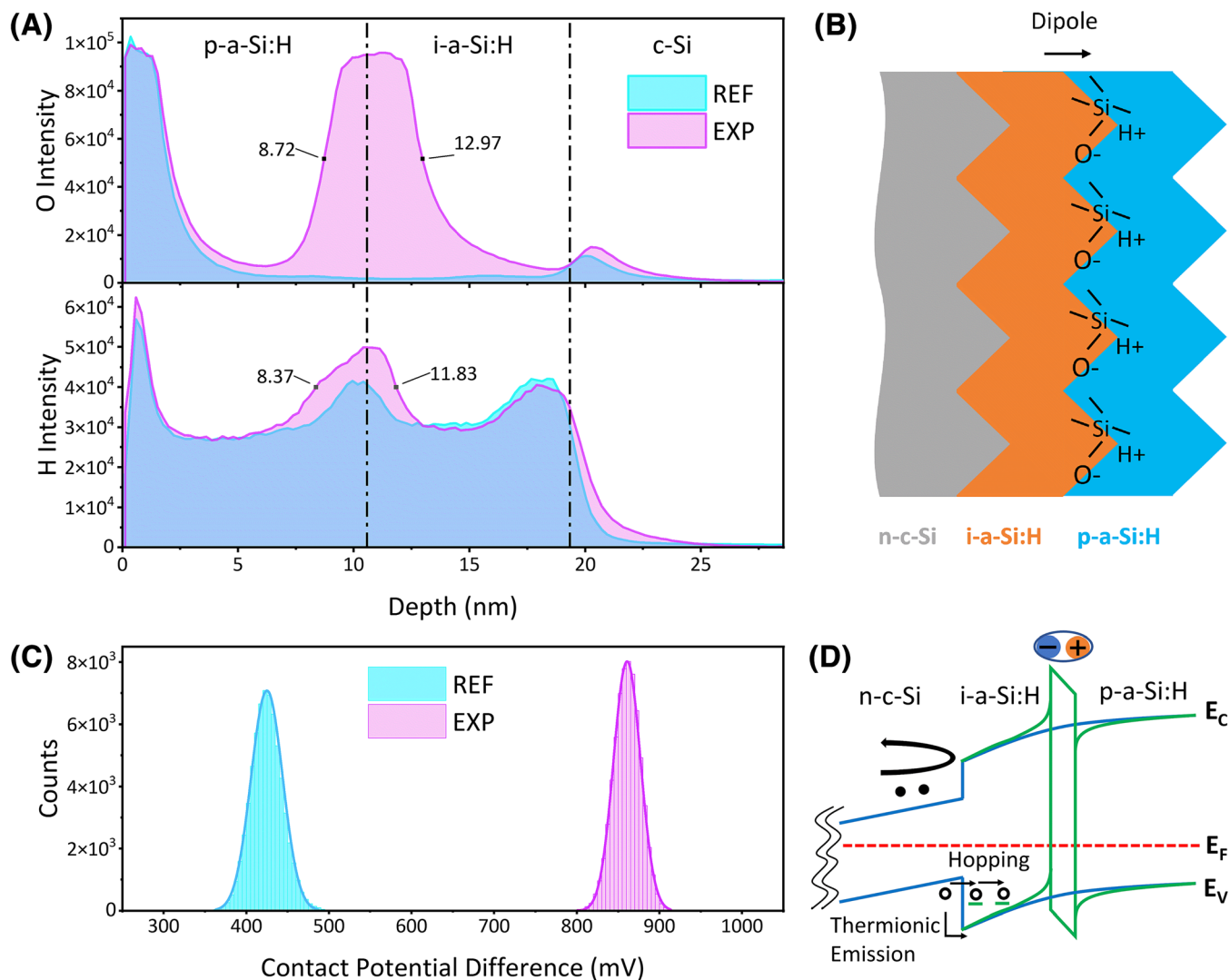


FIGURE 1 (A) Distributions of H and O in c-Si/i-a-Si:H/p-a-Si:H stacks with (EXP) and without (REF) a 20 s CO₂ PT process after the i-a-Si:H deposition. (B) Schematic of hole-selective contacts with a CO₂ PT process. (C) Contact potential difference (CPD) of i-a-Si:H films before (REF) and after (EXP) a 20 s CO₂ PT process and a 30 s H₂ PT process. (D) Schematics of charge carrier collection at the hole contact without (blue lines) and with (green lines) dipoles.

