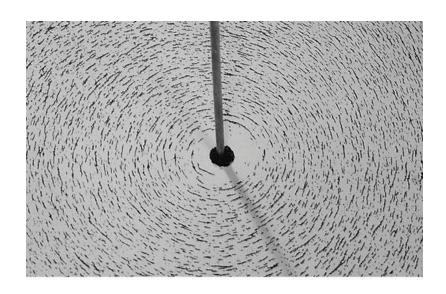
Chapter 28

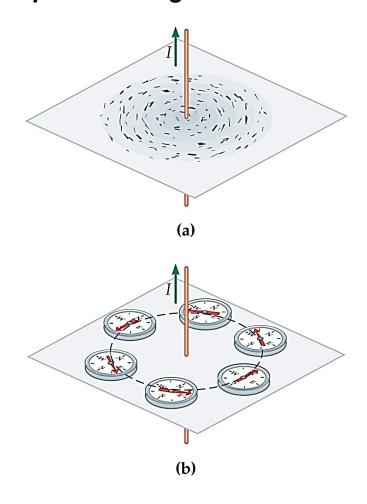
Sources of Magnetic Fields



*Electric Currents and Magnetism

Experimental observation:

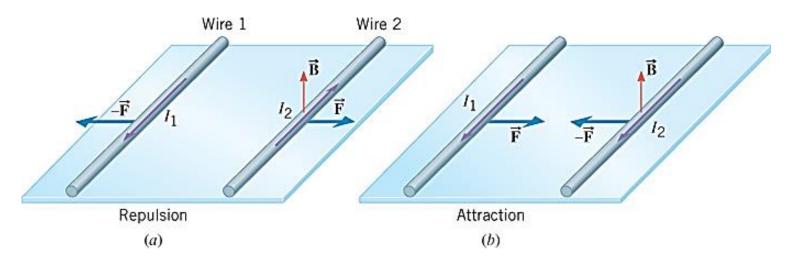
• Electric currents can produce magnetic fields.



*Electric Currents and Magnetism

Experimental observation:

• Current carrying wires can exert forces on each other.



The force between parallel wires

$$F = 2k_m \frac{I_1 I_2}{d} L = I_1 L \left(2k_m \frac{I_2}{d} \right) = I_2 L \left(2k_m \frac{I_1}{d} \right)$$

$$B = 2k_m \frac{I}{d} ?$$

Magnetic Field of a Moving Charge

For a single moving point charge q

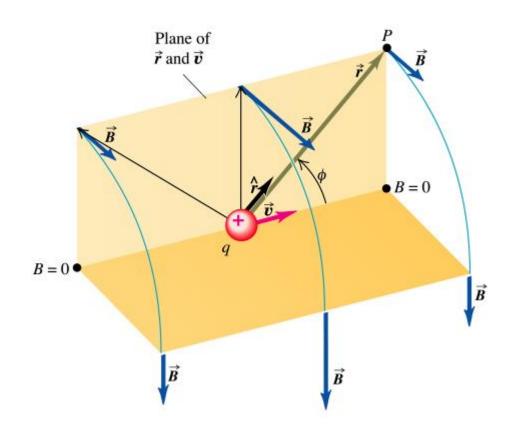
Magnetic field produced at P.

Force on parallel wires

$$B = k_m \frac{|q|v \sin\phi}{r^2}$$

In vector form

$$\vec{B} = k_m \frac{q \vec{v} \times \vec{r}}{r^3}$$



The permeability constant:

$$k_m = \frac{\mu_0}{4\pi} = 1 \times 10^{-7} \text{ T} \cdot \text{m/A}$$
 $\mu_0 = 4\pi \times 10^{-7} \text{ T} \cdot \text{m/A}$

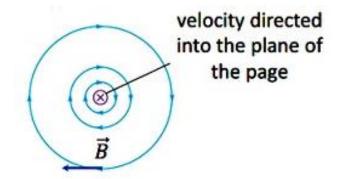
Magnetic Field of a Moving Charge

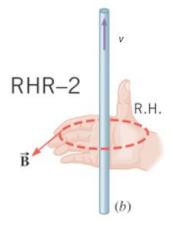
For a single moving point charge q

Vector magnetic field

$$\vec{B} = k_m \frac{q \, \vec{v} \times \vec{r}}{r^3}$$

Direction by right-hand rule No. 2





Magnetic Field of a Current Element

In a conductor, moving charges gives rise to the current.

