# 06-convection-diffusion

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# 1 Finite Differences for the Convection-Diffusion Equation

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The Finite Difference Method (FDM) is a numerical method to solve Partial Differential Equations (PDEs) approximately. We use second-order central finite differences for the diffusion term and investigate stability issue related to the discretization of the first-order advection term.

#### 1.1 Convection-Diffusion Problem

We want to solve the Convection-Diffusion equation

$$\begin{split} -\epsilon \Delta u(x,y) + v \cdot \nabla u(x,y) &= f(x,y), & (x,y) \in [0,L_x] \times [0,L_y] \\ u(x,0) &= b(x), & x \in [0,L] \\ u(x,L) &= t(x), & x \in [0,L] \\ u(L,y) &= r(y), & y \in [0,L] \\ u(0,y) &= l(y), & y \in [0,L] \end{split}$$

on a rectangle with  $\epsilon>0$  and  $v\in\mathbb{R}^2$ . The domain is discreitized structurally with  $N_x$  and  $N_y$  grid-points per dimension leading to  $h_x=\frac{1}{N_x-1}$  and  $h_y=\frac{1}{N_y-1}$ .

[1]: using LinearAlgebra
using SparseArrays
using PyPlot
using Plots

## 1.2 Domain and Boundary Conditions

The domain  $\bar{\Omega}$   $[0, L_x] \times [0, L_y]$  is split into the nodes  $\{(x_i, y_j)\}_{i,j=0...N} \in \bar{\Omega}_h$  with

$$x_i = \frac{L_x}{N_x - 1}, \quad y_j = \frac{L_y}{N_y - 1}$$

The interior nodes are  $\{(x_i, y_j)\}_{i,j=1...N-1} \in \Omega_h$ . We further need boundary conditions for all four sides of the unit square.

[2]: struct Rectangle
Lx::Float64

```
Ly::Float64
    Nx::Int64
    Ny:: Int64
    hx::Float64
    hy::Float64
    xh::Array{Float64,1}
    yh::Array{Float64,1}
    function Rectangle(Lx, Ly, Nx, Ny)
        hx = Lx/(Nx-1)
        hy = Ly/(Ny-1)
        xh = range(0, Lx, step=hx)
        yh = range(0, Ly, step=hy)
        N = new(Lx, Ly, Nx, Ny, hx, hy, xh, yh)
    end
end
struct UnitSquareUniform
    function UnitSquareUniform(N)
        h = 1/(N-1)
        xh = range(0, 1, step=h)
        yh = range(0, 1, step=h)
        N = Rectangle(1., 1., N, N)
    end
end
test = UnitSquareUniform(10)
display(test)
struct RectangleBCs
    bot
    right
    top
    left
end
```

### 1.3 Discretization Laplacian and First-Order Advection Term

We discretize the Laplacian using central finite differences with second order as

$$\Delta u(x,y) = \frac{-u(x+h,y) - u(x-h,y) + 4u(x,y) - u(x,y+h) - u(x,y-h)}{h^2} + \mathcal{O}(h^2)$$

such that the stencil for node values reads

$$u_{i,j} \approx \frac{-u_{i+1,j} - u_{i-1,j} + 4u_{i,j} - u_{i,j+1} - u_{i,j-1}}{h^2}$$

$$\Leftrightarrow [-\Delta u_h]_{\xi} = \frac{1}{h^2} \begin{bmatrix} 0 & -1 & 0 \\ -1 & 4 & -1 \\ 0 & -1 & 0 \end{bmatrix}, \quad \xi \in \Omega_h$$

We either use central finite differences or the upwind finite differences for the convection term (with  $v_1^+ := \max(0, v_1)$  and  $v_1^- := \min(0, v_1)$ ):

$$[v \cdot \nabla u_h^{up}]_{\xi} = \begin{bmatrix} 0 & \frac{v_2^-}{h_y} & 0\\ \frac{-v_1^+}{h_x} & \frac{|v_1|}{h_x} + \frac{|v_2|}{h_y} & \frac{v_1^-}{h_x}\\ 0 & \frac{-v_2^+}{h_y} & 0 \end{bmatrix}$$

$$[v \cdot \nabla u_h^{central}]_{\xi} = \frac{1}{h^2} \begin{bmatrix} 0 & v_2 & 0 \\ -v_1 & 0 & v_1 \\ 0 & -v_2 & 0 \end{bmatrix}$$
(FIXME: only for  $h_x = h_y$ )

