

P112A:Electromagnetic Theory

Vector Analysis (Chapter 1)

Physics 112A Approximate Schedule, Winter 2024

Week	Reading Assignment	Homework “(0)” indicate problem numbers in 4 th edition & Due date	Example & Practices	Concepts
Week 1	Ch1	HW1, 1/17 Ch1: 41, 50, 54, 55, 63	Ch1: 37, 38, 39, 40, 42	Vector field Curvilinear coordinates
Week 2	Ch2	HW2, 1/24 Ch2: 7, 12(11), 17(16), 23(22)	Ch2: 12, 14, 21, 23, 24, 28	Electric field Electrostatic potential
Week 3	Ch2	HW3, 1/31 Ch2: 35(34), 39(38), 43(42), 44(43)	Ch2: 25, 36, 39, 44, 47, 48, 50	Boundary conditions Conductors
Week 4	Ch3	HW4, 2/7 Ch3: 10(9), 11(10), 12(11), 16(14), 18(16)	Ch3: 3, 8, 14, 15, 21, 24	Method of images Laplace's Equation
Week 5	Ch3	HW5, 2/14 Ch3: 20(18), 21(19), 32(30), 34(32), 37(35)	Ch3: 34, 36, 43	Separation of variables Multipole expansion
Week 6	Midterm Ch4	HW6, 2/21 Ch4: 8, 10	Ch4: 7, 18, 19, 20, 26, 35	Midterm, 2/13 Tuesday, Ch1-3 Polarization
Week 7	Ch4 Ch5	HW7, 2/28 Ch4: 12, 15, 20, 21, 24 Ch5: 3, 6	Ch4: 22, 33, 34 Ch5: 7, 10, 11	Electric Displacement, Linear dielectric Lorentz force
Week 8	Ch5	HW8, 3/6 Ch5: 12, 13, 14	Ch5: 15	Ampere's Law Vector potential
Week 9 3/	Ch5	HW9, 3/13 Ch5: 25(23), 26(24), 27(25), 28(26), 37(35), 39(37), 44(41)	Ch5: 27, 34, 35, 36, 37, 57	Vector potential Multipole expansion
Week 10	Ch6	HW10, 3/20 Ch6: 8, 10, 12, 16	Ch6: 4, 7, 15, 17	Magnetization Ampere's Law in magnetized materials
Week 11 Final				Final, 3/21 Thursday 8:00-10:00am Ch1-6

Read textbook before class

Discussion/Quiz on Friday

Electricity and Magnetism

- Branch of physics that studies how charges and currents interact and the resulting electromagnetic phenomena
- In P112, Solve for Maxwell's equations:

Integral form (7D)

$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{\text{encl}}}{\epsilon_0} \quad (\text{Gauss's law for } \vec{E})$$

Area vector

$$\oint \vec{B} \cdot d\vec{A} = 0 \quad (\text{Gauss's law for } \vec{B})$$

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 \left(i_C + \epsilon_0 \frac{d\Phi_E}{dt} \right)_{\text{encl}} \quad (\text{Ampere's law})$$

$$\oint \vec{E} \cdot d\vec{l} = -\frac{d\Phi_B}{dt} \quad (\text{Faraday's law})$$

Differential form

$$\nabla \cdot \mathbf{E} = \rho / \epsilon_0$$

$$\nabla \cdot \mathbf{B} = 0$$

Displacement current

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{d\mathbf{E}}{dt}$$

$$\nabla \times \mathbf{E} = -\frac{d\mathbf{B}}{dt}$$

E field from charges is radial:

$$\nabla \times \mathbf{E} = 0$$

Forces on charge:

$$\mathbf{F} = Q(\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

Current Density:

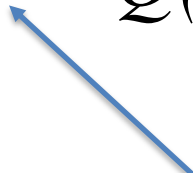
$$\nabla \cdot \mathbf{J} = -d\rho/dt$$

Math

- Vector algebra
 - Dot product: \bullet , cross-product: \times , triple products
- Vector calculus
 - Gradient: ∇F , Divergence: $\nabla \bullet F$, Curl: $\nabla \times F$
 - Integral calculus: Gradient theorem, Divergence theorem, Curls theorem
- Coordinate systems
 - Cartesian, Curvilinear: polar, cylindrical, and spherical
- Dirac Delta function (1D and 3D)
- Vector Fields
 - Scalar and vector potentials

Vector

- Vector fields are physical quantities, independent of the coordinate system
- Vector equations have same form in any coordinate system, e.g, $\mathbf{F} = Q(\mathbf{E} + \mathbf{v} \times \mathbf{B})$ Lorentz Force Law



Vector notation in our class: **Bold** or
error: \mathbf{F} or \vec{F}

Vector Operators in Cartesian Coordinates

Two vectors **F** and **G**: $\mathbf{F} = F_x \hat{x} + F_y \hat{y} + F_z \hat{z}$, $\mathbf{G} = G_x \hat{x} + G_y \hat{y} + G_z \hat{z}$

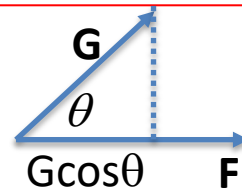
$\hat{x}, \hat{y}, \hat{z}$: unit vectors in x, y and z directions

1. Dot Product

$$\mathbf{F} \cdot \mathbf{G} = FG \cos \theta \quad \theta: \text{angle between } \mathbf{F} \text{ and } \mathbf{G}$$

$$\mathbf{F} \cdot \mathbf{G} = F_x G_x + F_y G_y + F_z G_z$$

Geometrically, it is projection of one vector on the other



2. Cross Product

$$\mathbf{F} \times \mathbf{G} = FG \sin \theta \hat{n}$$

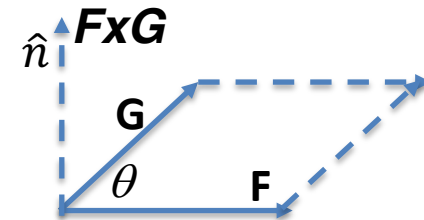
$$\mathbf{F} \times \mathbf{G} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ F_x & F_y & F_z \\ G_x & G_y & G_z \end{vmatrix}$$

determinant

θ : angle between **F** and **G**

\hat{n} : unit vector direction determined by right hand rule

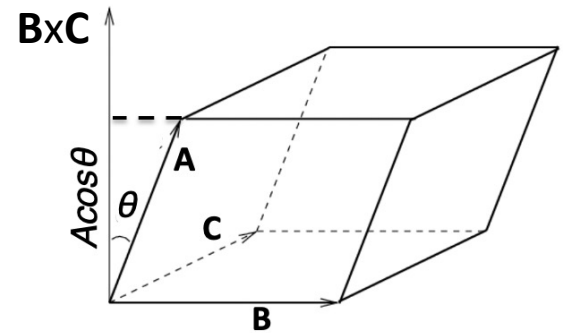
Geometrically, it is area of Parallelogram formed by **F** and **G**



$$\mathbf{F} \times \mathbf{G} = (F_y G_z - F_z G_y) \hat{x} - (F_x G_z - F_z G_x) \hat{y} + (F_x G_y - F_y G_x) \hat{z}$$