- The net magnetic field around the conductor is the superposition of the magnetic field of each charge.
- For a short segment dL

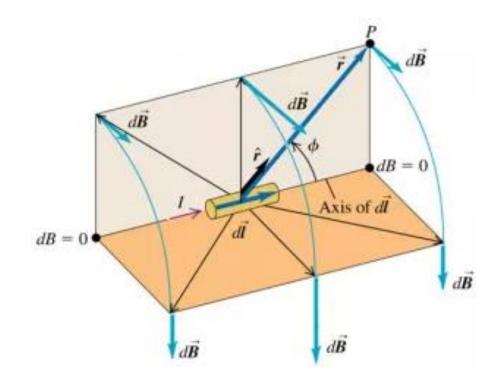
$$dQ = q n(AdL)$$

n =charge carrier density

• The field *dB*

$$dB = k_m \frac{q(nAvdL) \sin \phi}{r^2}$$

$$= k_m \frac{IdL \sin \phi}{r^2}$$
with $I = nqvA$



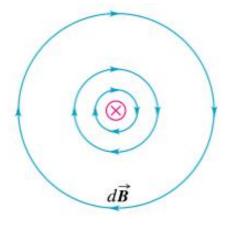
Magnetic Field of a Current Element

In a conductor, moving charges gives rise to the current.

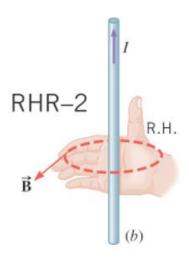
• The field $d\vec{B}$

$$dB = k_m \frac{IdL \sin\phi}{r^2}$$

$$d\vec{B} = k_m \frac{I \ d\vec{L} \times \vec{r}}{r^3}$$



Current into page



Magnetic Field of a Current Element

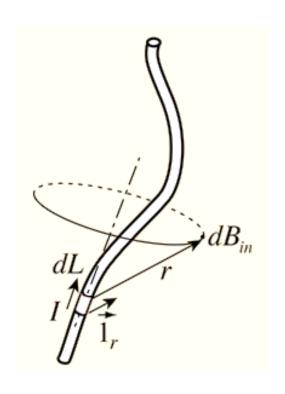
Biot-Savart Law

• The field $d\vec{B}$ due to a length element

$$d\vec{B} = k_m \frac{I \ d\vec{L} \times \vec{r}}{r^3}$$

For a finite length of wire

$$\vec{B} = k_m \int \frac{I \ d\vec{L} \times \vec{r}}{r^3}$$



Using Biot-Savart Law

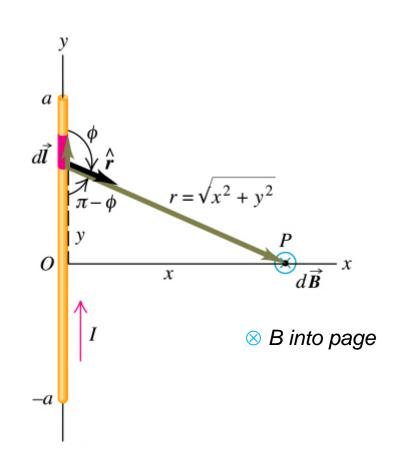
Magnetic field produced by a straight current-carrying conductor of length 2a.

$$\vec{B} = k_m \int \frac{I \ d\vec{L} \times \vec{r}}{r^3}$$

For a finite length of wire

$$dL = dy r = \sqrt{x^2 + y^2}$$
$$\sin \phi = x/\sqrt{x^2 + y^2}$$

$$B = k_m I \int \frac{x \, dy}{(x^2 + y^2)^{3/2}}$$



Using Biot-Savart Law

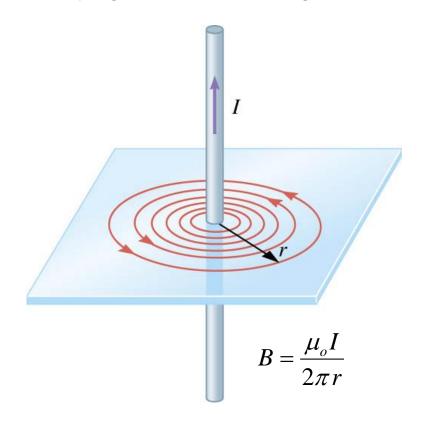
Magnetic field produced by a straight current-carrying conductor of length 2a.

On the perpendicular bisector

$$B = k_m I \frac{2a}{x\sqrt{(x^2 + a^2)}}$$

• In the limit $a \gg x$

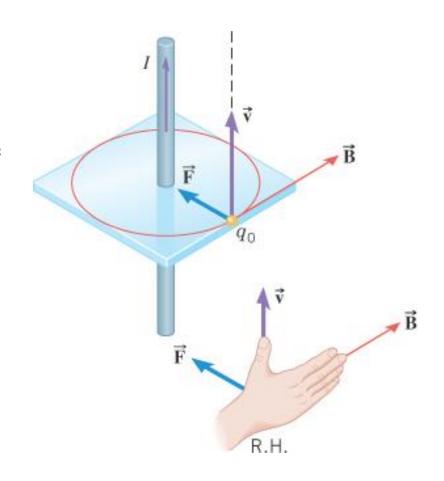
$$B = 2k_m \frac{I}{x}$$



Example: A Current Exerts a Force on a Moving Charge

The long straight wire carries a current of 3.0 A. A particle has a charge of +6.5x10⁻⁶ C and is moving parallel to the wire at a distance of 0.050 m. The speed of the particle is 280 m/s.

