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# Temperature dependent device characteristics of graphene/h-BN/Si heterojunction

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**Temperature Dependent Device Characteristics of Graphene/h-BN/Si Heterojunction**

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Graphene-on-semiconductor heterojunction solar cell is an emerging class of photovoltaics with potential for efficient and reliable energy conversion systems. The interfaces between graphene and lightly-doped semiconductor play a key role in charge-carrier separation and recombination dynamics. Owing to the low Schottky barrier height-induced interfacial charge carrier recombination, the graphene-on-silicon (Si) heterojunction solar cells suffer from instability in power conversion efficiency over time. Therefore, it is critical to engineer the interface to enhance the barrier height by interfacing a chemically-stable, insulating, and atomically-thin layer. Further, the temperature dependent photovoltaic characteristics of such stacked architectures are unknown, and temperature dependent behavior is critical to understand the MIS junction behavior and photovoltaic phenomenon. Here, we have introduced hexagonal boron nitride (h-BN) as a tunneling interlayer in graphene-on-Si heterojunction solar cells, which enables the passivation of the chemical dangling bonds on the Si surface. The effect of temperature on the performance of graphene/h-BN/Si PV cell is examined. Thin films of h-BN are directly synthesized on lightly-doped Si surface *via* a bottom-up chemical-surface-adsorption strategy followed by the transfer of a graphene monolayer. The 2D layer-on-2D layer-on-3D bulk semiconductor nanoarchitecture of graphene/h-BN/Si forms a metal-insulator-semiconductor (MIS)-type junction, where the h-BN acts as an electron-blocking layer to avoid interfacial charge carrier recombination. A 4-fold increase in open-circuit voltage ( $V_{oc}$ ) is found for graphene/h-BN/Si heterojunction cell (0.52 V) in contrast to the graphene/Si cell (0.13 V), which is due to the increase in the Schottky barrier height and hence built-in electric potential. Interestingly, the  $V_{oc}$  linearly decreases by only ~4% with every 10 K increase in temperature. This work will lead to an evolution of new 2D/2D/3D nanoarchitectures for mechanically-robust, high performance, and durable optoelectronic functionalities.

**INTRODUCTION**

Owing to the high optical transparency and ultra-high charge carrier mobility,  $sp^2$ -hybridized graphene<sup>1-3</sup>, which can function both as a transparent electrode and an active layer, is a potential candidate two-dimensional (2D) material for photonics, optoelectronics, and photovoltaics. A rectifying Schottky junction can be constructed by interfacing graphene with lightly doped Si, which can be employed as a photovoltaic cell<sup>4</sup> under the irradiation of light. The graphene-on-semiconductor heterostructure, which benefits from cost-effective fabrication process, has shown promising power conversion efficiency of over 10% at AM 1.5G irradiation with functionalized or doped graphene<sup>5-11</sup> and with an anti-reflection layer<sup>12-15</sup>. However, the graphene/Si-based 2D layer-on-3D bulk semiconductor solar cell suffers from serious instability in the photovoltaic performance.<sup>16</sup> This stems from the existence of an astronomically high dark-current in the graphene/Si Schottky barrier junctions, which is caused by the thermionic-emission

based dark current.<sup>17</sup> What's more, mid-gap surface states can be formed by dangling bonds on the Si surface, which will provide additional pathway for the charge carrier recombination.<sup>18</sup>

To reduce the thermionic-emission based dark current and charge-carrier recombination, a non-reactive and stable passivation layer is required to separate metal and semiconductor by forming a metal/insulator/semiconductor (MIS) architecture.<sup>19</sup> Previously, thin film insulators like transferred h-BN<sup>20</sup>, silicon dioxide (SiO<sub>2</sub>)<sup>21</sup>, Al<sub>2</sub>O<sub>3</sub><sup>22</sup> and fluorographene (FG)<sup>23</sup> have been reported as an electron blocking layer in MIS structure for photovoltaic application. Unlike SiO<sub>2</sub> and fluorographene,<sup>21</sup> the h-BN layer does not have trap-charges and surface states.<sup>24,25</sup> Hexagonal boron nitride (h-BN) is a sp<sup>2</sup>-hybridized 2D insulator with wide energy bandgap of 5.97 eV and ultra-smooth surface with no surface-dangling-bonds<sup>26–28</sup>. Further, the h-BN is isostructural and isoelectronic to graphene with lattice mismatching of only 1.7%, thus enabling compatible graphene/h-BN heterostructure.<sup>29</sup> Similar MIS structure has been built by transferring the h-BN layer<sup>30</sup> on Si surface, which can introduce heteroatom contaminations during chemical-transfer process.<sup>31</sup> Therefore, direct growth of h-BN on Si surfaces is critical to avoid the defects and for scalability. The temperature dependent diode and photovoltaic characteristics of the graphene/h-BN/n-Si heterojunction solar cell can be crucial to understand the charger carrier separation and recombination dynamics at the MIS structure interface. Previously, Kalita et al. investigated temperature dependent study on metal-semiconductor (MS)-type graphene-GaN heterojunction and MIS-type graphene/h-BN/GaN in the range of 298 to 373 K.<sup>32,33</sup> Graphene-GaN heterojunction showed photovoltaic effect and its V<sub>oc</sub> decreased nonlinearly with the temperature. Graphene/h-BN/GaN heterojunction showed increased Schottky barrier height. Temperature dependent study of Si based MIS solar cell has been explored in many previous studies.<sup>34</sup> However, no temperature dependent photovoltaic characteristic of graphene based MIS structure has been reported. It is therefore critical to investigate the temperature dependent photovoltaic characteristics of a graphene/h-BN/n-Si heterojunction device.