Triple Products

- *Scalar triple product is volume of parallelepiped since $\mathbf{B} \times \mathbf{C}$ is area of base, and $A \cos \theta$ is height:*



$$\mathbf{A} \cdot (\mathbf{B} \times \mathbf{C}) = \mathbf{B} \cdot (\mathbf{C} \times \mathbf{A}) = \mathbf{C} \cdot (\mathbf{A} \times \mathbf{B}) = \begin{bmatrix} A_x & A_y & A_z \\ B_x & B_y & B_z \\ C_x & C_y & C_z \end{bmatrix}$$

- *Note that dot and cross can be interchanged:*

$$\mathbf{A} \cdot (\mathbf{B} \times \mathbf{C}) = (\mathbf{A} \times \mathbf{B}) \cdot \mathbf{C}$$

Triple Products

- *Vector triple product*

$$\mathbf{A} \times (\mathbf{B} \times \mathbf{C}) = \mathbf{B} (\mathbf{A} \cdot \mathbf{C}) - \mathbf{C} (\mathbf{A} \cdot \mathbf{B})$$

- *If \mathbf{B} parallel to \mathbf{C} , or \mathbf{A} is perpendicular to \mathbf{B} and \mathbf{C}*

$$\mathbf{A} \times (\mathbf{B} \times \mathbf{C}) = \mathbf{0}$$

- *Vector triple products are not associative and Commutative*

$$\mathbf{A} \times (\mathbf{B} \times \mathbf{C}) \neq (\mathbf{A} \times \mathbf{B}) \times \mathbf{C}$$

$$(\mathbf{A} \times \mathbf{B}) \times \mathbf{C} = -\mathbf{C} \times (\mathbf{A} \times \mathbf{B})$$

First Derivatives - Gradient

∇ operator - “Del”

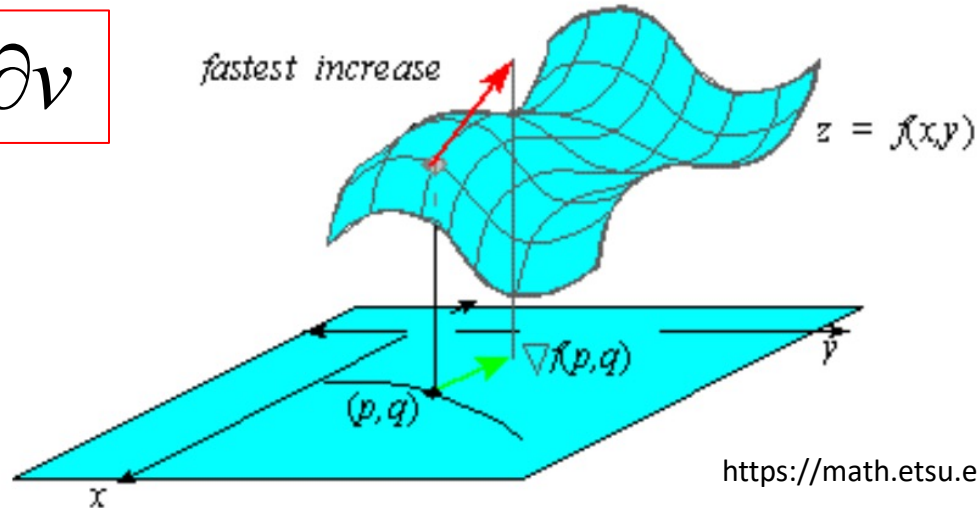
$$\nabla \equiv \hat{x} \partial/\partial x + \hat{y} \partial/\partial y + \hat{z} \partial/\partial z$$

- Gradient of scalar function $f(x,y,z)$:

$$\nabla f = \hat{x} \partial f / \partial x + \hat{y} \partial f / \partial y + \hat{z} \partial f / \partial z \quad \text{Or } \text{grad } f$$

- Gradient of f is a vector, gives the direction and the rate of fastest increase of f at a given point
- Changing rate of f respect to direction \mathbf{v} :

$$\nabla f \cdot \hat{\mathbf{v}} = \partial f / \partial v$$



Proof see A.3

Divergence

- Divergence of vector function F

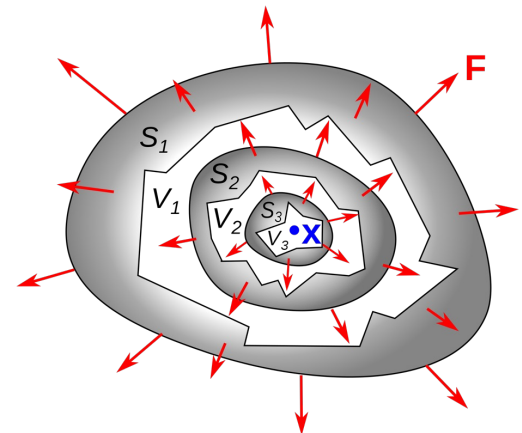
$$\nabla \cdot F = \partial F_x / \partial x + \partial F_y / \partial y + \partial F_z / \partial z, \text{ or } \text{div} F$$

- Divergence at point X is the ratio of the flux of F out of the closed surface of a volume V enclosing point X to the volume of V

$$\nabla \cdot F = \lim_{V \rightarrow 0} \frac{\Phi(S)}{V}$$

$$\Phi(S) = \oiint_{S(V)} \mathbf{F} \cdot \hat{\mathbf{n}} dS$$

Proof see A.4



Curl

- Curl of vector function \mathbf{F}

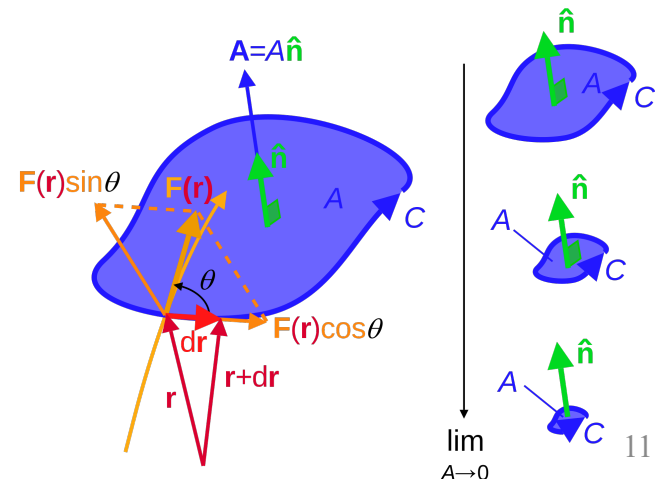
$$\nabla \times \mathbf{F} = \begin{bmatrix} \hat{x} & \hat{y} & \hat{z} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ F_x & F_y & F_z \end{bmatrix} \quad \text{Or } \text{curl } \mathbf{F}$$

$$\nabla \times \mathbf{F} = \left(\frac{\partial F_z}{\partial y} - \frac{\partial F_y}{\partial z} \right) \hat{x} + \left(\frac{\partial F_x}{\partial z} - \frac{\partial F_z}{\partial x} \right) \hat{y} + \left(\frac{\partial F_y}{\partial x} - \frac{\partial F_x}{\partial y} \right) \hat{z}$$