Determine the magnitude and direction of the magnetic force on the particle.



Example: A Current Exerts a Force on a Moving Charge

The magnetic field.

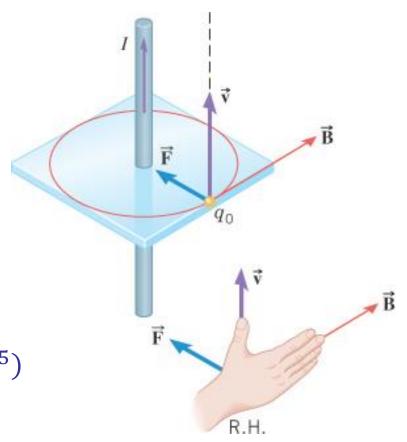
$$B = 2k_m \frac{I}{r}$$

$$= \frac{(2 \times 10^{-7})(3.0)}{(0.050)}$$

$$= 1.2 \times 10^{-5} \text{ T}$$

The magnetic force

$$F = qvB \sin 90^{\circ}$$
= $(6.5 \times 10^{-6})(280)(1.2 \times 10^{-5})$
= $2.18 \times 10^{-8} \text{ N}$



Force Between Parallel Currents

The magnetic force parallel current carrying wires

• Magnetic field due to I_1 :

$$B_1 = 2k_m \frac{I_1}{r}$$

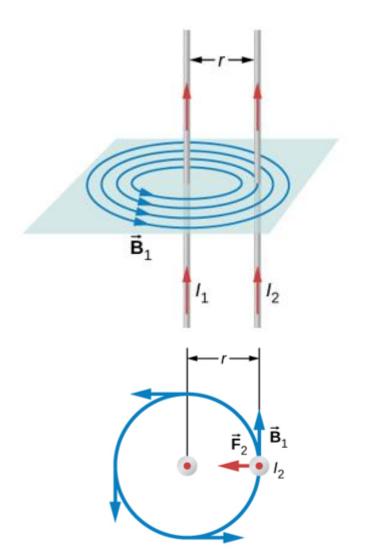
• The force on I_2 :

$$F_2 = I_2 L B_1$$
$$= 2k_m \frac{I_1 I_2}{r} L$$

• The force on I_1 :

$$F_1 = I_1 L B_2$$
$$= 2k_m \frac{I_1 I_2}{r} L$$

By Newton's 3rd law



Force Between Parallel Currents

Example:

Two straight parallel wires 4.5 mm apart carry equal currents of 15.0 A in opposite directions.

The magnetic force:

$$F = 2k_m \frac{I_1 I_2}{r} L = L(2 \times 10^{-7})(15.0)^2 / (0.0045)$$
$$= 0.01 L$$

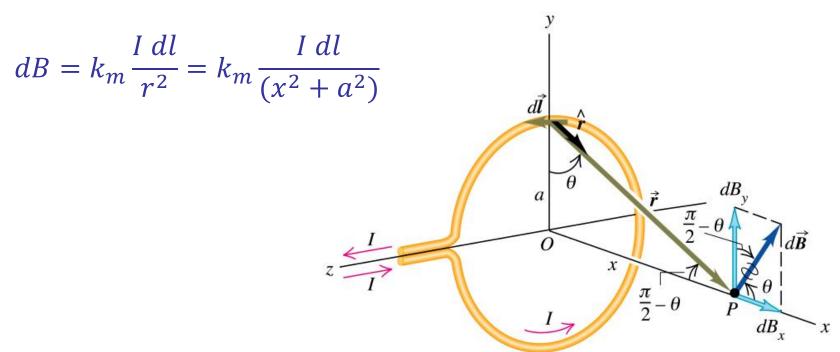
The force per unit length:

$$F/L = 0.01 \text{ N/m}$$

Magnetic field on the axis of circular loop of radius a.

• The field $d\vec{B}$ due to element dL

$$d\vec{B} = k_m \frac{I \ d\vec{L} \times \vec{r}}{r^3}$$



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Magnetic field on the axis of circular loop of radius a.

• The components of $d\vec{B}$

$$dB_{\chi} = dB\cos\theta = k_{m}Idl\frac{a}{(x^{2} + a^{2})^{3/2}}$$

$$dB_{y} = dB\sin\theta = k_{m}Idl\frac{x}{(x^{2} + a^{2})^{3/2}}$$

$$dB_{y} = dB\sin\theta = k_{m}Idl\frac{x}{(x^{2} + a^{2})^{3/2}}$$

Magnetic field on the axis of circular loop of radius a.

