

Force to Torque - Natural Orthogonal Complement Mapping Analysis

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1 Newton-Euler analysis

The analysis is performed base on the *Recursive Newton-Euler algorithm* Algorithm 1 equations. The objective of this document is to find a relation between the reaction force at the foot ($[\mathbf{f}_F]_0$) and the 3-dimensional torque at the hip ($\boldsymbol{\tau}_H = [\tau_x \ \tau_y \ \tau_z]^T$).

Algorithm 1 Recursive Newton-Euler (Inverse dynamics)

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1: function IDNE( $model, \dot{\mathbf{q}}, \ddot{\mathbf{q}}, \mathbf{Q}|_1^N, \boldsymbol{\delta}|_1^N, \mathbf{f}, \mathbf{n}$ )
2:
3:    $[\mathbf{c}_0]_1 \leftarrow \mathbf{0}$ 
4:    $[\boldsymbol{\omega}_0]_1 \leftarrow \mathbf{0}$ 
5:    $[\dot{\mathbf{c}}_0]_1 \leftarrow \mathbf{0}$ 
6:    $[\dot{\boldsymbol{\omega}}_0]_1 \leftarrow \mathbf{0}$ 
7:    $[\ddot{\mathbf{c}}_0]_1 \leftarrow [-g]_1$ 
8:    $[\boldsymbol{\delta}_0]_1 \leftarrow \mathbf{0}$ 
9:   for  $i = 1 : N_B$  do
10:     $[\mathbf{c}_i]_{i+1} \leftarrow \mathbf{Q}_i^T [\mathbf{c}_{i-1} + \boldsymbol{\delta}_{i-1}]_i + [\boldsymbol{\rho}_i]_{i+1}$ 
11:     $[\boldsymbol{\omega}_i]_{i+1} \leftarrow \mathbf{Q}_i^T [\boldsymbol{\omega}_{i-1} + \theta_i \mathbf{e}_i]_i$ 
12:     $[\mathbf{u}_{i-1}]_i \leftarrow [\boldsymbol{\omega}_{i-1} \times \boldsymbol{\delta}_{i-1}]_i$ 
13:     $[\mathbf{v}_i]_{i+1} \leftarrow [\boldsymbol{\omega}_i \times \boldsymbol{\rho}_i]_{i+1}$ 
14:     $[\dot{\mathbf{c}}_i]_{i+1} \leftarrow \mathbf{Q}_i^T [\dot{\mathbf{c}}_{i-1} + \mathbf{u}_{i-1}]_i + [\mathbf{v}_i]_{i+1}$ 
15:     $[\dot{\boldsymbol{\omega}}_i]_{i+1} \leftarrow \mathbf{Q}_i^T [\dot{\boldsymbol{\omega}}_{i-1} + \boldsymbol{\omega}_{i-1} \times \theta_i \mathbf{e}_i + \ddot{\theta}_i \mathbf{e}_i]_i$ 
16:     $[\ddot{\mathbf{c}}_i]_{i+1} \leftarrow \mathbf{Q}_i^T [\ddot{\mathbf{c}}_{i-1} + \dot{\boldsymbol{\omega}}_{i-1} \times \boldsymbol{\delta}_{i-1} + \boldsymbol{\omega}_{i-1} \times \mathbf{u}_{i-1}]_i + [\dot{\boldsymbol{\omega}}_i \times \boldsymbol{\rho}_i + \boldsymbol{\omega}_i \times \mathbf{v}_i]_{i+1}$ 
17:  end for
18:
19:   $[\mathbf{f}_{N_B}^P]_{N_B+1} \leftarrow [m_{N_B} \ddot{\mathbf{c}}_{N_B} - \mathbf{f}]_{N_B+1}$ 
20:   $[\mathbf{n}_{N_B}^P]_{N_B+1} \leftarrow [\mathbf{I}_{N_B} \dot{\boldsymbol{\omega}}_{N_B} + \boldsymbol{\omega}_{N_B} \times \mathbf{I}_{N_B} \boldsymbol{\omega}_{N_B} - \mathbf{n} + \boldsymbol{\rho}_{N_B} \times \mathbf{f}_{N_B}^P]_{N_B+1}$ 
21:   $[\boldsymbol{\tau}_{N_B}]_{N_B} \leftarrow (\mathbf{Q}_{N_B} [\mathbf{n}_{N_B}^P]_{N_B+1})_z$ 
22:  for  $i = N_B - 1 : 1$  do
23:     $[\phi_{i+1}]_{i+1} \leftarrow \mathbf{Q}_{i+1} [\mathbf{f}_{i+1}^P]_{i+2}$ 
24:     $[\mathbf{f}_i^P]_{i+1} \leftarrow [m_i \ddot{\mathbf{c}}_i + \phi_{i+1}]_{i+1}$ 
25:     $[\mathbf{n}_i^P]_{i+1} \leftarrow [\mathbf{I}_i \dot{\boldsymbol{\omega}}_i + \boldsymbol{\omega}_i \times \mathbf{I}_i \boldsymbol{\omega}_i + \boldsymbol{\rho}_i \times \mathbf{f}_i^P + \boldsymbol{\delta}_i \times \phi_{i+1}]_{i+1} + \mathbf{Q}_{i+1} [\mathbf{n}_{i+1}^P]_{i+2}$ 
26:     $[\boldsymbol{\tau}_i]_i \leftarrow (\mathbf{Q}_i [\mathbf{n}_i^P]_{i+1})_z$ 
27:  end for
28: end function

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1.1 Dynamic Computations - Inward Recursion

Let

$$\mathbf{e} = [0 \ 0 \ 1]^T$$

This vector denotes the rotation orientation of each generalized coordinate. For the case of the Newton-Euler algorithm and the architecture of the spatial pendulum this orientation is always parallel to the Z axis of the body-coordinate frame in the *six bodies model*.

The *six bodies model* for the spatial pendulum consists in one inertial frame (F_0) and 6 body-coordinate frames

