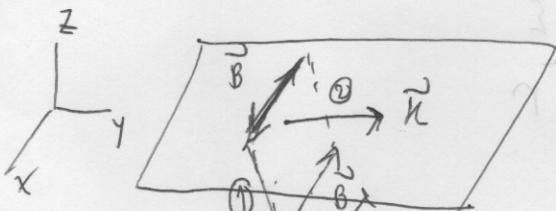


HW 1.

1. (a) \vec{B} should be parallel to the plane but perpendicular to \vec{n} due to symmetry



$$\vec{n} \rightarrow \hat{y}$$

$$\vec{B} \rightarrow \hat{x}$$

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I$$

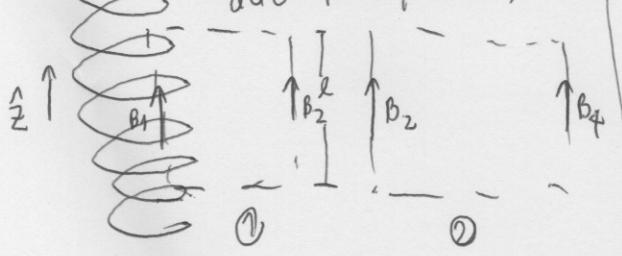
$$B l + 0 + B l + 0 = \mu_0 n l$$

\downarrow

$\vec{B} \perp d\vec{l}$ at ② & ①

$$\vec{B} = \frac{\mu_0 n l \hat{x}}{2} = \frac{\mu_0 n l \hat{x}}{2}$$

2. \vec{B} should be along \hat{z} due to symmetry | loop 2



$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I$$

$$B_2 l - B_4 l = 0 \quad \text{no } I$$

but since $B_4 = 0$ at ∞ (far away)

$$B_2 = 0$$

loop ①

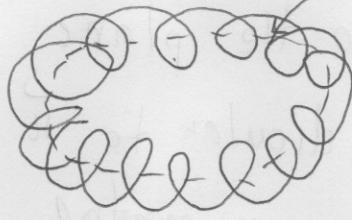
$$B_2 = 0$$

$$B_1 l - B_2 l = \mu_0 n I l$$

$$\vec{B}_1 = \mu_0 n I \hat{z}$$

(C)

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I$$



$$B 2\pi r = \mu_0 N I$$

$$B = \frac{\mu_0 N I}{2\pi r}$$

interesting \vec{B} is the same as in f
outside

$$B 2\pi r = 0$$

I pointing up is the same as
pointing down

$$B = 0$$

& good

$$I_{\text{eff}} = I_{\text{fd}} - I_{\text{sd}}$$

I or

$$I_{\text{sd}} = I_{\text{fd}}$$

(point rot) \Rightarrow to $I_{\text{sd}} = I_{\text{fd}}$ same field

$$I_{\text{sd}} = I_{\text{fd}}$$

① good

$$I_{\text{m},y} = I_{\text{fd}} - I_{\text{sd}}$$

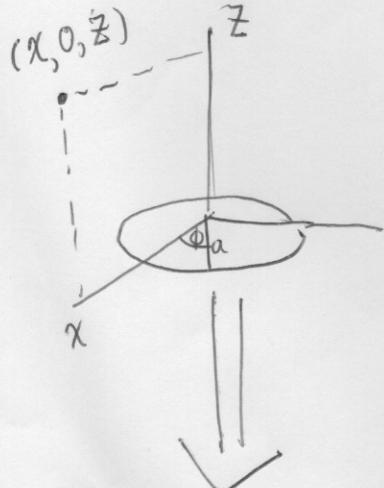
$$\therefore I_{\text{m},y} = I_{\text{fd}}$$

HW 2

$$(b) \vec{A} = \frac{\mu_0 I}{4\pi} \int \frac{d\vec{l}}{r}$$

$$A_x = \frac{\mu_0 I}{4\pi} \int_0^{2\pi} \frac{-a \sin \phi d\phi}{r} \quad A_y = \frac{\mu_0 I}{4\pi} \int_0^{2\pi} \frac{a \cos \phi d\phi}{r}$$

$$r = \sqrt{x^2 + a^2 + z^2 - 2az \cos \phi}$$



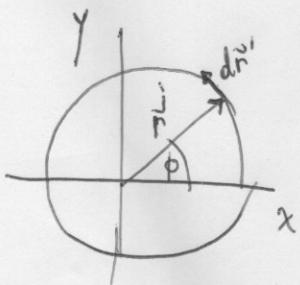
$$\vec{r} = (x, 0, z)$$

$$\vec{r}' = (a \cos \phi, a \sin \phi, 0)$$

$$\vec{r}_c = (x - a \cos \phi, -a \sin \phi, z)$$

$$|\vec{r}_c| = \sqrt{(x - a \cos \phi)^2 + a^2 \sin^2 \phi + z^2}$$

$$= \sqrt{x^2 - 2xa \cos \phi + a^2 + z^2}$$



$$d\vec{l}' = dr' = (-a \sin \phi d\phi, a \cos \phi d\phi, 0)$$

$$(c) A_x = -\frac{\mu_0 I}{4\pi} a \int_0^{2\pi} \frac{\sin \phi d\phi}{\sqrt{x^2 + a^2 + z^2 - 2az \cos \phi}}$$

$$= +\frac{\mu_0 I a}{4\pi} \int_0^{2\pi} \frac{d \cos \phi}{\sqrt{x^2 + a^2 + z^2 - 2az \cos \phi}}$$

$$= \frac{\mu_0 I a}{4\pi} \left[\frac{1}{(-2z)} \right] \left(x^2 + a^2 + z^2 - 2za \cos \phi \right) \Bigg|_{\cos \phi = 0}$$

$$= 0$$

(d)

$$A_y = \frac{\mu_0 I}{4\pi} \int_0^{2\pi} \frac{a \cos \phi}{\sqrt{x^2 + a^2 + z^2 - 2x a \cos \phi}} d\phi$$

$$\text{set } \phi = \pi + 2\beta$$

$$A_y = -\frac{\mu_0 I}{4\pi} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \frac{a \cos(2\beta)(2) d\beta}{\sqrt{x^2 + a^2 + z^2 + 2x a \cos(2\beta)}}$$

$$= -\frac{\mu_0 I}{4\pi} \int_0^{\frac{\pi}{2}} \frac{2a \cos 2\beta d\beta}{\sqrt{x^2 + a^2 + z^2 + 2x a \cos 2\beta}}$$

$$\text{use } \cos 2\beta = 1 - 2 \sin^2 \beta$$

$$= -\frac{\mu_0 I a}{4\pi} \int_0^{\frac{\pi}{2}} \frac{(1 - 2 \sin^2 \beta)}{\sqrt{x^2 + a^2 + z^2 + 2x a (1 - 2 \sin^2 \beta)}} d\beta$$

make a substitution

$$k^2 = \frac{4ax}{(a+x)^2 + z^2}$$

$$= -\frac{\mu_0 I a k}{4\pi} \int_0^{\frac{\pi}{2}} \frac{(1 - 2 \sin^2 \beta)}{\sqrt{4xa (\sqrt{1 - k^2 \sin^2 \beta})}} d\beta$$

$$= -\frac{\mu_0 I h}{2\pi} \sqrt{\frac{a}{x}} \left[\int_0^{\frac{\pi}{2}} \frac{d\beta}{\sqrt{1 - h^2 \sin^2 \beta}} - \int_0^{\frac{\pi}{2}} \frac{2 \sin^2 \beta d\beta}{\sqrt{1 - h^2 \sin^2 \beta}} \right]$$

$$\underbrace{\int_0^{\frac{\pi}{2}} \frac{d\beta}{\sqrt{1 - h^2 \sin^2 \beta}}}_{K(h)}$$

$$\underbrace{\int_0^{\frac{\pi}{2}} \frac{2 \sin^2 \beta d\beta}{\sqrt{1 - h^2 \sin^2 \beta}}}_{(*)}$$

(*) is related to $\frac{\partial \tilde{E}}{\partial k}$

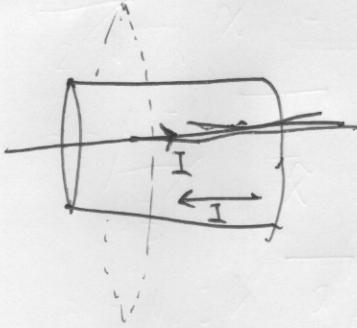
$$\tilde{E}(h) = \int_0^{\frac{\pi}{2}} \frac{1}{\sqrt{1 - h^2 \sin^2 \beta}} d\beta$$

$$\frac{\partial \tilde{E}}{\partial k} = \int_0^{\frac{\pi}{2}} \frac{1}{2 \sqrt{1 - h^2 \sin^2 \beta}} (2h) (-\sin^2 \beta) d\beta = -h \int_0^{\frac{\pi}{2}} \frac{\sin^2 \beta d\beta}{\sqrt{1 - h^2 \sin^2 \beta}}$$

$$\begin{aligned}
 H_y &= -\frac{\mu_0 I h}{2\pi} \sqrt{\frac{a}{x}} \left[\tilde{\kappa}(k) + \frac{2}{k} \frac{\partial \tilde{E}}{\partial k} \right] \\
 &= -\frac{\mu_0 I h}{2\pi} \sqrt{\frac{a}{x}} \left[\tilde{\kappa}(\cancel{k}) + \frac{2}{k} \left(\frac{\tilde{E}}{k} - \frac{\tilde{\kappa}}{k} \right) \right] \\
 &= -\frac{\mu_0 I h}{2\pi} \sqrt{\frac{a}{x}} \left[\tilde{\kappa} \left(1 - \frac{2}{k^2} \right) + \frac{2\tilde{E}}{k^2} \right] \\
 &= -\frac{\mu_0 I}{\pi k} \sqrt{\frac{a}{x}} \left[\tilde{\kappa} \left(\frac{k^2}{2} - 1 \right) + \tilde{E} \right] \\
 &= \frac{\mu_0 I}{\pi k} \sqrt{\frac{a}{x}} \left[\tilde{\kappa}(k) \left(1 - \frac{k^2}{2} \right) - \tilde{E}(k) \right]
 \end{aligned}$$

HW 3.

(a)



outside

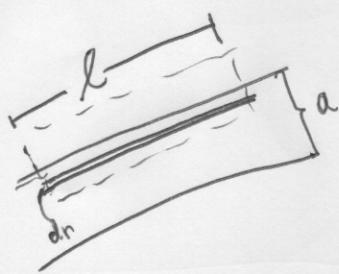
$$B 2\pi r = \mu_0 (I - I) \\ B = 0$$

inside

$$B 2\pi r = \mu_0 I_0 \cos \omega t$$

$$\vec{B} = \frac{\mu_0 I_0 \cos \omega t}{2\pi r} \hat{\phi}$$

$$(b) \Phi_B = \int \vec{B} \cdot d\vec{A}$$



$$= \int_r^a \frac{\mu_0 I_0 \cos \omega t l}{2\pi r} dr \quad \left. \begin{array}{l} \text{outside} \\ B = 0 \end{array} \right\}$$

$$= \frac{\mu_0 I_0 \cos \omega t l}{2\pi} \ln \left(\frac{a}{r} \right)$$

$$(c) \oint \vec{E} \cdot d\vec{l} = - \frac{\partial \Phi_B}{\partial t}$$

$$E_z = + \frac{\mu_0 I_0 (\omega) (\sin \omega t)}{2\pi} \ln \frac{a}{r}$$

$$E = \frac{\mu_0 I_0 \omega \sin \omega t}{2\pi} \ln \frac{a}{r} \hat{z}$$

(d) displacement current density

$$\vec{J}_D = \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$= \frac{\mu_0 \epsilon_0 I_0 \omega^2 \cos \omega t}{2\pi} \ln \frac{a}{r} \hat{z}$$

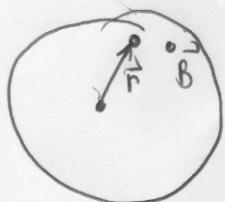
$$(e) I_D = \int \vec{J}_D \cdot d\vec{A} = \frac{\mu_0 \epsilon_0 I_0 \omega^2 \cos \omega t}{2\pi} \int_0^a (2\pi r) r \ln \frac{a}{r} dr$$

HW 4.

(4) (a)

$$\vec{F}_{\text{mag}} = q(\vec{v} \times \vec{B})$$

$$\vec{\tau} = \vec{r} \times \vec{F}_{\text{mag}} = q[\vec{r} \times (\vec{v} \times \vec{B})]$$



$$= q \left[\vec{v} (\vec{r} \cdot \vec{B}) - \vec{B} (\vec{r} \cdot \vec{v}) \right]$$

field into the page

$$\vec{\tau} = -q\vec{B}(\vec{r} \cdot \vec{v})$$

$$(b) \quad \vec{\tau} = \frac{d\vec{L}}{dt} = -q\vec{B}(\vec{r} \cdot \frac{d\vec{l}}{dt})$$

$$\text{put } d\vec{l} = dr \hat{r} + d\phi \hat{\phi}$$

$$\frac{d\vec{l}}{dt} = \vec{\tau} = -q\vec{B}(r \frac{dr}{dt})$$

$$\frac{dL}{dt} = -qBr \frac{dr}{dt} \Rightarrow L - L_0 = - \int_0^R qBr \frac{dr}{dt} dt$$

$$L - L_0 = -\frac{q\Phi_B}{2\pi} \quad \text{where } \Phi_B = 2\pi \int_0^R Br dr$$

(c) when $\Phi_B = 0$, then $L = L_0$ (i.e. conserve angular momentum)

If the particle comes out from the center,

$$\vec{L}_0 = 0 \text{ (since } \vec{p} = 0\text{)}$$

then it must exist with $L = 0$ and only possibility is when it travels radially (~~$\vec{L} = \vec{r} \times \vec{p} = 0$~~)

must be 0

$$(\vec{L} = \vec{r} \times \vec{p} = 0)$$

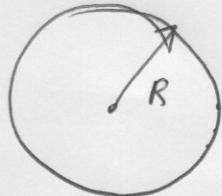
$\vec{v} \text{ must } \parallel \text{ to } \vec{r}$

(d)

again $L_0 = 0$ since it starts from center

$$L = \frac{q\Phi}{2\pi}$$

$$L = m\omega R^2 = \frac{qB}{2\pi} \cancel{\pi} R^2$$



$$\omega = \frac{qB}{2m}$$