Integrate

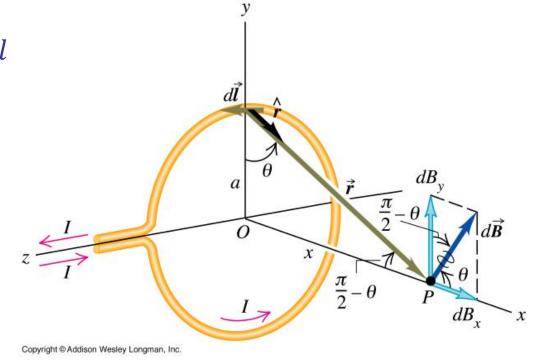
$$B_{x} = k_{m} Ia \int \frac{dl}{(x^{2} + a^{2})^{3/2}}$$

$$= k_{m} \frac{Ia}{(x^{2} + a^{2})^{3/2}} \int dl$$

$$= 2\pi k_{m} \frac{Ia^{2}}{(x^{2} + a^{2})^{3/2}}$$

$$B_{y} = \int dB_{y} = 0$$

(By symmetry)

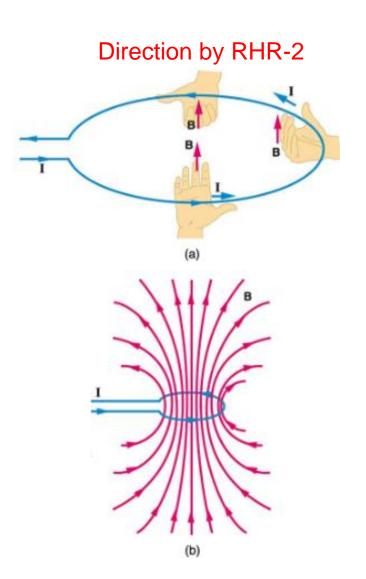


Magnetic field on the axis of circular loop of radius a.

$$B_{x} = 2\pi k_{m} \frac{Ia^{2}}{(x^{2} + a^{2})^{3/2}}$$

At the center (*x*=0)

$$B_0 = 2\pi k_m \frac{I}{a}$$

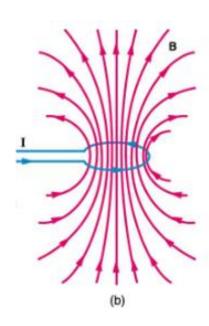


Example:

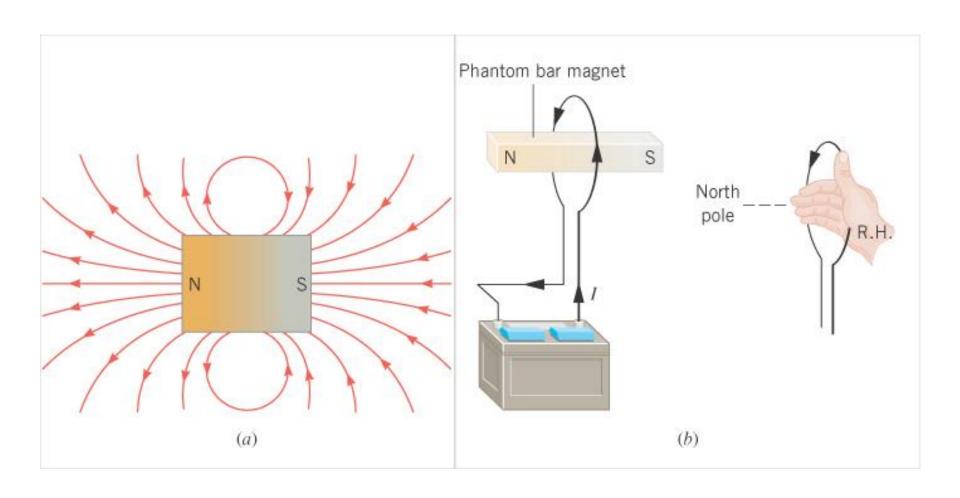
A coil consisting of 100 circular loops with radius 0.60 m carries a 5.0-A current.