$(F_2 - F_7)$. Now we define the equations for the forces and moments using this model and Algorithm 1

$$\begin{aligned}
[\mathbf{f}_6]_7 &= m_6 \ddot{\mathbf{c}}_6 \\
[\mathbf{n}_6]_7 &= \mathbf{I}_6 \dot{\boldsymbol{\omega}}_6 + \boldsymbol{\omega}_6 \times \mathbf{I}_6 \boldsymbol{\omega}_6 + \boldsymbol{\rho}_6 \times \mathbf{f}_6 \\
\boldsymbol{\tau}_6 &= \mathbf{e}^T \mathbf{Q}_6 \\
[\mathbf{f}_5]_6 &= \mathbf{Q}_6 [\mathbf{f}_6]_7 \\
[\mathbf{n}_5]_6 &= \mathbf{Q}_6 [\mathbf{n}_6]_7 \\
\boldsymbol{\tau}_5 &= \mathbf{e}^T \mathbf{Q}_5 [\mathbf{n}_5]_6 = \mathbf{e}^T \mathbf{Q}_5 \mathbf{Q}_6 [\mathbf{n}_6]_7 \\
[\mathbf{f}_4]_5 &= \mathbf{Q}_5 [\mathbf{f}_5]_6 = \mathbf{Q}_5 \mathbf{Q}_6 [\mathbf{f}_6]_7 \\
[\mathbf{n}_4]_5 &= \mathbf{Q}_5 [\mathbf{n}_5]_6 = \mathbf{Q}_5 \mathbf{Q}_6 [\mathbf{n}_6]_7 \\
\boldsymbol{\tau}_4 &= \mathbf{e}^T \mathbf{Q}_4 [\mathbf{n}_4]_5 = \mathbf{e}^T \mathbf{Q}_4 \mathbf{Q}_5 \mathbf{Q}_6 [\mathbf{n}_6]_7 \\
[\mathbf{f}_3]_4 &= m_3 \ddot{\mathbf{c}}_3 + \mathbf{Q}_4 [\mathbf{f}_4]_5 \\
&= m_3 \ddot{\mathbf{c}}_3 + \mathbf{Q}_4 \mathbf{Q}_5 \mathbf{Q}_6 [\mathbf{f}_6]_7 \\
[\mathbf{n}_3]_4 &= \mathbf{I}_3 \dot{\boldsymbol{\omega}}_3 + \boldsymbol{\omega}_3 \times \mathbf{I}_3 \boldsymbol{\omega}_3 + \boldsymbol{\rho}_3 \times \mathbf{f}_3 + \delta_3 \times \mathbf{Q}_4 [\mathbf{f}_4]_5 + \mathbf{Q}_4 [\mathbf{n}_4]_5 \\
&= \mathbf{I}_3 \dot{\boldsymbol{\omega}}_3 + \boldsymbol{\omega}_3 \times \mathbf{I}_3 \boldsymbol{\omega}_3 + \boldsymbol{\rho}_3 \times \mathbf{f}_3 + \delta_3 \times \mathbf{Q}_4 \mathbf{Q}_5 \mathbf{Q}_6 [\mathbf{f}_6]_7 + \mathbf{Q}_4 \mathbf{Q}_5 \mathbf{Q}_6 [\mathbf{n}_6]_7 \\
\boldsymbol{\tau}_3 &= \mathbf{e}^T \mathbf{Q}_3 [\mathbf{n}_3]_4 \\
[\mathbf{f}_2]_3 &= \mathbf{Q}_3 [\mathbf{f}_3]_4 \\
[\mathbf{n}_2]_3 &= \mathbf{Q}_3 [\mathbf{n}_3]_4 \\
\boldsymbol{\tau}_2 &= \mathbf{e}^T \mathbf{Q}_2 [\mathbf{n}_2]_3 = \mathbf{e}^T \mathbf{Q}_2 \mathbf{Q}_3 [\mathbf{n}_3]_4 \\
[\mathbf{f}_1]_2 &= \mathbf{Q}_2 [\mathbf{f}_2]_3 = \mathbf{Q}_2 \mathbf{Q}_3 [\mathbf{f}_3]_4 \\
[\mathbf{n}_1]_2 &= \mathbf{Q}_2 [\mathbf{n}_2]_3 = \mathbf{Q}_2 \mathbf{Q}_3 [\mathbf{n}_3]_4 \\
\boldsymbol{\tau}_1 &= \mathbf{e}^T \mathbf{Q}_1 [\mathbf{n}_1]_2 = \mathbf{e}^T \mathbf{Q}_1 \mathbf{Q}_2 \mathbf{Q}_3 [\mathbf{n}_3]_4
\end{aligned}$$

Note that \mathbf{I}_3 and \mathbf{I}_6 denote 3×3 inertia matrices corresponding to the lower and upper link, respectively, and should not be confused with the generalized inertia matrix $\mathbf{I}(\mathbf{q})$. From these equations we conclude the system can be reduced from 7 coordinate frames to 3 coordinate frames: inertial frame (F_0) and 2 body-coordinate frames (F_1 and F_2)

$$[\mathbf{f}_H]_2 = m_2 \ddot{\mathbf{c}}_2 \quad (1)$$

$$[\mathbf{n}_H]_2 = \mathbf{I}_2 \dot{\boldsymbol{\omega}}_H + \boldsymbol{\omega}_H \times \mathbf{I}_2 \boldsymbol{\omega}_H + \boldsymbol{\rho}_2 \times \mathbf{f}_H \quad (2)$$

$$[\mathbf{f}_F]_1 = m_1 \ddot{\mathbf{c}}_1 + \mathbf{R}_2 [\mathbf{f}_H]_2 \quad (3)$$

$$[\mathbf{n}_F]_1 = \mathbf{I}_1 \dot{\boldsymbol{\omega}}_F + \boldsymbol{\omega}_F \times \mathbf{I}_1 \boldsymbol{\omega}_F + \boldsymbol{\rho}_1 \times \mathbf{f}_F + \delta_1 \times \mathbf{R}_2 [\mathbf{f}_H]_2 + \mathbf{R}_2 [\mathbf{n}_H]_2 \quad (4)$$

and the expressions for $\boldsymbol{\tau}_F$ and $\boldsymbol{\tau}_H$ becomes

$$\boldsymbol{\tau}_F = \begin{bmatrix} \tau_{Fx} \\ \tau_{Fy} \\ \tau_{Fz} \end{bmatrix} = \begin{bmatrix} \mathbf{e}^T \mathbf{R}_1 \\ \mathbf{e}^T \mathbf{Q}_2 \mathbf{Q}_3 \\ \mathbf{e}^T \mathbf{Q}_3 \end{bmatrix} [\mathbf{n}_F]_1 \quad ; \quad \boldsymbol{\tau}_H = \begin{bmatrix} \tau_{Hx} \\ \tau_{Hy} \\ \tau_{Hz} \end{bmatrix} = \begin{bmatrix} \mathbf{e}^T \mathbf{R}_2 \\ \mathbf{e}^T \mathbf{Q}_5 \mathbf{Q}_6 \\ \mathbf{e}^T \mathbf{Q}_6 \end{bmatrix} [\mathbf{n}_H]_2 \quad (5)$$

where

$$\mathbf{R}_1 = \mathbf{Q}_1 \mathbf{Q}_2 \mathbf{Q}_3 \quad ; \quad \mathbf{R}_2 = \mathbf{Q}_4 \mathbf{Q}_5 \mathbf{Q}_6 \quad (6)$$

are rotation matrices that establish the following relation between coordinate frames

$$[\mathbf{f}_F]_2 = \mathbf{R}_2^T [\mathbf{f}_F]_1 \quad ; \quad [\mathbf{f}_F]_1 = \mathbf{R}_1^T [\mathbf{f}_F]_0 \quad (7)$$

Note that the notation $[\bullet]_i$ is omitted in those vectors and matrices expressed in the same frame that the vector on the LHS of the equation. The subscripts H and F denotes *hip* and *foot* respectively, and they are used to identify the vectors related to the first spherical joint at the foot or the second at the hip.