shunting path within the p-n junction. Figure 2E shows typical Nyquist plots of the SHJ solar cells with different CO₂ PT duration, recorded in the dark at a bias voltage of 0.25 V. The values of the shunt resistivity (R_{sh}) obtained from EC-lab analysis of the EIS data are shown in the inset. R_{sh} decreases with the increase of CO₂ PT duration. This demonstrates longer CO₂ PT increases defect states, carrier recombination, and shunting currents, which will decrease R_{sh} of the device and thus V_{OC} and FF to some extent.

3.2 | Performance of SHJ solar cells

To our surprise, these defect states on the other hand also show a positive role. Pradyumna Muralidharan et al.^{18,29} claimed that for a SHJ solar cell, charge transport through the i-a-Si:H layer depends

on the defect/phonon assisted transport. Holes hop from defect to defect within the i-a-Si:H layer before they are collected. Longer transit times are indicative of photocurrent suppression, which leads to an S-shape behavior in the J - V curve of SHJ solar cells. More available defect states reduce the number of phonons required for hole injection into the barrier, thereby reducing the transit time associated with injection and thus reducing the series resistance (R_s). In addition, as shown in Figure 1D, the valence band offset does not seem to change at the c-Si/i-a-Si:H interface, which means the thermionic emission and the direct tunneling are almost unaffected, but the hopping distance reduces as the valence band bends upward, which is also beneficial for hole to hop through the i-a-Si:H layer. The difference between FF and pFF is a misleading measure of R_s losses.³⁰ In Figure 3A,B, We plot $pFF - FF$ and R_s of solar cells prepared with various CO₂ PT duration. They show the same trend. R_s of SHJ solar cells