• At the center (*x*=0)

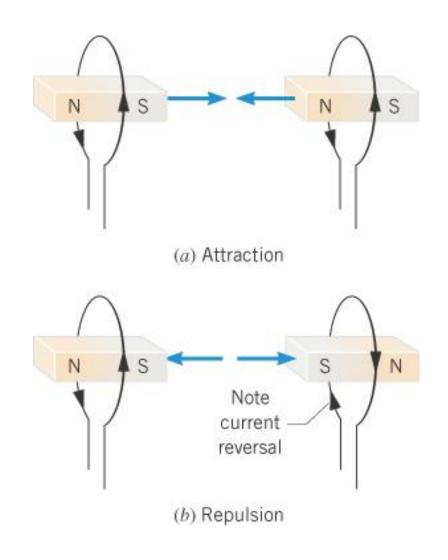
$$B_0 = N2\pi k_m \frac{I}{a} = \frac{100(2\pi \times 10^{-7})(5.0)}{(0.060)}$$
$$= 5.2 \times 10^{-3} \text{ T}$$



Magnetic field lines around a *circular loop* resemble those around a *bar magnet*.



Two circular loop interact the same way as two bar magnets would.

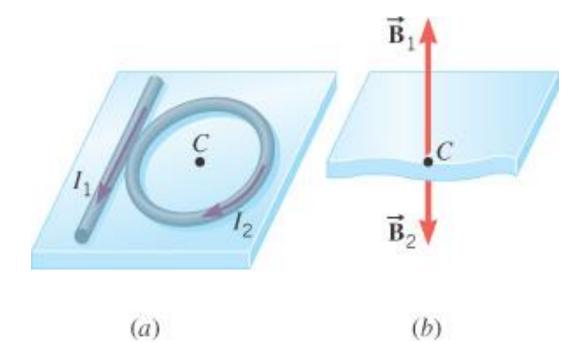


*Magnetic Fields - Superposition

Example: Superposition of Magnetic Fields

A long straight wire carries a current of 8.0 A and a circular loop of wire carries a current of 2.0 A and has a radius of 0.030 m. Find the magnitude and direction of the magnetic field at the center of the loop.

At the center C:



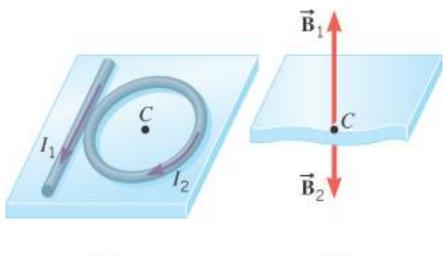
*Magnetic Fields - Superposition

Example: Superposition of Magnetic Fields

At the center C:

$$B = \frac{\mu_o I_1}{2\pi r} - \frac{\mu_o I_2}{2R} = \frac{\mu_o}{2} \left(\frac{I_1}{\pi r} - \frac{I_2}{R} \right)$$

$$B = \frac{\left(4\pi \times 10^{-7} \text{ T} \cdot \text{m/A}\right)}{2} \left(\frac{8.0 \text{ A}}{\pi (0.030 \text{ m})} - \frac{2.0 \text{ A}}{0.030 \text{ m}} \right) = 1.1 \times 10^{-5} \text{ T}$$

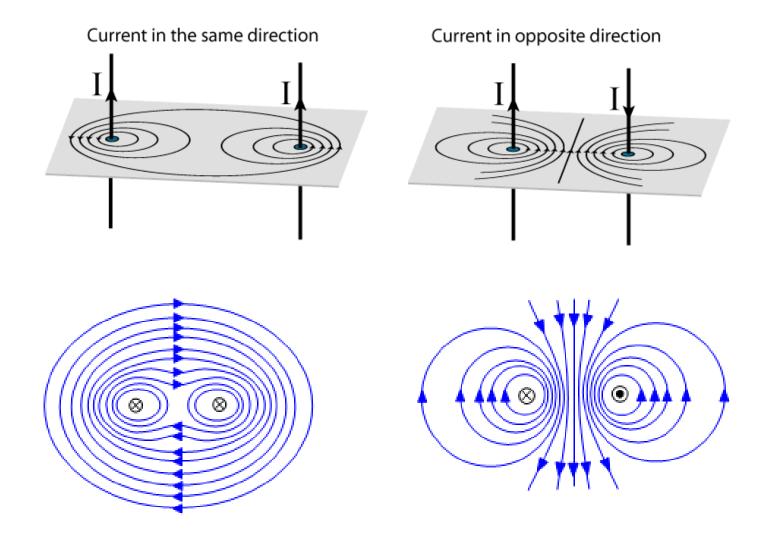


(a)

(b)

*Magnetic Fields - Superposition

The magnetic field due to two parallel currents.

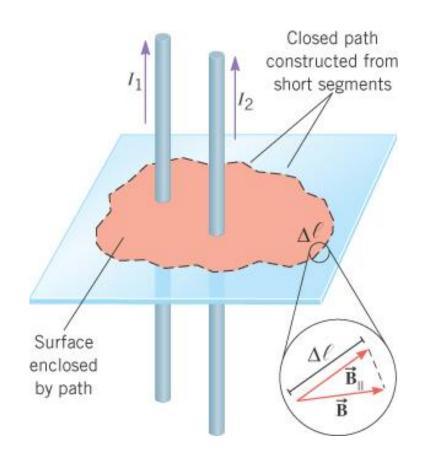


The line integral of the magnetic field.

$$\sum (\vec{\boldsymbol{B}} \cdot \Delta \vec{\boldsymbol{l}}) = \sum (B_{\parallel} \Delta l)$$

For infinitesimal line segments.

$$\oint (\vec{\mathbf{B}} \cdot d\vec{\mathbf{l}}) = \oint (B_{\parallel} dl)$$



The line integral of the magnetic field.

