# Equivalent Circuit of Quantum-Dot LED and Acquisition of Carrier Lifetime in Active Layer

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Abstract—We propose an equivalent circuit for a quantum-dot LED (QLED), where the resistances and the capacitances are expressed in terms of the physical parameters of the QLED. The validity of the equivalent circuit is verified by measurement results. From the measured frequency response of the QLED and the calculated frequency response of the equivalent circuit, we are able to deduce the carrier lifetime in the active layer of the QLED. The availability of the equivalent circuit can facilitate the analysis of the electrical characteristics of QLEDs.

Index Terms—Quantum-dot light-emitting diode, equivalent circuit, frequency response, carrier lifetime.

## I. INTRODUCTION

Electroluminescent (EL) quantum-dot (QD) LEDs (QLEDs) are promising light sources for display and illumination for their inkjet processability, high color tunability, and high color purity [1]. Their large bandwidths [2], compared with those of other solution-processed light sources, also render them suitable for communication applications. In this paper, we propose an equivalent circuit for the analysis of QLEDs and demonstrate the evaluation of the carrier lifetime in the active layer of the QLED with the help of the equivalent circuit.

Equivalent circuits have been proposed for semiconductor LEDs [3] and organic LEDs (OLEDs) [4]–[6] to facilitate the analysis of the electrical characteristics of these devices. These equivalent circuits are mostly developed by fitting the values of the circuit components with measured electrical characteristics of the devices. As such, the circuits are not directly related to the physical structures and the materials of the devices. Besides, some DC models [5] do not allow the evaluation of the modulation characteristics of the device. There is no equivalent circuit specifically developed for QLEDs.

The equivalent circuit we propose for a QLED reflects the physical structure of the device, where the component values are expressed in terms of the physical parameters of the device. Compared with reported equivalent circuits of OLEDs [4]–[6], our circuit allows a more direct and accurate evaluation of the roles of different layers of the device on its electrical characteristics. The validity of our circuit is verified with measured electrical characteristics. By measuring the frequency

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response of the QLED and with the effect of the equivalent circuit deducted, we are able to obtain the carrier lifetime in the active layer of the QLED. Carrier lifetime in the active layer, namely the lifetime of carriers that exist in the QD layer before recombination, is an important parameter of QLED, as it directly reflects the recombination process and the electron-hole dynamics in the active layer [7]. Existing methods for measuring carrier lifetimes in semiconductor structures, such as those based on the measurements of photoluminescence decay [8] and open-circuit voltage decay [9], could be applied to a QLED, but the result includes all the effects from the injection and transport layers of the QLED. There is no method reported for direct measurement of the carrier lifetime in the active layer of a QLED. Our approach provides a simple way to find this important parameter.

# II. DEVELOPMENT OF EQUIVALENT CIRCUIT

Figure 1(a) shows the structure of a 4-mm<sup>2</sup> CdSe/ZnS QLED considered in our study, which consists of the following layers: indium tin oxide (ITO, 100 nm) / poly(3.4-ethylenedioxythiophene) polystyrene sulfonate (PEDOT:PSS, 40 nm) / poly (9,9-dioctylfluorene-co-N-(4-(3-methylpropyl)) diphenylamine) (TFB, 29 nm) / CdSe/ZnS QDs (18 nm) / Zinc Oxide (ZnO, 33 nm) / Al (100 nm). The ITO, PEDOT:PSS, TFB, QD, ZnO, and Al layers serve as the anode, the hole injection layer (HIL), the hole transport layer (HTL), the emitting layer, the electron transport layer (ETL), and the cathode, respectively. The thicknesses were measured with a Bruker Stylus Profiler. This QLED emitted red light at a wavelength of 634 nm with a full-width-at-half-maximum (FWHM) of 22 nm. Details on the measurement of the optical properties and the frequency response of the QLED are given in [2]. Figure 1(b) shows the energy level diagram of the QLED [2]. In our experiments, a DC bias voltage (or an input voltage) varying from 3 to 10 V was applied to the QLED. For the measurement of the frequency response of the device, a sinusoidal signal with an amplitude of 1V was superimposed on the DC bias.