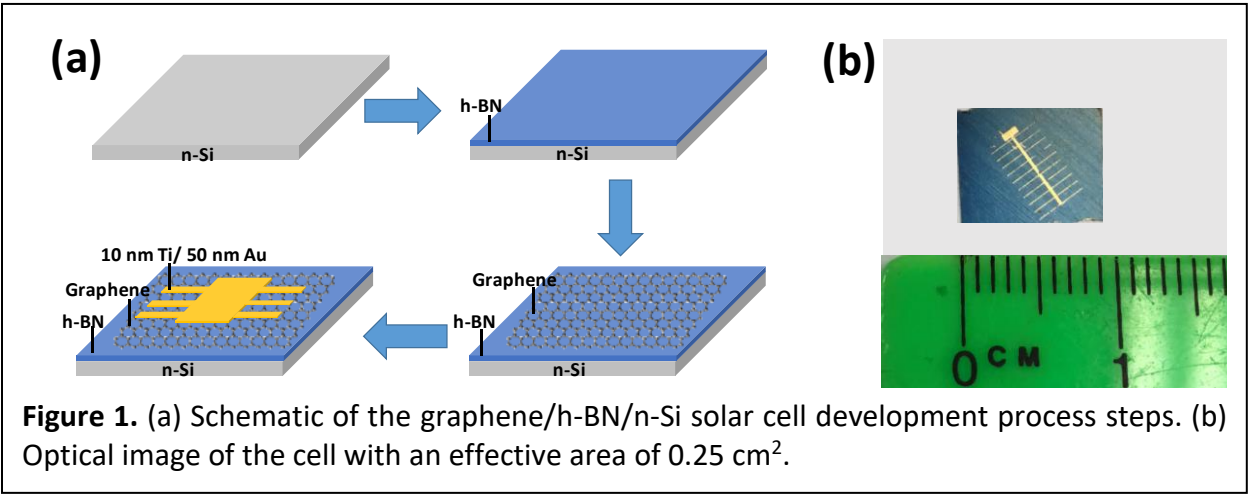
In this work, we directly nucleated h-BN layer on n-type silicon (n-Si) surface *via* surface-chemical-interaction mechanism<sup>25,35</sup>. A chemical-vapor deposited single-layer graphene was transferred on top of h-BN to construct a graphene/h-BN/n-Si MIS-type heterostructure. Temperature dependent photovoltaic characteristics is investigated. The open-circuit voltage (V<sub>oc</sub>) of maximum of 0.52 V is achieved merely by introducing h-BN interlayer, which is of 4-fold increase in comparison to the graphene/n-Si heterojunction cell. An energy-band model is developed to determine the effective Schottky barrier height ( $\phi'_{SBH}$ ) of graphene/h-BN/n-Si heterostructure after incorporating the h-BN layer. The temporal-stability of the photovoltaic parameters is investigated to illustrate the robustness and durability of the cell.

