

# Lecture 15: Slater Determinants

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# Hartree Product

- Configuration State Functions (CSF): (Simple) symmetry adapted basis state for  $N$ -electron systems
- Early attempt: Hartree product

$$\Phi_H(x_1, \dots, x_N) = \phi_1(x_1) \cdots \phi_N(x_N),$$

where the orbitals  $\{\phi_p\}$  are orthonormal

- Resulting  $N$  electron probability amplitude is simple product of (statistically independent) single-electron orbital probabilities,

$$|\Phi_H(x_1, \dots, x_N)|^2 = |\phi_1(x_1)|^2 \cdots |\phi_N(x_N)|^2.$$

- Unsuitable because electrons are indistinguishable, violates spin-statistics theorem

# Determinants

$f : \mathbf{A} \in \mathbf{C}^{n \times n} \rightarrow f(\mathbf{A}) \in \mathbf{C}$  is called *determinant* of  $\mathbf{A} = (\mathbf{a}_1 \cdots \mathbf{a}_n)$ ,  
i.e.,  $f(\mathbf{A}) = \det(\mathbf{A}) = |\mathbf{A}|$ , if and only if

(i)  $f$  is linear in each column, i.e.,

$$f(\cdots \mathbf{c}\mathbf{a}_k + \mathbf{b} \cdots) = cf(\cdots \mathbf{a}_k \cdots) + f(\cdots \mathbf{b} \cdots)$$

$$\forall k, 1 \leq k \leq n, c \in \mathbf{C}, \mathbf{b} \in \mathbf{C}^n.$$

(ii)  $f$  is antisymmetric under column exchange, i.e.,

$$f(\cdots \mathbf{a}_k \cdots \mathbf{a}_l \cdots) = -f(\cdots \mathbf{a}_l \cdots \mathbf{a}_k \cdots)$$

(iii)  $f(\mathbf{1}) = 1$

# Leibniz Formula

- Alternative definition using  $S(N)$ :

$$f(\mathbf{A}) = n! \hat{\mathcal{A}} a_{11} \cdots a_{nn},$$

- Projection operator on the totally antisymmetric IRREP of  $S(N)$  (“antisymmetrizer”):

$$\hat{\mathcal{A}} = \frac{1}{n!} \sum_{P \in S_n} \text{sgn}(P) \hat{P}$$

- $\hat{P}$ : permutations of  $n$  columns,  $\text{sgn}(P)$ : Antisymmetric IRREP characters,  $+1$  for even and  $-1$  for odd permutations

# Slater Determinant

- Antisymmetrized (and renormalized) Hartree product:

$$\Phi(x_1, \dots, x_N) = \sqrt{N!} \hat{\mathcal{A}} \Phi_H = \frac{1}{\sqrt{N!}} \begin{vmatrix} \phi_1(x_1) & \phi_1(x_2) & \dots & \phi_1(x_N) \\ \phi_2(x_1) & \phi_2(x_2) & \dots & \phi_2(x_N) \\ \dots & \dots & \vdots & \dots \\ \phi_N(x_1) & \phi_N(x_2) & \dots & \phi_N(x_N) \end{vmatrix}$$

- Satisfies antisymmetry requirement
- Can be CSF for closed-shell singlet and high-spin open shell states
- The simplest admissible many-electron wavefunction
- Describes  $N$  “noninteracting” electrons subject to Pauli principle:  
Only “Fermi”, no “Coulomb” correlation

## Example: Two Electrons in Two Orbitals

- Slater determinant:

$$\Phi(x_1, x_2) = \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_1(x_1) & \phi_1(x_2) \\ \phi_2(x_1) & \phi_2(x_2) \end{vmatrix} = \frac{1}{\sqrt{2}} [\phi_1(x_1)\phi_2(x_2) - \phi_1(x_2)\phi_2(x_1)]$$

- Triplet ( $M_s = 1$ ) case: Both orbitals have  $\alpha$  spin, but different spatial parts:

$$\phi_1(x_1) = \phi_1(\mathbf{r}_1)\delta_{\sigma_1\alpha}, \quad \phi_2(x_2) = \phi_2(\mathbf{r}_2)\delta_{\sigma_2\alpha}$$

- Triplet Slater determinant:

$$\Phi_t(x_1, x_2) = \frac{1}{\sqrt{2}} [\phi_1(\mathbf{r}_1)\phi_2(\mathbf{r}_2) - \phi_1(\mathbf{r}_2)\phi_2(\mathbf{r}_1)] \delta_{\sigma_1\alpha}\delta_{\sigma_2\alpha}$$

Separates into antisymmetric spacial and symmetric spin parts

- The two electrons must occupy spatially different orbitals
- Spatial triplet wavefunction has nodes for  $\mathbf{r}_1 = \mathbf{r}_2$  ("Fermi hole")
- This is a CSF with  $\hat{S}^2$  eigenvalue 1

