# Lab 1 Report

March 24, 2025

## 1 Lab 1: Implementing FEM for 1D Diffusion Equation

We implement the finite elements method (FEM) with piecewise linear  $(\mathcal{P}_1)$  basis functions to solve the one-dimensional diffusion equation:

(D): 
$$\begin{cases} -(D(x)u'(x))' = f(x) & \text{for } x \in (a,b) \\ u(0) = u(1) = 0 \end{cases}$$
.

In the equation (D),  $f:[a,b] \to \mathbf{R}$  is the forcing or source function,  $D:[a,b] \to \mathbf{R}$  is the diffusion (typically constant, in which case the equation reads -Du''=f), and we search for a suitable  $u:[a,b] \to \mathbf{R}$  in the trial space

$$\mathcal{V}([a,b]) = \{v : [a,b] \to \mathbf{R} : v \in C^0([a,b]), v' \text{ is piecewise continuous and bounded on } [a,b]\}$$

The constraints u(0) = u(1) = 0 are so-called *Dirichlet* boundary conditions.

#### 1.0.1 Importing the code

Please refer to https://github.com/labounad/comp-diffeq/blob/main/oneD/oneD\_galerkin.py for the full source code.

```
[1]: import os
  import sys
  module_path = os.path.abspath(os.path.join('...'))
  if module_path not in sys.path:
       sys.path.append(module_path)

from oneD.oneD_galerkin import *
```

### 1.1 Theory

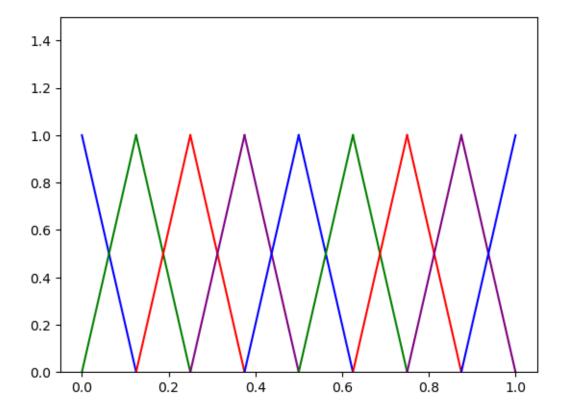
The idea is to discretize the interval [a, b] into pieces (elements) of width h, and approximate u by a sequence  $u_h$  converging to u, where each  $u_h$  lives in the finite-dimensional space  $\mathcal{V}_h = \text{span}\{\phi_0, \ldots, \phi_n\}$ . The basis functions  $\phi_i$  are chosen to be a piecewise linear partition of unity, such that, with  $x_j$ , j = [0:n] denoting the j-th node,  $\phi_i(x_j) = \delta_{ij}$ . Explicitly:

$$\phi_i(x) = \begin{cases} \frac{x - x_{i-1}}{x_i - x_{i-1}} & x \in [x_{i-1}, x_i] \\ \frac{x_{i+1} - x}{x_{i+1} - x_i} & x \in [x_i, x_{i+1}] \\ 0 & \text{else} \end{cases}$$

An example with n = 8 is depicted below:

```
[2]: # plot basis functions:
    n_elems = 8
    nodes = n_elems + 1
    x_start = TestParams.X_START
    x_end = TestParams.X_END

x_coords = np.linspace(x_start, x_end, nodes)
    colors = cycle(['blue', 'green', 'red', 'purple'])
    for n in range(nodes):
        col = next(colors)
        if n > 0:
            plt.plot([x_coords[n-1], x_coords[n]], [0,1], color=col)
        if n < n_elems:
            plt.plot([x_coords[n], x_coords[n+1]], [1,0], color=col)
    plt.ylim((0,1.5))
    plt.show()</pre>
```



Then, one solves the weak formulation in  $\mathcal{V}_h$ : Find  $u_h \in \mathcal{V}_h$  such that

$$(W_h): (u'_h, v') = (f, v) \text{ for all } v \in \mathcal{V}_h.$$

Since  $u_h \in \mathcal{V}_h$  we have  $u_h = \sum_j u_j \phi_j$  and since  $v \in \mathcal{V}_h$  must be arbitrary, we must at least have

 $(W_h)$  hold for the  $\phi_j$ , hence for all j:

$$(u'_h, \phi'_j) = \left(\sum_i u_i \phi'_i, \phi'_j\right) = \sum_i u_i(\phi'_i, \phi'_j) = (f, \phi_j).$$

By defining  $\mathbf{A} :=: \{a_{ij}\} := \{(\phi'_i, \phi'_j)\}$  and  $\mathbf{b} :=: \{b_j\} := \{(f, \phi_j)\}$ , we must solve  $\mathbf{A}\mathbf{u} = \mathbf{b}$ . Such a solution  $\mathbf{u} = \{u_i\}$  will then by linearity solve  $(W_h)$  for all  $v \in \mathcal{V}_h$ .