## EXPERIMENTAL

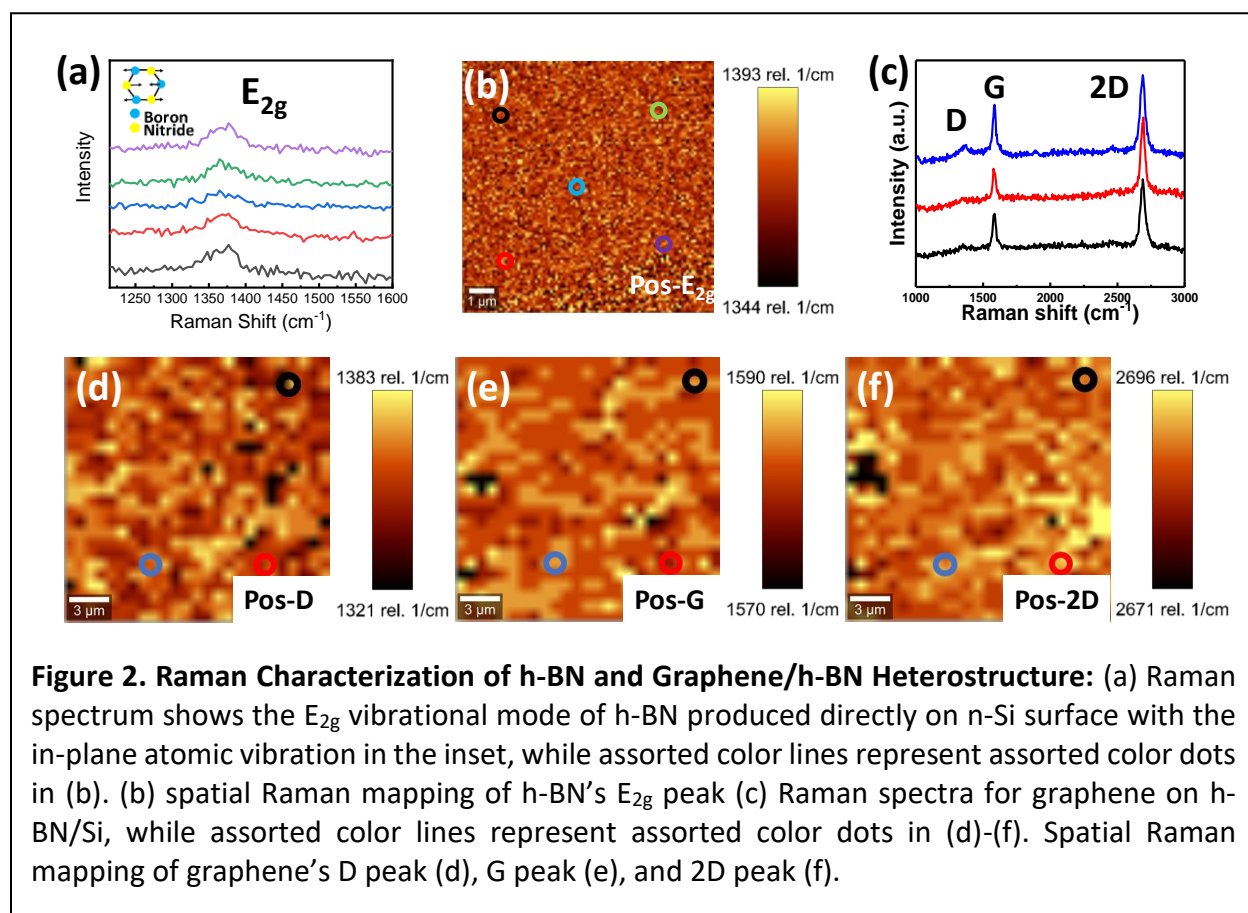
The fabrication process for a graphene/h-BN/Si solar cell is schematically illustrated in Figure 1 (a) Thin films of h-BN was directly synthesized on 1X1 cm<sup>2</sup> solar grade lightly doped n-Si by low-pressure chemical vapor deposition (LPCVD) as reported in the literature.<sup>25</sup> The process steps are presented in the Materials and Methods section. Monolayer graphene was synthesized on copper foil (99.8%, Alfa-Aesar) using thermal LPCVD and transferred onto h-BN/n-Si chip using the standard PMMA transfer method as reported previously.<sup>36</sup> The copper was removed by immersing PMMA coated graphene/copper into acid solution consisted of 1 part of HNO<sub>3</sub> and 1

part of DI water under ambient temperature for 10 minutes. The PMMA was removed by immersing the graphene-coated sample in the 60 °C-heated-acetone bath for 10 minutes. Both the process of synthesis of graphene on metal catalyst substrates and its chemical transfer are presented in Materials and Methods section. For solar cell fabrications, the Titanium (10 nm)/Gold (50 nm) pattern was made by standard e-beam evaporation and lift-off lithography strategy. The gallium-indium eutectic (99.99% Aldrich) was used as the back-contact metal-electrode for the graphene/h-BN/n-Si cell. Raman spectra were acquired by a confocal Raman microscope (Raman-AFM, WITec alpha 300 RA, laser wavelength of 532 nm and beam size of 721 nm). The light and dark current density/voltage (J-V) characteristic profiles were acquired under AM1.5G illumination (100 mW/cm<sup>2</sup>) and dark condition using a Keithley 2612 source-measure-unit. Quantum efficiency measurement (wavelength range: 200-1100 nm) facility was created in-house with the monochromator provided by ORIEL Instrument. The temperature dependent studies were conducted in Janis ST-100 setup.

