

# Lecture 3

niceguy

September 14, 2023

## 1 Potential Barriers

Let's say we have a long wave chain  $\psi(x, 0)$  that is approaching a potential barrier  $V(x)$ . We start by approximating  $\psi$  to have only one frequency. Then we can write

$$\psi(x, t) = \phi(x)e^{-i\frac{E}{\hbar}t} \quad (1)$$

Substituting into Schrödinger's Equation, we get (see complement H<sub>1</sub> in textbook)

$$\frac{d^2}{dx^2}\phi(x) + \frac{2m}{\hbar^2}(E - V(x))\phi(x) = 0 \quad (2)$$

Now if  $E > V$ , we put

$$E - V = \frac{\hbar^2 k^2}{2m} \quad (3)$$

This substituted into 2 gives

$$\phi(x) = Ae^{ix} + A'e^{-ikx}, |A| = |A'| \quad (4)$$

In quantum mechanics, in fact there is a nonzero probability that the wave is reflected.

$$\phi(x) = A_2 e^{ik_2 x} \quad (5)$$

If  $E < V$ , then we put

$$V - E = \frac{\hbar^2 \rho^2}{2m} \quad (6)$$

Plugging into 2 again, we get

$$\phi(x) = Be^{\rho x} + B'e^{-\rho x} \quad (7)$$

Classically, there is total reflection. However, in quantum mechanics, for  $x > 0$  we have

$$\phi(x) = B'e^{-\rho x} \quad (8)$$

## 2 Math facts of Quantum

*Note: This corresponds to the second chapter of the textbook.*

In section A of the textbook, we discuss the space of a 1 particle wave functions, where the integral of norm squared is 1. We define a vector space  $\mathcal{F}$  with elements being functions that are square integrable and satisfies our physical assumptions (boundary conditions). This is an infinite-dimension vector space.

In sections B - F, we talk about bracket notation. We link  $\psi(z) \in \mathcal{F}$  with  $|\psi\rangle \in \mathcal{E}_{\mathcal{F}}$ . We will look at more general kets later.

For kets  $|\phi\rangle \neq |\psi\rangle$  in  $\mathcal{E}$ , introduce a scalar product that is a complex number. We define in  $E_{\vec{r}}$

$$(|\phi\rangle, |\psi\rangle) = \int \phi^*(\vec{r})\psi(\vec{r})d\vec{r} \quad (9)$$

This is an inner product.

We also introduce a dual space  $\mathcal{E}^*$ . Recall a linear functional  $\chi$  maps from a ket to a complex number. The dual space is the set of all  $\chi$ s. In fact, we can prove that all  $\chi$  can be written uniquely as

$$\chi(|\psi\rangle) = \langle\chi|\psi\rangle \quad (10)$$