A topological three-dimensional quantum field theory of gravity.

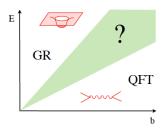
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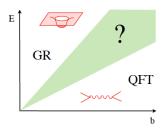
Plan

- Motivation and background
- Axiomatic formulation of a topological quantum field theory
- State sum invariants as a TQFT of gravity

Quantum gravity (QG) seeks to describe gravitational field through the principles of quantum mechanics.



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General relativity is the discovery that spacetime and the gravitational field are the same physical entity. All fields we know exhibit quantum properties at some scale, therefore it is believe spacetime to have quantum properties as well.



Definition

A differential manifold of dimension n is a topological space M with a family of bijective maps $\phi_{\alpha}:U_{\alpha}\subset M\to V_{\alpha}\subset \mathbb{R}^n$ such that

- i) $\bigcup_{\alpha} U_{\alpha} = M$
- ii) for all α, β such that $U_{\alpha\beta} := U_{\alpha} \cap U_{\beta} \neq \emptyset$, the sets $\phi_{\alpha}(U_{\alpha\beta})$ and $\phi_{\beta}(U_{\alpha\beta})$ are open in \mathbb{R}^n an the maps $\phi_{\beta} \circ \phi_{\alpha}^{-1}$ are smooths.

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The pairs $(U_{\alpha}, \phi_{\alpha})$ are called coordinate charts. The set of all charts is called an atlas. We will also require that any two points have disjoint neighborhoods and the atlas of M is countable.

Definition

Let M and N be differentiable manifolds. A continuous map $f: M \to N$ is called differentiable in p if for every chart (V, ϕ) with $f(p) \in V$ there exists a chart (U, ψ) with $p \in U$ such that $f(U) \subset V$ and $\psi \circ f \circ \phi^{-1}$ is differentiable in $\phi(p)$.

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A differentiable bijective map $f: M \to N$ with f^{-1} differentiable is called a diffeomorphism.

From the mathematical point of view we expect that a QG theory keeps

- general features of quantum theory (probabilistic)
- topological framework of spacetime

Both features are embraced by the precise notion of a topological quantum field theory.

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These data are subject to the following axioms:

1) Z is functorial with respect to orientation preserving diffeomorphisms of Σ and M. That is, given an orientation preserving diffeormorphism $f:\Sigma\to\Sigma'$ induces an isomorphism

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If f extends to an o. p. diff. $M \to M'$, with $\partial M = \Sigma$. $\partial M' = \Sigma'$. then Z(f) take Z(M) to Z'(M).

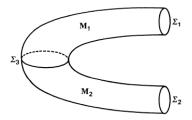
2) $Z(\Sigma^*) = Z(\Sigma)^*$ where Σ^* is Σ with opposite orientation and $Z(\Sigma)^*$ denotes de dual module. Note that theres is nothing said about the orientation of M.

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- 3) Z is multiplicative.

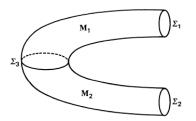
Thus, for disjoint unions

$$Z(\Sigma_1 \cup \Sigma_2) = Z(\Sigma_1) \otimes Z(\Sigma_2).$$

Moreover if $\partial M_1 = \Sigma_1 \cup \Sigma_3$, $\partial M_2 = \Sigma_2 \cup \Sigma_3^*$ and $M = M_1 \cup_{\Sigma_3} M_2$ is the manifold obtained by identification of the common Σ_3 – component.



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Then we require

$$Z(M) = < Z(M_1), Z(M_2) >$$

where <, > denotes the natural pairing

 $Z(\Sigma_1)\otimes Z(\Sigma_3)\otimes Z(\Sigma_3)^*\otimes Z(\Sigma_2) o Z(\Sigma_1)\otimes Z(\Sigma_2)$

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Similarly when $M = \emptyset$, $Z(M) \in R$ is an idempotent element.

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Note that when M is a closed d-dimensional manifold so that $\partial M = \emptyset$, then

$$Z(M) \in Z(\emptyset) = R$$



Einstein equations (movement equations) can be derived through minimizing a function known as the action S:

$$S = rac{1}{16\pi} \int R \sqrt{\det(g_{\mu
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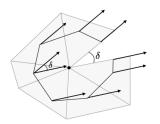
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1961 Regge gave a discrete formulation of general relativity using triangulated manifolds.



Regge showed that the curvature of a triangulated n dimensional manifold is encoded in the (n-2) skeleton.



The discrete form of the action S from which Einstein field equations are derived is given by

$$S_R = \sum |\sigma^i(T)| \, \varepsilon_i$$



In quantum mechanics the action is replaced by a partition function Z:

$$Z=\int e^{iS}D(x)$$

Given a manifold M, Z(M) obtained from the TQFT will have the physical meaning of the partition function.