Figure 1(c) shows the proposed circuit of the forward-biased QLED and Fig. 1(d) shows the correspondence between the circuit components and the physical layers of the device.

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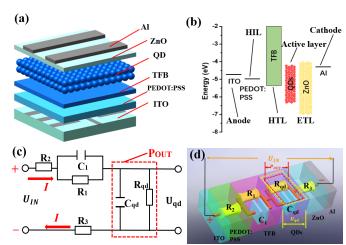


Fig. 1. (a) Schematic diagram of the structure of a CdSe/ZnS QLED; (b) energy level diagram of the QLED; (c) the proposed equivalent circuit of the QLED, and (d) the correspondence between the circuit components and the physical layers of the QLED.

The QD layer is equivalent to a parallel circuit of a forwardcurrent-caused capacitance  $C_{ad}$  and a resistance  $R_{ad}$ . Because of the low accessibility of holes from the TFB layer into the QD layer, the stack of holes at the surface between the HTL layer and the QD layer forms an accumulated-charge-caused barrier capacitance  $C_1$ . Therefore, the TFB layer is equivalent to a parallel circuit of the barrier capacitance  $C_1$  and a resistance  $R_1$ . These two parallel circuits are in series of the resistances of the PEDOT:PSS layer and the ZnO layer,  $R_2$  and  $R_3$ . Because the energy gaps across Al/ZnO and ZnO/QDs are small, as shown in Fig. 1(b), electrons can be easily injected from the Al layer into the ZnO layer and the QD layer. Electron accumulation in the ZnO/QDs interface and hence the capacitance in the ZnO layer can be ignored. If necessary, electron accumulation in the ZnO/QDs interface, which may occur in blue QLEDs [10], can be taken into account by placing an accumulated-charge-caused barrier capacitance across  $R_3$ . Similarly, the capacitance in the PEDOT:PSS layer can be ignored for its high carrier mobility and good electrode contact. The resistances of the Al cathode and the ITO anode are negligible and can also be ignored.

The transfer function of the QLED,  $Q(\omega)$ , is given by the ratio of the emitted optical power  $P_{OUT}(\omega)$  and the input voltage  $U_{IN}(\omega)$ , where  $\omega$  is the angular modulation frequency. According to Fig. 1(c) and (d), the optical power  $P_{OUT}$  should be generated from the voltage  $U_{qd}$  across the QD layer. The electrical-optical conversion in the QLED can then be described as the conversion from  $U_{IN}$  to  $U_{qd}$  followed by the conversion from  $U_{qd}$  to  $P_{OUT}$ .

The transfer function that describes the conversion from  $U_{IN}$  to  $U_{qd}$ , denoted as  $H(\omega)$ , depends solely on the components in the equivalent circuit. These components consume electric power and slow down the transportation of carriers in the QLED. The conversion from  $U_{qd}$  to  $P_{OUT}$ , denoted as  $F(\omega)$ , takes place in the active layer of the QLED and depends on the rate of electron-hole recombination. Similar to the case of a luminescent diode, the response in the active layer depends on the carrier lifetime  $\tau$  as  $1/\sqrt{1+\omega^2\tau^2}$  [11], which consists of both radiative and non-radiative contributions. Therefore, with

the help of the equivalent circuit, the transfer function of the QLED is obtained as

$$Q(\omega) = H(\omega)F(\omega), \tag{1}$$

with

$$H(\omega) = \frac{U_{qd}(\omega)}{U_{IN}(\omega)} = \frac{\frac{R_{qd} - j\omega R_{qd}^2 C_{qd}}{1 + (\omega R_{qd} C_{qd})^2}}{\frac{R_{qd} - j\omega R_{qd}^2 C_{qd}}{1 + (\omega R_{qd} C_{qd})^2} + \frac{R_{1} - j\omega R_{1}^2 C_{1}}{1 + (\omega R_{1} C_{1})^2} + R_{2} + R_{3}},$$
(2)

$$F(\omega) = \frac{P_{OUT}(\omega)}{U_{qd}(\omega)} = \frac{A}{\sqrt{1 + \omega^2 \tau^2}},$$
 (3)

where A is a constant.