Apply to a single straight current (a)

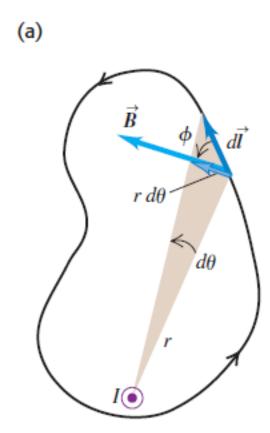
$$B = 2k_m \frac{I}{r}$$

$$\vec{B} \cdot d\vec{l} = Bdl \cos\phi = B(rd\theta)$$

$$\oint (\vec{\mathbf{B}} \cdot d\vec{\mathbf{l}}) = \oint \left(2k_m \frac{I}{r}\right) (rd\theta)$$

$$= (2k_m I) \int_0^{2\pi} d\theta$$

$$=4\pi k_m I = \mu_0 I$$



The line integral of the *magnetic field*.

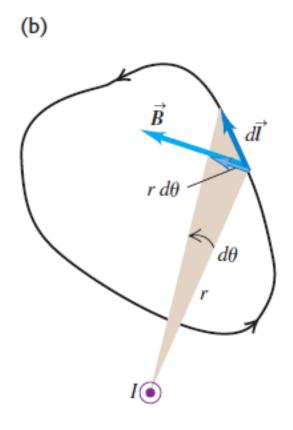
Apply to a single straight current (b)

$$\oint (\vec{B} \cdot d\vec{l}) = \oint \left(2k_m \frac{I}{r}\right) (rd\theta)$$

$$= (2k_m I) \int_{\theta_1}^{\theta_1} d\theta = 0$$

Ampere's law

$$\oint (\vec{\mathbf{B}} \cdot d\vec{\mathbf{l}}) = \mu_0 I_{enclosed}$$

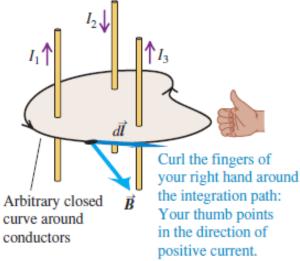


The line integral of the *magnetic field*.

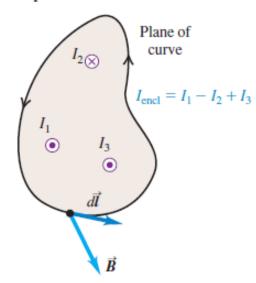
Apply to a system of three straight currents.

$$\oint (\vec{\mathbf{B}}_{net} \cdot d\vec{l}) = \mu_0 I_{encl}$$

$$I_{encl} = I_1 - I_2 + I_3$$



Top view

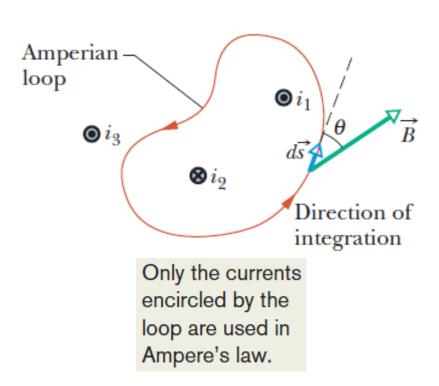


The line integral of the *magnetic field*.

Apply to any system of currents.

$$I_{encl} = \sum_{k} I_{k}$$

Net current through area bounded by the Amperian loop



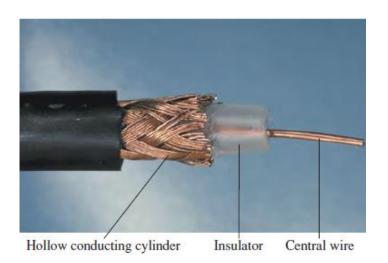
Example: Co-axial cable

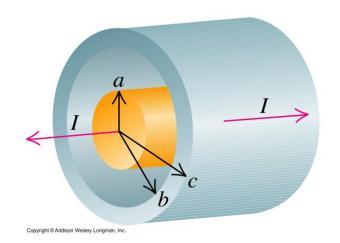
Circular Amperian loop outside the cable.

$$B = 0$$
 (for $r > c$)
(Net current = 0)

Circular Amperian loop inside the cable.