1.2 Kinematic Computations - Outwards Recursion

Find an expression for $\dot{\omega}_3$ and \ddot{e}_3 using the *six bodies model* and Algorithm 1 equations

$$\begin{aligned}
[\dot{\omega}_1]_2 &= Q_1^T [\ddot{q}_1 e]_1 \\
[\ddot{e}_1]_2 &= Q_1^T [\ddot{e}_0]_1 \quad \text{where} \quad \ddot{e}_0 = g \\
[\dot{\omega}_2]_3 &= Q_2^T [\dot{\omega}_1 + \omega_1 \times \dot{q}_2 e + \ddot{q}_2 e]_2 \\
&= (Q_1 Q_2)^T [\ddot{q}_1 e]_1 + Q_2^T [\omega_1 \times \dot{q}_2 e]_2 + Q_2^T [\ddot{q}_2 e]_2 \\
[\ddot{e}_2]_3 &= Q_2^T [\ddot{e}_1]_2 = (Q_1 Q_2)^T [\ddot{e}_0]_1 \\
[\dot{\omega}_3]_4 &= Q_3^T [\dot{\omega}_2 + \omega_2 \times \dot{q}_3 e + \ddot{q}_3 e]_3 \\
&= Q_3^T [\omega_2 \times \dot{q}_3 e]_3 + (Q_2 Q_3)^T [\omega_1 \times \dot{q}_2 e]_2 + (Q_1 Q_2 Q_3)^T [\ddot{q}_1 e]_1 + (Q_2 Q_3)^T [\ddot{q}_2 e]_2 + Q_3^T [\ddot{q}_3 e]_3 \\
[\ddot{e}_3]_4 &= Q_3^T [\ddot{e}_2]_3 + \dot{\omega}_3 \times \rho_3 + \omega_3 \times (\omega_3 \times \rho_3) \\
&= (Q_1 Q_2 Q_3)^T [\ddot{e}_0]_1 + \omega_3 \times (\omega_3 \times \rho_3) - \rho_3 \times \dot{\omega}_3 \\
[\dot{\omega}_4]_5 &= Q_4^T [\dot{\omega}_3 + \omega_3 \times \dot{q}_4 e + \ddot{q}_4 e]_4 \\
[\ddot{e}_4]_5 &= Q_4^T [\ddot{e}_3 + \dot{\omega}_3 \times \delta_3 + \omega_3 \times (\omega_3 \times \delta_3)]_4 \\
[\dot{\omega}_5]_6 &= Q_5^T [\dot{\omega}_4 + \omega_4 \times \dot{q}_5 e + \ddot{q}_5 e]_5 \\
&= (Q_4 Q_5)^T [\dot{\omega}_3]_4 + (Q_4 Q_5)^T [\omega_3 \times \dot{q}_4 e]_4 + (Q_4 Q_5)^T [\ddot{q}_4 e]_4 + Q_5^T [\omega_4 \times \dot{q}_5 e]_5 + Q_5^T [\ddot{q}_5 e]_5 \\
[\ddot{e}_5]_6 &= (Q_4 Q_5)^T [\ddot{e}_3 + \dot{\omega}_3 \times \delta_3 + \omega_3 \times (\omega_3 \times \delta_3)]_4 \\
[\dot{\omega}_6]_7 &= Q_6^T [\dot{\omega}_5 + \omega_5 \times \dot{q}_6 e + \ddot{q}_6 e]_6 \\
&= (Q_4 Q_5 Q_6)^T [\dot{\omega}_3]_4 + (Q_4 Q_5 Q_6)^T [\omega_3 \times \dot{q}_4 e]_4 + (Q_4 Q_5 Q_6)^T [\ddot{q}_4 e]_4 + (Q_5 Q_6)^T [\omega_4 \times \dot{q}_5 e]_5 + (Q_5 Q_6)^T [\ddot{q}_5 e]_5 \\
&\quad + Q_6^T [\omega_5 \times \dot{q}_6 e]_6 + Q_6^T [\ddot{q}_6 e]_6 \\
[\ddot{e}_6]_7 &= (Q_4 Q_5 Q_6)^T [\ddot{e}_3 + \dot{\omega}_3 \times \delta_3 + \omega_3 \times (\omega_3 \times \delta_3)]_4 + \dot{\omega}_6 \times \rho_6 + \omega_6 \times (\omega_6 \times \rho_6)
\end{aligned}$$

Then, the system can be reduced to 3 coordinate frames: 2 body-coordinate frames (F_1, F_2) and an inertial frame (F_0), given by

$$[\dot{\omega}_F]_1 = k_F + A_F \ddot{q}_F \quad (8)$$

$$[\ddot{e}_1]_1 = k_1 - CPM(\rho_1) \dot{\omega}_F \quad (9)$$

$$[\dot{\omega}_H]_2 = R_2^T [\dot{\omega}_F]_1 + k_H + A_H \ddot{q}_H \quad (10)$$

$$[\ddot{e}_2]_2 = k_2 + R_2^T [\ddot{e}_1 - CPM(\delta_1) \dot{\omega}_F]_1 - CPM(\rho_2) \dot{\omega}_H \quad (11)$$

where

$$\begin{aligned}
k_F(q, \dot{q}) &= Q_3^T [\omega_2 \times \dot{q}_3 e]_3 + (Q_2 Q_3)^T [\omega_1 \times \dot{q}_2 e]_2 \\
k_1(q, \dot{q}, g) &= R_1^T [\ddot{e}_0]_0 + \omega_F \times (\omega_F \times \rho_1) \\
k_H(q, \dot{q}) &= (Q_4 Q_5 Q_6)^T [\omega_3 \times \dot{q}_4 e]_4 + (Q_5 Q_6)^T [\omega_4 \times \dot{q}_5 e]_5 + Q_6^T [\omega_5 \times \dot{q}_6 e]_6 \\
k_2(q, \dot{q}) &= R_2^T [\omega_F \times (\omega_F \times \delta_1)]_1 + \omega_H \times (\omega_H \times \rho_2) \\
A_F(q) &= [R_1^T e \quad (Q_2 Q_3)^T e \quad Q_3^T e] \quad ; \quad A_H(q) = [R_2^T e \quad (Q_5 Q_6)^T \quad Q_6^T e] \\
\rho \times &= CPM(\rho) = \begin{bmatrix} 0 & -\rho_z & \rho_y \\ \rho_z & 0 & -\rho_x \\ -\rho_y & \rho_x & 0 \end{bmatrix} \quad \text{given that} \quad \rho = \begin{bmatrix} \rho_x \\ \rho_y \\ \rho_z \end{bmatrix}
\end{aligned}$$

