Lecture 26

niceguy

April 21, 2024

1 Plane Wave Incidence on Material Boundaries

For simple media, the wavenumber is in general

$$k = \omega \sqrt{\varepsilon \mu} = \omega \sqrt{\varepsilon_0 \varepsilon_r \mu_0 \mu_r} = \omega \sqrt{\varepsilon_0 \mu_0} \sqrt{\varepsilon_r \mu_r} = k_0 n$$

where we define the **refractive index** to be

$$n = \sqrt{\varepsilon_r \mu_r}$$

If the incident waves are

$$\tilde{\vec{E}}_i = E_i e^{-jk_0 n_1 z} \hat{x}$$
$$\tilde{\vec{H}}_i = \frac{E_i}{\eta_1} e^{-jk_0 n_1 z} \hat{y}$$

The reflected wave will be

$$\tilde{\vec{E}}_r = \Gamma E_i e^{jk_0 n_1 z} \hat{x}$$
$$\tilde{\vec{H}}_r = -\Gamma \frac{E_i}{\eta_1} e^{jk_0 n_1 z} \hat{y}$$

where Γ is the reflection coefficient. The transmitted wave is

$$\begin{split} \tilde{\vec{E}}_t &= \tau E_i e^{-jk_0 n_2 z} \hat{x} \\ \tilde{\vec{H}}_t &= \tau \frac{E_i}{\eta_2} e^{-jk_0 n_2 z} \hat{y} \end{split}$$

By continuity of tangential \vec{E} at z=0, we have

$$\vec{E}_r + \vec{E}_i = \vec{E}_t$$

$$\Gamma E_i \hat{x} + E_i \hat{x} = \tau E_i \hat{x}$$

$$1 + \Gamma = \tau$$

If there is no surface current, the same applies to \vec{H} , so

$$\frac{E_i}{\eta_1} - \Gamma \frac{E_i}{\eta_1} = \tau \frac{E_i}{\eta_2}$$
$$\frac{1 - \Gamma}{\eta_1} = \frac{\tau}{\eta_2}$$

Solving,

$$\Gamma = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1}, \tau = \frac{2\eta_2}{\eta_1 + \eta_2}$$

In a perfect conductor, $\Gamma = -1, \tau = 0$. In the first medium,

$$\tilde{\vec{E}}(z<0) = E_i e^{-jk_0 n_1 z} - E_i e^{jk_0 n_1 z} \hat{x}$$

$$= -2j E_0 \sin(k_0 n_1 z) \hat{x}$$

and similarly

$$\tilde{\vec{H}}(z<0) = \frac{E_i}{\eta_1} e^{-jk_0 n_1 z} + \frac{E_i}{\eta_1} e^{jk_0 n_1 z} = \frac{2E_i}{\eta_1} \cos(k_0 n_1 z)\hat{y}$$

At z = 0,

$$\begin{split} \hat{n} \times (\vec{H}_2 - \vec{H}_1) &= \vec{J}_s \\ \hat{z} \times \left(0 - \frac{2E_i}{\eta_1} \hat{y}\right) &= \vec{J}_s \\ \vec{J}_s &= \frac{2E_i}{\eta_1} \hat{x} \end{split}$$