$$B = \frac{\mu_0 I}{2\pi r}$$
(for $b > r > a$)



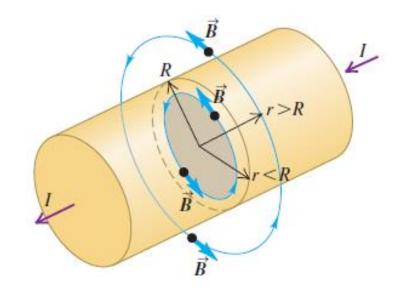


Example: Solid cylindrical conductor

 Circular Amperian loop outside the cable.
 (Symmetry requires B to be uniform along the circle)

$$\oint (\vec{B} \cdot d\vec{l}) = B(2\pi r) = \mu_0 I$$

$$B = \frac{\mu_0 I}{2\pi r} \quad \text{(for } r > R\text{)}$$



Example: Solid cylindrical conductor

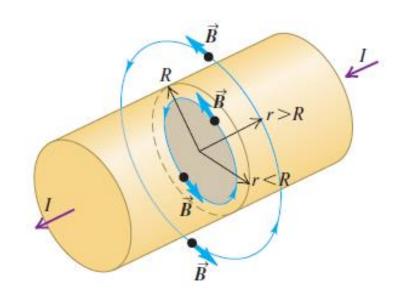
 Circular Amperian loop inside the cable. (Symmetry requires B to be uniform along the circle)

Current density =
$$I/(\pi R^2)$$

$$\oint (\vec{B} \cdot d\vec{l}) = B(2\pi r)$$

$$= \mu_0 I \left(\frac{\pi r^2}{\pi R^2}\right)$$

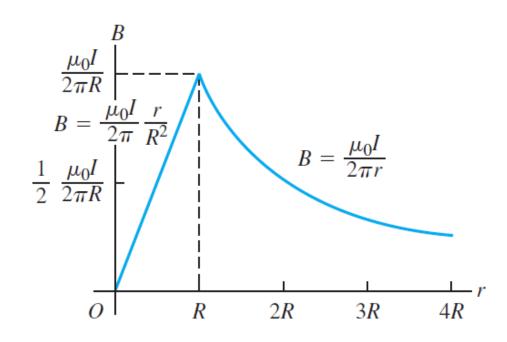
$$B = \frac{\mu_0 I}{2\pi} \frac{r}{R^2} \qquad \text{(for } r < R\text{)}$$



Example: Solid cylindrical conductor

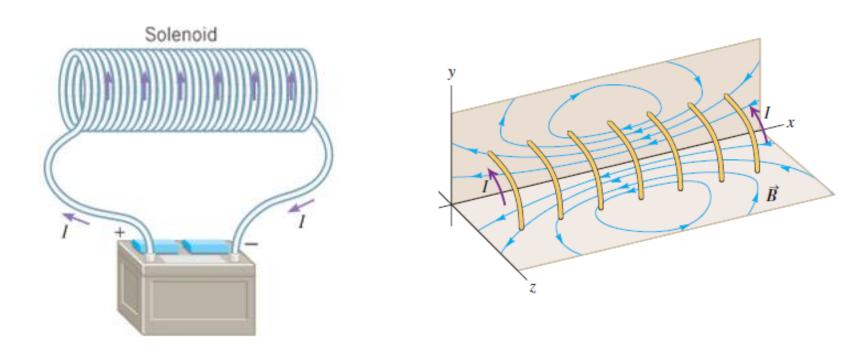
Solution

$$B = \begin{cases} \frac{\mu_0 I}{2\pi} \frac{r}{R^2} & \text{(for } r < R) \\ \frac{\mu_0 I}{2\pi r} & \text{(for } r > R) \end{cases} \qquad \frac{\frac{\mu_0 I}{2\pi R}}{\frac{1}{2\pi R}} \frac{R}{R^2}$$



Example: Ideal solenoid

N turns of wire around a cylinder of length L



n = N/L = number of turns per unit length

Example: Ideal solenoid

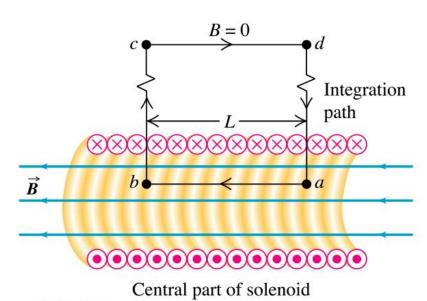
 Amperian loop is rectangular path abcd. (Symmetry requires B inside is uniform.)