Applying Eq. (8) in Eq. (9) gives

$$\begin{aligned}
[\ddot{e}_1]_1 &= k_1 - CPM(\rho_1) [k_F + A_F \ddot{q}_F] \\
&= k_3 - CPM(\rho_1) A_F \ddot{q}_F
\end{aligned} \quad (12)$$

Substituting Eq. (8) in Eq. (10) yields

$$\begin{aligned} [\dot{\omega}_H]_2 &= \mathbf{R}_2^T [\mathbf{k}_F + \mathbf{A}_F \ddot{\mathbf{q}}_F]_1 + \mathbf{k}_H + \mathbf{A}_H \ddot{\mathbf{q}}_H \\ &= \mathbf{k}_{H2} + \mathbf{R}_2^T \mathbf{A}_F \ddot{\mathbf{q}}_F + \mathbf{A}_H \ddot{\mathbf{q}}_H \end{aligned} \quad (13)$$

Applying this equation to Eq. (11) gives

$$[\ddot{\mathbf{c}}_2]_2 = \mathbf{k}_2 + \mathbf{R}_2^T [\ddot{\mathbf{c}}_1 - CPM(\delta_1) \dot{\omega}_F]_1 - CPM(\rho_2) [\mathbf{k}_{H2} + \mathbf{R}_2^T \mathbf{A}_F \ddot{\mathbf{q}}_F + \mathbf{A}_H \ddot{\mathbf{q}}_H] \quad (14)$$

Substituting Eq. (9) and Eq. (8) yields

$$\begin{aligned} [\ddot{\mathbf{c}}_2]_2 &= \mathbf{k}_2 + \mathbf{R}_2^T [\mathbf{k}_1 - CPM(\rho_1) \dot{\omega}_F - CPM(\delta_1) \dot{\omega}_F]_1 - CPM(\rho_2) [\mathbf{k}_{H2} + \mathbf{R}_2^T \mathbf{A}_F \ddot{\mathbf{q}}_F + \mathbf{A}_H \ddot{\mathbf{q}}_H] \\ &= \mathbf{k}_2 + \mathbf{R}_2^T [\mathbf{k}_1 - CPM(\rho_1 + \delta_1) \dot{\omega}_F]_1 - CPM(\rho_2) [\mathbf{k}_{H2} + \mathbf{R}_2^T \mathbf{A}_F \ddot{\mathbf{q}}_F + \mathbf{A}_H \ddot{\mathbf{q}}_H] \\ &= \mathbf{k}_2 + \mathbf{R}_2^T [\mathbf{k}_1 - CPM(\rho_1 + \delta_1) (\mathbf{k}_F + \mathbf{A}_F \ddot{\mathbf{q}}_F)]_1 - CPM(\rho_2) [\mathbf{k}_{H2} + \mathbf{R}_2^T \mathbf{A}_F \ddot{\mathbf{q}}_F + \mathbf{A}_H \ddot{\mathbf{q}}_H] \\ &= \mathbf{k}_4 - \mathbf{R}_2^T CPM(\rho_1 + \delta_1) \mathbf{A}_F \ddot{\mathbf{q}}_F - CPM(\rho_2) [\mathbf{R}_2^T \mathbf{A}_F \ddot{\mathbf{q}}_F + \mathbf{A}_H \ddot{\mathbf{q}}_H] \\ &= \mathbf{k}_4 - \left[\mathbf{R}_2^T CPM(\rho_1 + \delta_1) + CPM(\rho_2) \mathbf{R}_2^T \right] \mathbf{A}_F \ddot{\mathbf{q}}_F - CPM(\rho_2) \mathbf{A}_H \ddot{\mathbf{q}}_H \end{aligned} \quad (15)$$

Recall the expression for \mathbf{f}_F , Eq. (3)

$$\begin{aligned} [\mathbf{f}_F]_1 &= m_1 \ddot{\mathbf{c}}_1 + \mathbf{R}_2 [\mathbf{f}_H]_2 \\ &= m_1 \ddot{\mathbf{c}}_1 + m_2 \mathbf{R}_2 [\ddot{\mathbf{c}}_2]_2 \end{aligned} \quad (16)$$

Expressed at the inertia frame becomes

$$[\mathbf{f}_F]_0 = m_1 \mathbf{R}_1 [\ddot{\mathbf{c}}_1]_1 + m_2 \mathbf{R}_1 \mathbf{R}_2 [\ddot{\mathbf{c}}_2]_2 \quad (17)$$

Substituting Eq. (12) yields

$$\begin{aligned} [\mathbf{f}_F]_0 &= m_1 \mathbf{R}_1 [\mathbf{k}_3 - CPM(\rho_1) \mathbf{A}_F \ddot{\mathbf{q}}_F]_1 + m_2 \mathbf{R}_1 \mathbf{R}_2 [\mathbf{k}_4 - \left[\mathbf{R}_2^T CPM(\rho_1 + \delta_1) + CPM(\rho_2) \mathbf{R}_2^T \right] \mathbf{A}_F \ddot{\mathbf{q}}_F \\ &\quad - CPM(\rho_2) \mathbf{A}_H \ddot{\mathbf{q}}_H]_2 \\ &= \mathbf{k}_5 + \mathbf{F} \mathbf{A}_F \ddot{\mathbf{q}}_F + \mathbf{G} \mathbf{A}_H \ddot{\mathbf{q}}_H \end{aligned} \quad (18)$$

where

$$\begin{aligned} \mathbf{F}(\mathbf{q}) &= - \left[m_1 \mathbf{R}_1 CPM(\rho_1) + m_2 \mathbf{R}_1 CPM(\rho_1 + \delta_1) + m_2 \mathbf{R}_1 \mathbf{R}_2 CPM(\rho_2) \mathbf{R}_2^T \right] \\ \mathbf{G}(\mathbf{q}) &= - m_2 \mathbf{R}_1 \mathbf{R}_2 CPM(\rho_2) \end{aligned}$$

In the next section we find an expression for $\ddot{\mathbf{q}}_F$ and $\ddot{\mathbf{q}}_H$ in terms of τ_H