- Curl of \mathbf{F} is a vector, projection of the curl of \mathbf{F} at point p onto axis $\hat{u} = \hat{x}, \hat{y}, \hat{z}$ is a closed line integral in a plane orthogonal to \hat{u} divided by the area enclosed when $A \rightarrow 0$, \mathbf{A} oriented via the right-hand rule.

$$(\nabla \times \mathbf{F})(p) \cdot \hat{u} \stackrel{\text{def}}{=} \lim_{A \rightarrow 0} \frac{1}{|A|} \oint_C \mathbf{F} \cdot d\mathbf{r}$$

Proof see A.5



2nd Derivatives - Laplacian

Laplace operator, Laplacian: $\nabla^2 = \nabla \cdot \nabla$

- Laplacian of scalar function $f = \text{div of } \nabla f$

$$\nabla^2 f = \nabla \cdot (\nabla f) = \partial^2 f / \partial x^2 + \partial^2 f / \partial y^2 + \partial^2 f / \partial z^2$$

Laplace operator: $\Delta f = \nabla^2 f$

- Laplacian of vector function $\vec{v} = v_x \hat{x} + v_y \hat{y} + v_z \hat{z}$:

$$\nabla^2 \vec{v} = \nabla^2 v_x \hat{x} + \nabla^2 v_y \hat{y} + \nabla^2 v_z \hat{z}$$

Other 2nd Derivatives

- Curl of Grad

$$\nabla \times (\nabla f) = 0$$

- Div of curl

$$\nabla \cdot (\nabla \times \mathbf{A}) = 0$$

- Curl of curl

$$\nabla \times (\nabla \times \mathbf{A}) = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$$

Vector Identities

Triple Products

$$(1) \mathbf{A} \cdot (\mathbf{B} \times \mathbf{C}) = \mathbf{B} \cdot (\mathbf{C} \times \mathbf{A}) = \mathbf{C} \cdot (\mathbf{A} \times \mathbf{B})$$

$$(2) \mathbf{A} \times (\mathbf{B} \times \mathbf{C}) = \mathbf{B}(\mathbf{A} \cdot \mathbf{C}) - \mathbf{C}(\mathbf{A} \cdot \mathbf{B})$$

Product Rules

$$(3) \nabla(fg) = f(\nabla g) + g(\nabla f)$$

$$(4) \nabla(\mathbf{A} \cdot \mathbf{B}) = \mathbf{A} \times (\nabla \times \mathbf{B}) + \mathbf{B} \times (\nabla \times \mathbf{A}) + (\mathbf{A} \cdot \nabla)\mathbf{B} + (\mathbf{B} \cdot \nabla)\mathbf{A}$$

$$(5) \nabla \cdot (f\mathbf{A}) = f(\nabla \cdot \mathbf{A}) + \mathbf{A} \cdot (\nabla f)$$

$$(6) \nabla \cdot (\mathbf{A} \times \mathbf{B}) = \mathbf{B} \cdot (\nabla \times \mathbf{A}) - \mathbf{A} \cdot (\nabla \times \mathbf{B})$$

$$(7) \nabla \times (f\mathbf{A}) = f(\nabla \times \mathbf{A}) - \mathbf{A} \times (\nabla f)$$

$$(8) \nabla \times (\mathbf{A} \times \mathbf{B}) = (\mathbf{B} \cdot \nabla)\mathbf{A} - (\mathbf{A} \cdot \nabla)\mathbf{B} + \mathbf{A}(\nabla \cdot \mathbf{B}) - \mathbf{B}(\nabla \cdot \mathbf{A})$$

Second Derivatives

$$(9) \nabla \cdot (\nabla \times \mathbf{A}) = 0$$

$$(10) \nabla \times (\nabla f) = 0$$

$$(11) \nabla \times (\nabla \times \mathbf{A}) = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$$

- Prove by writing out in cartesian coordinates
- Provide in midterm and final :)

Integral Calculus

- Gradient Theorem

$$\int_a^b (\nabla f) \cdot d\mathbf{l} = f(b) - f(a)$$

- Fundamental Theorem for Divergence (Gauss's Thm)

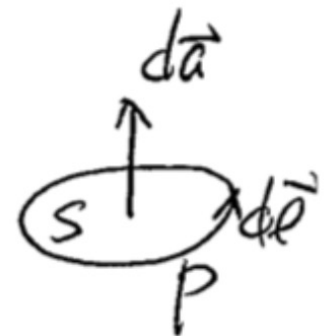
$$\int_{Vol} (\nabla \cdot \mathbf{F}) d\tau = \int_{Sur} \mathbf{F} \cdot d\mathbf{a}$$

(over closed surface)



- Fundamental Theorem for Curl (Stoke's Thm)

$$\int_{Sur} (\nabla \times \mathbf{F}) \cdot d\mathbf{a} = \int_{loop} \mathbf{F} \cdot d\mathbf{l}$$



Proof see Appendix A

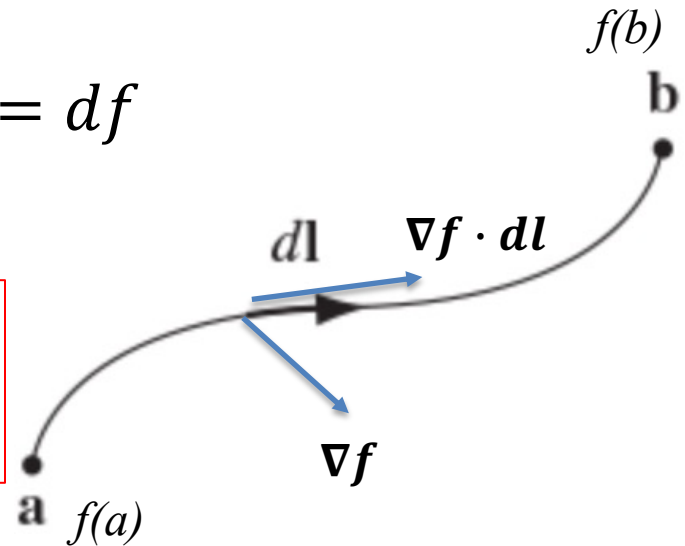
Gradient Theorem

$$d\mathbf{l} = dx\hat{\mathbf{x}} + dy\hat{\mathbf{y}} + dz\hat{\mathbf{z}}$$

$$\nabla f = \frac{\partial f}{\partial x}\hat{\mathbf{x}} + \frac{\partial f}{\partial y}\hat{\mathbf{y}} + \frac{\partial f}{\partial z}\hat{\mathbf{z}}$$

$$\nabla f \cdot d\mathbf{l} = \frac{\partial f}{\partial x}dx + \frac{\partial f}{\partial y}dy + \frac{\partial f}{\partial z}dz = df$$

$$\int_a^b \nabla f \cdot d\mathbf{l} = \int_a^b df = f(b) - f(a)$$



Divergence Theorem (Gauss's)

