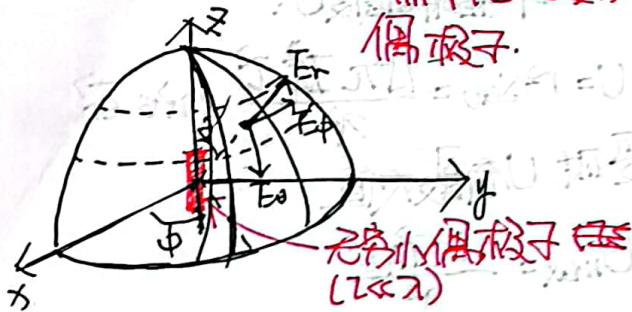


Chapter II ~~电磁波~~ 偶极子天线

三、偶极子天线 ~~← 重点讨论无穷小偶极子~~



~~偶极子~~

1. 无穷小偶极子 (infinitesimal dipole)

假设电流为: $i(t) = I_0 \cos \omega t$

沿 z 轴的电流为: $I_z = \vec{e}_z I$

(1) 矢位 (vector potential)

$$\begin{aligned} \vec{A} &= \frac{\mu_0}{4\pi} \int_V \vec{J} \frac{e^{-jkr}}{r} dV \\ &= \frac{\mu_0 I l}{4\pi r} e^{-jkr} \vec{e}_z \\ &= A_z \vec{e}_z \end{aligned}$$

$$\vec{A} = \vec{e}_r A_r + \vec{e}_\theta A_\theta + \vec{e}_\phi A_\phi$$

转换为球坐标

$$\begin{cases} A_r = A_z \cos \theta \\ A_\theta = -A_z \sin \theta \\ A_\phi = 0 \end{cases}$$

由 A 垂直 r 向外, 故 A 有分量

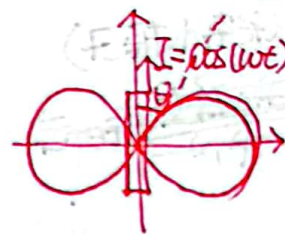
(2) 磁场 (magnetic field)

$$\begin{aligned} \vec{H} &= \frac{1}{\mu_0} \nabla \times \vec{A} \\ &= \frac{1}{\mu_0 r^2 \sin \theta} \begin{vmatrix} \vec{e}_r & r\vec{e}_\theta & r\sin \theta \vec{e}_\phi \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ A_r & rA_\theta & 0 \end{vmatrix} \end{aligned}$$

$$\begin{aligned} &= \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (0) - \frac{\partial}{\partial \phi} (rA_\theta) \right] \vec{e}_r \\ &+ \frac{1}{r} \left[\frac{\partial}{\partial \phi} A_r - \frac{\partial}{\partial r} (0) \right] \vec{e}_\theta + \frac{1}{r} \left[\frac{\partial}{\partial r} (rA_\theta) - \frac{\partial}{\partial \theta} A_r \right] \vec{e}_\phi \\ &= \frac{I l}{4\pi r} \sin \theta \left(j\frac{k}{r} + \frac{1}{r^2} \right) e^{-jkr} \vec{e}_\phi \end{aligned}$$

~~与 r 无关~~

$\theta: 0 \rightarrow \pi$ 过程中, H 先变大后变小



$$\begin{cases} H_r = H_\theta = 0 \\ H_\phi = \frac{I l}{4\pi r} \sin \theta \left(j\frac{k}{r} + \frac{1}{r^2} \right) e^{-jkr} \end{cases}$$

(3) 电场 (electric field)

$$\vec{E} = \frac{1}{j\omega \epsilon} \nabla \times \vec{H}$$

$$\vec{E} = E_r \vec{e}_r + E_\theta \vec{e}_\theta$$

E 和 H 正交, 在 phi 上有 H 就没有 E

$$\begin{cases} E_r = \frac{I l}{4\pi} \frac{2}{\omega \epsilon_0} \cos \theta \left(\frac{k}{r^3} - j\frac{1}{r^2} \right) e^{-jkr} \\ E_\theta = \frac{I l}{4\pi} \frac{1}{\omega \epsilon_0} \sin \theta \left(j\frac{k^2}{r} + \frac{k}{r^2} - j\frac{1}{r^3} \right) e^{-jkr} \end{cases}$$