1.3 Equation of motion of the spatial pendulum

The equation of motion of the spatial pendulum is

$$\mathbf{I}(\mathbf{q}) \ddot{\mathbf{q}} + \mathbf{C}(\mathbf{q}, \dot{\mathbf{q}}) \dot{\mathbf{q}} - \gamma(\mathbf{q}) = \tau \quad (19)$$

$$\begin{bmatrix} \mathbf{I}_{11} & \mathbf{I}_{12} \\ \mathbf{I}_{21} & \mathbf{I}_{22} \end{bmatrix} \begin{bmatrix} \ddot{\mathbf{q}}_F \\ \ddot{\mathbf{q}}_H \end{bmatrix} + \begin{bmatrix} \mathbf{C}_{11} & \mathbf{C}_{12} \\ \mathbf{C}_{21} & \mathbf{C}_{22} \end{bmatrix} \begin{bmatrix} \dot{\mathbf{q}}_F \\ \dot{\mathbf{q}}_H \end{bmatrix} - \begin{bmatrix} \gamma_F \\ \gamma_H \end{bmatrix} = \begin{bmatrix} \tau_F \\ \tau_H \end{bmatrix} \quad (20)$$

where $\mathbf{I}(\mathbf{q})$ is the 6×6 generalized inertia matrix and $\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}})$ is the Coriolis coefficient matrix. We are looking for an expression for $\ddot{\mathbf{q}}_F$ without considering gravity and only one active joint at the hip. Thus, the equation of motion becomes

$$\begin{bmatrix} \mathbf{I}_{11} & \mathbf{I}_{12} \\ \mathbf{I}_{21} & \mathbf{I}_{22} \end{bmatrix} \begin{bmatrix} \ddot{\mathbf{q}}_F \\ \ddot{\mathbf{q}}_H \end{bmatrix} + \begin{bmatrix} \mathbf{C}_{11} & \mathbf{C}_{12} \\ \mathbf{C}_{21} & \mathbf{C}_{22} \end{bmatrix} \begin{bmatrix} \dot{\mathbf{q}}_F \\ \dot{\mathbf{q}}_H \end{bmatrix} = \begin{bmatrix} \mathbf{0} \\ \tau_H \end{bmatrix} \quad (21)$$

So

$$\ddot{\mathbf{q}}_F = \bar{\mathbf{I}}_{12} \tau_H + \mathbf{k}_{MF} \quad (22)$$

$$\ddot{\mathbf{q}}_H = \bar{\mathbf{I}}_{22} \tau_H + \mathbf{k}_{MH} \quad (23)$$

where

$$\begin{aligned}\bar{\mathbf{I}}_{12}(\mathbf{q}) &= \begin{bmatrix} \mathbf{1} & \mathbf{0} \end{bmatrix} \mathbf{I}(\mathbf{q})^{-1} \begin{bmatrix} \mathbf{0} \\ \mathbf{1} \end{bmatrix} \quad ; \quad \mathbf{k}_{MF}(\mathbf{q}, \dot{\mathbf{q}}) = - \begin{bmatrix} \mathbf{1} & \mathbf{0} \end{bmatrix} \mathbf{I}(\mathbf{q})^{-1} \mathbf{C}(\mathbf{q}, \dot{\mathbf{q}}) \dot{\mathbf{q}} \\ \bar{\mathbf{I}}_{22}(\mathbf{q}) &= \begin{bmatrix} \mathbf{0} & \mathbf{1} \end{bmatrix} \mathbf{I}(\mathbf{q})^{-1} \begin{bmatrix} \mathbf{0} \\ \mathbf{1} \end{bmatrix} \quad ; \quad \mathbf{k}_{MH}(\mathbf{q}, \dot{\mathbf{q}}) = - \begin{bmatrix} \mathbf{0} & \mathbf{1} \end{bmatrix} \mathbf{I}(\mathbf{q})^{-1} \mathbf{C}(\mathbf{q}, \dot{\mathbf{q}}) \dot{\mathbf{q}}\end{aligned}$$

Substituting the result in Eq. (19) yields

$$\begin{aligned}[\mathbf{f}_F]_0 &= \mathbf{k}_5 + \mathbf{F} \mathbf{A}_F [\bar{\mathbf{I}}_{12} \boldsymbol{\tau}_H + \mathbf{k}_{MF}] + \mathbf{G} \mathbf{A}_H [\bar{\mathbf{I}}_{22} \boldsymbol{\tau}_H + \mathbf{k}_{MH}] \\ &= \mathbf{b} + [\mathbf{F} \mathbf{A}_F \bar{\mathbf{I}}_{12} + \mathbf{G} \mathbf{A}_H \bar{\mathbf{I}}_{22}] \boldsymbol{\tau}_H \\ &= \mathbf{b} + \mathbf{M} \boldsymbol{\tau}_H\end{aligned}\tag{24}$$

2 Summary

2.1 Reduced Newton-Euler equations

$$[\mathbf{f}_H]_2 = m_2 \ddot{\mathbf{c}}_2 \tag{25}$$

$$[\mathbf{n}_H]_2 = \mathbf{I}_2 \dot{\boldsymbol{\omega}}_H + \boldsymbol{\omega}_H \times \mathbf{I}_2 \boldsymbol{\omega}_H + \boldsymbol{\rho}_2 \times \mathbf{f}_H \tag{26}$$

$$[\mathbf{f}_F]_1 = m_1 \ddot{\mathbf{c}}_1 + \mathbf{R}_2 [\mathbf{f}_H]_2 \tag{27}$$

$$[\mathbf{n}_F]_1 = \mathbf{I}_1 \dot{\boldsymbol{\omega}}_F + \boldsymbol{\omega}_F \times \mathbf{I}_1 \boldsymbol{\omega}_F + \boldsymbol{\rho}_1 \times \mathbf{f}_F + \boldsymbol{\delta}_1 \times \mathbf{R}_2 [\mathbf{f}_H]_2 + \mathbf{R}_2 [\mathbf{n}_H]_2 \tag{28}$$

$$[\dot{\boldsymbol{\omega}}_F]_1 = \mathbf{k}_F + \mathbf{A}_F \ddot{\mathbf{q}}_F \tag{29}$$

$$[\ddot{\mathbf{c}}_1]_1 = \mathbf{k}_1 - \text{CPM}(\boldsymbol{\rho}_1) \dot{\boldsymbol{\omega}}_F \tag{30}$$

$$[\dot{\boldsymbol{\omega}}_H]_2 = \mathbf{R}_2^T [\dot{\boldsymbol{\omega}}_F]_1 + \mathbf{k}_H + \mathbf{A}_H \ddot{\mathbf{q}}_H \tag{31}$$

$$[\ddot{\mathbf{c}}_2]_2 = \mathbf{k}_2 + \mathbf{R}_2^T [\ddot{\mathbf{c}}_1 - \text{CPM}(\boldsymbol{\delta}_1) \dot{\boldsymbol{\omega}}_F]_1 - \text{CPM}(\boldsymbol{\rho}_2) \dot{\boldsymbol{\omega}}_H \tag{32}$$