$$\oint (\vec{B} \cdot d\vec{l}) = B \int_{a}^{b} dl = B(L)$$

$$= \mu_0 NI$$

$$B = \mu_0 \left(\frac{N}{L}\right) I = \mu_0 n I$$

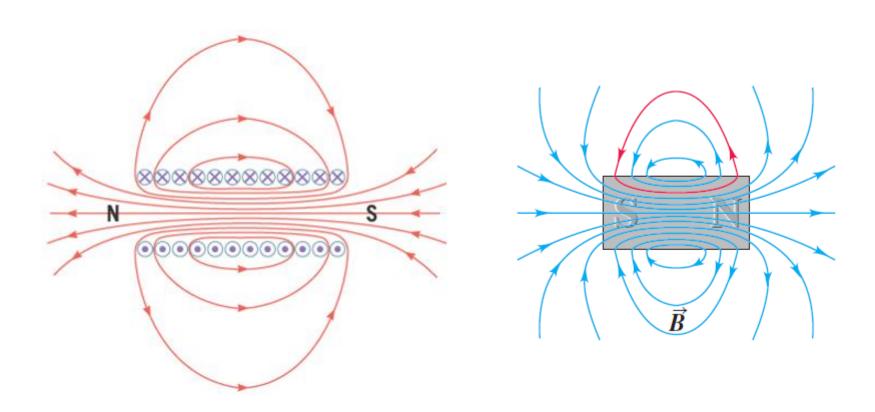
Magnetic field inside the solenoid is *uniform*!



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Example: Ideal solenoid

Real solenoid and bar magnet



Three general types of magnetism:

- Diamagnetism: A diamagnetic material placed in an external field B_{ext} develops a magnetic dipole moment directed opposite B_{ext}.
 - When the field is non-uniform, the material is *repelled* from a region of *greater* field toward region of *lesser* field.
- Paramagnetism: A paramagnetic material placed in an external field B_{ext} develops a magnetic dipole moment in the direction of B_{ext}.
 - If the field is non-uniform, the material is *attracted* toward a region of *greater* field from a region of *lesser* field.

Three general types of magnetism:

- Ferromagnetism: A ferromagnetic material placed in an external field B_{ext} develops a strong magnetic dipole moment in the direction of B_{ext}.
 - In a non-uniform field, the material is *attracted* toward a region of *greater* field from a region of *lesser* field.

Origins of magnetism:

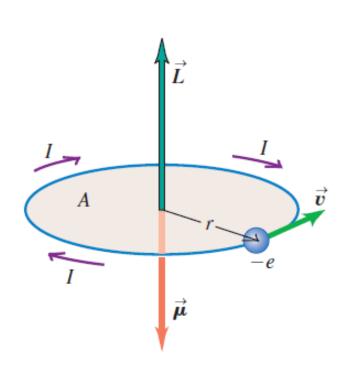
Orbital spin:

In the atom, an electron has an additional angular momentum \mathbf{L}_{orb} . Associated with \mathbf{L}_{orb} is an **orbital magnetic dipole moment** μ_{orb} .

$$\vec{\mu}_{\text{orb}} = -\frac{e}{2m} \vec{\mathbf{L}}_{\text{orb}}$$

Fundamental unit:

$$\mu_{orb} = \mu_B = 9.274 \times 10^{-24} \; \; \text{J/T}$$
 the **Bohr magneton**



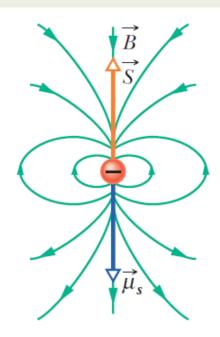
Origins of magnetism:

Intrinsic spin:

An electron has an *intrinsic* angular momentum (or spin) **S**. Associated with spin is an intrinsic **spin magnetic dipole moment** μ_s .

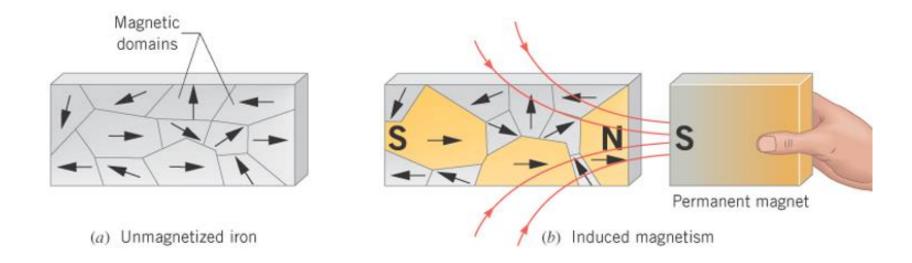
The *intrinsic "spin"* and *orbital motion* of electrons gives rise to the magnetic properties of materials.

For an electron, the spin is opposite the magnetic dipole moment.



Ferromagnetism

• In *ferromagnetic materials* (iron, nickel, cobalt) groups of neighboring atoms, forming *magnetic domains*, the spins of electrons are naturally aligned with each other.



END