The resistivity of each layer in the QLED is given by  $\rho=\rho_0/\kappa(1-In\kappa)$  [12], where  $\rho_0=1/ne\mu$  is the bulk resistivity,  $n=J/e\mu E$  is the electron concentration with J being the current density, e the electron charge,  $\mu$  the mobility of the material, and E the electric field in the layer, and  $\kappa$  is given by  $d/2l_{bulk}$  with d being the layer thickness and  $l_{bulk}$  the bulk mean free path. The resistance of the layer is then given by  $R=\rho d/S$ , where S=4 mm² is the emission area of the QLED. The values of  $l_{bulk}$  for PEDOT:PSS [13], TFB [14], QDs [15], ZnO [16] are 2.7, 3, 20, and 2.6 nm, respectively, and the values of the mobilities for ITO [17], PEDOT:PSS [18], TFB [19], QDs [20], ZnO [21], and Al are 20,  $3\times10^{-2}$ ,  $1.2\times10^{-2}$ ,  $1.0\times10^{-2}$ ,  $2.3\times10^{-3}$ , and  $10^3$  cm² $V^{-1}s^{-1}$ , respectively.

The injection current density J can be measured experimentally [2], but the electric field E in each layer cannot be measured. Instead, we calculate E with the commercial software Setfos [22] – the software is widely used for the analysis of optoelectronic devices. The calculation takes into account all the QLED parameters, namely the layer thicknesses, the mobilities, the dielectric constants, and the energy levels of different layers shown in Fig. 1(b), as well as the applied voltage. Given the electric field E in each layer and the measured injection current density J, we can calculate the bulk resistivity  $\rho_0 = E/J$  in each layer and hence the corresponding resistance.

The barrier capacitance caused by the space charges is given by  $C_1 = (\varepsilon \varepsilon_0/\sqrt{2}L_d) \exp{(-eV_{d1}/2kT)}$  [23], where  $\varepsilon$  is the dielectric constant of the material,  $\varepsilon_0$  is the permittivity of the vacuum,  $L_d = \sqrt{\varepsilon kT/(e^2n)}$  is the debye length, k is the Boltzmann coefficient, T (= 298 K) is the temperature, and  $V_{d1}$  is the potential difference between the TFB layer and the QD layer. The values of  $\varepsilon$  for PEDOT:PSS [24], TFB [24], QDs [25], ZnO [26] are 3, 3, 6, and 3.5, respectively. The amount of residual charges  $Q_R$  in the QD layer is the amount of charges  $Q_I$  in the current I minus the amount of charges  $Q_P$  consumed for photon conversion. Therefore,  $C_{ad}$  can be expressed as

$$C_{qd} = S \frac{dQ_R}{dU_{IN}} = S \frac{dQ_I}{dU_{IN}} - S \frac{dQ_P}{dU_{IN}}, \tag{4}$$

with  $Q_I = It$  and  $Q_P = N_e et = \frac{et}{hc} \int \frac{P(\lambda)\lambda}{\eta_{EQE}} d\lambda$ , where t is the unit time,  $N_e$  is the electron number, h is Planck's constant, c is the speed of light in vacuum,  $\eta_{EQE}$  is the external quantum efficiency, and  $P(\lambda)$  is the spectral power density of the emitted light from a packaged QLED. The effects of both radiative and non-radiative recombinations are implicitly taken into account by  $\eta_{EQE}$ .  $C_{qd}$  can be calculated from Eq. (4) with the

knowledge of  $dI/dU_{IN}$ ,  $P(\lambda)$ , and  $\eta_{EQE}$ , which can be found experimentally [2].

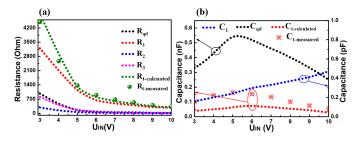


Fig. 2. (a) Values of  $R_{qd}$ ,  $R_1$ ,  $R_2$ , and  $R_3$  calculated at different applied voltages and comparison between calculated and measured input resistance  $R_t$ ; (b) values of  $C_{qd}$  and  $C_1$  calculated at different applied voltages and comparison of calculated and measured input capacitance  $C_t$ .