where

$$\begin{aligned}\mathbf{k}_F(\mathbf{q}, \dot{\mathbf{q}}) &= \mathbf{Q}_3^T [\boldsymbol{\omega}_2 \times \dot{\mathbf{q}}_3 \mathbf{e}]_3 + (\mathbf{Q}_2 \mathbf{Q}_3)^T [\boldsymbol{\omega}_1 \times \dot{\mathbf{q}}_2 \mathbf{e}]_2 \\ \mathbf{k}_1(\mathbf{q}, \dot{\mathbf{q}}, g) &= \mathbf{R}_1^T [\ddot{\mathbf{c}}_0]_0 + \boldsymbol{\omega}_F \times (\boldsymbol{\omega}_F \times \boldsymbol{\rho}_1) \\ \mathbf{k}_H(\mathbf{q}, \dot{\mathbf{q}}) &= (\mathbf{Q}_4 \mathbf{Q}_5 \mathbf{Q}_6)^T [\boldsymbol{\omega}_3 \times \dot{\mathbf{q}}_4 \mathbf{e}]_4 + (\mathbf{Q}_5 \mathbf{Q}_6)^T [\boldsymbol{\omega}_4 \times \dot{\mathbf{q}}_5 \mathbf{e}]_5 + \mathbf{Q}_6^T [\boldsymbol{\omega}_5 \times \dot{\mathbf{q}}_6 \mathbf{e}]_6 \\ \mathbf{k}_2(\mathbf{q}, \dot{\mathbf{q}}) &= \mathbf{R}_2^T [\boldsymbol{\omega}_F \times (\boldsymbol{\omega}_F \times \boldsymbol{\delta}_1)]_1 + \boldsymbol{\omega}_H \times (\boldsymbol{\omega}_H \times \boldsymbol{\rho}_2) \\ \mathbf{A}_F(\mathbf{q}) &= [\mathbf{R}_1^T \mathbf{e} \quad (\mathbf{Q}_2 \mathbf{Q}_3)^T \mathbf{e} \quad \mathbf{Q}_3^T \mathbf{e}] \quad ; \quad \mathbf{A}_H(\mathbf{q}) = [\mathbf{R}_2^T \mathbf{e} \quad (\mathbf{Q}_5 \mathbf{Q}_6)^T \quad \mathbf{Q}_6^T \mathbf{e}]\end{aligned}$$

and

$$\boldsymbol{\tau}_H = \begin{bmatrix} \tau_{Hx} \\ \tau_{Hy} \\ \tau_{Hz} \end{bmatrix} = \begin{bmatrix} \mathbf{e}^T \mathbf{R}_2 \\ \mathbf{e}^T \mathbf{Q}_5 \mathbf{Q}_6 \\ \mathbf{e}^T \mathbf{Q}_6 \end{bmatrix} [\mathbf{n}_H]_2 \tag{33}$$

2.2 Mapping

$$\begin{aligned}\mathbf{f}_F &= \mathbf{b} + [\mathbf{F} \mathbf{A}_F \bar{\mathbf{I}}_{12} + \mathbf{G} \mathbf{A}_H \bar{\mathbf{I}}_{22}] \boldsymbol{\tau}_H \\ &= \mathbf{b} + \mathbf{M} \boldsymbol{\tau}_H\end{aligned}\tag{34}$$

where

$$\mathbf{b} = \mathbf{f}_0$$

meaning \mathbf{b} is the reaction force vector at the foot when $\boldsymbol{\tau}_H = \mathbf{0}$. The expression for matrices \mathbf{F} , \mathbf{A}_F , \mathbf{G} and \mathbf{A}_H is given in next section.

3 Coefficients analysis

$$\mathbf{A}_F(\mathbf{q}_F) = [\mathbf{R}_1^T \mathbf{e} \quad (\mathbf{Q}_2 \mathbf{Q}_3)^T \mathbf{e} \quad \mathbf{Q}_3^T \mathbf{e}] = \begin{bmatrix} \cos \bar{q}_3 \sin \bar{q}_2 & \sin \bar{q}_3 & 0 \\ -\sin \bar{q}_2 \sin \bar{q}_3 & \cos \bar{q}_3 & 0 \\ -\cos \bar{q}_2 & 0 & 1 \end{bmatrix} \quad (35)$$

$$\mathbf{A}_H(\mathbf{q}_H) = [\mathbf{R}_2^T \mathbf{e} \quad (\mathbf{Q}_5 \mathbf{Q}_6)^T \mathbf{e} \quad \mathbf{Q}_6^T \mathbf{e}] = \begin{bmatrix} \cos \bar{q}_6 \sin \bar{q}_5 & \sin \bar{q}_6 & 0 \\ -\sin \bar{q}_5 \sin \bar{q}_6 & \cos \bar{q}_6 & 0 \\ -\cos \bar{q}_5 & 0 & 1 \end{bmatrix} \quad (36)$$

$$\mathbf{F} \mathbf{A}_F = -\left[m_1 \mathbf{R}_1 \text{CPM}(\boldsymbol{\rho}_1) - m_2 \mathbf{R}_1 \text{CPM}(\boldsymbol{\rho}_1 + \boldsymbol{\delta}_1) - m_2 \mathbf{R}_1 \mathbf{R}_2 \text{CPM}(\boldsymbol{\rho}_2) \mathbf{R}_2^T \right] \mathbf{A}_F = \begin{bmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ 0 & a_{32} & a_{33} \end{bmatrix} \quad (37)$$

$$\begin{aligned} a_{11} &= \beta_1 \left[\cos \bar{q}_1 \sin \bar{q}_3 - \sin \bar{q}_1 \cos \bar{q}_2 \cos \bar{q}_3 \right] + \beta_2 \left[\cos \bar{q}_1 \sin(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 - \sin \bar{q}_1 \cos \bar{q}_2 \cos(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 \right] \\ a_{12} &= -\cos \bar{q}_1 \left[\beta_1 \cos \bar{q}_3 \sin \bar{q}_2 + \beta_2 \left(\cos \bar{q}_2 \cos \bar{q}_5 + \sin \bar{q}_2 \cos(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 \right) \right] \\ a_{13} &= \beta_1 \left[\sin \bar{q}_1 \cos \bar{q}_3 - \cos \bar{q}_1 \cos \bar{q}_2 \sin \bar{q}_3 \right] + \beta_2 \left[\sin \bar{q}_1 \cos(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 - \cos \bar{q}_1 \cos \bar{q}_2 \sin(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 \right] \\ a_{21} &= \beta_1 \left[\sin \bar{q}_1 \sin \bar{q}_3 + \cos \bar{q}_1 \cos \bar{q}_2 \cos \bar{q}_3 \right] + \beta_2 \left[\sin \bar{q}_1 \sin(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 + \cos \bar{q}_1 \cos \bar{q}_2 \cos(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 - \cos \bar{q}_1 \sin \bar{q}_2 \cos \bar{q}_5 \right] \\ a_{22} &= -\sin \bar{q}_1 \left[\beta_1 \sin \bar{q}_2 \cos \bar{q}_3 + \beta_2 \left(\cos \bar{q}_2 \cos \bar{q}_5 + \sin \bar{q}_2 \cos(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 \right) \right] \\ a_{23} &= -\beta_1 \left[\cos \bar{q}_1 \cos \bar{q}_3 + \sin \bar{q}_1 \cos \bar{q}_2 \sin \bar{q}_3 \right] - \beta_2 \left[\cos \bar{q}_1 \cos(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 + \sin \bar{q}_1 \cos \bar{q}_2 \sin(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 \right] \\ a_{31} &= 0 \\ a_{32} &= \beta_1 \cos \bar{q}_2 \cos \bar{q}_3 + \beta_2 \left[\cos \bar{q}_2 \cos(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 - \sin \bar{q}_2 \cos \bar{q}_5 \right] \\ a_{33} &= -\sin \bar{q}_2 \left[\beta_1 \sin \bar{q}_3 + \beta_2 \sin(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 \right] \end{aligned}$$