2. ~~无穷小偶极子产生近场, 远场的~~ 电场

(1) 近场: $\frac{1}{kr} \ll \frac{1}{kr^2} \ll \frac{1}{kr^3}$

$$\begin{cases} H_\phi = \frac{I l}{4\pi r^2} \sin \theta \\ E_r = -j \frac{I l}{4\pi r^3} \frac{2}{\omega \epsilon_0} \cos \theta \\ E_\theta = -j \frac{I l}{4\pi r^3} \frac{1}{\omega \epsilon_0} \sin \theta \\ E_\phi = H_r = H_\theta = 0 \end{cases}$$



12) 远场

$$\frac{1}{kr} \gg \frac{1}{4\pi r^2} \gg \frac{1}{4\pi r^3} \quad k = \frac{2\pi}{\lambda}$$

$$\begin{cases} H_\varphi = j \frac{I L}{2\pi r} \sin\theta e^{-jkr} = \frac{k I L}{4\pi r} \sin\theta e^{-jkr} \\ E_\theta = j \frac{60\pi I L}{\lambda r} \sin\theta e^{-jkr} = j \frac{k I L}{4\pi r} e^{-jkr} \sin\theta \\ E_r = 0 \\ H_r = H_\theta = E_\varphi = 0 \end{cases}$$

$E_\theta = j H_\varphi$

3) 远场下的辐射特性 (无小振子)

$$\vec{S}_{av} = \frac{1}{2} \text{Re}[\vec{E} \times \vec{H}^*] = \frac{15\pi I_A^2 \lambda^2}{\lambda^2 r^2} \sin^2\theta \vec{e}_r$$

(1) 辐射序变量

$$\vec{S}_{av} = \frac{1}{2} \text{Re}[\vec{E} \times \vec{H}^*] = \frac{15\pi I_A^2 \lambda^2}{\lambda^2 r^2} \sin^2\theta \vec{e}_r$$

12) 辐射功率

$$\begin{aligned} P_r &= \oint_S \vec{S}_{av} \cdot d\vec{S} = \oint_S \frac{1}{2} \text{Re}[\vec{E} \times \vec{H}^*] \cdot d\vec{S} \\ &= \int_0^{2\pi} d\varphi \int_0^\pi \frac{15\pi I^2 \lambda^2}{\lambda^2 r^2} \sin^2\theta d\theta \\ &= 40\pi^2 I^2 \left(\frac{\lambda}{2}\right)^2 \end{aligned}$$

$$P_{rad} = \oint_S \vec{u}_{rad} \cdot \vec{r}^2 \sin\theta d\theta d\varphi$$

13) 辐射阻抗

由于辐射的能量不再返回波源，
认为天线的辐射功率完全被等效电阻
吸收，称为辐射阻抗 (radiation impedance)

$$P_r = \frac{1}{2} I^2 R_r$$

$$\text{故 } R_r = \frac{2P_r}{I^2} = 80\pi^2 \left(\frac{\lambda}{2}\right)^2$$

(4) 方向性

由 S_{av} 到辐射强度 U :

$$U = r^2 S_{av} = \frac{15\pi I_A^2 \lambda^2}{\lambda^2} \sin^2\theta \vec{e}_r$$

$\theta = \frac{\pi}{2}$ 时 U 有最大值

$$U_{max} = \frac{15\pi I_A^2 \lambda^2}{\lambda^2}$$

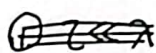
$$P_r = 40\pi^2 I^2 \left(\frac{\lambda}{2}\right)^2$$

$$D_0 = 4\pi \frac{U_{max}}{P_r} = \frac{3}{2}$$

4. 有限长振子

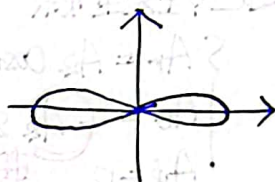
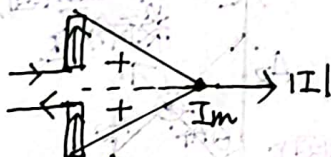
(1) 指定电流 I 和磁电流密度 M
对于很细的振子 (直径为 a), 电流分布:

$$I_e = \begin{cases} e z I_0 \sin[k(\frac{L}{2} - z')] & 0 \leq z' \leq \frac{L}{2} \\ e z I_0 \sin[k(\frac{L}{2} + z')] & -\frac{L}{2} \leq z' \leq 0 \end{cases}$$