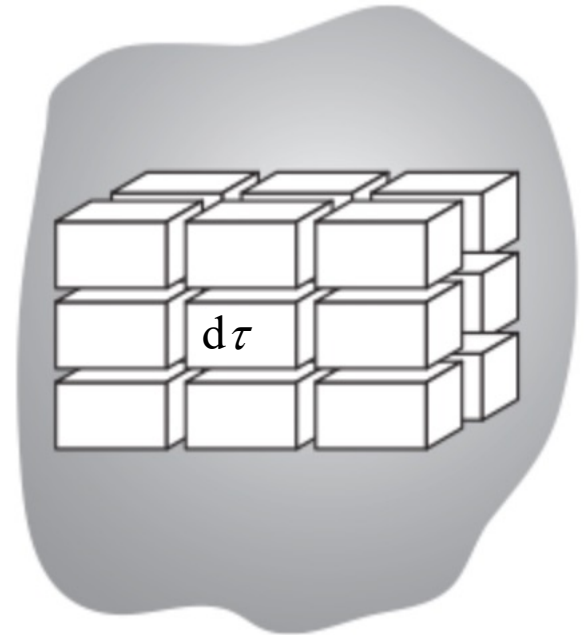
$$\nabla \cdot \mathbf{F} = \lim_{V \rightarrow 0} \frac{\Phi(S)}{V}$$



for each $d\tau$

$$(\nabla \cdot \mathbf{F})d\tau = \oint \mathbf{F} \cdot d\mathbf{a}$$

Internal contributions cancel in pairs for each internal boundary of two adjacent $d\tau$, since $d\mathbf{a}$ always opposite



$$\int_{Vol} (\nabla \cdot \mathbf{F}) d\tau = \int_{Sur} \mathbf{F} \cdot d\mathbf{a}$$

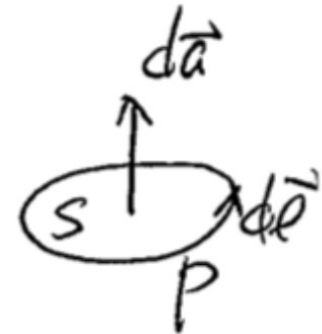
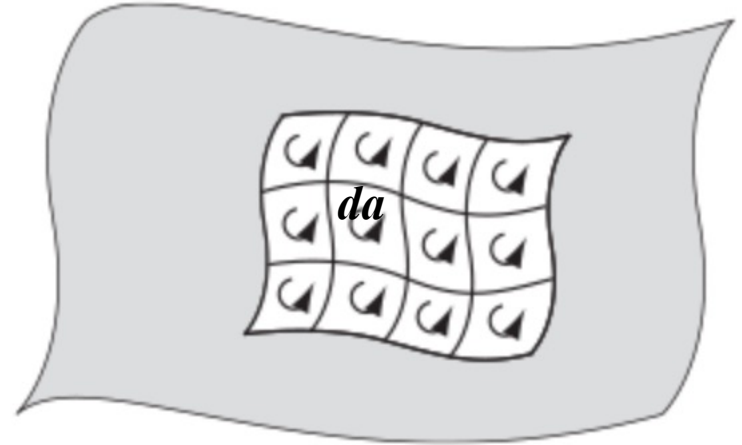
Curl Theorem (Stoke's)

$$(\nabla \times \mathbf{F})(p) \cdot \hat{\mathbf{u}} \stackrel{\text{def}}{=} \lim_{A \rightarrow 0} \frac{1}{|A|} \oint_C \mathbf{F} \cdot d\mathbf{r}$$

for each $d\mathbf{a}$

$$(\nabla \times \mathbf{F}) \cdot d\mathbf{a} = \oint \mathbf{F} \cdot d\mathbf{l}$$

Internal contributions cancel in pairs, because every internal line is the edge of two adjacent loops running in opposite directions



$$\int_{Sur} (\nabla \times \mathbf{F}) \cdot d\mathbf{a} = \int_{loop} \mathbf{F} \cdot d\mathbf{l}$$

Electricity and Magnetism

- Branch of physics that studies how charges and currents interact and the resulting electromagnetic phenomena
- In P112, Solve for Maxwell's equations:

Integral form (7D)

$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{\text{encl}}}{\epsilon_0} \quad (\text{Gauss's law for } \vec{E})$$

Area vector

$$\oint \vec{B} \cdot d\vec{A} = 0 \quad (\text{Gauss's law for } \vec{B})$$

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 \left(i_C + \epsilon_0 \frac{d\Phi_E}{dt} \right)_{\text{encl}} \quad (\text{Ampere's law})$$

$$\oint \vec{E} \cdot d\vec{l} = -\frac{d\Phi_B}{dt} \quad (\text{Faraday's law})$$

Differential form

$$\nabla \cdot \vec{E} = \rho / \epsilon_0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{d\vec{E}}{dt}$$

Displacement current

$$\nabla \times \vec{E} = -\frac{d\vec{B}}{dt}$$

E field from charges is radial:

$$\nabla \times \vec{E} = 0$$

Forces on charge:

$$\vec{F} = Q(\vec{E} + \vec{v} \times \vec{B})$$

Current Density:

$$\nabla \cdot \vec{J} = -d\rho/dt$$

Convert Maxwell's Equations from Integral Form into Differential Form

Gauss' Law

$$\oint_S \mathbf{E} \cdot d\mathbf{S} = \frac{Q}{\epsilon_0}$$

$$\oint_S \mathbf{E} \cdot d\mathbf{S} = \int_V \frac{\rho}{\epsilon_0} dV$$

Invoke divergence theorem:

$$\int_V (\nabla \cdot \mathbf{E}) dV = \int_V \frac{\rho}{\epsilon_0} dV$$

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

Gauss' Law for Magnetism

$$\oint_S \mathbf{B} \cdot d\mathbf{S} = 0$$

divergence theorem:

$$\int_V (\nabla \cdot \mathbf{B}) dV = 0$$

$$\nabla \cdot \mathbf{B} = 0$$

Convert Maxwell's Equations from Integral Form into Differential Form

Faraday's Law

$$\oint_C \mathbf{E} \cdot d\mathbf{l} = -\frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{S}$$

Stokes' theorem:

$$\int_S (\nabla \times \mathbf{E}) \cdot d\mathbf{S} = -\frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{S}$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

Ampere's Law

$$\oint_C \mathbf{B} \cdot d\mathbf{l} = \mu_0 \int_S \mathbf{J} \cdot d\mathbf{S} + \mu_0 \epsilon_0 \frac{d}{dt} \int_S \mathbf{E} \cdot d\mathbf{S}$$

