## Hörmander's $L^2$ -estimate in dimension one

Following original work by Lars Hörmander we establish a result about the  $\bar{\partial}$ -operator in planar domains. Thus we restrict the attention to  $\mathbf{C}$  where z = x + iy is the complex coordinate.

**Remark.** The full strength of  $L^2$ -estimates appears in dimension  $n \geq 2$  where one works with plurisubharmonic functions and impose conditions on strictly pesudo-convex subsets of  $\mathbb{C}^n$  where one seeks solutions of inhomogeneous  $\bar{\partial}$ -equations for differential forms of every bi-degree (p,q) where  $0 \leq p, q \leq n$ . See Hörmander's text-boook in several complex variables for details.

The Cauchy-Riemann operator sends a differentiable function f into

$$\bar{\partial}(f) = \frac{1}{2}(\partial_x(f) + i \cdot \partial_y(f))$$

The Hilbert space  $\mathcal{H}_{\phi}(\Omega)$ . Let  $\Omega$  be an open set in  $\mathbf{C}$ . A real-valued continuous and non-negative function  $\phi$  on  $\Omega$  gives the Hilbert space  $\mathcal{H}_{\phi}$  whose elements are complex-valued Lebesgue measurable functions f in  $\Omega$  such that

$$\int_{\Omega} |f|^2 \cdot e^{-\phi} \, dx dy < \infty$$

The square root yields norm denoted by  $||f||_{\phi}$ . Let  $\psi$  be another continuous and non-negative function which gives the Hilbert space  $\mathcal{H}_{\psi}$  where the norm of an element g is denoted by  $||g||_{\psi}$ . We consider the  $\bar{\partial}$ -operator which sends a function  $f \in \mathcal{H}_{\phi}$  to  $\bar{\partial}(f) = df/d\bar{z}$  and study the equation

(2) 
$$\bar{\partial}(f) = g : g \in H_{\psi}$$

One seeks conditions for the pair  $(\phi, \psi)$  in order that there exists a constant C such that (2) has a solution f for every g where

$$(3) ||f||_{\phi} \le C \cdot ||g||_{\psi}$$

Notice that (2) does not have a unique solution since f can be replaced by f + a(z) where a is a holomorphic function which belongs to  $\mathcal{H}_{\phi}$ . For example, non-uniqueness fails when  $\Omega$  is a bounded open set and the function  $e^{-\phi}$  is bounded in  $\Omega$ . For then f can be replaced by f + p for an arbitrary polynomial p(z). We shall find a sufficient condition in order that (2-3) above hold.

**Hörmander's condition.** The pair  $\phi, \psi$  satisfies the Hörmander condition if  $\psi$  is a  $C^2$ -function and  $\phi$  is at least a  $C^1$ -function, and there exists a positive constant  $c_0$  such that the following pointwise inequality holds in  $\Omega$ :

(\*) 
$$\Delta(\psi) - 2 \cdot |\nabla(\psi)|^2 + \psi_x \phi_x + \psi_y \phi_y \ge 2 \cdot c_0^2 \cdot e^{\psi(z) - \phi(z)}$$

where we have put  $|\nabla(\psi)|^2 = \psi_x^2 + \psi_y^2$ .

**Main Theorem.** If the pair  $(\phi, \psi)$  satisfies (\*) the equation  $\bar{\partial}(f) = g$  has a solution for every  $g \in \mathcal{H}_{\psi}$  where

$$||f||_{\phi} \leq \frac{1}{c_0} \cdot ||q||_{\psi}$$

Before the proof starts we recall some facts shout linear operators between Hilbert spaces. In general, let  $\mathcal{H}_0$  and  $\mathcal{H}_1$  be a pair of complex Hilbert spaces and  $T \colon \mathcal{H}_0 \to \mathcal{H}_1$  is a densely defined linear operator. Following Torsten Carlemsn's famous monograph about unbounded operators on Hilbert spaces published by Uppsala university in 1923, we recall the construction of an adjoint. Namely, a vector  $y \in \mathcal{H}_1$  belongs to the domain of definition for the adjoint operator  $T^*$  if and only if there exists a constant C such that

(i) 
$$\left| \langle Tx, y \rangle_1 \right| \le C \cdot |x|_0 \quad : x \in \mathcal{D}(T)$$

where  $|x|_0$  is the norm of the vector x taken in  $\mathcal{H}_0$ , and in the left hand side we considered the hermitiain inner product on  $\mathcal{H}_1$ . Since  $\mathcal{D}(T)$  is dense and Hilbert spaces are self-dual, each y for which (i) holds yields a unique vector  $T^*(y) \in \mathcal{H}_0$  such that

(ii) 
$$\langle Tx, y \rangle_1 = \langle x, T^*y \rangle_0$$

In general  $\mathcal{D}(T^*)$  is not a dense subspace of  $\mathcal{H}_1$ . But let us add this as an hypothesis on T. Moreover, assume that the two densely defined operators T and  $T^*$  both are closed, i.e. their graphs are closed in the product of the two Hilbert spaces.