RESULTS AND DISCUSSION

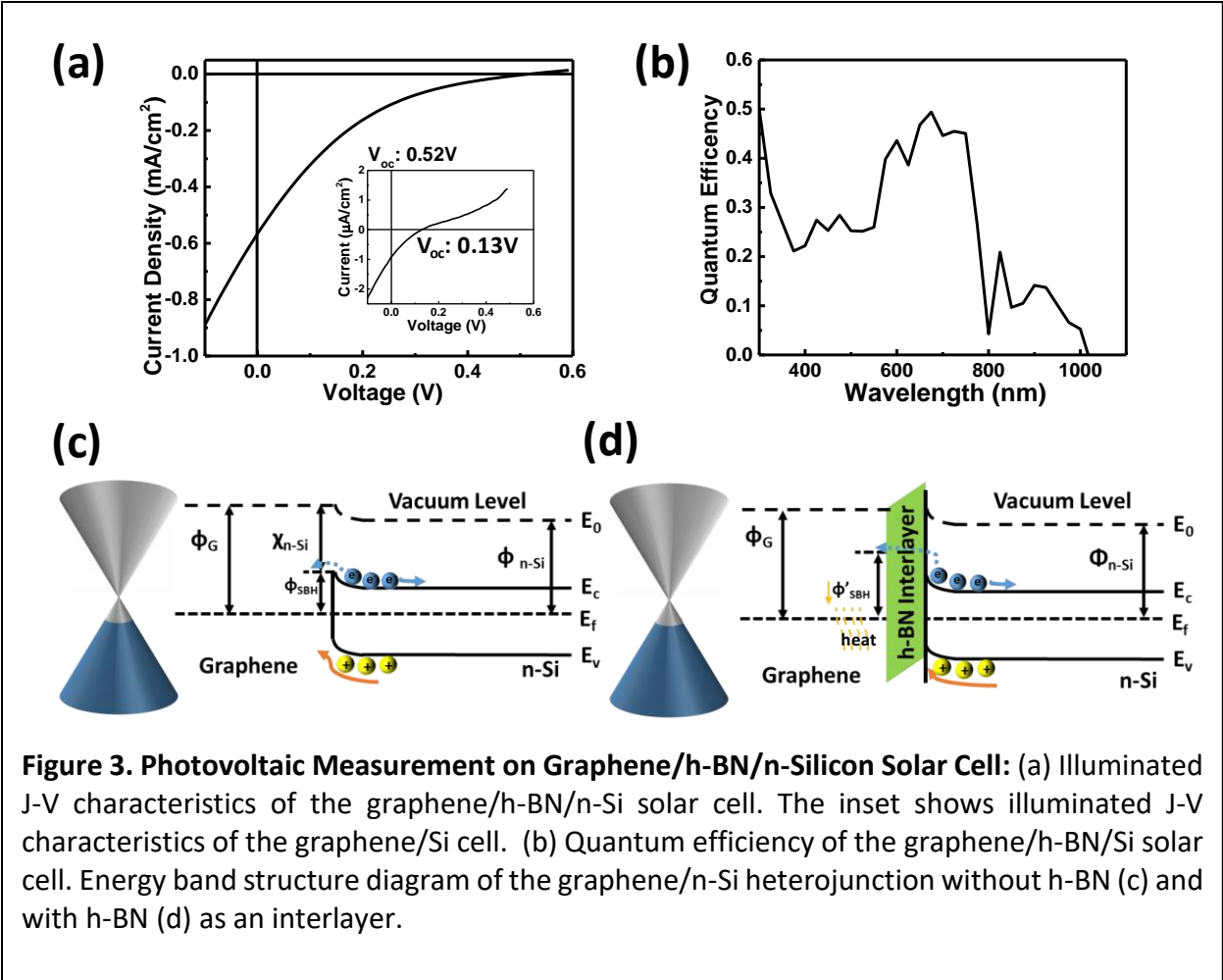


Raman vibrational spectroscopy is typically employed to identify chemical and electronic structure by observing vibrational, rotational, and other low-frequency modes in a material system.<sup>18</sup> Raman spectroscopy has been extensively utilized as an inelastic-scattering based finger-print to characterize the 2D nanomaterials.<sup>37–39</sup> Scanning Raman spectra and mapping results are presented in Figure 2 to confirm the nucleation of h-BN on n-Si and formation of graphene/h-BN heterostructure. The presence of a weak peak at 1368.8 cm<sup>-1</sup> in h-BN/n-Si Raman spectrum (Figure 2 (a)) corresponds to the E<sub>2g</sub> in-plane vibration of boron and nitrogen atoms in the molecular building unit of h-BN as shown in the inset.<sup>40</sup> There are also no competing Raman peaks such as 1304 cm<sup>-1</sup> and 1055 cm<sup>-1</sup> which correspond to the longitudinal optical (LO) phonon and the transversal optical (TO) phonon of c-BN.<sup>39</sup> The color circles (black, green, blue, red and purple) at different area of Figure 2(b) corresponds to the Raman spectrum of the same color in Figure 2(a). The homogeneous color contrast in Figure 2 (b) and presence 1368 cm<sup>-1</sup> Raman peaks in all the Raman spectra in Figure 1(a) clearly shows a continuous and uniform h-BN film formation on solar grade light p-doped n-Si surfaces.<sup>35</sup> Further, for graphene/h-BN/n-Si hybrid



system, two characteristic Raman modes at  $1583\text{ cm}^{-1}$  (G-peak) and  $2687\text{ cm}^{-1}$  (2D peak) are found, which demonstrates the continuous interfacing of graphene on h-BN surface. It is noticed that the Raman shift of this graphene's G-peak is less than  $1584.5\text{ cm}^{-1}$ , which indicates lightly doped graphene.<sup>41</sup> To further confirm the uniformity of graphene/h-BN heterostructure, spatial Raman mapping for graphene's D peak ( $1344\text{ cm}^{-1}$ ), G peak ( $1583\text{ cm}^{-1}$ ), and 2D peak ( $2687\text{ cm}^{-1}$ ) are presented in Figure 2 (d)–(f). The colored circles (blue, red, and black) at different areas of Figure 2 (d-f) correspond to the Raman spectra of the same color in Figure 2 (c). The homogenous color contrast (Figure 2 (d-f)) clearly shows a continuous and uniform transferred graphene on the h-BN/n-Si surface.

Photo-electrical measurements are essential to investigate the performance of photovoltaic cells. Figure 3 (a) shows the photovoltaic profile for graphene/h-BN/n-Si solar cell. As expected, the performance of the device with h-BN as an interlayer increased in comparison to the reference graphene/n-Si solar cell. The photovoltaic characteristics such as: open-circuit voltage ( $V_{oc}$ ), short-circuit current density ( $J_{sc}$ ) and fill-factor (FF) of the graphene/h-BN/n-Si device are  $0.52\text{ V}$ ,  $0.57\text{ mA/cm}^2$  and  $25\%$ . The measured  $V_{oc}$  for graphene/n-Si solar cell is about  $0.13\text{ V}$  as shown in the inset. It is interesting to observe that by inserting a thin-film of h-BN layer between graphene and n-Si enhances the  $V_{oc}$  by 4-fold. Further, it is critical to understand the recombination dynamics in the graphene/h-BN/n-Si photovoltaic cell. The external quantum efficiency (EQE) is a spectro-electronic tool and it is the ratio of the number of charge carriers collected by the solar cell to the number of incident photons of a given wavelength. As presented

