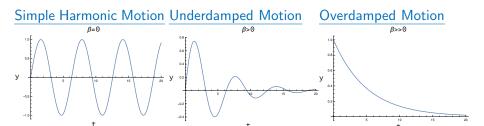


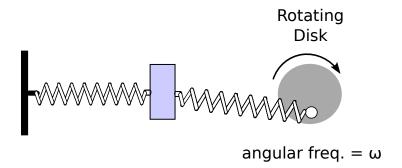
x(t) = displacement from rest position

Newton's 2nd Law:

$$F = ma$$
$$x'' + \beta x' + kx = 0$$



Spring oscillators with periodic forcing: $f(t) = F_0 \cos(\omega t)$



$$\underbrace{x'' + \omega_0^2 x}_{\frac{1}{m}(mx'' + kx) \text{ with } \omega_0 = \sqrt{k/m}} = \frac{f(t)}{m} = \frac{F_0}{m} \cos(\omega t) \quad \text{with} \quad \frac{\omega \neq \omega_0}{x(0) = x'(0) = 0}$$

$$x_h = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) \qquad x_p = A \cos(\omega t) + B \sin(\omega t)$$

$$x_p'' = -\omega^2 A \cos(\omega t)$$

$$-\omega^2 A \cos(\omega t) + \omega_0^2 A \cos(\omega t) = \frac{F_0}{m} \cos(\omega t)$$

$$\frac{\sin(t):}{\cos(t):} -\omega^2 B + \omega_0^2 B = 0 \qquad \Rightarrow B = 0$$

$$\frac{\cos(t):}{\cos(t):} -\omega^2 A + \omega_0^2 A = \frac{F_0}{m} \qquad \Rightarrow A = \frac{F_0}{m(\omega_0^2 - \omega^2)}$$

$$x_p = \frac{F_0}{m(\omega_0^2 - \omega^2)} \cos(\omega t)$$

$$x'' + \omega_0^2 x = \frac{F_0}{T} \cos(\omega t) \quad \text{with} \quad$$

 $\begin{array}{cc} \omega \neq \omega_0 \\ x(0) = x'(0) = 0 \end{array}$

$$x_h = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t)$$

$$x_p = \frac{F_0}{m(\omega_0^2 - \omega^2)} \cos(\omega t)$$

Initial Conditions:

$$x(0) = 0 = c_1 + \frac{F_0}{m(\omega_0^2 - \omega^2)}$$

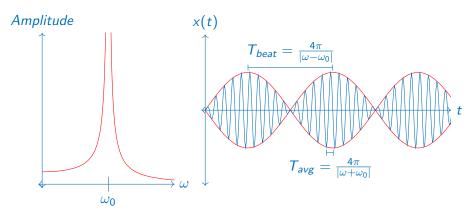
$$c_1 = \frac{F_0}{m(\omega^2 - \omega_0^2)}$$

$$x'(0) = 0 = c_2$$

$$x(t) = \frac{F_0}{m(\omega_0^2 - \omega^2)} \underbrace{(\cos(\omega t) - \cos(\omega_0 t))}_{\sin(\frac{\omega_0 + \omega}{2} t) \sin(\frac{\omega_0 - \omega}{2} t)}$$

Beat phenomena

$$x(t) = \frac{F_0}{m(\omega_0^2 - \omega^2)} \sin\left(\frac{\omega_0 + \omega}{2}t\right) \sin\left(\frac{\omega_0 - \omega}{2}t\right)$$



What happens to x(t) as $\omega \to \omega_0$?

Resonance ($\omega \to \omega_0$)

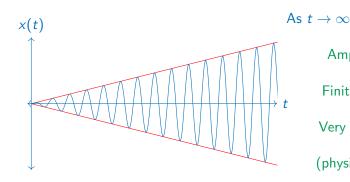
$$x'' + \omega_0^2 x = f(t) = \frac{F_0}{m} \cos(\omega_0 t)$$

$$x_h = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) \qquad x_p = ?$$

$$y_p \neq A \cos(\omega_0 t) + B \sin(\omega_0 t) \qquad \text{(mathematical resonance)}$$

$$y_p = At \cos(\omega_0 t) + Bt \sin(\omega_0 t)$$

 $B = \frac{F_0}{2m\omega_0}$



Amplitude $\rightarrow \infty$

Finite input power

Very large response

(physical resonance)

Damped oscillator with forcing

$$mx'' + \beta x' + kx = F_0 \cos(\omega t)$$
$$x(t) = x_h(t) + x_p(t)$$

Three cases:

$$x_h(t) = \begin{cases} e^{-\frac{\beta}{2m}t} \left(c_1 \cos(\omega_1 t) + c_2 \sin(\omega_1 t) \right) & \beta^2 < 4km \\ c_1 e^{r_1 t} + c_2 e^{r_2 t} & \beta^2 > 4km \\ c_1 e^{rt} + c_2 t e^{rt} & \beta^2 = 4km \end{cases}$$

Since $\beta > 0$, all exponents are negative

$$\Rightarrow$$
 $x_h \to 0$ as $t \to \infty$

Damped oscillator with forcing

$$mx'' + \beta x' + kx = F_0 \cos(\omega t)$$

$$x(t) = x_h(t) + x_p(t)$$

Three cases:

$$x_h(t) = \begin{cases} e^{-\frac{\beta}{2m}t} \left(c_1 \cos(\omega_1 t) + c_2 \sin(\omega_1 t) \right) & \beta^2 < 4km \\ c_1 e^{r_1 t} + c_2 e^{r_2 t} & \beta^2 > 4km \\ c_1 e^{rt} + c_2 t e^{rt} & \beta^2 = 4km \end{cases}$$