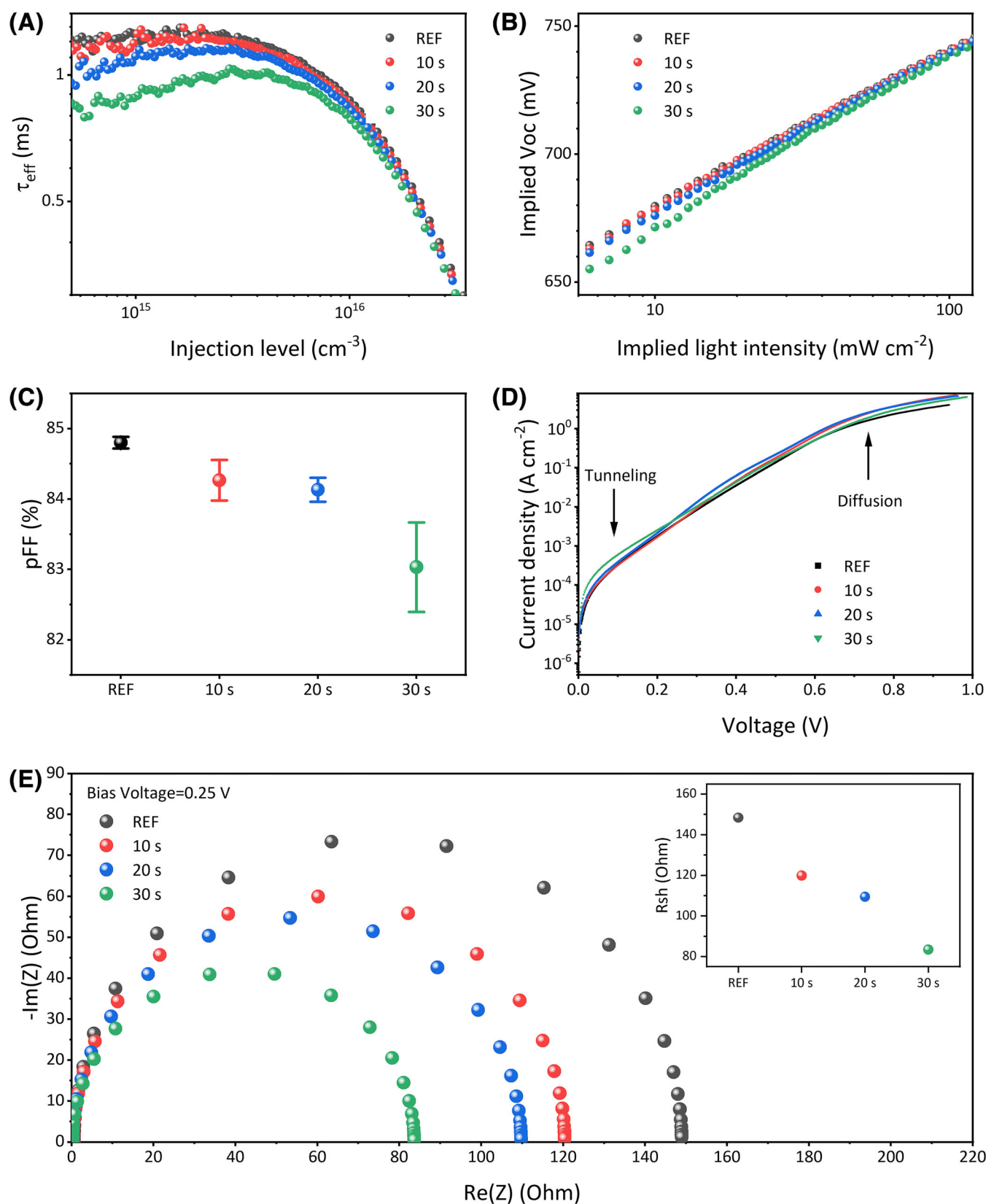


FIGURE 2 (A) Injection-dependent τ_{eff} and (B) implied V_{OC} (iV_{OC}) dependence on implied light intensity of IWO/n-a-Si:H/i-a-Si:H/n-c-Si/i-a-Si:H/p-a-Si:H/IWO with different CO₂ PT duration. (C) Pseudo-FF plots and (D) dark J-V curve of SHJ solar cells with different CO₂ PT duration. (E) Nyquist plots of electrochemical impedance spectroscopy (EIS) for devices with different CO₂ PT duration in the dark at a bias of 0.25 V. The inset shows the shunt resistance (R_{sh}) at a bias of 0.25 V of SHJ solar cells with different CO₂ PT duration obtained from EC-lab analysis of the EIS data.

