

# Lecture 1

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## 1 Course Information

Office hours Fridays 14:00 - 17:00 MP408, including reading week. Problem set questions appear on quizzes.

## 2 Waves and Particles

$$E = h\omega \tag{1}$$

$$p = \frac{h}{\lambda} \tag{2}$$

The LHS of both equations are usually linked with *particles*, while the RHS are linked with *waves*.

Relativistically, we have

$$E = \frac{mc^2}{\sqrt{1 - \frac{v^2}{c^2}}} \tag{3}$$

$$p = \frac{mv}{\sqrt{1 - \frac{v^2}{c^2}}} \tag{4}$$

And the nonrelativistic counterparts are trivial.

### 2.1 Electrons

Electrons are used to study surfaces, since they don't penetrate too deeply.

## 2.2 Wavefunctions

In 1 dimension, we have  $\cos(kx - \omega t)$  or  $\sin(kx - \omega t)$ . We also know that these functions are complex in quantum, hence we use  $e^{i(kx - \omega t)}$ .

At  $t = 0$ , we form a wavepacket

$$\psi(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(k) e^{ikx} dk \quad (5)$$

A simpler expression is to sum over  $ks$  instead of integration.

**Example 2.1.** Let

$$g(k) = \begin{cases} 1 & i = k_0 \\ \frac{1}{2} & i = k_0 \pm \frac{\Delta}{2} \end{cases}$$

Then

$$\psi(x) = \frac{e^{ik_0x}}{\sqrt{2\pi}} \left[ 1 + \cos\left(\frac{\Delta k_0 x}{2}\right) \right]$$

If we introduct time,  $e^{ikx}$  becomes  $e^{i(kx - \omega(k)t)}$ .

In a nonrelativistic case,

$$E = \frac{p^2}{2m} = \frac{\hbar^2 k^2}{2m} \quad (6)$$

Assuming superposition,

$$\begin{aligned} \psi(x, t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(k) e^{i(kx - \omega(k)t)} dk \\ &= \frac{1}{\sqrt{2\pi}} e^{i(k_0x - \omega(k_0)t)} \int_{-\infty}^{\infty} g(k) e^{i(k - k_0)x} e^{-i(\omega(k) - \omega_0)t} dk \\ &\approx \frac{1}{\sqrt{2\pi}} e^{i(k_0x - \omega_0t)} \int_{-\infty}^{\infty} g(k) e^{i(k - k_0)(x - Vt)} dk \end{aligned}$$

where we define

$$V = \left( \frac{d\omega}{dk} \right)_{k_0} \quad (7)$$

Letting the final integral be  $F(x, t)$ , we see  $F$  moves at a phase velocity  $V$ .

For our nonrelativistic particle,

$$\left(\frac{d\omega}{dk}\right)_{k_0} = \frac{\hbar k_0}{m} = \frac{p_0}{m} \quad (8)$$

For a relativistic particle,

$$\frac{\omega}{k} = \omega\lambda = \frac{E}{p} = \frac{c^2}{v} \quad (9)$$