where

$$\begin{aligned} \beta_1 &= m_1 \frac{l_1}{2} + m_2 l_1 \\ \beta_2 &= m_2 \frac{l_2}{2} \end{aligned}$$

$$\mathbf{G} \mathbf{A}_H = -\left[m_2 \mathbf{R}_1 \mathbf{R}_2 \text{CPM}(\boldsymbol{\rho}_2) \right] \mathbf{A}_H = \begin{bmatrix} b_{11} & b_{12} & 0 \\ b_{21} & b_{22} & 0 \\ b_{31} & b_{32} & 0 \end{bmatrix} \quad (38)$$

$$\begin{aligned} b_{11} &= \beta_2 \sin \bar{q}_5 \left[\sin \bar{q}_1 \cos(\bar{q}_3 + \bar{q}_4) - \cos \bar{q}_1 \cos \bar{q}_2 \sin(\bar{q}_3 + \bar{q}_4) \right] \\ b_{12} &= \beta_2 \left[\cos \bar{q}_1 \sin \bar{q}_2 \sin \bar{q}_5 + \sin \bar{q}_1 \sin(\bar{q}_3 + \bar{q}_4) \cos \bar{q}_5 + \cos \bar{q}_1 \cos \bar{q}_2 \cos(\bar{q}_3 + \bar{q}_4) \cos \bar{q}_5 \right] \\ b_{13} &= 0 \\ b_{21} &= -\beta_2 \sin \bar{q}_5 \left[\cos \bar{q}_1 \cos(\bar{q}_3 + \bar{q}_4) + \sin \bar{q}_1 \cos \bar{q}_2 \sin(\bar{q}_3 + \bar{q}_4) \right] \\ b_{22} &= \beta_2 \left[\sin \bar{q}_1 \sin \bar{q}_2 \sin \bar{q}_5 - \cos \bar{q}_1 \sin(\bar{q}_3 + \bar{q}_4) \cos \bar{q}_5 + \sin \bar{q}_1 \cos \bar{q}_2 \cos(\bar{q}_3 + \bar{q}_4) \cos \bar{q}_5 \right] \\ b_{23} &= 0 \\ b_{31} &= -\beta_2 \sin \bar{q}_2 \sin(\bar{q}_3 + \bar{q}_4) \sin \bar{q}_5 \\ b_{32} &= \beta_2 \left[\sin \bar{q}_2 \cos(\bar{q}_3 + \bar{q}_4) \cos \bar{q}_5 - \cos \bar{q}_2 \sin \bar{q}_5 \right] \\ b_{33} &= 0 \end{aligned}$$

In this section the computations of each rotation matrix ($\mathbf{Q}_1 - \mathbf{Q}_6$) were omitted, but they are available in the Appendix of this document.

4 Appendix: Rotation matrices

For this analysis we are using the rotation matrix definition of the Natural Orthogonal methodology. In this methodology the rotation matrices are defined as

$$\mathbf{Q}_i = \begin{bmatrix} \cos \bar{q}_i & -\lambda_i \sin \bar{q}_i & \mu_i \sin \bar{q}_i \\ \sin \bar{q}_i & \lambda_i \cos \bar{q}_i & -\mu_i \cos \bar{q}_i \\ 0 & \mu_i & \lambda_i \end{bmatrix} \quad (39)$$

where

$$\lambda_i = \cos \alpha_i$$

$$\mu_i = \sin \alpha_i$$

$$\bar{q}_i = q_i + q_{ini}$$

For the spatial pendulum under consideration

$$\mathbf{q}_{ini} = \begin{bmatrix} \pi & -\frac{\pi}{2} & \pi & 0 & \frac{\pi}{2} & 0 \end{bmatrix}^T \quad (40)$$

$$\boldsymbol{\alpha} = \begin{bmatrix} \frac{\pi}{2} & \frac{\pi}{2} & 0 & \frac{\pi}{2} & \frac{\pi}{2} & 0 \end{bmatrix}^T \quad (41)$$

Note that both vectors are constant and they depend on the architecture of the spatial pendulum. Thus, the rotation matrix for each generalized coordinate is as follows

$$\begin{aligned} \lambda_1 &= \cos\left(\frac{\pi}{2}\right) = 0 \quad ; \quad \mu_1 = \sin\left(\frac{\pi}{2}\right) = 1 & \lambda_4 &= \cos\left(\frac{\pi}{2}\right) = 0 \quad ; \quad \mu_4 = \sin\left(\frac{\pi}{2}\right) = 1 \\ \mathbf{Q}_1 &= \begin{bmatrix} \cos \bar{q}_1 & 0 & \sin \bar{q}_1 \\ \sin \bar{q}_1 & 0 & -\cos \bar{q}_1 \\ 0 & 1 & 0 \end{bmatrix} & \mathbf{Q}_4 &= \begin{bmatrix} \cos \bar{q}_4 & 0 & \sin \bar{q}_4 \\ \sin \bar{q}_4 & 0 & -\cos \bar{q}_4 \\ 0 & 1 & 0 \end{bmatrix} \end{aligned} \quad (42)$$

$$\begin{aligned} \lambda_2 &= \cos\left(\frac{\pi}{2}\right) = 0 \quad ; \quad \mu_2 = \sin\left(\frac{\pi}{2}\right) = 1 & \lambda_5 &= \cos\left(\frac{\pi}{2}\right) = 0 \quad ; \quad \mu_5 = \sin\left(\frac{\pi}{2}\right) = 1 \\ \mathbf{Q}_2 &= \begin{bmatrix} \cos \bar{q}_2 & 0 & \sin \bar{q}_2 \\ \sin \bar{q}_2 & 0 & -\cos \bar{q}_2 \\ 0 & 1 & 0 \end{bmatrix} & \mathbf{Q}_5 &= \begin{bmatrix} \cos \bar{q}_5 & 0 & \sin \bar{q}_5 \\ \sin \bar{q}_5 & 0 & -\cos \bar{q}_5 \\ 0 & 1 & 0 \end{bmatrix} \end{aligned} \quad (43)$$

$$\begin{aligned} \lambda_3 &= \cos(0) = 1 \quad ; \quad \mu_3 = \sin(0) = 0 & \lambda_6 &= \cos(0) = 1 \quad ; \quad \mu_6 = \sin(0) = 0 \\ \mathbf{Q}_3 &= \begin{bmatrix} \cos \bar{q}_3 & -\sin \bar{q}_3 & 0 \\ \sin \bar{q}_3 & \cos \bar{q}_3 & 0 \\ 0 & 0 & 1 \end{bmatrix} & \mathbf{Q}_6 &= \begin{bmatrix} \cos \bar{q}_6 & -\sin \bar{q}_6 & 0 \\ \sin \bar{q}_6 & \cos \bar{q}_6 & 0 \\ 0 & 0 & 1 \end{bmatrix} \end{aligned} \quad (44)$$