## Example: Two Electrons in Two Orbitals

- Closed-shell singlet ( $M_s = 0$ ) case: Both orbitals have identical spatial parts:

$$\phi_1(x_1) = \phi(\mathbf{r}_1)\delta_{\sigma_1\alpha}, \quad \phi_2(x_2) = \phi(\mathbf{r}_2)\delta_{\sigma_2\beta}$$

- Closed-shell singlet Slater determinant:

$$\Phi_s(x_1, x_2) = \phi(\mathbf{r}_1)\phi(\mathbf{r}_2)\frac{1}{\sqrt{2}} [\delta_{\sigma_1\alpha}\delta_{\sigma_2\beta} - \delta_{\sigma_2\alpha}\delta_{\sigma_1\beta}]$$

- No Fermi hole
- This is a CSF with  $\hat{S}^2$  eigenvalue 0

## Example: Two Electrons in Two Orbitals

- Closed-shell singlet ( $M_s = 0$ ) case: Both orbitals have identical spatial parts:

$$\phi_1(x_1) = \phi(\mathbf{r}_1)\delta_{\sigma_1\alpha}, \quad \phi_2(x_2) = \phi(\mathbf{r}_2)\delta_{\sigma_2\beta}$$

- Closed-shell singlet Slater determinant:

$$\Phi_s(x_1, x_2) = \phi(\mathbf{r}_1)\phi(\mathbf{r}_2)\frac{1}{\sqrt{2}} [\delta_{\sigma_1\alpha}\delta_{\sigma_2\beta} - \delta_{\sigma_2\alpha}\delta_{\sigma_1\beta}]$$

- No Fermi hole
- This is a CSF with  $\hat{S}^2$  eigenvalue 0



## Example: Two Electrons in Two Orbitals

- Open-shell singlet ( $M_s = 0$ ) case: Both orbitals have different spatial parts:

$$\phi_1(x_1) = \phi_1(\mathbf{r}_1)\delta_{\sigma_1\alpha}, \quad \phi_2(x_2) = \phi_2(\mathbf{r}_2)\delta_{\sigma_2\beta}$$

- Open-shell singlet Slater determinant (one of two):

$$\Phi_{\text{oss}}(x_1, x_2) = \frac{1}{\sqrt{2}} [\phi_1(\mathbf{r}_1)\phi_2(\mathbf{r}_2)\delta_{\sigma_1\alpha}\delta_{\sigma_2\beta} - \phi_1(\mathbf{r}_1)\phi_2(\mathbf{r}_2)\delta_{\sigma_2\alpha}\delta_{\sigma_1\beta}]$$

- This is *not* a CSF, because it is not an  $\hat{S}^2$  eigenstate!
- OSS singlet CSF requires linear combination of (at least) 2 determinants:

$$\begin{aligned} \Psi_{\text{oss}}(x_1, x_2) = \frac{1}{\sqrt{2}} [\phi_1(\mathbf{r}_1)\phi_2(\mathbf{r}_2) + \phi_1(\mathbf{r}_1)\phi_2(\mathbf{r}_2)] \\ \times \frac{1}{\sqrt{2}} [\delta_{\sigma_1\alpha}\delta_{\sigma_2\beta} - \delta_{\sigma_2\alpha}\delta_{\sigma_1\beta}] \end{aligned}$$

## Example: Two Electrons in Two Orbitals

- Triplet ( $M_s = 0$ ) case: Both orbitals have different spatial parts:

$$\phi_1(x_1) = \phi_1(\mathbf{r}_1)\delta_{\sigma_1\alpha}, \quad \phi_2(x_2) = \phi_2(\mathbf{r}_2)\delta_{\sigma_2\beta}$$

- Open-shell singlet Slater determinant (one of two):

$$\Phi_{\text{oss}}(x_1, x_2) = \frac{1}{\sqrt{2}} [\phi_1(\mathbf{r}_1)\phi_2(\mathbf{r}_2)\delta_{\sigma_1\alpha}\delta_{\sigma_2\beta} + \phi_1(\mathbf{r}_1)\phi_2(\mathbf{r}_2)\delta_{\sigma_2\alpha}\delta_{\sigma_1\beta}]$$

- This is *not* a CSF, because it is not an  $\hat{S}^2$  eigenstate!
- $M_s = 0$  triplet CSF requires linear combination of (at least) 2 determinants:

$$\begin{aligned} \Psi_{t,0}(x_1, x_2) = & \frac{1}{\sqrt{2}} [\phi_1(\mathbf{r}_1)\phi_2(\mathbf{r}_2) - \phi_1(\mathbf{r}_1)\phi_2(\mathbf{r}_2)] \\ & \times \frac{1}{\sqrt{2}} [\delta_{\sigma_1\alpha}\delta_{\sigma_2\beta} + \delta_{\sigma_2\alpha}\delta_{\sigma_1\beta}] \end{aligned}$$