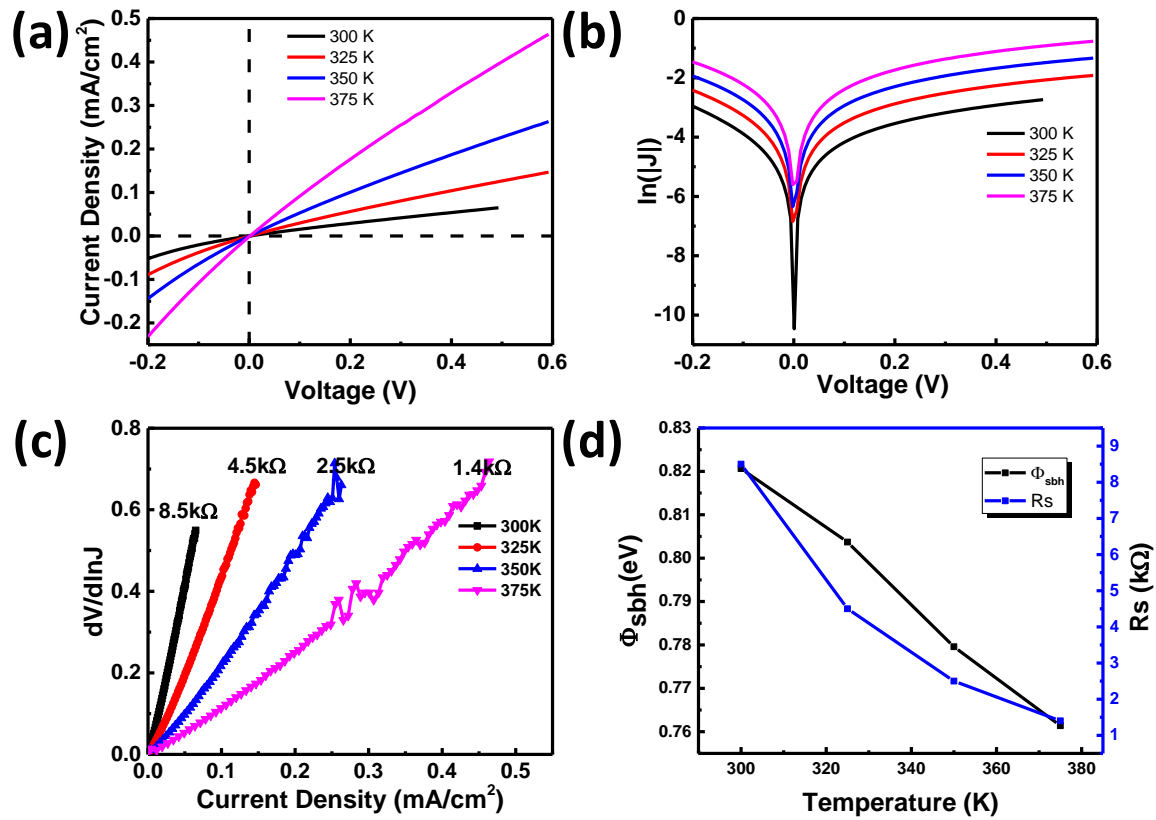


**Figure 3. Photovoltaic Measurement on Graphene/h-BN/n-Silicon Solar Cell:** (a) Illuminated J-V characteristics of the graphene/h-BN/n-Si solar cell. The inset shows illuminated J-V characteristics of the graphene/Si cell. (b) Quantum efficiency of the graphene/h-BN/Si solar cell. Energy band structure diagram of the graphene/n-Si heterojunction without h-BN (c) and with h-BN (d) as an interlayer.

in Figure 3 (b), it is found that the cell exhibits higher EQE (30 to 50%) in the longer wavelength range of 600-800 nm. Lower EQE indicates higher recombination in the short wavelength range: 300 - 600 nm. As illustrated in Figure 3 (c), the photo-excited charge carriers are separated by the built-in electric potential ( $V_{bi}$ ), where holes are pulled to the graphene electrode and the electrons towards n-Si. Here, the sum of  $V_{bi}$  and  $\Delta E$  ( $\Delta E = E_c - E_f$ ) is known as the Schottky Barrier Height ( $\phi_{SBH}$ ), which blocks the transport of electrons from n-Si back to graphene and holes from graphene back to n-Si. However, the  $\phi_{SBH}$  is relatively low for graphene/Si heterojunction, so that electrons may drift back to graphene, and that may cause a current leakage. To reduce the current leakage, a h-BN layer is introduced between graphene and n-Si to increase the  $\phi_{SBH}$  and thus create an extra barrier for electrons drifting back to graphene/h-BN/n-Si interface. Effective Schottky Barrier Height ( $\phi'_{SBH}$ ) in Figure 3(d) is larger than  $\phi_{SBH}$  due to the insert of h-BN layer.

