

Generalizing D'Alembert's Solution

The solution method is agnostic to boundary conditions:

$$\begin{aligned} y_{tt} &= c^2 y_{xx} & a < x < b \\ y(x, 0) &= f(x) & y_t(x, 0) &= g(x) \end{aligned}$$

$$y(x, t) = A(x - ct) + B(x + ct)$$

$$A(z) = \frac{1}{2} \left[F(z) - \frac{1}{c} \int_a^z G(x) dx \right] \quad B(z) = \frac{1}{2} \left[F(z) + \frac{1}{c} \int_a^z G(x) dx \right]$$

$$\begin{aligned} y(x, t) &= \frac{F(x - ct) + F(x + ct)}{2} + \frac{\int_a^{x+ct} G(s) ds - \int_a^{x-ct} G(s) ds}{2c} \\ &= \left[\frac{F(x - ct) + F(x + ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} G(s) ds \right] \end{aligned}$$

Generalizing D'Alembert's Solution

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$$y(x, t) = \frac{F(x - ct) + F(x + ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} G(s) ds$$

Pick $F(z)$ and $G(z)$ to match the boundary conditions.

- $y(a, t) = y(b, t) = 0$
 $\Rightarrow F$ & G are odd periodic extensions of f and g
- $y_x(a, t) = y_x(b, t) = 0$
 $\Rightarrow F$ & G are even periodic extensions of f and g

Separation of variables with zero-derivative boundaries

$$\begin{aligned}
 y_{tt} &= c^2 y_{xx} & 0 < x < L & & y_x(0, t) = y_x(L, t) = 0 \\
 y(x, 0) &= f(x) & & & y_t(x, 0) = g(x)
 \end{aligned}$$

Perform even periodic extensions of the initial conditions

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos\left(\frac{n\pi}{L}x\right)$$

$$g(x) = \frac{d_0}{2} + \sum_{n=1}^{\infty} d_n \cos\left(\frac{n\pi}{L}x\right)$$

$$y(x, t) = \frac{d_0}{2}t + \frac{a_0}{2} + \sum_{n=1}^{\infty} \cos\left(\frac{n\pi}{L}x\right) \left[d_n \frac{L}{n\pi c} \sin\left(\frac{n\pi c}{L}t\right) + a_n \cos\left(\frac{n\pi c}{L}t\right) \right]$$

Suppose you are given

$$\begin{aligned}y_{tt} &= c^2 y_{xx} & 0 < x < L & & y_x(0, t) = y_x(L, t) = 0 \\ y(x, 0) &= f(x) & & & y_t(x, 0) = g(x)\end{aligned}$$

Compare the two methods for computing the solution to this problem

Separation of variables

1. Find the even periodic extension of $f(x)$

- a_n = an integral
- a_0 = another integral

2.

$$y(x, t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos\left(\frac{n\pi}{L}x\right) \cos\left(\frac{n\pi c}{L}t\right)$$

Infinite sums are too costly, in practice we only get approximate solutions.

<https://www.desmos.com/calculator/auxhik8lsh>

Suppose you are given

$$\begin{aligned} y_{tt} &= c^2 y_{xx} & 0 < x < L & & y_x(0, t) = y_x(L, t) = 0 \\ y(x, 0) &= f(x) & & & y_t(x, 0) = g(x) \end{aligned}$$

Compare the two methods for computing the solution to this problem

D'Alembert

1. Define an even extension of $f(x)$ onto $[-L, L]$

$$f^{\text{even}}(x) = \begin{cases} f(x) & 0 \leq x < L \\ -f(-x) & -L \leq x < 0 \end{cases}$$

- 2.

$$\begin{aligned} y(x, t) = \frac{1}{2} [& f^{\text{even}}(\text{mod}(x - ct + L, 2L) - L) \\ & + f^{\text{even}}(\text{mod}(x + ct + L, 2L) - L)] \end{aligned}$$

Exact solution, no infinite sums or integrals required.

<https://www.desmos.com/calculator/auxhik8lsh>

Other boundary conditions?

Suppose we have $y(0, t) = y_0 \neq 0$ and $y(L, t) = y_L \neq 0$

Proceed like with the heat equation

$$y_p(x) = y_0 + \frac{y_L - y_0}{L}x$$

define $w(x, t) = y(x, t) - y_p(x)$

$$\begin{aligned} w_{tt} &= c^2 w_{xx} & 0 < x < L & & w(0, t) = w(L, t) = 0 \\ w(x, 0) &= f(x) - y_p(x) & & & y_t(x, 0) = g(x) \end{aligned}$$

Same strategy applies for non-zero derivatives.

