

D'Alembert's Method: $y_{tt} = c^2 y_{xx}$

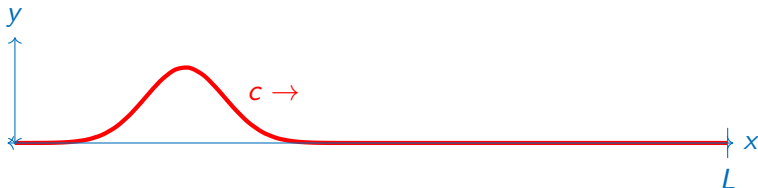
The Separation of Variables/Fourier method can be used to analyze the spectrum of frequencies present in the wave.

Useful to understand music - not useful to visualize a wave.

D'Alembert's Method: an alternative approach based on change of coordinates:

$$\xi = x - ct, \quad \eta = x + ct$$

Basic idea, solutions are waves travelling with speed $\pm c$.



D'Alembert's Method

$$\xi = x - ct, \quad \eta = x + ct$$

Let's compute $\frac{\partial^2}{\partial x^2}$ in this new coordinate system.

Apply the chain rule:

$$\begin{aligned} \frac{\partial}{\partial x} &= \frac{\partial \xi}{\partial x} \frac{\partial}{\partial \xi} + \frac{\partial \eta}{\partial x} \frac{\partial}{\partial \eta} = \frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta} \\ \frac{\partial^2}{\partial x^2} &= \frac{\partial}{\partial x} \frac{\partial}{\partial x} = \left(\frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta} \right) \left(\frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta} \right) \\ &= \frac{\partial^2}{\partial \xi^2} + 2 \frac{\partial^2}{\partial \xi \partial \eta} + \frac{\partial^2}{\partial \eta^2} \end{aligned}$$

D'Alembert's Method

$$\xi = x - ct, \quad \eta = x + ct$$

Let's compute $\frac{\partial^2}{\partial t^2}$ in this new coordinate system.

Apply the chain rule:

$$\begin{aligned} \frac{\partial}{\partial t} &= \frac{\partial \xi}{\partial t} \frac{\partial}{\partial \xi} + \frac{\partial \eta}{\partial t} \frac{\partial}{\partial \eta} = \overset{-c}{\cancel{\frac{\partial \xi}{\partial t}}} \frac{\partial}{\partial \xi} + \overset{+c}{\cancel{\frac{\partial \eta}{\partial t}}} \frac{\partial}{\partial \eta} = -c \frac{\partial}{\partial \xi} + c \frac{\partial}{\partial \eta} \\ \frac{\partial^2}{\partial t^2} &= \frac{\partial}{\partial t} \frac{\partial}{\partial t} = \left(-c \frac{\partial}{\partial \xi} + c \frac{\partial}{\partial \eta} \right) \left(-c \frac{\partial}{\partial \xi} + c \frac{\partial}{\partial \eta} \right) \\ &= c^2 \left(\frac{\partial^2}{\partial \xi^2} - 2 \frac{\partial^2}{\partial \xi \partial \eta} + \frac{\partial^2}{\partial \eta^2} \right) \end{aligned}$$

D'Alembert's Method

$$\frac{\partial^2}{\partial x^2} = \frac{\partial^2}{\partial \xi^2} + 2\frac{\partial^2}{\partial \xi \partial \eta} + \frac{\partial^2}{\partial \eta^2} \qquad \frac{\partial^2}{\partial t^2} = c^2 \frac{\partial^2}{\partial \xi^2} - 2c^2 \frac{\partial^2}{\partial \xi \partial \eta} + c^2 \frac{\partial^2}{\partial \eta^2}$$

Write down the wave equation in the new coordinate system.

$$\begin{aligned} y_{tt} &= c^2 y_{xx} \\ c^2 \frac{\partial^2 y}{\partial \xi^2} - 2c^2 \frac{\partial^2 y}{\partial \xi \partial \eta} + c^2 \frac{\partial^2 y}{\partial \eta^2} &= c^2 \frac{\partial^2 y}{\partial \xi^2} + c^2 2 \frac{\partial^2 y}{\partial \xi \partial \eta} + c^2 \frac{\partial^2 y}{\partial \eta^2} \\ 0 &= 4c^2 \frac{\partial^2 y}{\partial \xi \partial \eta} \end{aligned}$$

D'Alembert's Method

$$0 = \frac{\partial^2 y}{\partial \xi \partial \eta}$$

Find the general solution to the wave equation:

Drop the multiplicative constants, integrate w.r.t. ξ

$$y_\eta(\xi, \eta) = \int \frac{\partial^2 y}{\partial \xi \partial \eta} d\xi = \int 0 d\xi = C(\eta)$$

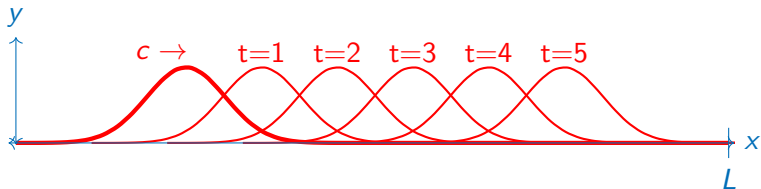
$$y(\xi, \eta) = \int C(\eta) d\eta = A(\eta) + B(\xi)$$

where A and B are single variable functions.