Stokes' theorem:

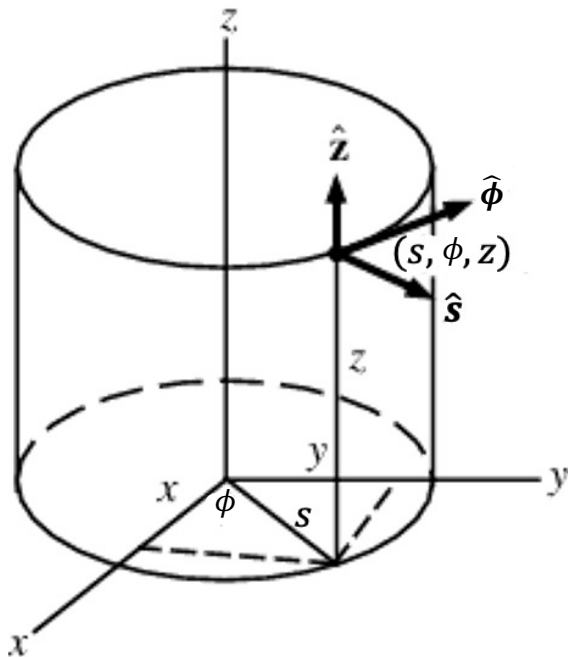
$$\int_S (\nabla \times \mathbf{B}) \cdot d\mathbf{S} = \mu_0 \int_S \mathbf{J} \cdot d\mathbf{S} + \mu_0 \epsilon_0 \frac{d}{dt} \int_S \mathbf{E} \cdot d\mathbf{S}$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

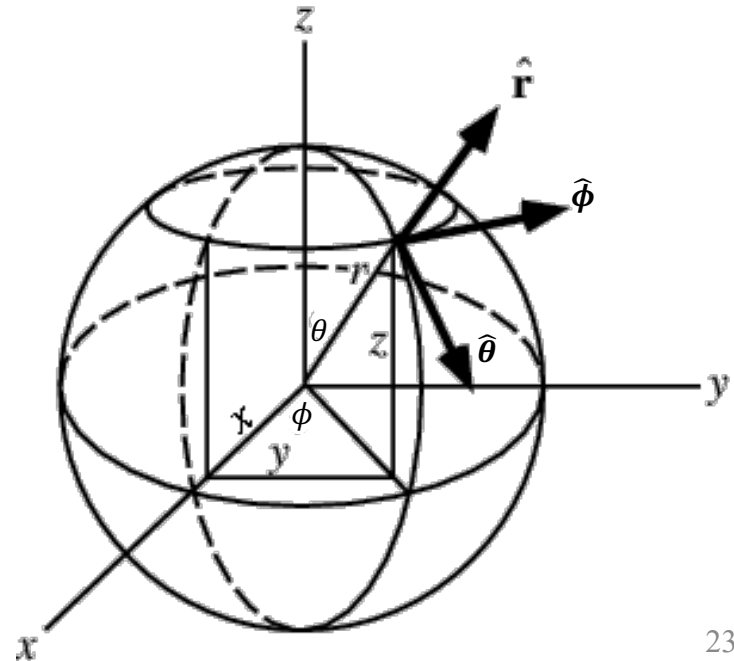
Curvilinear Coordinates

- Vector identities were developed irrespective of coordinate system
- Three coordinate systems are commonly used: Cartesian, Cylindrical, and Spherical (in 2D, polar system, which is special case of cylindrical)

Cylindrical coordinate



Spherical coordinate



Curvilinear – Cartesian Transformation

Spherical

$$\begin{cases} x = r \sin \theta \cos \phi \\ y = r \sin \theta \sin \phi \\ z = r \cos \theta \end{cases}$$

$$\begin{cases} \hat{\mathbf{x}} = \sin \theta \cos \phi \hat{\mathbf{r}} + \cos \theta \cos \phi \hat{\boldsymbol{\theta}} - \sin \phi \hat{\boldsymbol{\phi}} \\ \hat{\mathbf{y}} = \sin \theta \sin \phi \hat{\mathbf{r}} + \cos \theta \sin \phi \hat{\boldsymbol{\theta}} + \cos \phi \hat{\boldsymbol{\phi}} \\ \hat{\mathbf{z}} = \cos \theta \hat{\mathbf{r}} - \sin \theta \hat{\boldsymbol{\theta}} \end{cases}$$

$$\begin{cases} r = \sqrt{x^2 + y^2 + z^2} \\ \theta = \tan^{-1}(\sqrt{x^2 + y^2}/z) \\ \phi = \tan^{-1}(y/x) \end{cases}$$

$$\begin{cases} \hat{\mathbf{r}} = \sin \theta \cos \phi \hat{\mathbf{x}} + \sin \theta \sin \phi \hat{\mathbf{y}} + \cos \theta \hat{\mathbf{z}} \\ \hat{\boldsymbol{\theta}} = \cos \theta \cos \phi \hat{\mathbf{x}} + \cos \theta \sin \phi \hat{\mathbf{y}} - \sin \theta \hat{\mathbf{z}} \\ \hat{\boldsymbol{\phi}} = -\sin \phi \hat{\mathbf{x}} + \cos \phi \hat{\mathbf{y}} \end{cases}$$

Cylindrical

$$\begin{cases} x = s \cos \phi \\ y = s \sin \phi \\ z = z \end{cases}$$

$$\begin{cases} \hat{\mathbf{x}} = \cos \phi \hat{\mathbf{s}} - \sin \phi \hat{\boldsymbol{\phi}} \\ \hat{\mathbf{y}} = \sin \phi \hat{\mathbf{s}} + \cos \phi \hat{\boldsymbol{\phi}} \\ \hat{\mathbf{z}} = \hat{\mathbf{z}} \end{cases}$$

$$\begin{cases} s = \sqrt{x^2 + y^2} \\ \phi = \tan^{-1}(y/x) \\ z = z \end{cases}$$

$$\begin{cases} \hat{\mathbf{s}} = \cos \phi \hat{\mathbf{x}} + \sin \phi \hat{\mathbf{y}} \\ \hat{\boldsymbol{\phi}} = -\sin \phi \hat{\mathbf{x}} + \cos \phi \hat{\mathbf{y}} \\ \hat{\mathbf{z}} = \hat{\mathbf{z}} \end{cases}$$