Periodic extensions of functions defined for $a < x < b$

We have seen how to make odd/even periodic extensions of functions defined for $0 < x < L$

- Graphically
- Fourier Series with period $2L$
 - even=Fourier Cosine Series
 - odd=Fourier Sine Series

How does this work for a function defined for $a < x \leq b$?

1. Define a new coordinate $\tilde{x} = x - a$
2. Take a Fourier series in \tilde{x} with $L = b - a$

Hint: You need to do this for Assignment Q5c

What happens as $a \rightarrow -\infty$ and $b \rightarrow \infty$

$$\begin{aligned}y_{tt} &= c^2 y_{xx} & -\infty < x < \infty \\y(x, 0) &= f(x) & y_t(x, 0) = g(x)\end{aligned}$$

$$y(x, t) = \frac{f(x - ct) + f(x + ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} g(s) ds$$

The wave extends over all space... eventually.

However, the wave travels at finite speed.

Information about the initial condition at a point x_0 is always contained in the interval $[x_0 - ct, x_0 + ct]$.

What type of solutions do we get for $-\infty < x < \infty$?

For $0 < x < L$, we got

$$y(x, t) \approx \sum_{n=1}^{\infty} \left[a_n \cos \left(\frac{2\pi t}{T_n} \right) + b_n \sin \left(\frac{2\pi t}{T_n} \right) \right] \left[\sin \left(\frac{2\pi x}{\lambda_n} \right) + \cos \left(\frac{2\pi x}{\lambda_n} \right) \right]$$

a_n and b_n are found from periodic extensions of the initial conditions

$$T_n = \frac{2L}{cn} \quad \lambda_n = \frac{2L}{n} \quad (\text{periods of the vibrational modes})$$

alternatively, we can rewrite this as

$$y(x, t) = \sum_{n=-\infty}^{\infty} c_n e^{i\frac{2\pi}{T_n}t} e^{i\frac{2\pi}{\lambda_n}x} = \sum_{n=-\infty}^{\infty} c_n e^{2\pi i \left(\frac{x}{\lambda_n} + \frac{t}{T_n} \right)}$$

$$c_n = \begin{cases} a_0/2 & n = 0 \\ \frac{1}{2}(a_n + b_n) & n > 0 \\ \frac{1}{2}(a_{|n|} - b_{|n|}) & n < 0 \end{cases}$$

What type of solutions do we get for $-\infty < x < \infty$?

For $-\infty < x < \infty$, we get

Inspired by the finite case, we try separation of variables with

$$y(x, t) = e^{\alpha x} e^{\beta t}$$

$$y_{tt} = c^2 y_{xx}$$

$$\beta^2 e^{\alpha x} e^{\beta t} = c^2 \alpha^2 e^{\alpha x} e^{\beta t}$$

$$\beta^2 = c^2 \alpha^2$$

$$\beta = \pm c\alpha$$

For general initial conditions, β and α are purely imaginary. Let $\alpha = ik$

$$y(x, t) = \int_{-\infty}^{\infty} w(k) e^{2\pi i k(x+ct)} dk$$

The function $w(k)$ is found from a Fourier Transform of the initial conditions

NOT ON EXAM

What type of solutions do we get for $-\infty < x < \infty$?

For $-\infty < x < \infty$, we get

For some very special initial conditions, Fourier Transforms are not necessary.

ex: $y_{tt} = c^2 y_{xx}$ with $y(x, 0) = e^{-2x}$

Try $y_{1,2}(x, t) = e^{\alpha x \pm c\alpha t}$

Initial condition: $y(x, 0) = e^{-2x} = e^{\alpha x} \Rightarrow \alpha = -2$

$$y_{1,2}(x, t) = e^{-2x \mp 2ct}$$

Summary

- D'Alembert solution to the wave equation is complementary to the Fourier series solution:
 - Same solution, different perspective
- Two waves one moving left and another moving right
 - The two waves are fully determined by the initial conditions.
 - The initial conditions propagate through the domain at finite speed
- Finite Domain: use appropriate periodic extension + wave interference to match boundary conditions.
- Infinite Domain: No need to use periodic extensions, there are no boundary conditions.