## III. RESULTS AND DISCUSSIONS

Figure 2 shows the values of the resistances and the capacitances in the equivalent circuit calculated at different input voltages. As shown in Fig. 2(a),  $R_1$  has the largest value compared with other resistances and all the resistances decrease with an increase in the applied voltage. The reason is that the electric potential changes severely between the TFB layer and the QD layer, which directly influences the values of E and  $R_1$ . As shown in Fig. 2(b),  $C_{qd}$  has a peak value at a voltage of 5.2 V and the value of  $C_1$  increases monotonically with the applied voltage. The peak of  $C_{qd}$  is due to the balance between the speed of charge injection and the speed of charge conversion into light in the QD layer. As  $C_1$  is inversely proportional to the carrier density, its value increases with the applied voltage. The input impedance of the circuit can be written as a resistance  $R_t$  in parallel with a capacitance  $C_t$ , where  $R_t$  is the sum of  $R_{qd}$ ,  $R_1$ ,  $R_2$ , and  $R_3$ .

We calculate the transfer function  $H(\omega)$  of the equivalent circuit and the corresponding bandwidth B. With the knowledge of the bandwidth B and the resistance  $R_t$ , we deduce  $C_t$  from the expression  $B = 1/(2\pi R_t C_t)$ . We measured the values of  $R_t$  and  $C_t$  of the QLED at different DC bias voltages with a semiconductor parameter analyzer [2] and compare the measurement data with the calculated values. As shown in Fig. 2, the measured values agree well with the calculated values, which thus verifies the validity of the equivalent circuit. We measured the frequency response of the QLED, i.e.,  $Q(\omega)$ , by varying the modulating frequency of the input sinusoidal signal at different DC bias voltages. The results are shown in Fig. 3(a). We also calculate the frequency response of the equivalent circuit  $H(\omega)$  and the results are shown in Fig. 3(b).

The bandwidth of  $H(\omega)$  is wider than that of  $Q(\omega)$  at the same voltage. For example, at a bias voltage of 10 V,  $H(\omega)$  has a bandwidth of ~20 MHz, while  $Q(\omega)$  has a bandwidth of only ~8 MHz. We obtain  $F(\omega)$  by dividing  $Q(\omega)$  by  $H(\omega)$ . The results are shown by the discrete points in Fig. 3(c). By fitting the discrete points with curves in the form of  $F(\omega) = 1/\sqrt{1 + \omega^2 \tau^2}$ , we find the values of the carrier lifetime  $\tau$  at different applied voltages, which are shown by the discrete points in Fig. 3(d). The carrier lifetime decreases with an increase in the applied voltage, which indicates that the

response of the QLED is faster at a higher applied voltage. This finding agrees with our previous experimental study [2], where the bandwidth of a QLED is shown to be mainly limited by the carrier lifetime at a large injection current.

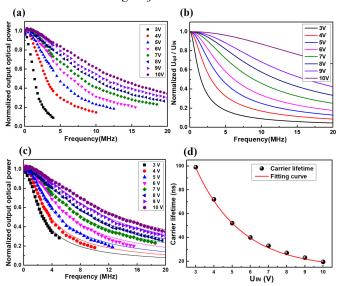


Fig. 3. (a) Measured frequency response of the QLED,  $Q(\omega)$ , (b) calculated frequency response of the equivalent circuit,  $H(\omega)$ ; (c)  $F(\omega)$  obtained by dividing  $Q(\omega)$  by  $H(\omega)$ , and (d) carrier lifetime obtained by fitting the results in (c) with  $F(\omega) = 1/\sqrt{1 + \omega^2 \tau^2}$ .

## IV. CONCLUSION

We have proposed an equivalent circuit for QLEDs, where the resistances and the capacitances are expressed in terms of the physical parameters of the device. We have measured the input resistance and capacitance of a specific QLED and the measurement data agree well with the results calculated from the equivalent circuit. We are able to deduce the carrier lifetime in the active layer of the QLED by dividing the frequency response of the device by that of the equivalent circuit. Our proposed equivalent circuit provides a convenient model for the analysis of the electrical characteristic of QLEDs within the linear operation regimes of the devices.

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