Compute the rotation matrix $\mathbf{Q}_1\mathbf{Q}_2$

$$\begin{aligned} \mathbf{Q}_1\mathbf{Q}_2 &= \begin{bmatrix} \cos \bar{q}_1 & 0 & \sin \bar{q}_1 \\ \sin \bar{q}_1 & 0 & -\cos \bar{q}_1 \\ 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} \cos \bar{q}_2 & 0 & \sin \bar{q}_2 \\ \sin \bar{q}_2 & 0 & -\cos \bar{q}_2 \\ 0 & 1 & 0 \end{bmatrix} \\ &= \begin{bmatrix} \cos \bar{q}_1 \cos \bar{q}_2 & \sin \bar{q}_1 & \cos \bar{q}_1 \sin \bar{q}_2 \\ \sin \bar{q}_1 \cos \bar{q}_2 & -\cos \bar{q}_1 & \sin \bar{q}_1 \sin \bar{q}_2 \\ \sin \bar{q}_2 & 0 & -\cos \bar{q}_2 \end{bmatrix} \end{aligned} \quad (45)$$

Compute the rotation matrix $\mathbf{Q}_2\mathbf{Q}_3$

$$\begin{aligned} \mathbf{Q}_2\mathbf{Q}_3 &= \begin{bmatrix} \cos \bar{q}_2 & 0 & \sin \bar{q}_2 \\ \sin \bar{q}_2 & 0 & -\cos \bar{q}_2 \\ 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} \cos \bar{q}_3 & -\sin \bar{q}_3 & 0 \\ \sin \bar{q}_3 & \cos \bar{q}_3 & 0 \\ 0 & 0 & 1 \end{bmatrix} \\ &= \begin{bmatrix} \cos \bar{q}_2 \cos \bar{q}_3 & -\cos \bar{q}_2 \sin \bar{q}_3 & \sin \bar{q}_2 \\ \sin \bar{q}_2 \cos \bar{q}_3 & -\sin \bar{q}_2 \sin \bar{q}_3 & -\cos \bar{q}_2 \\ \sin \bar{q}_3 & \cos \bar{q}_3 & 0 \end{bmatrix} \end{aligned} \quad (46)$$

Using Eq. (45) and the definition for \mathbf{Q}_3 , we find an expression for the rotation matrix from the inertial frame to the body frame of the first link

$$\begin{aligned}\mathbf{Q}_1\mathbf{Q}_2\mathbf{Q}_3 &= \mathbf{Q}_1\mathbf{Q}_2 \begin{bmatrix} \cos \bar{q}_3 & -\sin \bar{q}_3 & 0 \\ \sin \bar{q}_3 & \cos \bar{q}_3 & 0 \\ 0 & 0 & 1 \end{bmatrix} \\ &= \begin{bmatrix} \cos \bar{q}_1 \cos \bar{q}_2 \cos \bar{q}_3 + \sin \bar{q}_1 \sin \bar{q}_3 & -\cos \bar{q}_1 \cos \bar{q}_2 \sin \bar{q}_3 + \sin \bar{q}_1 \cos \bar{q}_3 & \cos \bar{q}_1 \sin \bar{q}_2 \\ \sin \bar{q}_1 \cos \bar{q}_2 \cos \bar{q}_3 - \cos \bar{q}_1 \sin \bar{q}_3 & -\sin \bar{q}_1 \cos \bar{q}_2 \sin \bar{q}_3 - \cos \bar{q}_1 \cos \bar{q}_3 & \sin \bar{q}_1 \sin \bar{q}_2 \\ \sin \bar{q}_2 \cos \bar{q}_3 & -\sin \bar{q}_2 \sin \bar{q}_3 & -\cos \bar{q}_2 \end{bmatrix} \quad (47)\end{aligned}$$

Compute the rotation matrix $\mathbf{Q}_5\mathbf{Q}_6$

$$\begin{aligned}\mathbf{Q}_5\mathbf{Q}_6 &= \begin{bmatrix} \cos \bar{q}_5 & 0 & \sin \bar{q}_5 \\ \sin \bar{q}_5 & 0 & -\cos \bar{q}_5 \\ 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} \cos \bar{q}_6 & -\sin \bar{q}_6 & 0 \\ \sin \bar{q}_6 & \cos \bar{q}_6 & 0 \\ 0 & 0 & 1 \end{bmatrix} \\ &= \begin{bmatrix} \cos \bar{q}_5 \cos \bar{q}_6 & -\cos \bar{q}_5 \sin \bar{q}_6 & \sin \bar{q}_5 \\ \sin \bar{q}_5 \cos \bar{q}_6 & -\sin \bar{q}_5 \sin \bar{q}_6 & -\cos \bar{q}_5 \\ \sin \bar{q}_6 & \cos \bar{q}_6 & 0 \end{bmatrix} \quad (48)\end{aligned}$$

Using Eq. (48) and the definition for \mathbf{Q}_4 , we find an expression for the rotation matrix from the body frame of the first link to the body frame of the second link

$$\begin{aligned}\mathbf{Q}_4\mathbf{Q}_5\mathbf{Q}_6 &= \begin{bmatrix} \cos \bar{q}_4 & 0 & \sin \bar{q}_4 \\ \sin \bar{q}_4 & 0 & -\cos \bar{q}_4 \\ 0 & 1 & 0 \end{bmatrix} \mathbf{Q}_5\mathbf{Q}_6 \\ &= \begin{bmatrix} \cos \bar{q}_4 \cos \bar{q}_5 \cos \bar{q}_6 + \sin \bar{q}_4 \sin \bar{q}_6 & -\cos \bar{q}_4 \cos \bar{q}_5 \sin \bar{q}_6 + \sin \bar{q}_4 \cos \bar{q}_6 & \cos \bar{q}_4 \sin \bar{q}_5 \\ \sin \bar{q}_4 \cos \bar{q}_5 \cos \bar{q}_6 - \cos \bar{q}_4 \sin \bar{q}_6 & -\sin \bar{q}_4 \cos \bar{q}_5 \sin \bar{q}_6 - \cos \bar{q}_4 \cos \bar{q}_6 & \sin \bar{q}_4 \sin \bar{q}_5 \\ \sin \bar{q}_5 \cos \bar{q}_6 & -\sin \bar{q}_5 \sin \bar{q}_6 & -\cos \bar{q}_5 \end{bmatrix} \quad (49)\end{aligned}$$