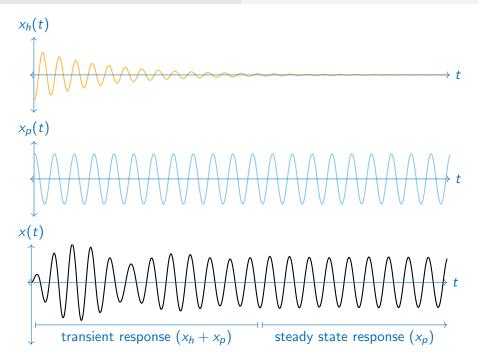
From the Method of Undetermined Coefficients:

$$x_p = A\cos(\omega t) + B\sin(\omega t) \quad \forall \ \omega \in \mathbb{R}$$

Impossible to have mathematical resonance with periodic forcing

As $t \to \infty$

$$x(t) \rightarrow x_p(t)$$



$$mx'' + \beta x' + kx = \cos(\omega t)$$

How does the amplitude of the steady state response vary with ω ?

$$x_p = A\cos(\omega t) + B\sin(\omega t)$$

$$x'_p = -\omega A\sin(\omega t) + \omega B\cos(\omega t)$$

$$x''_p = -\omega^2 A\cos(\omega t) - \omega^2 B\sin(\omega t)$$

$$m\left(-\omega^2 A \cos(\omega t) - \omega^2 B \sin(\omega t)\right) + \beta\left(-\omega A \sin(\omega t) + \omega B \cos(\omega t)\right) + k\left(A \cos(\omega t) + B \sin(\omega t)\right) = \cos(\omega t)$$

$$(Ak + B\beta\omega - mA\omega^{2})\cos(\omega t) + (Bk - A\beta\omega - mB\omega^{2})\sin(\omega t)$$
$$= \cos(\omega t)$$

Match coefficients

$$\frac{\cos :}{\sin :} (Ak + B\beta\omega - mA\omega^{2}) = 1$$

$$\frac{\sin :}{\sin :} (Bk - A\beta\omega - mB\omega^{2}) = 0$$

$$B = A\frac{\beta\omega}{k - m\omega^{2}}$$

$$\left(Ak + A\frac{(\beta\omega)^{2}}{k - m\omega^{2}} - mA\omega^{2}\right) = 1$$

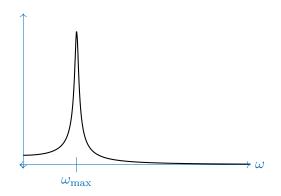
$$A = \frac{k - m\omega^{2}}{(\beta\omega)^{2} - (k - m\omega^{2})^{2}} \qquad B = \frac{\beta\omega}{(\beta\omega)^{2} - (k - m\omega^{2})^{2}}$$

Amplitude:

$$R = \sqrt{A^2 + B^2}$$

$$= \frac{\sqrt{(k - m\omega^2)^2 + (\beta\omega^2)}}{(\beta\omega)^2 - (k - m\omega^2)^2}$$

$$= \frac{1}{\sqrt{k^2 - 2km\omega^2 + m^2\omega^4 + \omega^2\beta^2}}$$



Amplitude vs Forcing Frequency

Steady-state response: $x_p = R \cos(\omega t - \phi)$ with

$$R(\omega) = \frac{1}{\sqrt{k^2 - 2km\omega^2 + m^2\omega^4 + \omega^2\beta^2}}$$
$$= \frac{1}{\sqrt{\beta^2\omega^2 + (k - m\omega^2)^2}}$$

What value of ω creates the largest amplitude response?

$$\frac{\mathrm{d}}{\mathrm{d}\omega}R = -\frac{\omega\left(\beta^2 - 2km + 2m^2\omega^2\right)}{\left(\beta^2\omega^2 + (k - m\omega^2)^2\right)^{3/2}} = 0$$

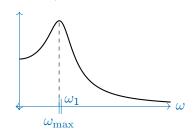
Critical points:
$$\omega^* = 0, \quad \pm \sqrt{\frac{k}{m} - \frac{1}{2} \left(\frac{\beta}{m}\right)^2}$$

Amplitude vs Forcing Frequency

Underdamped:

$$x_h = e^{-\frac{\beta}{2m}t}(c_1\cos(\omega_1t) + c_2\sin(\omega_1t))$$

$$\omega_1 = \sqrt{\frac{k}{m} - \left(\frac{\beta}{2m}\right)^2}$$

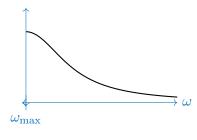


$$\omega_{\max} = \sqrt{\frac{k}{m} - \frac{1}{2} \left(\frac{\beta}{m}\right)^2} \approx \omega_1$$

Overdamped:

$$x_h = c_1 e^{r_1 t} + c_2 e^{r_2 t}$$

$$r_{1,2} = \frac{-\beta \pm \sqrt{\beta^2 - 4km}}{2m}$$



$$\omega_{\mathrm{max}} = 0$$

Resonance with damped oscillators?

$$mx'' + \beta x' + kx = F_0 \cos(\omega t)$$

What happens as $\omega \to \omega_{\rm max}$?

- Transient response $(x_h + x_p \text{ for } t \approx 0)$
 - amplitude and duration both grow larger
- Steady state response (x_p)
 - amplitude grows larger

This phenomenon is called quasi-resonance or practical resonance.

 ω_{max} is called the quasi-resonance frequency