TODO: Explain RHS vector and source vector

```
[3]: function \Delta (\Omega :: Rectangle)
              = kron
            dxx = spdiagm(-1=>ones(\Omega.Nx-3), 0=>-2ones(\Omega.Nx-2), 1=>ones(\Omega.Nx-3))
            dyy = spdiagm(-1=>ones(\Omega.Ny-3), 0=>-2ones(\Omega.Ny-2), 1=>ones(\Omega.Ny-3))
            return 1/\Omega.hx^2 * I(\Omega.Ny-2) dxx + 1/\Omega.hy^2 * dyy I(\Omega.Nx-2)
      end
      function vo (Ω::Rectangle, v, use_upwind)
              = kron
            if use_upwind
                 v1dxLoc = 1/\Omega.hx * spdiagm(
                       -1 \Rightarrow -\max(0, v[1]) * \operatorname{ones}(\Omega.Nx-3),
                       0 =  abs(v[1])*ones(\Omega.Nx-2),
                       1 = \min(0, v[1]) * \operatorname{ones}(\Omega.Nx-3)
                 v2dyLoc = 1/\Omega.hy * spdiagm(
                       -1 = > -\max(0, v[2]) * ones(\Omega.Ny-3),
                       0 =  abs (v[2]) * ones(\Omega.Ny-2),
                       1 = \min(0, v[2]) * \operatorname{ones}(\Omega.Ny-3)
                 )
            else
                  # central as default
                 v1dxLoc = 1/\Omega.hx^2 * spdiagm(
                       -1 \Rightarrow -v[1] * ones(\Omega.Nx-3),
                       1=>v[1]*ones(\Omega.Nx-3)
                 v2dyLoc = 1/\Omega.hy^2 * spdiagm(
                       -1 \Rightarrow -v[2] * ones(\Omega.Ny-3),
```

```
1=>v[2]*ones(\Omega.Ny-3)
        )
    return I(\Omega.Ny-2) v1dxLoc + v2dyLoc I(\Omega.Nx-2)
end
function b (Ω::Rectangle, f, bcs::RectangleBCs)
    Nx = \Omega.Nx
    Ny = \Omega.Ny
    xInt = \Omega.xh[2:end-1]
    yInt = \Omega.yh[2:end-1]
    fh = vec(f.(xInt,yInt'))
    bh = 1/\Omega.hy^2 .* vec(bcs.bot.(xInt))
    rh = 1/\Omega.hx^2 .* vec(bcs.right.(yInt))
    th = 1/\Omega.hy^2 .* vec(bcs.top.(xInt))
    lh = 1/\Omega.hx^2 .* vec(bcs.left.(yInt))
    bvec = zeros((Nx-2)*(Ny-2))
    bvec += fh
    bvec[1
                                      : Nx-2] += bh
                            : 1
    bvec[(Nx-2)*(Ny-2-1)+1 : 1 : end] += th
    bvec[(Nx-2) : (Nx-2) : end] += rh
    bvec[1
                             : (Nx-2) : end] += 1h
    return byec
end
function solveConvDiff(Ω::Rectangle, f, v, , use_upwind, bcs::RectangleBCs)
    A = - *\Delta (\Omega) + v \circ (\Omega, v, use\_upwind)
    b = b (\Omega, f, bcs)
    return (A) \ b
end
function plotSol(Ω::Rectangle, u, bcs::RectangleBCs, edgeAvg=true)
    pyplot()
    Nx = \Omega.Nx
    Ny = \Omega.Ny
    uMat = zeros(Nx,Ny)
    uMat[2:end-1,2:end-1] = reshape(u, (Nx-2, Ny-2))
    uMat[2:Nx-1,1] = vec(bcs.bot.(\Omega.xh[2:end-1]))
    uMat[2:Nx-1,Ny] = vec(bcs.top.(\Omega.xh[2:end-1]))
    uMat[Nx,2:Ny-1] = vec(bcs.right.(\Omega.yh[2:end-1]))
    uMat[1,2:Ny-1] = vec(bcs.left.(\Omega.yh[2:end-1]))
    if edgeAvg
        uMat[1,1] = 0.5 * (uMat[1,2] + uMat[2,1])
        uMat[1,Ny] = 0.5 * (uMat[1,Ny-1] + uMat[2,Ny])
```

[3]: plotSol (generic function with 2 methods)

## 1.4 Experiments

- Show failure of central diffs for convection-dominated cases
- Show success for upwind diffs for arbitrary cases

```
[4]: \Omega = \text{UnitSquareUniform}(50)

bcs = RectangleBCs(x -> 0, y -> 0, x -> 0, y -> 0)

f(x,y) = 10

v = [-0.2, -0.2]

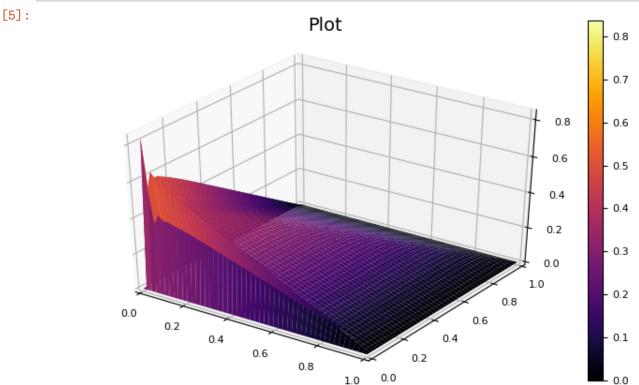
use_upwind = false

u_stable = \text{solveConvDiff}(\Omega, f, v, 1E0, use_upwind, bcs})

plotSol(\Omega, u_stable, bcs})
```

[4]: Plot 0.30 0.25 0.3 0.20 0.2 - 0.15 0.1 0.0 - 0.10 1.0 0.8 0.0 0.6 0.2 - 0.05 0.4 0.4 0.6 0.2 0.8 0.0 1.0 0.00

```
[5]: u_{instable} = solveConvDiff(\Omega, f, v, 1E-1, use_upwind, bcs) # lower epsilon/v_u 
 <math>\hookrightarrow ratio
 plotSol(\Omega, u_{instable}, bcs)
```



```
[6]:  \begin{aligned} \Omega &= \text{UnitSquareUniform(50)} \\ \text{bcs} &= \text{RectangleBCs(x } \rightarrow 0, \text{ y } \rightarrow 0, \text{ x } \rightarrow 0, \text{ y } \rightarrow 0) \\ \text{f(x,y)} &= 10 \\ \text{v} &= [-0.2, -0.2] \\ \text{use\_upwind} &= \text{true} \end{aligned}   \begin{aligned} \text{u\_stable} &= \text{solveConvDiff}(\Omega, \text{ f, v, } 1\text{E-2, use\_upwind, bcs)} \\ \text{plotSol}(\Omega, \text{ u\_stable, bcs)} \end{aligned}
```

[6]:

