

Eigenvalue-eigenfunctions: Legendre

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Objective: Obtain the eigenvalues and eigenfunctions of our linear differential operator

$$\mathcal{L} \phi_n(x) = \lambda_n \phi_n(x)$$

- * homogeneous linear differential equation
- * λ_n : eigenvalues \rightarrow to be determined
- * $\phi_n(x)$: eigenfunctions \rightarrow to be determined

Generally, we can write

$$\frac{d^2}{dx^2} \phi_n(x) + p(x) \frac{d}{dx} \phi_n(x) + q(x) \phi_n(x) = 0$$

Mode of solution \rightarrow power series (Frobenius method)
- infinite series in powers of $(x-x_0)$
 $x_0 \rightarrow$ expansion point

ordinary (expansion) point \rightarrow when $P(x)$ and $Q(x)$ are well defined at x_0 .
 $\xrightarrow{\text{no singularity}}$

Example: Legendre Equation

$$\mathcal{L} \phi(x) = \lambda \phi(x); -1 \leq x \leq 1$$

$$-(1-x^2) \frac{d^2}{dx^2} \phi(x) + 2x \frac{d}{dx} \phi(x) = \lambda \phi(x)$$

\rightarrow this ODE always appear in physics
e.g. angular part of \vec{r}^2 in spherical form ($x = \cos\theta$)

\rightarrow the differential operator is already self-adjoint

$$\mathcal{L} = -(1-x^2) \frac{d^2}{dx^2} + 2x \frac{d}{dx}$$

$$P_0(x) = -(-1)^0 x^0 \quad P_1(x) = \frac{dP_0}{dx}$$

We solve this using the Frobenius method (series solution)

D) Assume a general series form

$$\phi(x) = x^s \sum_{j=0}^{\infty} a_j x^j$$

$$\phi(x) = \sum_{j=0}^{\infty} a_j x^{s+j}$$

\Rightarrow get the needed derivatives

$$\frac{d\phi}{dx} = \sum_{j=0}^{\infty} a_j (s+j)x^{s+j-1}$$

$$\frac{d^2\phi}{dx^2} = \sum_{j=0}^{\infty} a_j (s+j)(s+j-1)x^{s+j-2}$$

$$\text{at } j=0 \quad (s+j)(s+j-1)=0 \quad \Rightarrow \quad s=0$$

$$\text{or } s(s-1)=0 \quad (\text{circular equation})$$

The circular equation tells us that there are two solutions to our 2nd order ODE.

$$\text{solution for } s=0 \quad \phi(x) = \sum_{j=0}^{\infty} a_j x^j$$

$$\text{solution for } s=1 \quad \phi(x) = \sum_{j=0}^{\infty} a_j x^{j+1}$$

3) obtain the recurrence relation for a_j for each solution

$$\text{for } s=0 \quad -(1-x^2) \sum_{j=2}^{\infty} a_j j(j-1)x^{j-2} + 2x \sum_{j=1}^{\infty} a_j jx^{j-1} = \sum_{j=0}^{\infty} a_j x^j$$

$$-\sum_{j=2}^{\infty} a_j j(j-1)x^{j-2} + \sum_{j=2}^{\infty} a_j j(j-1)x^{j-1} + \sum_{j=1}^{\infty} a_j 2jx^j = \sum_{j=0}^{\infty} a_j x^j$$

outlast in terms of powers of x

make this consistent with the other terms by shifting the index.

$$-\sum_{j=2}^{\infty} a_j j(j-1)x^{j-2} (j+2)(j+1)x^{j-2}$$

$$= -\sum_{j=0}^{\infty} a_{j+2} (j+2)(j+1)x^j$$

canceling same powers of x

$$-a_{j+2}(j+2)(j+1) + a_j j(j-1) + a_j z_j = a_j \lambda$$

$$a_j j(j-1) + a_j z_j - a_j \lambda = (j+2)(j+1)a_{j+2}$$

$$a_{j+2} = \frac{j(j-1) + 2j - \lambda}{(j+2)(j+1)} a_j$$

↳ recurrence relation
you can obtain a_{j+2} by knowing the previous a_j value

Notice from the recurrence relation

$$a_{j+2} = \frac{j(j-1) + 2j - \lambda}{(j+2)(j+1)} a_j \quad \begin{array}{l} j=0 \\ j=1 \\ j=2 \\ \vdots \\ j=S \end{array} \quad \begin{array}{l} a_0 = \square a_0 \\ a_1 = \square a_0 \\ a_2 = \frac{2-\lambda}{3} a_1 \\ \vdots \\ a_S = \square a_1 \end{array}$$

Solutions can be separated to odd and even solutions

even solutions: $a_0 \neq 0, a_1 = 0 \quad (s=0)$

odd solution: $a_0 = 0, a_1 \neq 0 \quad (s=1)$

The general solution to the legendre equation is

$$\phi(x) = a_0 \phi_{\text{even}}(x) + a_1 \phi_{\text{odd}}(x)$$

$\Rightarrow \phi_{\text{even}}(x)$ and $\phi_{\text{odd}}(x)$ are linearly independent

\Rightarrow Both $\phi_{\text{even}}(x)$ and $\phi_{\text{odd}}(x)$ are infinite series about $x=0$ with the same recurrence relation

The infinite series are convergent only when $|x| < 1$

Consider the j terms in the series

$$a_{j+2} = \frac{j(j-1) + 2j - \lambda}{(j+2)(j+1)} a_j$$

$$= \frac{j^2 + j - \lambda}{j^2 + 3j + 2} a_j$$

large j $a_{j+2} \sim \frac{j^2}{j^2} a_j$ or $a_{j+2} \sim a_j$

This means that

$$a_0 \phi_{\text{even}}(x) = a_0 + a_2 x^2 + a_4 x^4 + \dots \quad a_j \sim a_{j+2}$$

large j behavior: $a_{j+2} \sim a_j$

When we set $x \pm 1$

$$a_0 \phi_{\text{even}}(1) = a_0 + a_2 + a_4 + \dots$$

these numbers are NOT decreasing. This term will not converge!!!

④ Fix the eigenvalue λ by imposing the series to be finite (polynomial)

Terminate the infinite series! $\phi_\ell(x) = \sum_{j=0}^{\ell} a_j x^j \rightarrow \sum_{j=0}^{\ell} a_j x^j$

$\ell \Rightarrow$ max integer c_j

$$a_{\ell+2} = \frac{\ell(\ell-1) + 2\ell - \lambda}{(\ell+2)(\ell+1)} a_\ell$$

a_ℓ is now the coefficient of the highest power x^ℓ

at $j=\ell \Rightarrow$ the series will stop

$$\ell(\ell-1) + 2\ell - \lambda = 0$$

$$\ell^2 - \ell + 2\ell - \lambda = 0$$

$$\ell^2 + \ell - \lambda = 0$$

$$\lambda = \ell(\ell+1)$$

for $\ell=0, 1, 2, \dots$

We obtained the eigenvalue in our eigenvalue-eigenfunction problem

= Impose that our solution is finite at $x=\pm 1$

Let us list down some eigenvalues and their corresponding un-normalized eigenfunctions.

Note: for each eigenvalue we have the recurrence relation

$$\phi_\ell(x) = \sum_j a_j^{(\ell)} x^j$$

Note: $\lambda = \ell(\ell+1)$

$$\ell=0 \rightarrow \lambda=0$$

$$\ell=1 \rightarrow \lambda=1(1+1)=2$$

$$\ell=2 \rightarrow \lambda=2(2+1)=6$$

$$\vdots$$

$$\phi_0(x) = a_0^{(0)}$$

$$\phi_1(x) = a_1^{(1)} x$$

$$\phi_2(x) = a_0^{(0)} + a_2^{(0)} x^2$$

$$\vdots$$

use the recurrence formula

$$a_{j+2}^{(\ell)} = \frac{j(j-1) + 2j - \lambda}{(j+2)(j+1)} a_j^{(\ell)}$$

$$\phi_2(x) = a_0^{(0)} + \frac{0(0-1) + 2 \cdot 0 - 6}{(0+2)(0+1)} a_0^{(0)} x^2$$

$$= a_0^{(0)} (1 - 3x^2)$$

$$\ell=3 \rightarrow \lambda=3(3+1)=12$$

$$\phi_3(x) = a_1^{(1)} x + a_3^{(3)} x^3$$

$$\phi_3(x) = a_1^{(1)} x + \frac{1(1-1) + 2 \cdot 1 - 12}{(1+2)(1+1)} a_1^{(1)} x^3$$

$$= a_1^{(1)} (x - \frac{5}{3}x^3)$$

Notice that all $\phi_\ell(x)$ will have a common factor $a_0^{(\ell)}$ or $a_1^{(\ell)}$

the value can be fixed by normalization

Legendre ODE is already self-adjoint, so $\phi(x) = 1$ and

$$\langle f | g \rangle = \int_{-1}^1 dx f^*(x) g(x)$$

Q4. List down the $\ell=4$ & $\ell=5$ eigenvalues and obtain the corresponding eigenfunctions. Normalize $\phi_\ell(x)$ to $\phi(x)$.

Table 15.1 Legendre Polynomials

$P_0(x) = 1$	$P_1(x) = x$	$P_2(x) = \frac{1}{2}(3x^2 - 1)$	$P_3(x) = \frac{1}{2}(5x^3 - 3x)$	$P_4(x) = \frac{1}{8}(35x^4 - 30x^2 + 3)$	$P_5(x) = \frac{1}{8}(63x^5 - 70x^3 + 15x)$	$P_6(x) = \frac{1}{16}(231x^6 - 315x^4 + 105x^2 - 5)$	$P_7(x) = \frac{1}{16}(429x^7 - 693x^5 + 315x^3 - 35x)$	$P_8(x) = \frac{1}{128}(6435x^8 - 12012x^6 + 6930x^4 - 1260x^2 + 35)$
$\phi_\ell(x) = \boxed{P_\ell(x)}$								

Legendre polynomials are usually listed in their unnormalized form.

The orthogonality condition takes the form

$$\langle P_n | P_m \rangle = \int_{-1}^1 P_n(x) P_m(x) dx = \frac{2}{2n+1} S_m$$

Recall

$$\mathcal{L} \phi(x) = \ell(\ell+1) \phi(x); -1 \leq x \leq 1$$

$$-(1-x^2) \frac{d^2}{dx^2} \phi(x) + 2x \frac{d}{dx} \phi(x) = \ell(\ell+1) \phi(x)$$

For a given ℓ , the general solution to the legendre equation is of the 2nd kind

$\phi(x) = C_1 P_\ell(x) + C_2 Q_\ell(x)$

- Legendre function of the 2nd kind
- finite polynomial
- finite-valued for all x in $-1 \leq x \leq 1$
- not used in physics