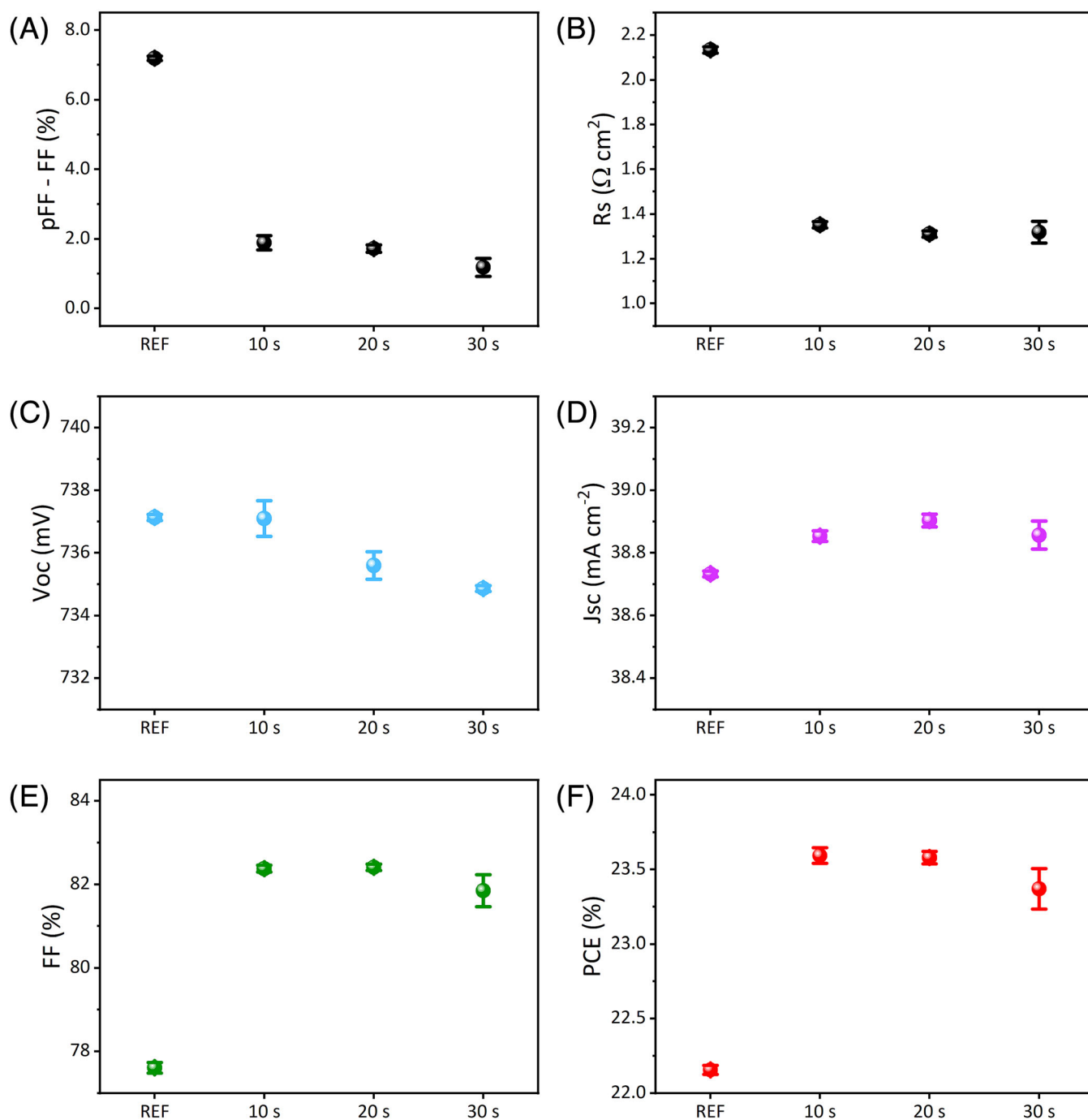


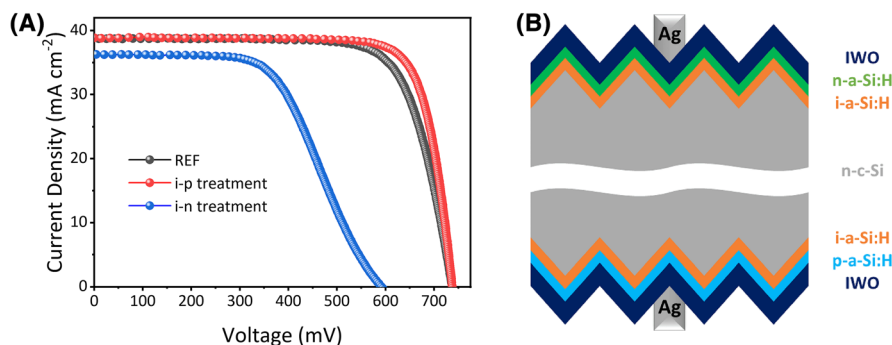
FIGURE 3 Performance of SHJ solar cells with different CO₂ PT duration, (A) pFF-FF, (B) R_s , (C) V_{oc} , (D) J_{sc} , (E) FF, and (F) PCE. The values are average values from three pieces of 6-in. rear-emitter SHJ solar cells (~ 244 cm²), and error bars represent the standard deviation.

with a 10 s CO₂ PT process significantly decreases but does not decrease further with additional CO₂ PT duration.