Curvilinear Expressions of Vector Derivatives

$$\text{Cartesian } d\mathbf{l} = dx \hat{\mathbf{x}} + dy \hat{\mathbf{y}} + dz \hat{\mathbf{z}}, d\tau = dx dy dz$$

$$\text{Gradient: } \nabla t = \frac{\partial t}{\partial x} \hat{\mathbf{x}} + \frac{\partial t}{\partial y} \hat{\mathbf{y}} + \frac{\partial t}{\partial z} \hat{\mathbf{z}}$$

$$\text{Divergence: } \nabla \cdot \mathbf{v} = \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z}$$

$$\text{Curl: } \nabla \times \mathbf{v} = \left(\frac{\partial v_z}{\partial y} - \frac{\partial v_y}{\partial z} \right) \hat{\mathbf{x}} + \left(\frac{\partial v_x}{\partial z} - \frac{\partial v_z}{\partial x} \right) \hat{\mathbf{y}} + \left(\frac{\partial v_y}{\partial x} - \frac{\partial v_x}{\partial y} \right) \hat{\mathbf{z}}$$

$$\text{Laplacian: } \nabla^2 t = \frac{\partial^2 t}{\partial x^2} + \frac{\partial^2 t}{\partial y^2} + \frac{\partial^2 t}{\partial z^2}$$

$$\text{Spherical } d\mathbf{l} = dr \hat{\mathbf{r}} + r d\theta \hat{\boldsymbol{\theta}} + r \sin \theta d\phi \hat{\boldsymbol{\phi}}, d\tau = r^2 \sin \theta dr d\theta d\phi$$

$$\text{Gradient: } \nabla t = \frac{\partial t}{\partial r} \hat{\mathbf{r}} + \frac{1}{r} \frac{\partial t}{\partial \theta} \hat{\boldsymbol{\theta}} + \frac{1}{r \sin \theta} \frac{\partial t}{\partial \phi} \hat{\boldsymbol{\phi}}$$

$$\text{Divergence: } \nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta v_\theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi}$$

$$\text{Curl: } \nabla \times \mathbf{v} = \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (\sin \theta v_\phi) - \frac{\partial v_\theta}{\partial \phi} \right] \hat{\mathbf{r}} + \frac{1}{r} \left[\frac{1}{\sin \theta} \frac{\partial v_r}{\partial \phi} - \frac{\partial}{\partial r} (r v_\phi) \right] \hat{\boldsymbol{\theta}} + \frac{1}{r} \left[\frac{\partial}{\partial r} (r v_\theta) - \frac{\partial v_r}{\partial \theta} \right] \hat{\boldsymbol{\phi}}$$

$$\text{Laplacian: } \nabla^2 t = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial t}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial t}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 t}{\partial \phi^2}$$

$$\text{Cylindrical } d\mathbf{l} = ds \hat{\mathbf{s}} + s d\phi \hat{\boldsymbol{\phi}} + dz \hat{\mathbf{z}}, d\tau = s ds d\phi dz$$

$$\text{Gradient: } \nabla t = \frac{\partial t}{\partial s} \hat{\mathbf{s}} + \frac{1}{s} \frac{\partial t}{\partial \phi} \hat{\boldsymbol{\phi}} + \frac{\partial t}{\partial z} \hat{\mathbf{z}}$$

$$\text{Divergence: } \nabla \cdot \mathbf{v} = \frac{1}{s} \frac{\partial}{\partial s} (s v_s) + \frac{1}{s} \frac{\partial v_\phi}{\partial \phi} + \frac{\partial v_z}{\partial z}$$

$$\text{Curl: } \nabla \times \mathbf{v} = \left[\frac{1}{s} \frac{\partial v_z}{\partial \phi} - \frac{\partial v_\phi}{\partial z} \right] \hat{\mathbf{s}} + \left[\frac{\partial v_s}{\partial z} - \frac{\partial v_z}{\partial s} \right] \hat{\boldsymbol{\phi}} + \frac{1}{s} \left[\frac{\partial}{\partial s} (s v_\phi) - \frac{\partial v_s}{\partial \phi} \right] \hat{\mathbf{z}}$$

$$\text{Laplacian: } \nabla^2 t = \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial t}{\partial s} \right) + \frac{1}{s^2} \frac{\partial^2 t}{\partial \phi^2} + \frac{\partial^2 t}{\partial z^2}$$

Proof see Appendix A

Dirac Delta Function

- Want to describe point charge Q as special case of an expression for charge density:

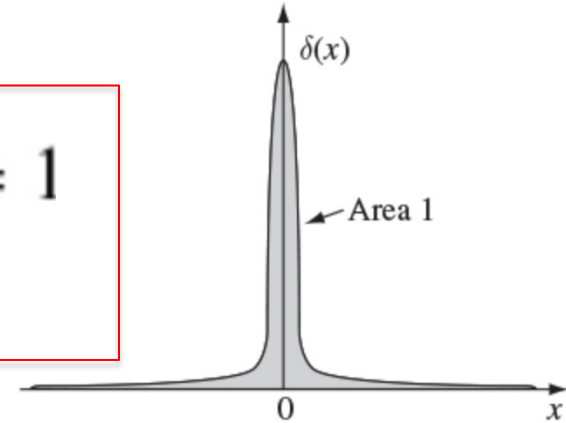
$$\rho(x) = Q\delta(x)$$

$$\text{where: } \delta(x) = 0 \text{ if } x \neq 0$$
$$= \infty \text{ if } x = 0$$

- Also need $\int \delta(x) dx = 1$
 - for any limits that include $x=0$, so total charge is Q

1D Dirac Delta Function

$$\delta(x) = \begin{cases} 0 & \text{if } x \neq 0, \\ \infty & \text{if } x = 0, \end{cases} \quad \int_{-\infty}^{\infty} \delta(x) dx = 1$$



If $f(x)$ is ordinary (continuous) function:

$$f(x)\delta(x) = f(0)\delta(x)$$

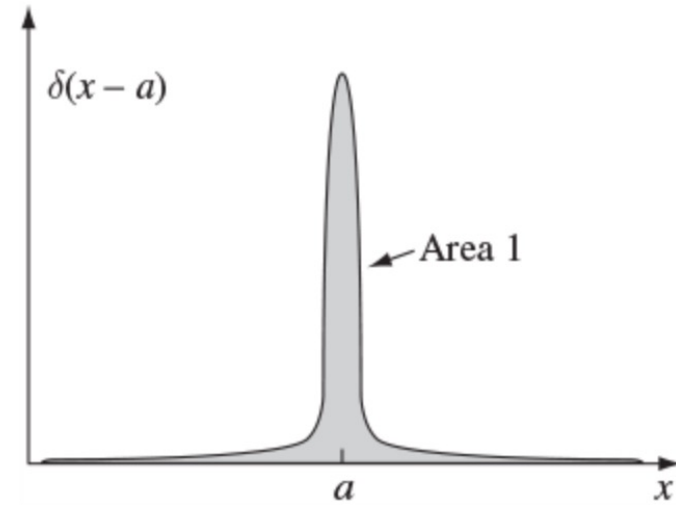
$$\int_{-\infty}^{\infty} f(x)\delta(x) dx = f(0) \int_{-\infty}^{\infty} \delta(x) dx = f(0)$$

1D Dirac Delta Function

Can shift the spike from $x=0$ to $x=a$

$$\delta(x-a) = \begin{cases} 0 & \text{if } x \neq a, \\ \infty & \text{if } x = a, \end{cases} \quad \text{with } \int_{-\infty}^{\infty} \delta(x-a) dx = 1$$

Typically, the limits of integration are from $x = -\infty$ to $x = +\infty$, but will give 1 for any limits that include $x=a$.