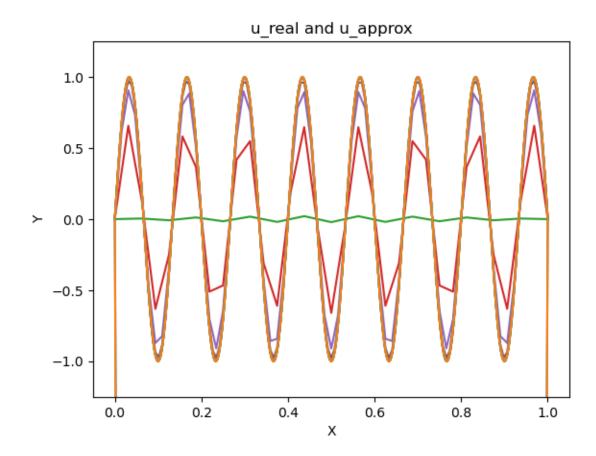
The matrix  $\mathbf{A}$  is called the *stiffness* matrix and the vector  $\mathbf{b}$  the *load* vector.

### 1.2 Implementation

To demonstrate the implementation, we take [a,b] = [0,1], constant diffusion D = 1, and pick a test function  $u_{\text{real}}(x) = \sin(k\pi x)$  and source  $f(x) = -u''_{\text{real}}(x) = (k\pi)^2 \sin(k\pi x)$ , with k = 15, so that  $u_{\text{real}}$  is a solution of (D).

We then approximate  $u_{\text{real}}$  by functions  $u_h$ , where  $h = n^{-1}$  for  $n = 2^2, 2^3, \dots, 2^{12}$ .

```
[7]: TestParams.K_CONST = 15
     def constant_diffusion(x):
         return np.ones(len(x))
     TestParams.diffusion_function = constant_diffusion
     n_{elems} = [2**i for i in range(2, 13)]
     x_start = 0
     x_end = 1
     residuals = []
     for n_elem in n_elems:
         x_coords, u_approx, u_real = galerkin(n_elem, x_start, x_end)
         plt.plot(x_coords, u_approx)
         residuals.append(calc_l2err(x_coords, u_approx, u_real))
     x_{fine} = np.linspace(x_{start}, x_{end}, 10 ** 4)
     plt.plot(x_fine, TestParams.u_real(x_fine))
     plt.ylim((-1.25, 1.25))
     plt.xlabel("X")
     plt.ylabel("Y")
     plt.title("u_real and u_approx")
     plt.show()
```



Here we have used k = 13, leading to instability for very small values of n, however, convergence is achieved.

In the following test, we let k = 4 and D(x) be the step function

$$D(x) = \begin{cases} 1 & x \in [0, 0.5) \\ 2 & x \in [0.5, 1] \end{cases}.$$

```
[33]: def step_diffusion(x):
    half_nodes = len(x) // 2
    other_half = len(x) - half_nodes
    return np.concatenate((np.ones(half_nodes), 2 * np.ones(other_half)))

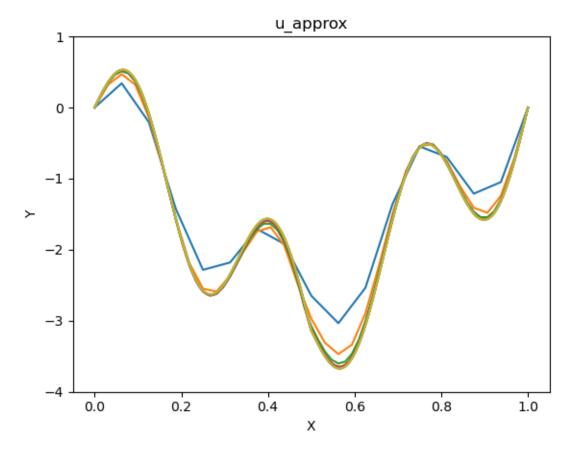
TestParams.K_CONST = 6
TestParams.diffusion_function = step_diffusion

n_elems = [2**i for i in range(4, 13)]

for n_elem in n_elems:
    x_coords, u_approx, u_real = galerkin(n_elem, x_start, x_end)
    plt.plot(x_coords, u_approx)
```

```
x_fine = np.linspace(x_start, x_end, 10 ** 4)
plt.ylim((-4, 1))

plt.xlabel("X")
plt.ylabel("Y")
plt.title("u_approx")
plt.show()
```



## 1.3 Convergence Analysis

In the first implementation above (k = 15 and D = 1), we calculated the residuals at each iteration. Theory (Nietsche-Aubin) states that the residual error is  $O(h^2)$ :

$$||u_{\text{real}} - u_h||_{L^2} \le h^2 ||u''||.$$

Thus if h is halved every time, the ratio of residuals  $e_i/e_{i+1}$  should tend to 4. Indeed, as the following plots demonstrate, this is precisely what is obtained:

```
[59]: residuals = np.array(residuals)
plt.plot(range(2,13), np.log(residuals))
plt.title('$\log \log$ plot of residuals vs number of elements $n$')
plt.xlabel('$\log_2 n$')
plt.show()

ratios = np.divide(residuals[:-1], residuals[1:])
plt.plot(ratios)
plt.title("Residual ratios")
plt.xlim((2,9))
plt.ylim((1.5,4.5))
plt.show()
```

 $\log \log plot$  of residuals vs number of elements n