Output parameters of 6-in. SHJ solar cells (~ 244 cm²) with different CO₂ PT duration are plotted in Figure 3C-F. With the increase of CO₂ PT duration, V_{oc} decreases gradually, in good agreement with iV_{oc} , dark $J-V$ and R_{sh} , stemming from increasing density states and passivation degradation. However, it has a gain about 0.15 mA cm⁻² in short-circuit current density (J_{sc}), and more importantly, FF boosts

from 77.61% to 82.37% (10 s) and 82.41% (20 s) due to a significant decrease in R_s . When CO₂ PT duration is added to 30 s, FF decreases to 81.85%. Since both R_s and R_{sh} take effects on FF of SHJ solar cells,³¹ this trend means that the gain of R_s reduction on FF is much greater than the loss of R_{sh} reduction, but the loss of FF cannot be ignored in case of overlong CO₂ PT. Thanks to improvement of FF, PCE increases from 22.16% to 23.59% or 23.58% with a 10 or 20 s CO₂ PT duration and then decreases to 23.37% with a 30 s CO₂ PT

FIGURE 4 (A) One-Sun J - V curve of SHJ solar cells with a CO_2 PT at different sides. (B) The schematic diagram of the 6-in. rear-emitter SHJ solar cell



process. The J - V curve under simulated AM 1.5G solar irradiation of the best device named i-p treatment is plotted in Figure 4A. The best device undergoing a 10 s CO_2 PT process shows the highest PCE of 23.64% with a J_{SC} of 38.87 mA cm^{-2} , a V_{OC} of 737.6 mV, and an FF of 82.45%.

Moreover, we also apply CO_2 PT to the interface of i-n stacks of SHJ solar cells. The J - V curve under simulated AM 1.5G solar irradiation of the best device named i-n treatment is plotted in Figure 4A. The J - V curve shows S-shape, and the device performance shows an overall degradation; that is, V_{OC} , J_{SC} , FF, and PCE are 591.9 mV, 36.29 mA cm^{-2} , 58.81%, and 12.63%, respectively. The poor performance indicates the selective electron transport of the device is seriously hindered. This further proves that the interface with dipoles and defect states caused by CO_2 PT helps hole transport but hinders electron transport.

4 | CONCLUSION

In conclusion, we demonstrate that the CO_2 PT after the deposition of i-a-Si:H film injects O atoms to amorphous silicon network and then getters H atoms from the plasma of SiH_4 and H_2 during the p-a-Si:H deposition. This leads to a longitudinal distribution of H and O atoms that creates a dipole moment. It reduces $W_{\text{F surface}}$ by 0.44 eV and causes the conduction band and the valence band to bend upwards and thus reduces the hole hopping distance and blocks electron transport. We also demonstrate that the CO_2 PT increases defect states. To some extent, more defect states reduce passivation performance, which causes the decrease of R_{sh} and V_{OC} , but are beneficial for hole to hop through the i-a-Si:H layer. Even though increasing the CO_2 PT duration from 10 to 30 s leads to some performance loss, the key thing is that even at 30 s, there is a clear benefit to the solar cell efficiency. As a whole consequence, a 10 s CO_2 PT process at the i/p interface of SHJ solar cells gains their FF by 6.13% and PCE by 6.45% relatively.

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DATA AVAILABILITY STATEMENT

The data that support the findings of this study are available from the corresponding author upon reasonable request.

CONFLICT OF INTEREST STATEMENT

The authors declare no conflict of interest.

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SUPPORTING INFORMATION

Additional supporting information can be found online in the Supporting Information section at the end of this article.

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