If $f(x)$ is ordinary (continuous) function:

$$f(x)\delta(x-a) = f(a)\delta(x-a) \quad \int_{-\infty}^{\infty} f(x)\delta(x-a) dx = f(a)$$

For delta functions $D_1(x)$ and $D_2(x)$

$$\text{if } \int_{-\infty}^{\infty} f(x)D_1(x) dx = \int_{-\infty}^{\infty} f(x)D_2(x) dx \quad \text{then } D_1(x)=D_2(x)$$

3D Dirac Delta Function

Generalize the delta function to three dimensions:

$$\delta^3(\mathbf{r}) = \delta(x) \delta(y) \delta(z)$$

$$\int_{\text{all space}} \delta^3(\mathbf{r}) d\tau = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \delta(x) \delta(y) \delta(z) dx dy dz = 1$$

If $f(\mathbf{r})$ is ordinary (continuous) function:

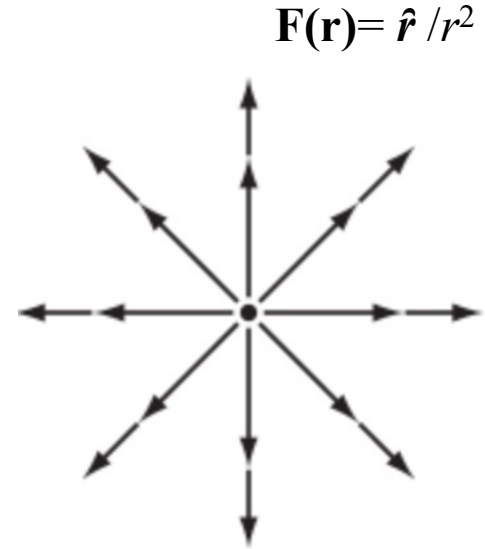
$$\int_{\text{all space}} f(\mathbf{r}) \delta^3(\mathbf{r} - \mathbf{a}) d\tau = f(\mathbf{a})$$

Dirac delta function: $\nabla \cdot \frac{\hat{\mathbf{r}}}{r^2}$

$$\nabla \cdot \frac{\hat{\mathbf{r}}}{r^2} = 4\pi\delta^3(\mathbf{r})$$

$$r \neq 0 \quad \nabla \cdot \frac{\hat{\mathbf{r}}}{r^2} = \frac{1}{r^2} \frac{\partial}{\partial r} \left[r^2 \left(\frac{1}{r^2} \right) \right] = 0$$

Spherical Divergence: $\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta v_\theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi}$



Consider a sphere around $r=0$

$$r=0 \quad \nabla \cdot \frac{\hat{\mathbf{r}}}{r^2} = \lim_{\text{volume} \rightarrow 0} \frac{\text{Total Flux}}{\text{Volume}} = \lim_{r \rightarrow 0} \frac{4\pi r^2 \left(\frac{1}{r^2} \right)}{\frac{4}{3}\pi r^3} = \lim_{r \rightarrow 0} \frac{4\pi}{\frac{4}{3}\pi r^3} = +\infty$$

$$\iiint \nabla \cdot \frac{\hat{\mathbf{r}}}{r^2} d^3r = \text{Total flux} = \oiint \frac{\hat{\mathbf{r}}}{r^2} \cdot d\mathbf{a} = 4\pi$$

Dirac delta function: $\nabla^2 \frac{1}{r}$

$$\nabla^2 \frac{1}{r} = -4\pi\delta^3(\mathbf{r})$$

$$\nabla(1/r) = \hat{\mathbf{r}}[\partial/\partial r](1/r) = \hat{\mathbf{r}}(-1/r^2) = -\hat{\mathbf{r}}/r^2$$

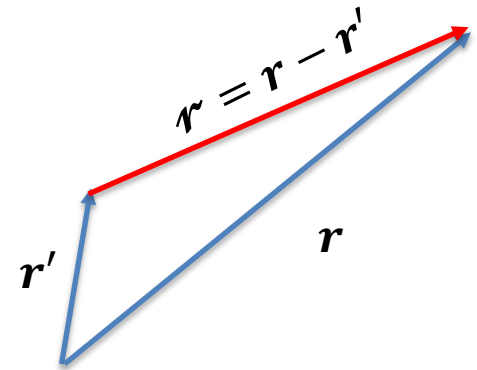
Spherical Gradient: $\nabla = \frac{\partial}{\partial r}\hat{\mathbf{r}} + \frac{1}{r}\frac{\partial}{\partial\theta}\hat{\boldsymbol{\theta}} + \frac{1}{r\sin\theta}\frac{\partial}{\partial\phi}\hat{\boldsymbol{\phi}}$

$$\nabla^2(1/r) = \nabla \cdot \nabla(1/r) = \nabla \cdot (-\hat{\mathbf{r}}/r^2) = -4\pi\delta^3(\mathbf{r})$$

Dirac delta function

More general:

Define: $\mathbf{r} = \mathbf{r} - \mathbf{r}'$



$$\nabla \cdot \frac{\hat{\mathbf{r}}}{r^2} = 4\pi\delta^3(\mathbf{r})$$

$$\nabla^2 \frac{1}{r} = -4\pi\delta^3(\mathbf{r})$$

Derivative of Dirac Delta Function

$$\delta[\delta(x)]/dx = -\delta(x)/x$$

Integration by parts: $\int u dv = uv - \int v du$ $u = xf(x), v = \delta(x)$

$$\begin{aligned} \int xf(x)[d\delta(x)/dx]dx &= \int xf(x)d\delta(x) = xf(x)\delta(x) - \int \delta(x)d[xf(x)] \\ &= xf(x)\delta(x) - \int \delta(x)(f(x)dx + xdf(x)) \\ &= xf(x)\delta(x) - \int \delta(x)f(x)dx - \int \delta(x)xf(x)'dx \\ &= \textcolor{red}{x*f(x)\delta(x)} - f(0) - 0*f(0)' = -f(0) = -\int f(x)\delta(x)dx \end{aligned}$$

Red term is zero evaluated at $x = -\infty$ to $x = +\infty$

$$xf(x)[d\delta(x)/dx]dx = -f(x)\delta(x)$$

$$xd\delta(x)/dx = -\delta(x)$$

$$d\delta(x)/dx = -\delta(x)/x$$

Example 1.15.