The current density-voltage (J-V) profile under dark condition is critical for understanding the electrical performance of a diode. Figure 4 (a) shows the dark J-V characteristics of the graphene/h-BN/n-Si MIS-type heterojunction cell with variation in temperature (300 K – 375 K). An increased forward current under forward bias, as well as an increased reverse current, were observed with the increased temperature for the applied voltage ranging from -0.2 V to 0.6 V. Similar phenomenon has been reported for Si-based MIS structure solar cell and graphene-based





**Figure 4. Temperature Dependency on Graphene/h-BN/n-Si Solar Cell without Light Illumination:** (a) J-V characteristics, (b) Log(|J|)-V plot under different temperatures. (c) Series resistance derived from dark J-V characteristics. (d) The effect of temperature on Schottky Barrier Height and series resistance.

MS structure heterojunction solar cell.<sup>42,43</sup> As shown in Equation 6, the  $V_{oc}$  decreases as dark saturation current increases. Intrinsic carrier concentration increases as the band gap of Si decreases with the increase in temperature. The result is in accordance with realistic dark J-V characteristics.<sup>44</sup> Figure 4 (b) shows the  $\log(|J|)$ -V plot in the voltage range of -0.2 V to 0.6 V with the increase of temperature. The series resistance was derived from the dark J-V curve as is shown in Figure 4 (c).<sup>45</sup> Detailed  $R_s$  values under different temperature are displayed in Figure 4 (d), the series resistance of graphene/h-BN/n-Si solar cell decreases as temperature increases. For every 25 K increase in temperature, the series resistance decreases by approximately 45%. The  $\phi_{SBH}$  can be calculated through combining Equation (1) and (6). The  $\phi_{SBH}$  versus temperature is plotted in Figure 4 (d), the effect of heat on  $\phi_{SBH}$  is plotted in Figure 3(d). The  $\phi_{SBH}$  decreases with temperature and with the decrease of  $\phi_{SBH}$ , electrons are easier to diffuse back to interface and recombine with holes and thus efficiency of solar cell decreased due to higher recombination.

The photovoltaic characteristics of graphene/h-BN/n-Si solar cell as a function of temperature is critical to understand the role of thermal energy in exciton carrier transport. A photovoltaic effect

was obtained in Figure 5 (a) for all the temperatures studied. As shown by the gray arrow in Figure 5 (a), the  $V_{oc}$  decreases as the temperature increases. The changes of  $V_{oc}$  with temperature is plotted in Figure 5 (b) and detailed calculation is presented in the Supporting Information section 2. A linear fit was performed on data points of  $V_{oc}$  versus temperature. The  $R^2$  is 0.999, which means the  $V_{oc}$  is linear to temperature and  $V_{oc}$  decreases by 2.0 mV for every 1 K increase in temperature. This linear property of  $V_{oc}$  versus temperature is exactly the same as the Si solar cell (2.0 mV)<sup>42</sup>. At the same time,  $J_{sc}$  stays almost the same (variance is  $4.1 \times 10^{-5}$ ) when temperature increases from 300 to 375 K (Figure 5 (b)).

In order to further analyze the effect of temperature on MIS structure graphene/h-BN/n-Si solar cell and the effect of h-BN layer, we analyzed their J-V relations as presented below:<sup>46</sup>

$$J = J_s \exp\left(\frac{V_a}{nkT}\right) - 1 \quad (1)$$

$$J_s = A^* T^2 \exp\left(-\frac{q\phi_{SBH}}{kT}\right) \quad (2)$$

Here  $J$ ,  $J_s$  are the current density and reverse saturation current density respectively<sup>17</sup>,  $A^*$  is the effective Richardson constant,  $T$  is temperature,  $k$  is the Boltzmann constant, and  $\phi_{SBH}$  is the Schottky barrier height.

The short red shift of graphene G peak indicates the lightly doped graphene, and photovoltaic effect of graphene/h-BN/n-silicon heterostructure further indicates the slightly p-doped graphene. Although graphene under ambient environment is slightly p-doped, it still behaves as a metallic electrode.<sup>47</sup> Adding a h-BN tunneling layer will add a tunneling probability factor  $\exp(-\sqrt{\Delta E_c} \delta)$ .<sup>48</sup> Now, the  $J_s$  can be expressed as:

$$J_s = A^* T^2 \exp\left(-\frac{q\phi_{SBH}}{kT}\right) \cdot \exp(-\sqrt{\Delta E_c} \delta) \quad (3)$$

$$J_s = A^* T^2 \exp\left(-\frac{q\left(\phi_{SBH} + \frac{kT}{q} \sqrt{\Delta E_c} \delta\right)}{kT}\right) \quad (4)$$

The new effective Schottky barrier height ( $\phi_{SBH}'$ ) by comparing Equations (1) & (3) is:

$$\phi_{SBH}' = \phi_{SBH} + \frac{kT}{q} \sqrt{\Delta E_c} \delta \quad (5)$$

As per the Equation 5, the effective Schottky barrier height is increased by  $1.01 \frac{kT}{q} \sqrt{\Delta E_c} \delta$  due to the insertion of h-BN layer. The conduction band offset ( $\Delta E_c$ ) represents the effective tunneling barrier height that h-BN layer creates for the electrons in the conduction band of n-Si. For graphene/h-BN/n-Si, the  $\Delta E_c$  is the work function of n-Si, which is approximately 4.1 eV and  $\delta$  is the thickness of the h-BN film. As is shown by the Equation 5,  $\Delta E_c$  and thickness affects the  $J_s$  exponentially.