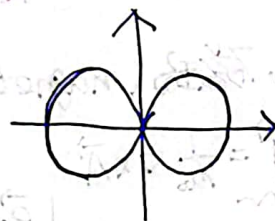
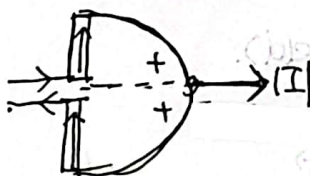


• 电流在天线上的分布和辐射方向图

① $L < \lambda$

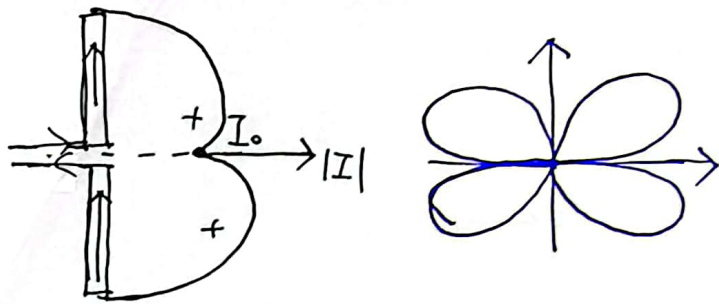


② $L = \lambda$

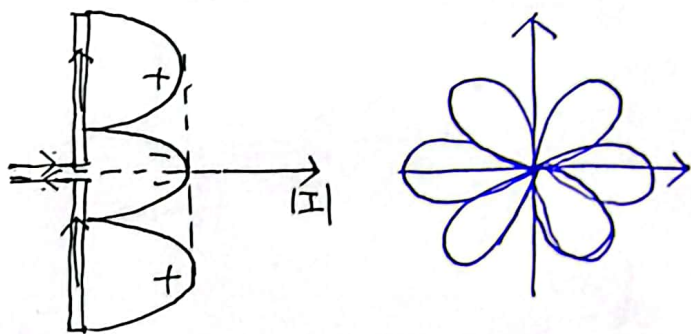


Chapter II 偶极子天线

③ $\frac{\lambda}{2} < L < \lambda$.



④ $\lambda < L < \frac{3}{2}\lambda$.



(2) 确定辐射阻抗 R_r 和输入阻抗 R_{in} .

$$R_{in} = \left[\frac{I_0}{I_{in}} \right]^2 R_r = \frac{R_r}{\sin^2\left(\frac{kL}{2}\right)}$$

~~当~~ $\sin^2\left(\frac{kL}{2}\right) = 1$. 即 $\frac{kL}{2} = \frac{2\pi}{\lambda} \cdot \frac{L}{2} = \frac{\pi}{2}$. $L = \frac{\lambda}{2}$
得 $L = \frac{\lambda}{2}$

此时 $R_{in} = R_r$.

$$\begin{array}{c} \text{---} \square \text{---} \square \text{---} \\ R_{in} \quad R_r \end{array} \quad \Gamma = \frac{Z_L - Z_0}{Z_L + Z_0}$$

故 $L = \frac{\lambda}{2}$ 时, $R_{in} = R_r$, 天线效率最高

但 $L = 1.5\lambda$, $L = 2.5\lambda$ 时 依然有 $R_{in} = R_r$

E (3) ~~确定辐射~~

(3) 最大有效面积.

① 极小偶极子天线的有效面积

$$A_{rm} = 0.119\lambda^2$$

② 各向同性天线有效面积

$$A_{tm} = \frac{A_{rm}}{D_{ro}} = \frac{0.119\lambda^2}{1.5} = \frac{\lambda^2}{4\pi}$$

③ 通常, 最大有效面积,

$$A_{em} = \frac{\lambda^2}{4\pi} D_o$$

• 比较: 两根天线的 $L = \frac{\lambda}{2}$, 但 λ 不同

$\lambda \uparrow$
f ↓
Arm ↑
Pr ↑

$\lambda \downarrow$
f ↑
Arm ↓
Pr ↓

波长的长度
越高, 接收到的功率越小