Show that

$k > 0$:

$$\delta(kx) = \frac{1}{|k|} \delta(x),$$

where k is any (nonzero) constant. (In particular, $\delta(-x) = \delta(x)$: it's an *even* function.)

set $y=kx$ Integration by substitution

$$k > 0: \int f(x) \delta(kx) dx = \int_{y=-\infty}^{+\infty} f(y/k) \delta(y) (dy/k) = \int f(y/k) \delta(y) (dy/k) = f(0)/k$$

$$k < 0: \int f(x) \delta(kx) dx = \int_{y=+\infty}^{-\infty} f(y/k) \delta(y) (dy/k) = - \int f(y/k) \delta(y) (dy/k) = -f(0)/k$$

$$\int f(x) \delta(kx) dx = f(0) / |k|$$

For delta functions $D_1(x)$ and $D_2(x)$

$$\text{if } \int_{-\infty}^{\infty} f(x) D_1(x) dx = \int_{-\infty}^{\infty} f(x) D_2(x) dx \quad \text{then } D_1(x) = D_2(x)$$

$$\delta(kx) = \frac{1}{|k|} \delta(x)$$

Helmholtz Theorem of Vector Fields

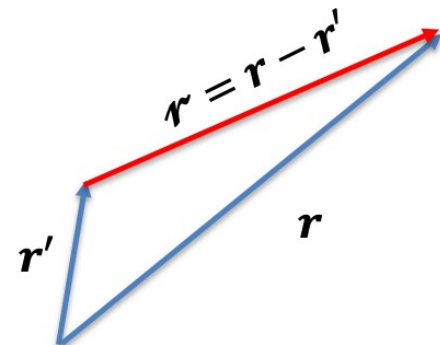
For a vector field \mathbf{F} (\mathbf{E} or \mathbf{B}), if you are given $\text{Div } \mathbf{F}$, $\text{Curl } \mathbf{F}$ and boundary condition

- $\nabla \cdot \mathbf{F} = D$
- $\nabla \times \mathbf{F} = \mathbf{C}$ ($\nabla \cdot \mathbf{C} = 0$ because $\nabla \cdot (\nabla \times \mathbf{F}) = 0$)
- $\mathbf{F} = 0$ at $r = \infty$ or \mathbf{F} values defined on boundary (boundary condition)

Then vector field \mathbf{F} is uniquely determined:


$$\mathbf{F}(\mathbf{r}) = \nabla \left(\frac{-1}{4\pi} \int \frac{\nabla' \cdot \mathbf{F}(\mathbf{r}')}{r} d\tau' \right) + \nabla \times \left(\frac{1}{4\pi} \int \frac{\nabla' \times \mathbf{F}(\mathbf{r}')}{r} d\tau' \right). \quad (\text{B.10})$$

\swarrow V , scalar potential
 \nwarrow \mathbf{A} , vector potential



Scalar Potential

- If $\nabla \times \mathbf{F} = \mathbf{0}$ (ie, curl-less field), then $\mathbf{F} = -\nabla V$, where V is called the scalar potential



(10) $\nabla \times (\nabla f) = \mathbf{0}$

- The following are equivalent conditions (divergence theorem)
 1. $\nabla \times \mathbf{F} = \mathbf{0}$, everywhere in space
 2. $\int \mathbf{F} \cdot d\mathbf{l}$ is independent of path
 3. $\int \mathbf{F} \cdot d\mathbf{l} = 0$ for a closed loop
- When boundary condition, \mathbf{F} is unique but V is not unique since a constant can be added

For example, in electrostatics, $\nabla \cdot \mathbf{E} = \rho/\epsilon_0$ and $\nabla \times \mathbf{E} = \mathbf{0}$, so

$$\mathbf{E}(\mathbf{r}) = -\nabla \left(\frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{r}')}{r} d\tau' \right) = -\nabla V$$

Vector Potential

- If $\nabla \cdot \mathbf{F} = 0$ (ie, divergence-less everywhere), then $\mathbf{F} = \nabla \times \mathbf{A}$, where \mathbf{A} is a vector potential

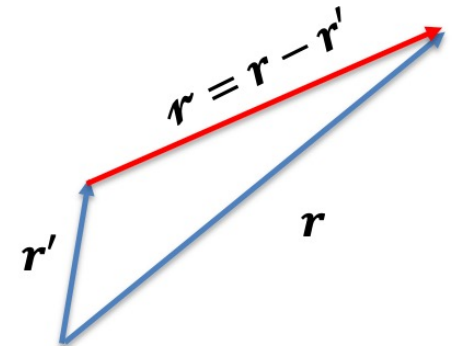
$$(9) \nabla \cdot (\nabla \times \mathbf{A}) = 0$$

- The following are equivalent conditions (curl theorem)

1. $\nabla \cdot \mathbf{F} = 0$
 2. $\int \mathbf{F} \cdot d\mathbf{a}$ is independent of surface shape, for any given boundary line on the surface
 3. $\oint \mathbf{F} \cdot d\mathbf{a} = 0$ for any closed surface
- \mathbf{A} is not unique. Gradient of any scalar function f can be added to \mathbf{A} because $\nabla \times (\nabla f) = 0$

in magnetostatics $\nabla \cdot \mathbf{B} = 0$ and $\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$, so

$$\mathbf{B}(\mathbf{r}) = \nabla \times \left(\frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}')}{r} d\tau' \right) = \nabla \times \mathbf{A}$$

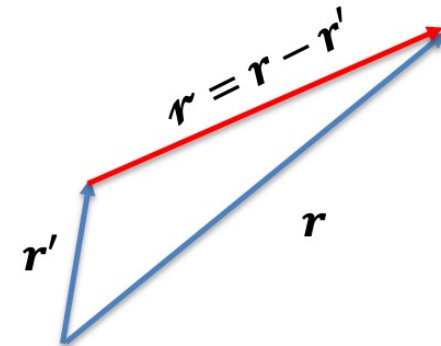


Helmholtz Decomposition

Any differentiable vector function \mathbf{F} that goes to zero faster than $1/r$ as $r \rightarrow \infty$ can be expressed as (gradient of a scalar) + (curl of vector):

$$\mathbf{F} = -\nabla V + \nabla \times \mathbf{A} \text{ (always true)}$$

$$\mathbf{F}(\mathbf{r}) = \nabla \left(\frac{-1}{4\pi} \int \frac{\nabla' \cdot \mathbf{F}(\mathbf{r}')}{r} d\tau' \right) + \nabla \times \left(\frac{1}{4\pi} \int \frac{\nabla' \times \mathbf{F}(\mathbf{r}')}{r} d\tau' \right). \quad (\text{B.10})$$



For example, in electrostatics, $\nabla \cdot \mathbf{E} = \rho/\epsilon_0$ and $\nabla \times \mathbf{E} = \mathbf{0}$, so

$$\mathbf{E}(\mathbf{r}) = -\nabla \left(\frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{r}')}{r} d\tau' \right) = -\nabla V \quad (\text{B.11})$$

(where V is the scalar potential), while in magnetostatics $\nabla \cdot \mathbf{B} = 0$ and $\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$, so

$$\mathbf{B}(\mathbf{r}) = \nabla \times \left(\frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}')}{r} d\tau' \right) = \nabla \times \mathbf{A} \quad (\text{B.12})$$

(where \mathbf{A} is the vector potential).