**Exercise.** Suppose that both T and  $T^*$  are closed with dense domains of definition. Assume in addition that there exists a positive constant c such that

$$|T^*y|_0 \ge |y|_1$$
 :  $y \in \mathcal{D}(T^*)$ 

Show that this implies that the range  $T^*(\mathcal{D}(T^*))$  is a closed subspace of  $\mathcal{H}_0$  which is equal to the orthogonal complement of the nullspace  $\mathrm{Ker}(T)$  Moreover, show that for each  $y \in \mathcal{H}_1$  we can find  $x \in \mathcal{D}(T)$  such that

$$Tx = y$$
 &  $|x|_0 \le c^{-1} \cdot |y|_1$ 

## Proof of the Main Theorem

Since  $C_0^{\infty}(\Omega)$  is a dense subspace of  $\mathcal{H}_{\phi}$  the linear operator  $T\colon f\mapsto \bar{\partial}(f)$  from  $\mathcal{H}_{\phi}$  to  $\mathcal{H}_{\psi}$  is densely defined and we leave as an exercise to the reader to check that T is closed. In fact, this relies upon a general fact about closedness of operators defined by differential operators. The reader may also check that Stokes Teorem entails that test-functions in  $\Omega$  belong to  $\mathcal{D}(T^*)$  and since  $C_0^{\infty}(\Omega)$  is dense in the Hilbert space  $\mathcal{H}_{\psi}$  the adjoint is also densely defined. Let us then consider some  $g \in \mathcal{D}(T^*)$ . For each  $f \in C_0^{\infty}(\Omega)$  Stokes theorem gives

(i) 
$$\langle T(f), g \rangle = \int \bar{\partial}(f) \cdot \bar{g} \cdot e^{-\psi} \, dx dy = -\int f \cdot \left[ \bar{\partial}(\bar{g}) - \bar{g} \cdot \bar{\partial}(\psi) \right] \cdot e^{-\psi} \, dx dy$$

Since  $\psi$  is real-valued,  $\bar{\partial}(\bar{w}) - \bar{w} \cdot \bar{\partial}(\psi)$  is equal to the complex conjugate of  $\partial(w) - w \cdot \partial(\psi)$ . We conclude that (i) gives

(ii) 
$$T^*(g) = -\left[\partial(g) - g \cdot \partial(\psi)\right] \cdot e^{\phi - \psi}$$

In particular  $T^*$  is defined via a differential operator and has therefore a closed graph. Taking the squared  $L^2$ -norm in  $\mathcal{H}_{\phi}$  we obtain

$$||T^*(g)||_{\phi}^2 = \int |\partial(g) - g \cdot \partial(\psi)|^2 \cdot e^{\phi - 2\psi} =$$

(iii) 
$$\int \left( |\partial(g)|^2 + |g|^2 \cdot |\partial(\psi)|^2 \right) \cdot e^{\phi - 2\psi} - 2 \cdot \Re\left( \int \partial(g) \cdot \bar{g} \cdot \bar{\partial}(\psi) \cdot e^{\phi - 2\psi} \right)$$

By partial integration the last integral in (iii) is equal to

$$\text{(iv)} \qquad 2 \cdot \mathfrak{Re} \Big( \int g \cdot [\partial(\bar{w}) \cdot \bar{\partial}(\psi) + \bar{g} \cdot \partial \bar{\partial}(\psi) - 2\bar{w} \cdot \bar{\partial}(\psi) \cdot \partial(\psi) + \bar{g} \cdot \bar{\partial}(\psi) \cdot \partial(\phi) \Big] \cdot e^{\phi - 2\psi} \Big)$$

Next, the Cauchy-Schwarz inequality gives

(v) 
$$|2 \cdot \int g \cdot \partial(\bar{g}) \cdot \bar{\partial}(\psi) \cdot e^{\phi - 2\psi} | \leq \int (|\partial(g)|^2 + |g|^2 \cdot |\partial(\psi)|^2) \cdot e^{\phi - 2\psi}$$

Together (iii-v) give

$$||T^*(g)||_\phi^2 \geq 2 \cdot \mathfrak{Re} \int |g|^2 \cdot \left[ \, \partial \bar{\partial}(\psi) - 2 \cdot \bar{\partial}(\psi) \cdot \partial(\psi) + \bar{\partial}(\psi) \cdot \partial(\phi) \, \right] \cdot e^{\phi - 2\psi} = 0$$

(vi) 
$$2 \cdot \mathfrak{Re} \int |g|^2 \cdot \frac{1}{4} \left[ \Delta(\psi) - 2 \cdot |\nabla(\psi)|^2 + \psi_x \phi_x + \psi_y \phi_y \right] \cdot e^{\phi - 2\psi}$$

where the last equality follows since  $\phi$  and  $\psi$  are real-valued. Funally, since (4.1) is assumed it follows that

(vi) 
$$||T^*(g)||_{\phi}^2 \ge c_0^2 \cdot \int |g|^2 \cdot e^{\psi - \phi} \cdot e^{\phi - 2\psi} = c_0^2 \cdot ||g||_{\psi}^2$$

Now we apply the Exercise above and the proof of the Main Theorem is finished.

**Remark.** Let  $\Omega$  be an open subset of a disc  $\{|z| < r\}$  for somne r < 1 which is centered at the origin. Consider the function

$$\phi(z) = \log(1 - |z|^2) = \log(1 - x^2 - y^2)$$

Now we can take  $\psi=\phi$  and Hörmander's condition (\*) is valid. To see this we notice that

(i) 
$$\Delta(\psi) = \frac{4}{(1 - x^2 - y^2)^2}$$

(ii) 
$$\psi_x^2 + \psi_y^2 = \frac{4x^2 + 4y^2}{(1 - x^2 - y^2)^2}$$

Since  $\phi = \psi$  we see that the right hand side in (\*) becomes

(iii) 
$$\frac{4}{(1-x^2-y^2)^2} - \frac{4x^2+4y^2}{(1-x^2-y^2)^2}$$

Inside the disc of radius r < 1 we notice that (iii) is

$$4 \cdot \frac{1 - r^2}{(1 - r^2)^2}$$

which can be taken as  $c_0$  in the Main Theorem.