The  $V_{oc}$  of solar cell can be expressed as<sup>49</sup>:



$$V_{oc} = \frac{nkT}{q} \ln \left[ 1 + \frac{J_L}{J_s} \right] \quad (6)$$

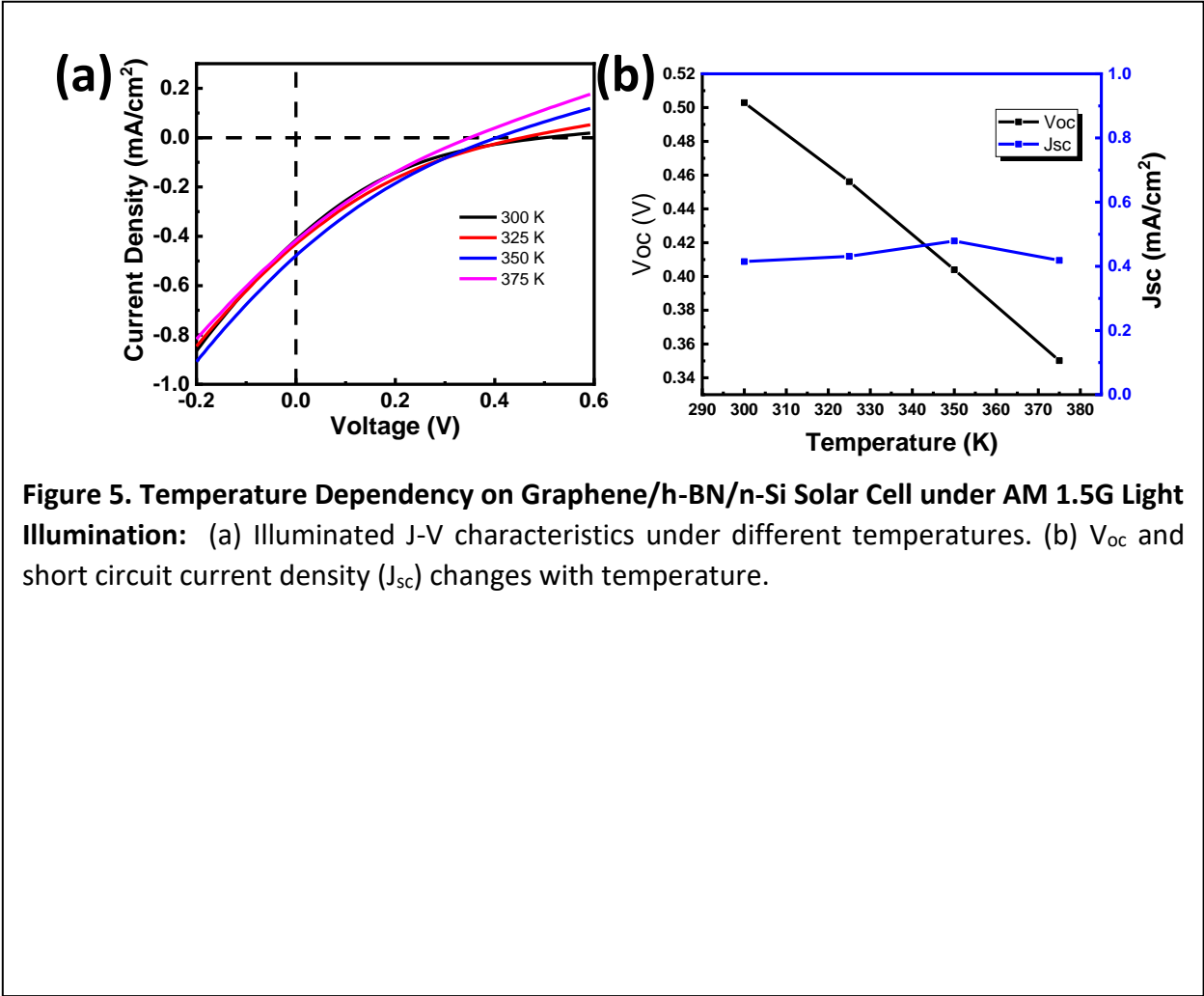
Here  $J_L$  is the light generated current density,  $J_s$  is the dark saturation current density, which is composed of  $J_{ms}$  (thermionic emission dark current),  $J_{rg}$  (depletion layer recombination-generation current density),  $J_d$  (injection-diffusion current density), and  $J_{ss}$  (the surface state current density due to the charge exchange between the metal and semiconductor band edges *via* surface states). We mainly focus on the effect of  $J_{ms}$  on the  $J_s$  since large  $J_{ms}$  is the main cause for the high dark current for the metal-semiconductor heterojunctions. According to Equation 4, the  $J_{ms}$  decreases with the increase of effective Schottky barrier height. Since  $J_L$  remains same for the same graphene/n-Si solar cell, the  $V_{oc}$  increases as  $J_s$  decreases with the insertion of h-BN layer in the graphene/n-Si solar cell.

The relation of saturation current density and diffusion length can be expressed as follows:<sup>50</sup>

$$J_s \approx \frac{qDn_i^2}{L_{diff}N} \quad (7)$$

D is diffusivity of the minority carrier,  $L_{diff}$  is the minority carrier diffusion length, N is the doping, and  $n_i$  is the intrinsic carrier concentration. According to Equation 7, the dark saturation current density is positive dependent on  $n_i$ , and  $n_i$  increases with temperature. By combining Equation 7 and Equation for intrinsic carrier concentration, we can get Equation 8 (detail derivation is presented in the Supporting Information in which  $V_{oc}$  is linearly decreased with the temperature. The theoretic analysis is in accordance with the experiment results.

$$\frac{dV_{oc}}{dT} = -\frac{\phi'_{SBH} - V_{oc} + 3\frac{kT}{q}}{T} \quad (8)$$



**Figure 5. Temperature Dependency on Graphene/h-BN/n-Si Solar Cell under AM 1.5G Light Illumination:** (a) Illuminated J-V characteristics under different temperatures. (b)  $V_{oc}$  and short circuit current density ( $J_{sc}$ ) changes with temperature.