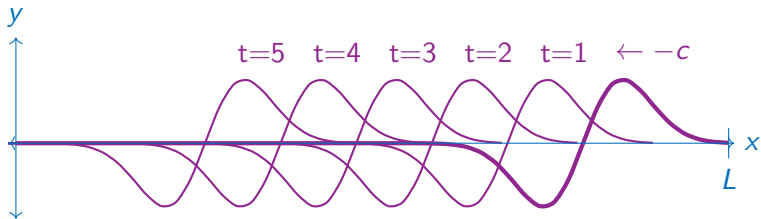
$$y(x, t) = A(x - ct) + B(x + ct)$$

$$y(x, t) = A(x - ct) + B(x + ct)$$

$A(x - ct)$: moves rightwards



$B(x + ct)$: moves leftwards



These are just two unrelated examples, the wave profiles A and B are not completely different like in this slide.

D'Alembert's Formula

General Solution to:

$$\begin{aligned} y_{tt} &= c^2 y_{xx} & 0 < x < L & & y(0) = y(L) = 0, \\ y(x, 0) &= f(x) & & & y_t(x, 0) = g(x) \end{aligned}$$

$$y(x, t) = A(x - ct) + B(x + ct)$$

Let F be the odd periodic extension of $f(x)$, and G be the odd periodic extension of $g(x)$. Then we have

$$A(z) = \frac{1}{2} \left[F(z) - \frac{1}{c} \int_0^z G(x) dx \right] \quad \text{and} \quad B(z) = \frac{1}{2} \left[F(z) + \frac{1}{c} \int_0^z G(x) dx \right]$$

Proof: DiffyQs §3.8.2

Interpreting D'Alembert's Formula

$$\begin{aligned}
 y_{tt} &= c^2 y_{xx} & 0 < x < L & & y(0) = y(L) = 0, \\
 y(x, 0) &= f(x) & & & y_t(x, 0) = g(x)
 \end{aligned}$$

$$y(x, t) = A(x - ct) + B(x + ct) \quad \text{with}$$

$$A(z) = \frac{1}{2} \left[F(z) - \frac{1}{c} \int_0^z G(x) dx \right] \quad \text{and} \quad B(z) = \frac{1}{2} \left[F(z) + \frac{1}{c} \int_0^z G(x) dx \right]$$

F and G are odd periodic extensions of $f(x)$ and $g(x)$, respectively.

Suppose $g(x) = 0$, then $G(x) = 0$ so

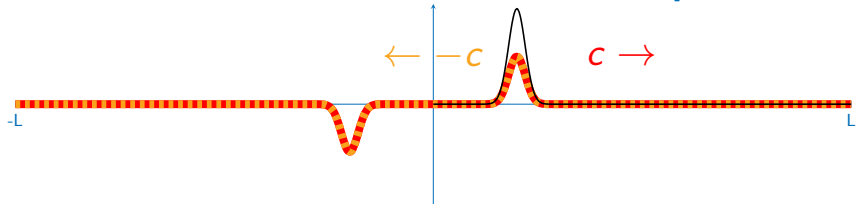
$$A(x) = B(x) = \frac{1}{2} F(z)$$

We have odd extensions of the initial condition moving left and right.

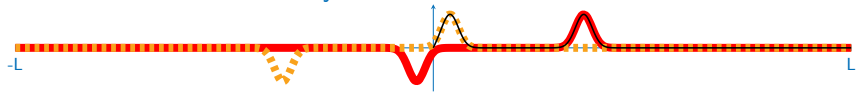
This leads to reflection of the wave at the domain boundaries.

$t = 0$: A and B sum to give the initial condition

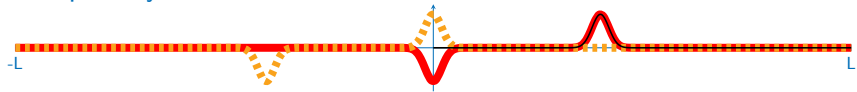
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$t = 0.8$: B hits the boundary at $x = 0$



$t = 1$: primary wave from B annihilates with the odd extension of A



$t = 1.8$: wave is reflected after hitting boundary