Conjointly, the interfacing of h-BN layer on n-Si passivates the surface dangling bonds. To reduce the saturation current density, longer diffusion length for the minority carriers is critically required. Diffusion length is affected by the recombination-active defects, intrinsic mobility limitation, and absence of percolation pathways. Due to the passivation of h-BN layer, dangling bonds on the Si surface will not cause the mid-gap surface states and thus recombination-active defects are reduced. Further, the MIS-type architecture is a minority carrier-based solar cell. The recombination is controlled by the number of minority carriers at the junction edge, how fast

they move away from the junction and how quickly they recombine. High recombination rate increases the forward bias diffusion current, and hence increases the reverse dark-saturation current. Less reverse dark-saturation current represents a smaller number of electrons diffusing backwards and thus the chance of electrons recombining with holes gets minimized. So, smaller reverse dark saturation current means less recombination.

## CONCLUSION

In summary, we have demonstrated the development of directly introduced h-BN film as a passivation and tunneling interlayer for graphene-on-Si heterojunction photovoltaic cell *via* chemical-surface-adsorption technique and investigated the temperature dependent photovoltaic characteristics of graphene/h-BN/n-Si MIS solar cell. This MIS-type heterojunction shows non-linear rectifying characteristics with increased forward current and reverse current with the increase of temperature under dark condition. Under AM 1.5G light illumination, a 4-fold increase in  $V_{oc}$  is found for the graphene/h-BN/n-Si cell compared to graphene/n-Si cell with a performance stability of 7-month. This may be due to the h-BN interlayer-induced increase in Schottky barrier height and decrease in surface mid-gap trap states. The  $V_{oc}$  was found to decrease linearly by 2.0 mV/K and  $J_{sc}$  was constant with increase in temperature. Our findings revealed the temperature dependent photovoltaic responses of the graphene/h-BN/n-Si MIS heterojunction. Further improvement can be achieved by tuning the Fermi level of graphene by doping, functionalization with nano particles and electrostatic gating. The nanoarchitecture and phenomena developed here may further provide a new avenue of designing stable, high-performance, and cost-effective tandem 2D-photovoltaic-cells.

## MATERIALS AND METHODS

**Direct Synthesis of h-BN on n-Si Surface.** The n-Si substrates were cleaned by DI water, acetone and isopropyl alcohol (IPA). Then, they were loaded into 1 inch quartz tube inside the furnace and heated up to 1100 °C in 25 minutes with  $H_2$  flow at 30 sccm. When the furnace temperature reached 900 °C, precursor chamber which held crucible with ammonia borane complex inside would be heated by thermal tape, whose setting temperature is 100 °C. After substrate temperature reached 1098 °C, valve connecting precursor chamber to furnace was open. Right after that, tube pressure was adjusted to 5 Torr using right angle valve near the mechanical pump. After an hour reaction, thermal tape was turned off, the valve connecting precursor chamber to furnace was closed and furnace cover was opened for fast cooling.  $H_2$  gas was turned off when furnace temperature dropped to 60 °C and vent valve was opened to unload the sample.

**Synthesis of Graphene.** Pretreated copper foil was loaded in 1-inch quartz tube. The reaction chamber was evacuated to 6 mTorr and flushed with 10 sccm of  $H_2$ . Then the system was heated up to 1050 °C in 25 minutes and hold for 40 minutes with the same  $H_2$  condition. Afterwards, the

flow rate of H<sub>2</sub> was increased to 22 sccm and 10 sccm CH<sub>4</sub> was introduced for 20 seconds. Finally, the furnace was cooled down to room temperature in 35 minutes with 22 sccm H<sub>2</sub>.

**Chemical Transfer of Graphene.** Copper foil with graphene was put on the spin coater (Headway Research, Inc. Model: PWM32). After its surface was covered by PMMA solution, copper foil rotated with spinner to 4000 rpm for 60 seconds. Then, copper foil was transferred to Nitric Acid Solution (30% by weight) to have copper etched away. After 10 minutes, graphene was picked by SiO<sub>2</sub>/Si chip (3 cm\*3 cm) and transferred to DI water. After cleaning in DI water for 5 minutes, graphene was transferred to second DI water and picked up by h-BN/n-Si chips from previous step. The chip was left in the hood for overnight drying. Then chip was put on hotplate which was 160 °C for 25 minutes. After that, chip was transferred to 60 °C acetone solution for 10 minutes. Finally, chip was cleaned by acetone and IPA to remove the PMMA on the top.

**ASSOCIATED CONTENT**

**Supporting Information**

Section 1: 1. Calculation of Raman spot size and pixel size. Section 2: Effect of temperature on the open-circuit voltage of graphene/h-BN/n-Si solar cell. Section 3: Derivation of V<sub>oc</sub> vs. Temperature Relationship

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**Notes**

The authors declare no competing financial interests.

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